

# Alfvén Eigenmode evolution in NBI-heated plasmas with dynamic magnetic configuration in the TJ-II stellarator



## magnetic configuration in the TJ-II stellarator

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### INTRODUCTION

Alfvén Eigenmodes (AEs) were studied in the low magnetic shear flexible heliac TJ-II ( $B_0=1$  T,  $\langle R \rangle=1.5$  m,  $\langle a \rangle=0.22$  m). The AE-modes were excited by hydrogen co-NBI in L-mode hydrogen plasmas ( $P_{NBI} \leq 0.56$  MW,  $E_{NBI}=32$  keV), and were diagnosed by HIBP [1], MPs and bolometers. Taking advantage of the unique TJ-II capabilities, a dynamic magnetic configuration experiment with iota variation during discharge was performed via inducing the net plasma current  $I_{pl}$ . Experiment has shown a strong effect of the iota value on the mode frequency. A drastic frequency change from  $\sim 50$  to  $\sim 250$  kHz was observed for some AEs, when plasma current as low as  $\pm 2$  kA was induced by small ( $\leq 10\%$ ) changes in the vertical field (VF). On top of the conventional linear link between  $f_{AE}$  and  $I_{pl}$ , which could explain the local extrema of  $f_{AE}$  coinciding with the extrema of  $I_{pl}$  via  $k_{||}$ , a new type of  $f_{AE}$  dependence on  $I_{pl}$  has been observed.

### EXPERIMENTAL SET-UP

Beam characteristics		Plasma parameters	
Energy	$\Delta E_{beam}$	$\phi \rightarrow E_r$	
Beam current	$\sim I_{beam}$	$\sim n_e$	
Displacement	$\sim Z_d$	$\sim B_{pol}$	

Poloidal orientation of sample volumes:

$$E_p = (\phi_1 - \phi_2)/x, \quad x \sim 1 \text{ cm}, \quad k_\theta < 3 \text{ cm}^{-1}$$

Turbulent particle flux

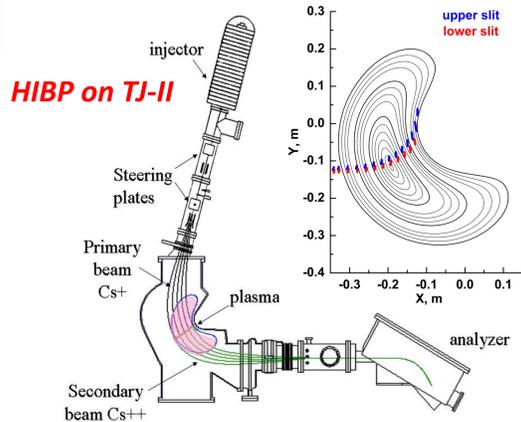
$$\Gamma_r = \Gamma_{EpolkBot} < n_e \sim v_r \rightarrow = \Gamma_{ExB}$$

Poloidal mode number

$$m = Lk_\theta / 2\pi$$

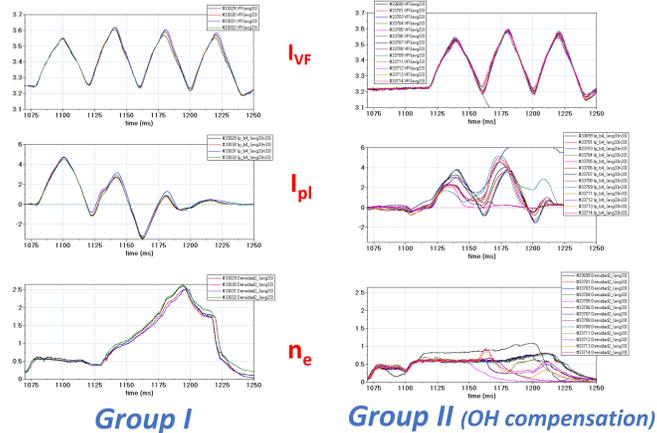
Poloidal propagation velocity

$$V_\theta = Lf/m$$

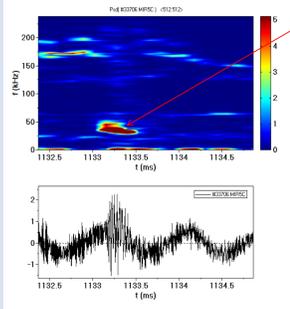


### SCENARIO: Variable Magnetic Configuration

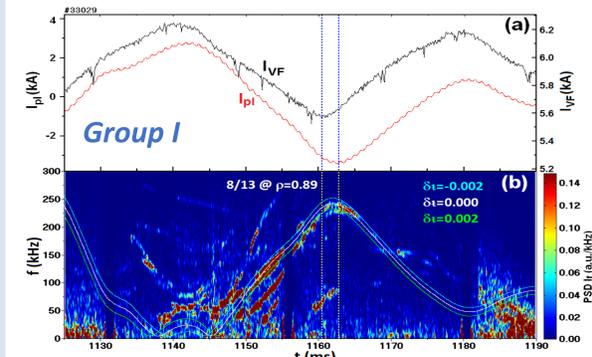
Vertical Field (VF) Modulation keeps vacuum iota the same, but causes the change of the plasma current  $I_{pl}$  and therefore changes the real iota.



Extrema of type II are reproducible and robust phenomenon. They are accompanied by splash of the frequencies 35-45 kHz and higher up to 300 kHz.



### EXPERIMENT: Extrema of Type I



Minima and maxima of AE frequency correspond to the extrema of  $I_{pl}$

#### Modelling

AE frequency follows the formula [2]

$$f_{AE}(\rho_{AE}) = \frac{1}{2\pi} k_{||} V_A$$

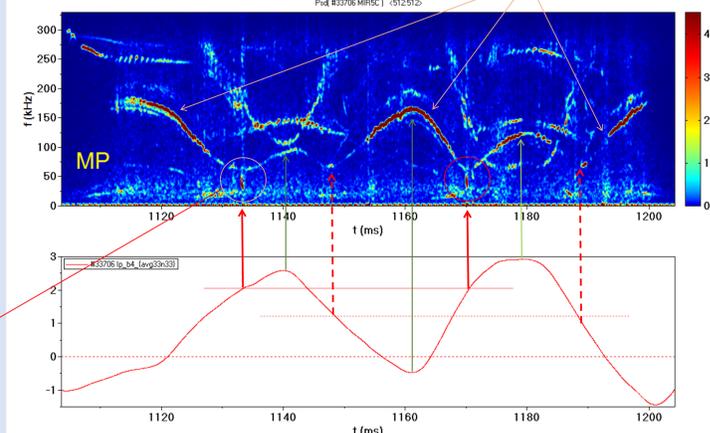
$$k_{||} = \frac{1}{R} |m - n|$$

$$\pm(\rho_{AE}, t) = \pm_{vac}(\rho_{AE}, t) + C(\rho_{AE}) I_{pl}(t)$$

- The first group of shots is characterized by:
- 1) Lifetime of AE is below a period of VF variation.
  - 2) AE frequency minima and maxima corresponds to extrema of actual iota caused by  $I_{pl}$ .

### EXPERIMENT: Extrema of Type II

Group II Mode under study lives as long as two periods of  $I_{pl}$  modulation – long-living mode (90 ms), very rare case for TJ-II



For the selected mode:

- 1) Local maxima of  $f_{AE}$  happen at the time of local extrema (maxima or minima) of  $I_{pl}$  - extrema of type I.
- 2) There are local minima of  $f_{AE}$ , which are located inside the intervals of linear  $I_{pl}$  growth/decay - extrema of type II.

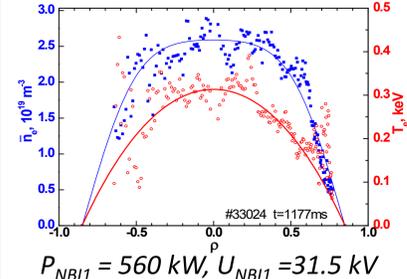
AE frequency minima of type II:  $I_{pl}$  doesn't have extrema in these points.

There is the experimental lower limit for  $f_{AE}$ :  $f_{AE} > 50$  kHz.

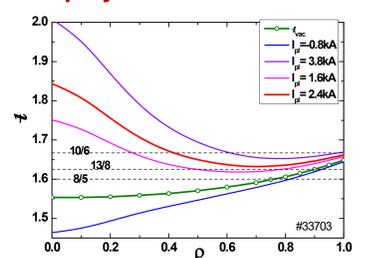
A theory predicts  $f_{AE} > f_{GAM}$

### PROFILES

#### TS: $T_e$ and $n_e$ profiles

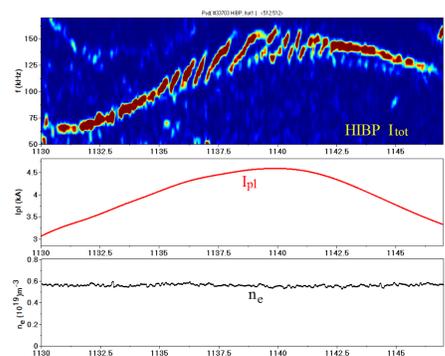


#### iota profiles at various currents

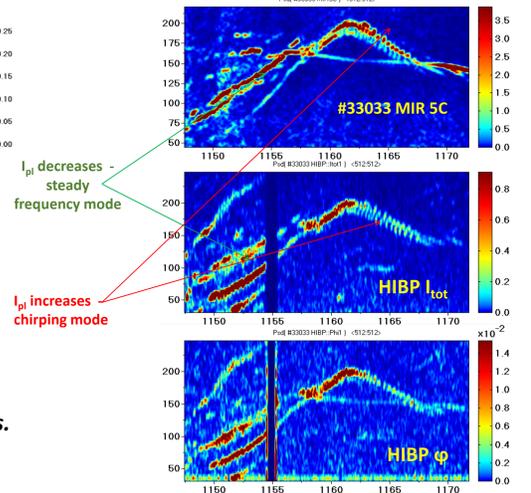


### STEADY TO CHIRPING TRANSITION

#### Transformation due to $I_{pl}$ variation caused by VF

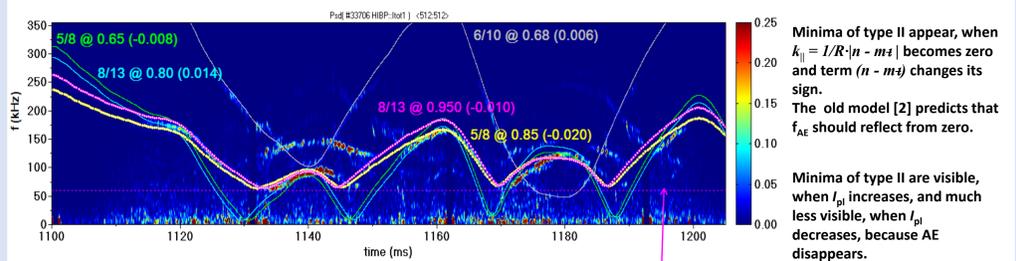


#### Transformation due to change sign of $I_{pl}$ time-derivative



$n_e$  and  $T_e$  are permanent, collisionality (slowing-down time) remains. How does iota cause the transformation?

### MODELLING: Modified model



Modified model:  $f^*(t) = f_{model}(I_{pl}(t)) + f_0$ ,  $f_0 = f_{min} = const$

If we shift  $f_{model}$  on  $f_{min} = const$ , we can obtain the curve  $f^*(t)$  very close to the experimental  $f_{AE}(t)$ .

Minima of type II in this case are result of change of sign of  $|n - \epsilon m|$  in  $k_{||} = 1/R * |n - \epsilon m|$ .

Minima of type II appear, when  $k_{||} = 1/R * |n - \epsilon m|$  becomes zero and term  $(n - \epsilon m)$  changes its sign. The old model [2] predicts that  $f_{AE}$  should reflect from zero.

Minima of type II are visible, when  $I_{pl}$  increases, and much less visible, when  $I_{pl}$  decreases, because AE disappears.

Minima of type II are smooth due to the iota smoothing. They should be sharp due to the modulus dependence in  $|n - \epsilon m|$ .

### SUMMARY

- The global minimum value for  $f_{AE}$  was observed, no AE exists with  $f_{AE} < f_{min}$
- On top of conventional linear link between  $f_{AE}$  and plasma current  $I_{pl}$ , a new type of  $f_{AE}$  dependence on  $I_{pl}$  has been observed in TJ-II: when  $\epsilon = n/m$  for specific mode,  $f_{AE}$  reaches  $f_{min}$ , then changes direction of evolution (decrease to raise).
- The evolution of AE from steady frequency mode to chirping mode and back takes place, when the plasma current reaches certain values or changes direction of evolution (raise to decrease).

### ACKNOWLEDGEMENTS / REFERENCES

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[1] A.V. Melnikov *et al*, Nuclear Fusion **57** (2017) 072004

[2] A.V. Melnikov *et al*, Nuclear Fusion **54** (2014) 123002