

## Hybrid model for Soledge2D:

Based on [1]: atom vdf  $f$  split into a kinetic part  $f_K$  and a fluid part  $f_F$ , such that  $f = f_F + f_K$

$$\frac{\partial f_{K,F}}{\partial t} + \vec{v} \cdot \vec{\nabla} f_{K,F} = S_{iz}^{K,F} + S_{rc}^{K,F} + S_{cx}^{K,F} + S_{mol}^{K,F} \pm S_{F \rightarrow K} \pm S_{K \rightarrow F}$$

Transition processes between kinetic and fluid phases introduced – evaporation & condensation

$$S_{F \rightarrow K} = -v_{evap} f_F \quad S_{K \rightarrow F} = -v_{cond} f_K \quad \text{At this point: choice of rates is arbitrary}$$

Aim : condensation  $\gg$  evaporation in highly collisionality regions (HCRs), i.e.  $f \sim f_F \sim \text{Maxw.}$ , evaporation  $\gg$  condensation in LCRs

Next step : collisional closure for moments of  $f_F$  with the idea that  $f_F/f_K \ll 1$  in regions where the later is not valid.

Here two moments assuming  $T_n = T_i$  in the HCRs – capture momentum exchanges

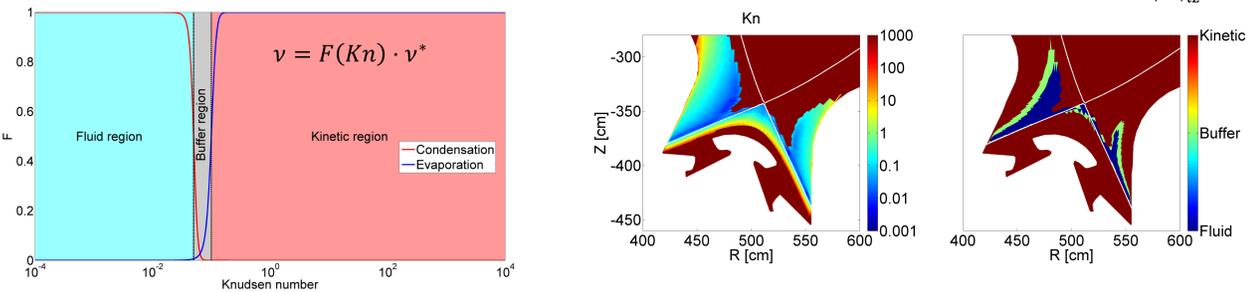
$$\frac{\partial n_F}{\partial t} + \vec{v} \cdot (n_F \vec{u}_F) = S_{iz,n}^F + S_{rc,n}^F + S_{evap,n}^F + S_{cond,n}^F$$

$$\frac{\partial n_F \vec{u}_F}{\partial t} + \vec{v} \cdot (n_F \vec{u}_F \vec{u}_F + \frac{p_F}{m} I) = \vec{S}_{iz,mom}^F + \vec{S}_{rc,mom}^F + \vec{S}_{cx,mom}^F + \vec{S}_{evap,mom}^F + \vec{S}_{cond,mom}^F$$

$f_F$  is assumed purely Maxwellian (appears explicitly in the equation for  $f_K$  +  $f_K = f - f_F$  – positive definiteness of  $f_K$  ?)

Fluid breakdown parameter  $\rightarrow$  local gradient length Knudsen number  $Kn = \frac{\lambda}{L} = \frac{v_{th}/(n_i \langle \sigma v \rangle_{cx})}{\max(|v_{||} n_e, \nabla_{\perp} n_e|)/n_e}$

Kn can be small in regions with no plasma or not dominated by cx  $\rightarrow$  kinetic zone forced where  $n_e = n_i \leq n_{min}$  and/or  $\frac{\langle \sigma v \rangle_{cx}}{\langle \sigma v \rangle_{iz}} \leq min$



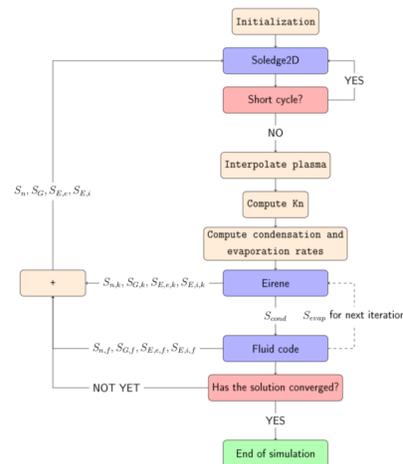
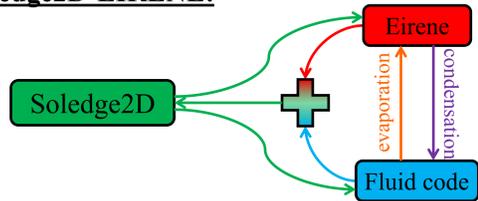
MC slow because of cx  $\rightarrow$  make condensation more probable than cx in fluid areas  $v_{cond}^* = M \cdot n_i \langle \sigma v \rangle_{cx}$  with  $M \gg 1$  (here  $v_{evap}^* = v_{cond}^*$ )

fluid approximation valid for  $Kn \ll 1$ , so the shape of the factor  $F(Kn)$  determines whether mesh region is fluid or kinetic

So essentially this approach amounts to a spatially hybrid model with a dynamic boundary between fluid and kinetic domain, enforced by an immersed boundary technique similar to the penalization scheme implemented in Soledge2D.

boundary conditions discussed in the literature (analysis ongoing). Buffer region added to prevent “fluid/kinetic” oscillations.

## Implementation in Soledge2D-EIRENE:



Specifically developed fluid code based on HDG scheme, on the same grid as EIRENE [4]

Absorbing BC for fluid atoms at the wall, recycled as kinetic atoms/mol.

$$\Gamma_{F \rightarrow K} = R \Gamma_F^- = R \int_{-} f_{Max}(\vec{u}_F, T_F) \vec{v}_F \cdot \hat{n} d\vec{v}_F$$

Molecules treated kinetically (hybrid may be needed too because of elastic collisions !)

Condensation implemented as an extra CX process with background fluid atoms

## Conclusions:

- The method initially proposed by Karney&Stotler can be seen as the implementation of a spatially hybrid scheme. Complementary with macro/micro approach developed by Horsten et al.
- First results encouraging but robustness with model parameters need to be investigated further.
- Could be seen as a convergence accelerator (subsequent fine tuning with kinetic only allowing one to assess modeling errors related to the hybrid model)

## References :

- H. Bufferand et al, Nuclear Fusion 55, 053025 (2015)
- D. Reiter et al, Fusion Science and Techn. 47, 172 (2005)
- C.F.F. Karney et al, Contrib. Plasma Phys. 38/ 1/ 2, 319(1998)
- M. Valentinuzzi et al, Contrib. Plasma Phys.(2018)
- M. Valentinuzzi NME (2019), Phd Thesis

## Motivation :

- Modelling of edge plasma through fluid plasma solver (like SoEdge2D [1]) coupled to kinetic Monte Carlo code for neutrals (like Eirene [2]).
- In most of the simulation domain Knudsen number  $Kn \equiv m \cdot f \cdot p/L$  for neutrals much larger than one, but may not be true near divertor targets in detached regime due to high plasma density ( $10^{20} - 10^{21} m^{-3}$ ) and low temperature (below 5eV)
- Neutrals undergo high number of charge-exchange and elastic collisions before getting ionized, so Monte Carlo codes inefficient/inaccurate
- Hybrid kinetic-fluid model for neutrals could improve performances of edge plasma codes

## Hybrid simulations: comparison with fully kinetic case

### ITER test case (P. Tamain et al., and Refs. [5])

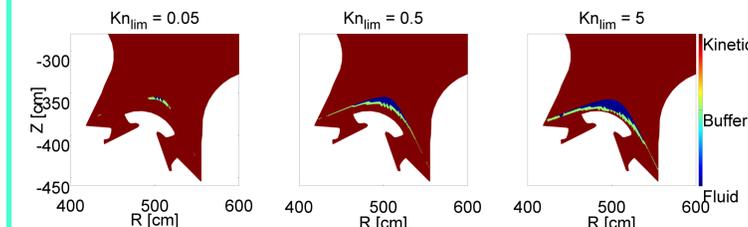
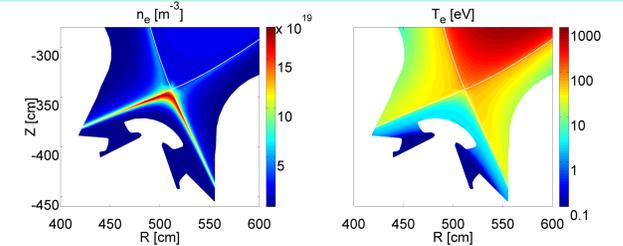
Pure H,  $P_{SOL} = 20MW$ , self-consistent core density BC

$D = 0.3m^2/s$ ,  $\chi_e = \chi_i = 1m^2/s$

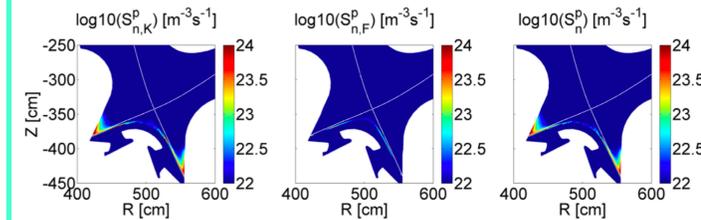
Be wall  $R = 1$

1 pump below dome with  $R = 0.9928$

2 puffs (top) with same rate  $5.65 \cdot 10^{21} e^-/s$



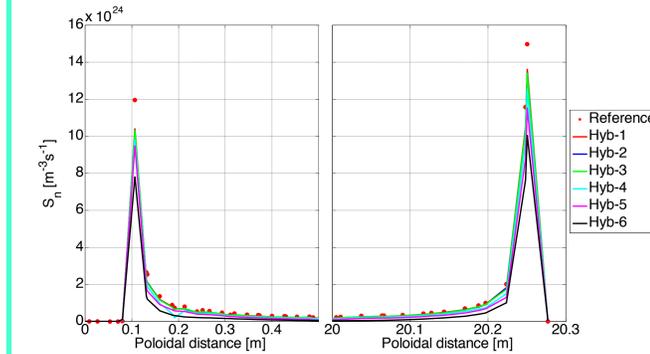
Total plasma sources = kinetic sources + fluid sources



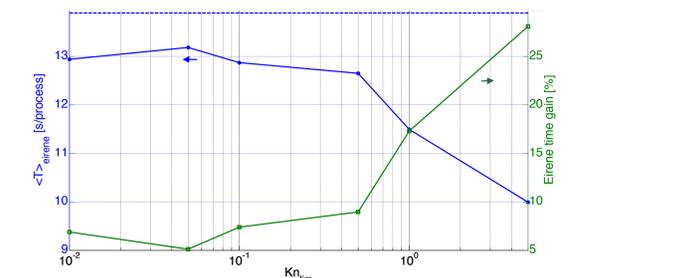
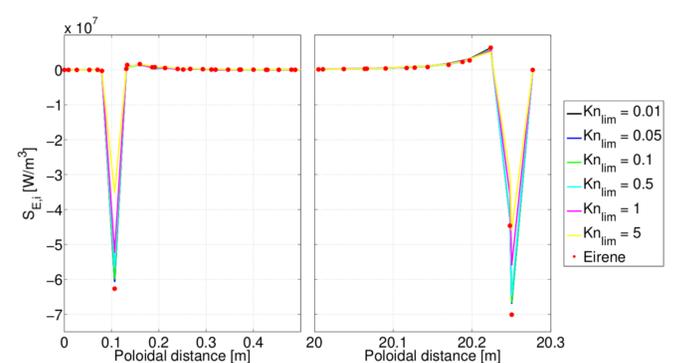
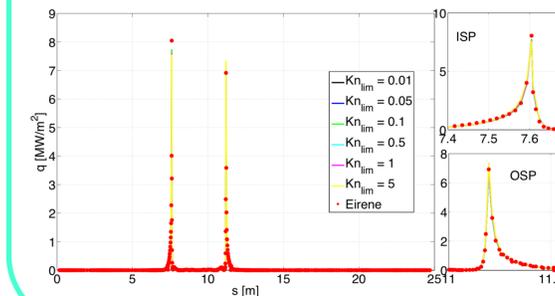
	Eirene	Fluid code	$Kn_{lim}$
Reference	✓	✗	✗
Hyb-1	✓	✓	0.01
Hyb-2	✓	✓	0.05
Hyb-3	✓	✓	0.1
Hyb-4	✓	✓	0.5
Hyb-5	✓	✓	1.
Hyb-6	✓	✓	5.

Code package converged with the hybrid model

$S_{Ei}$  approximated due to  $T_F = T_i$  assumption  $\rightarrow$  possible error introduced in ion energy equation in fluid areas. Here not critical but likely to be case dependent ... would be made more robust with an energy equation for neutrals



BUT negligible impact on the target profiles



reduction of 5 ÷ 30% average time for Eirene vs reference (note that the case considered is not troublesome when using EIRENE alone & that the neutral fluid code is not parallelized yet)

