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## Studies of the pedestal structure in JET with the ITER-like wall

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The H-mode pedestal structure is characterized in JET-ILW plasmas limited by type I ELMs. The pre-ELM pressure width  $\Delta p$  increases with the square root of  $\beta_{pol\_PED}$ , as assumed in EPED, in H-modes with low D gas injection. In dimensionless experiments  $\Delta p$  broadens at constant  $\beta_{pol\_PED}$  with increasing  $\nu_{star\_PED}$ . In power and gas scans  $\Delta p/\sqrt{\beta_{pol\_PED}}$  is constant with  $\nu_{star\_PED}$ , but is systematically wider at higher than at lower D rates.  $\Delta p$  may therefore depend also on other parameters, directly or indirectly connected with the D neutral content in the plasma, implying that atomic physics could contribute in setting the pedestal width.

The pedestal evolution during the ELM cycle is more complex than what would be expected if KBMs were to control the inter-ELM pressure gradient evolution. At high beta and low D gas injection rates  $p_{e\_PED}$  increases due to narrowing of the width and steepening of the gradient. The  $n_e$  width narrows and the gradient increases until the ELM, suggesting qualitative consistency with the neutral penetration model. The  $n_e$  pedestal structure evolves similarly at low and high D gas rates. At high beta  $T_{e\_PED}$  saturates half way through the ELM cycle at high D gas rates. This causes the reduction in  $p_{e\_PED}$  in higher beta plasmas when the D gas rate is increased at constant net input power. The edge bootstrap current  $J_{BS}$ , derived with NEO from the measured kinetic profiles and  $Z_{eff}$ , increases throughout the ELM cycle at low beta, while it saturates well in advance of the ELM crash at high beta.

Initial isotope experiments have investigated pedestal formation in H vs D. In the high density branch,  $P_{L-H}(H) = 2 \times P_{L-H}(D)$ . H and D plasmas have matched stored energy and diamagnetic edge  $E_r$ . The higher power in H required to achieve the same stored energy as in D is consistent with the lower L-mode energy confinement in H. Assuming that the L-H transition requires an  $E_r \times B$  shearing rate  $\sim$  to the growth rate of the most unstable mode,  $\gamma_{Ecrit} \sim \gamma_{turb}$ , the above result suggests  $\gamma_{turb}$  independent of mass in the high  $n_e$  branch. In the low  $n_e$  branch  $P_{L-H}(H) \gg 2 \times P_{L-H}(D)$ . Similar edge  $n_e$  profiles are observed in H and D, but H plasmas have a stronger  $T_e$  gradient, indicating the need for a larger  $\gamma_{Ecrit}$  to trigger the L-H transition and suggesting an isotope effect on  $\gamma_{turb}$ .

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**Primary author:** Dr MAGGI, Costanza (CCFE)

**Co-authors:** Ms LUNNISS, Amelia (York Plasma Institute, University of York, UK); Dr DELABIE, Ephrem (Oak Ridge National Laboratory); Dr CASSON, Francis (CCFE); Dr URANO, Hajime (JAEA); Prof. WILSON, Howard

(York Plasma Institute, University of York, UK); Dr LUPELLI, Ivan (CCFE); Mr HORVATH, Laszlo (York Plasma Institute, University of York, UK); FRASSINETTI, Lorenzo (KTH, Royal Institute of Technology); Dr LEYLAND, Matthew (York Plasma Institute, University of York UK); Dr STANISLAS, Pamela (CCFE); Dr SAARELMA, Samuli (CCFE)

**Presenter:** Dr MAGGI, Costanza (CCFE)

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