



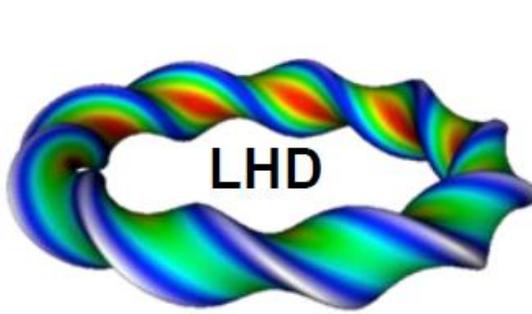
# Multi-machine analysis of turbulent transport in helical systems via gyrokinetic simulation

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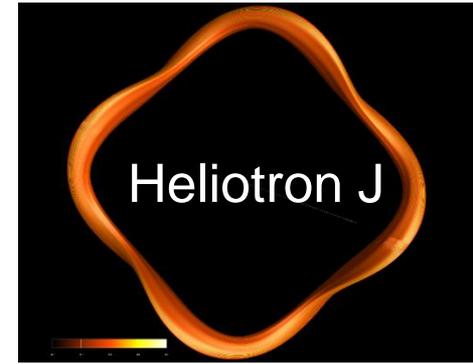
# Outline

- Introduction
- Turbulent transport in Large Helical Device (LHD)
- Turbulent transport in Heliotron J (HJ)
- Comparison of helical systems
- Summary

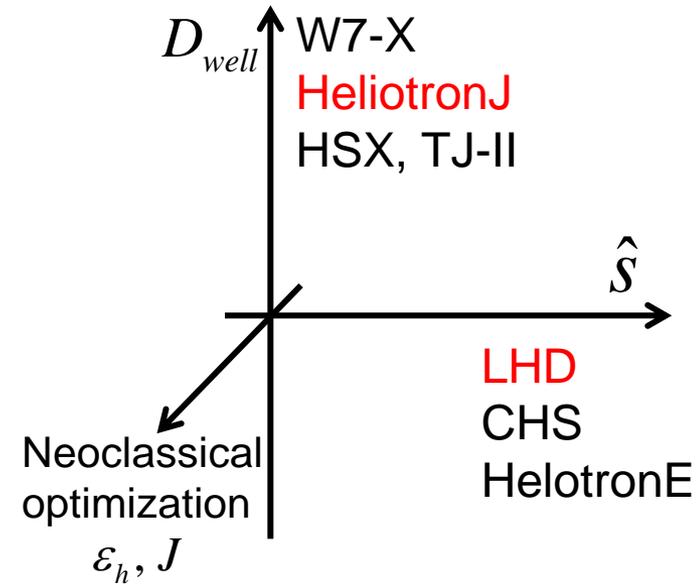
# Helical plasmas



D. Spong, APS2014



- Helical plasmas are studied for optimization by improving MHD stability and neoclassical transport.
  - W7-X: optimized against neoclassical transport as well as MHD stability.
  - Heliotron J : optimized against MHD stability by producing the magnetic well.
  - LHD: better neoclassical transport for the inward-shifted configuration, while it has better MHD stability for the outward-shifted one.
- Two strategies for stability
  - Utilizing magnetic shear: LHD, CHS, Heliotron-E
  - Utilizing magnetic well (Mercier well index): W7-X, Heliotron-J, HSX, TJ-II
- Recently, optimization against turbulent transport becomes a hot topic.



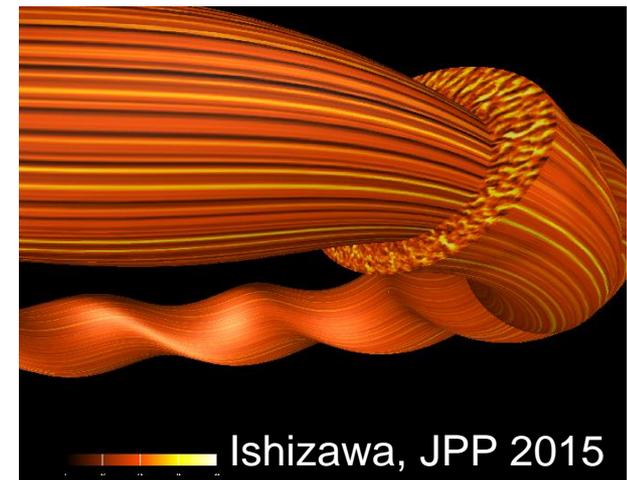
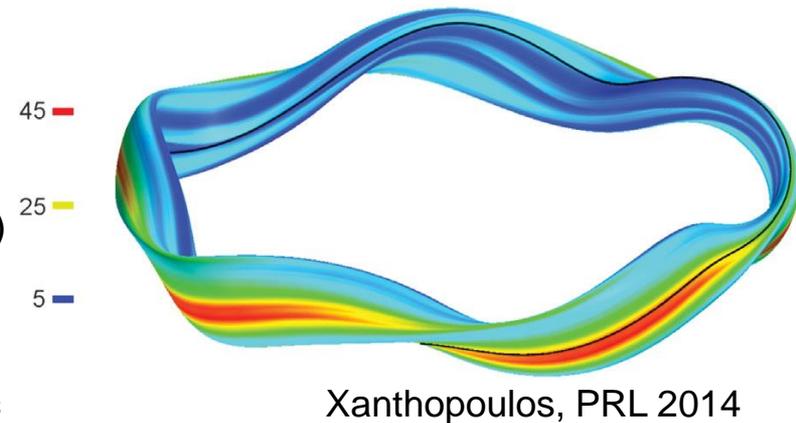
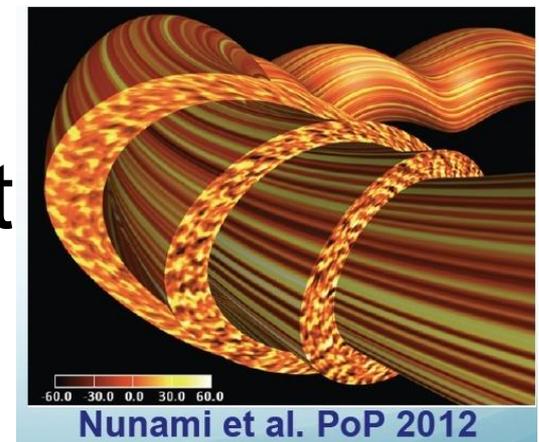
# Parameters of LHD and Heliotron J plasmas

- The aspect ratio  $R/a$ , safety factor  $q$ , normalized Larmor radius, and temperature ratio  $T_e/T_i$ , are similar. On the other hand, the normalized collision frequency and the density gradient length  $L_n$  are significantly different.
- The LHD is the inward-shifted configuration, and it is the magnetic hill with a moderate shear.
- The Heliotron J is the magnetic well with a very weak shear.

	LHD-L	HJ-ST
$R_0/a$	6.2	7.3
$\rho = r/a$	0.68	0.5
$q$	1.5	1.7
$\rho_* [10^{-3}]$	2.	4.5
$v_i^*$	0.083	3.2
$\beta [\%]$	0.2	0.05
$T_e/T_i$	0.96	1.3
$R_0/L_n$	2.7	9.3
$R_0/L_{Ti}$	8.7	13.
$R_0/L_{Te}$	9.1	17.
$\hat{s}$	1.2	0.023
$D_{\text{well}}$	-0.01	0.74

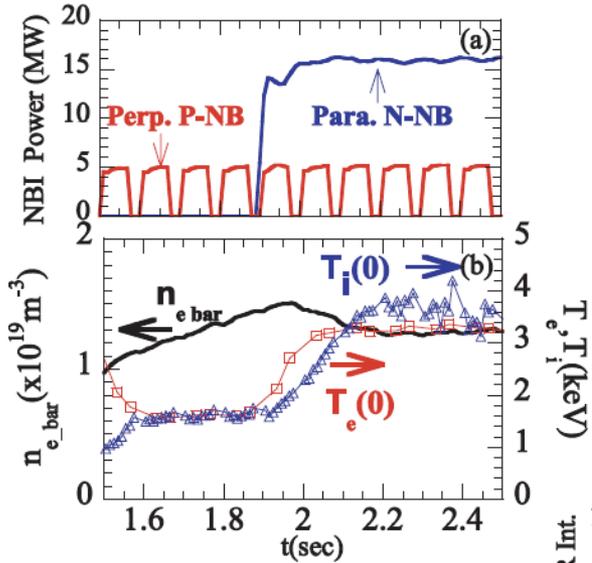
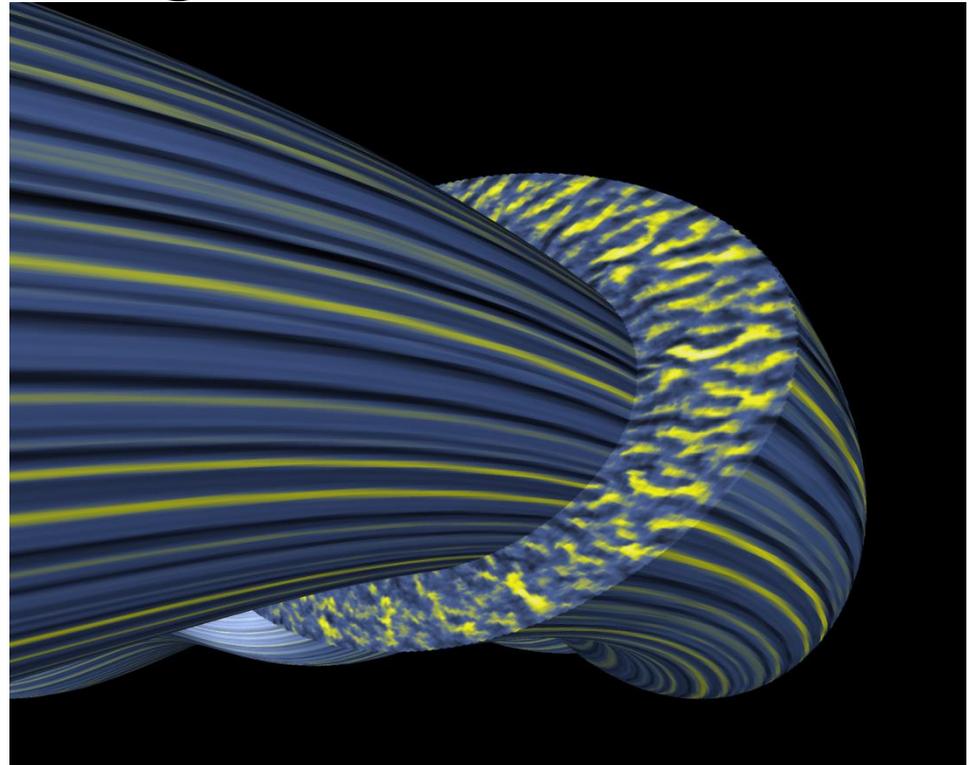
# Progress of gyro-kinetic analysis of turbulent transport in helical systems

- Adiabatic electron simulations
  - LHD: T.-H. Watanabe Phys. Rev. Lett. (2008)  
M. Nunami, Phys. Plasmas (2012, 2013)
  - W7-X: P. Xanthopoulos, Phys. Rev. Lett. (2007)  
P. Xanthopoulos, Phys. Rev. Lett. (2014)
  - NCSX: H. Mynick, Phys. Rev. Lett. (2010)  
H. Mynick, Plasma Phys. Cont. Fusion (2014)
- Kinetic electron simulations
  - Enable us to evaluate particle and electron heat fluxes.
  - Trapped electron effects enhance the growth rate of ITG mode.
  - HSX: B. J. Faber, Phys. Plasmas (2015)
  - LHD: A. Ishizawa, Nuclear Fusion (2013, 2015)  
A. Ishizawa, Phys. Plasmas (2014)  
A. Ishizawa, J. Plasma Phys. (2015)
- This conference
  - TH/P2-3, M. Nunami
  - TH/P4-10 D. Spong
- We use GKV code (Local flux tube code).

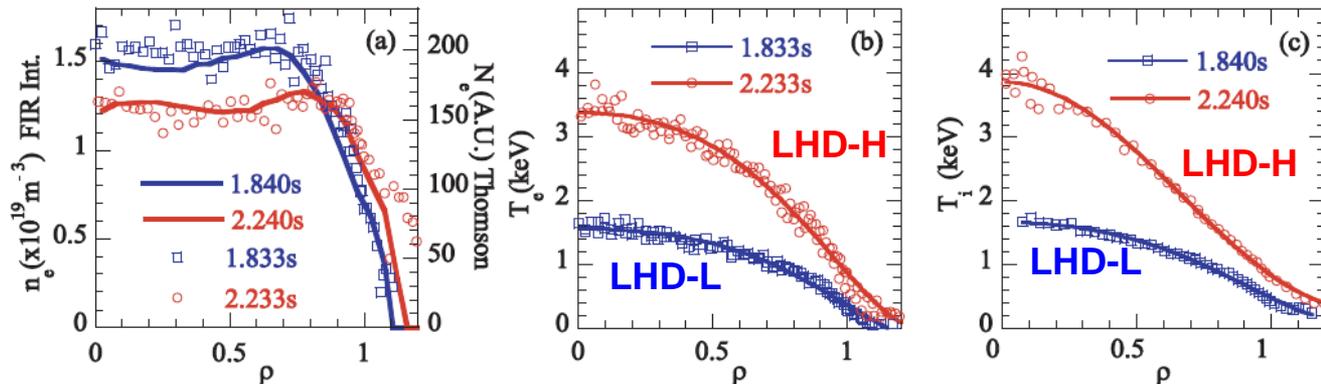


# LHD discharge #88343

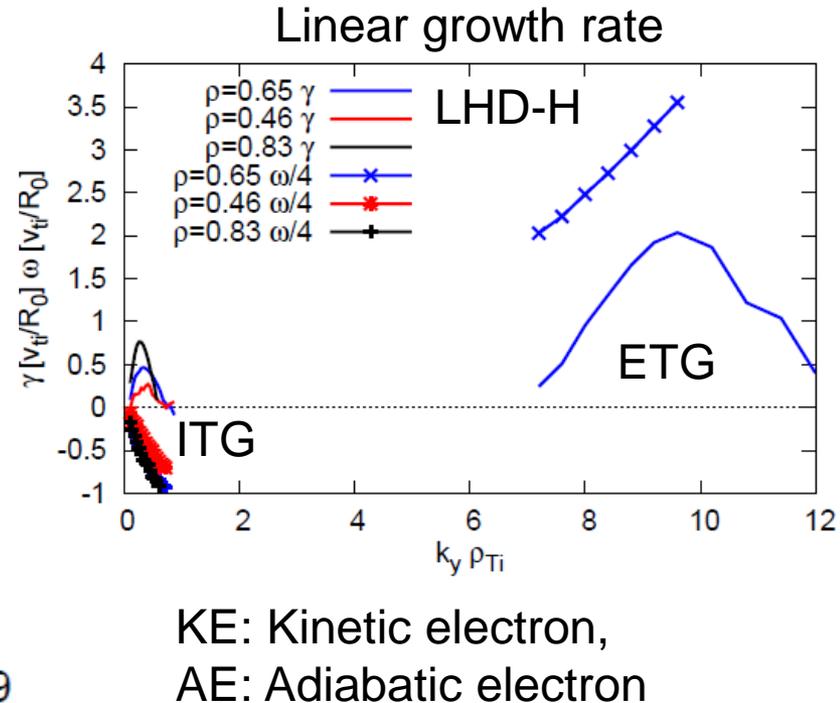
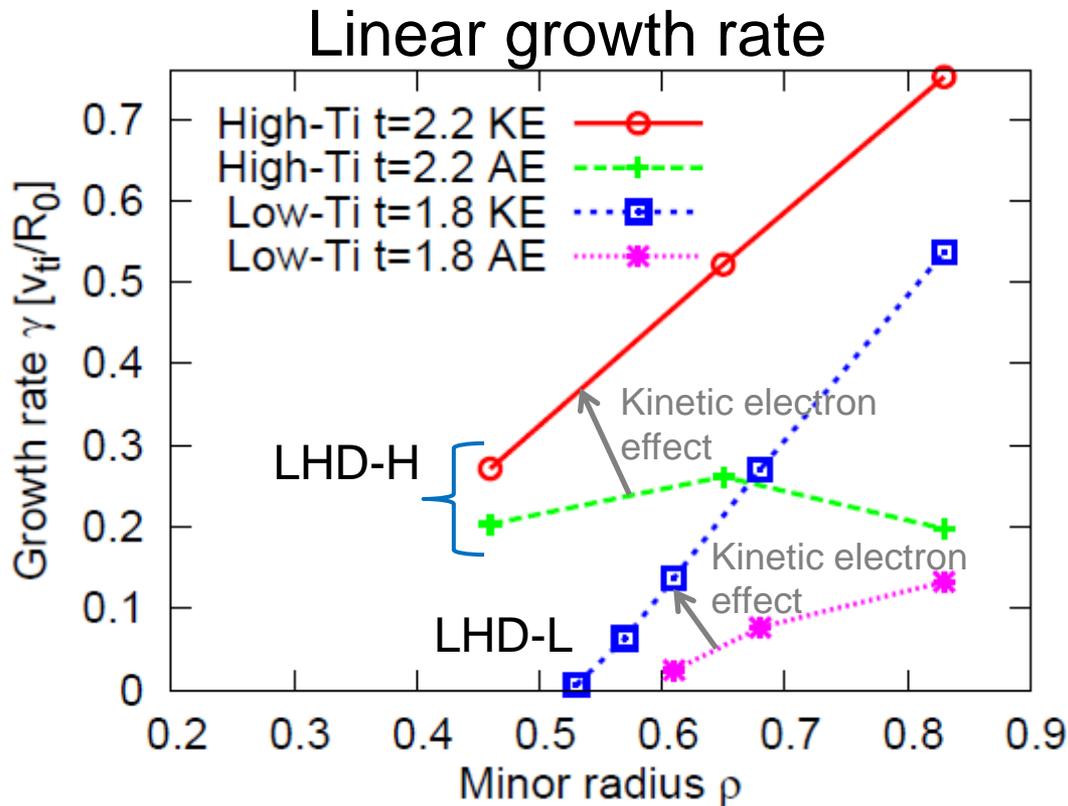
- $B=2.75\text{T}$ ,  $R=3.6\text{m}$
- Low-Ti phase:  $T_i=1.6\text{keV}$   $t=1.8\text{s}$
- High-Ti phase:  $T_i=3.9\text{keV}$   $t=2.2\text{s}$ 
  - $\text{Beta}(r/a=0.65)=0.3\%$
  - Collision:  $1/\nu$  regime



Tanaka, PFR 2010



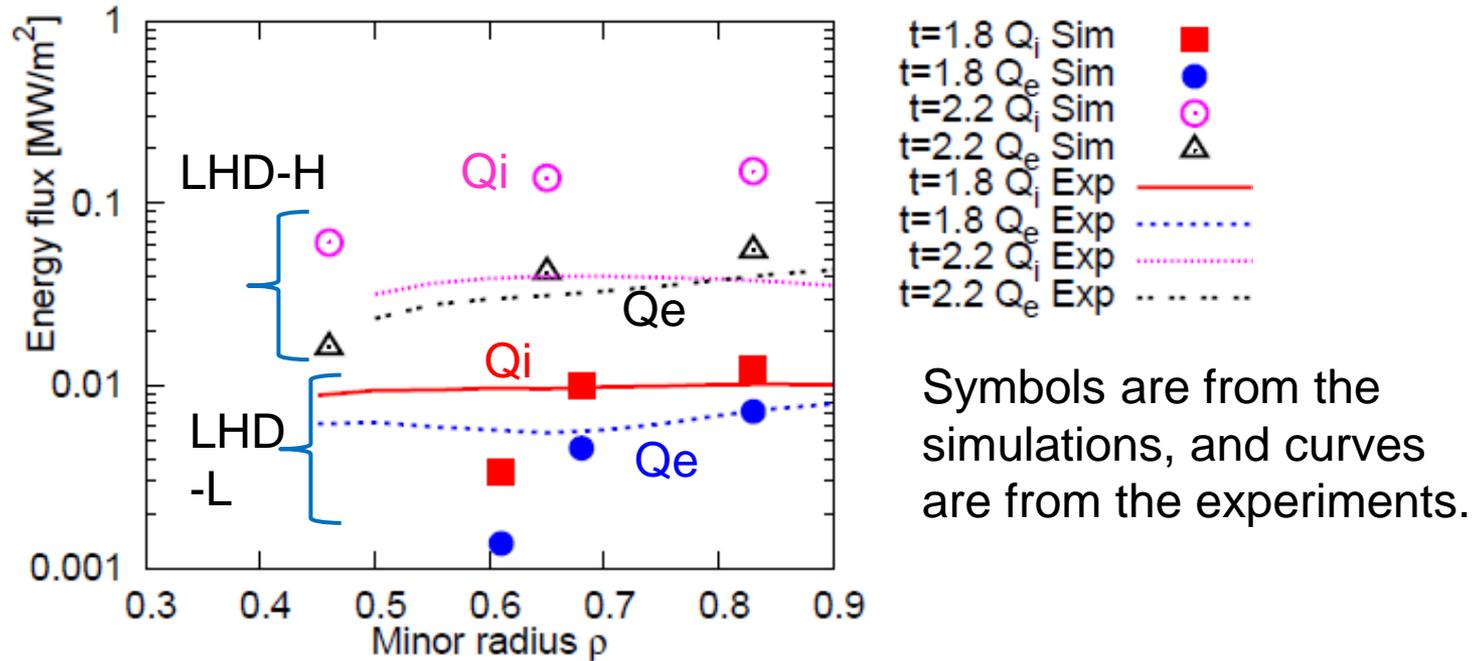
# Kinetic electron effects enhance ITG mode



- Ion temperature gradient (ITG) modes are unstable.
- Kinetic electron (KE) effects enhance the ITG mode.
- ETG modes are unstable at high wavenumber.

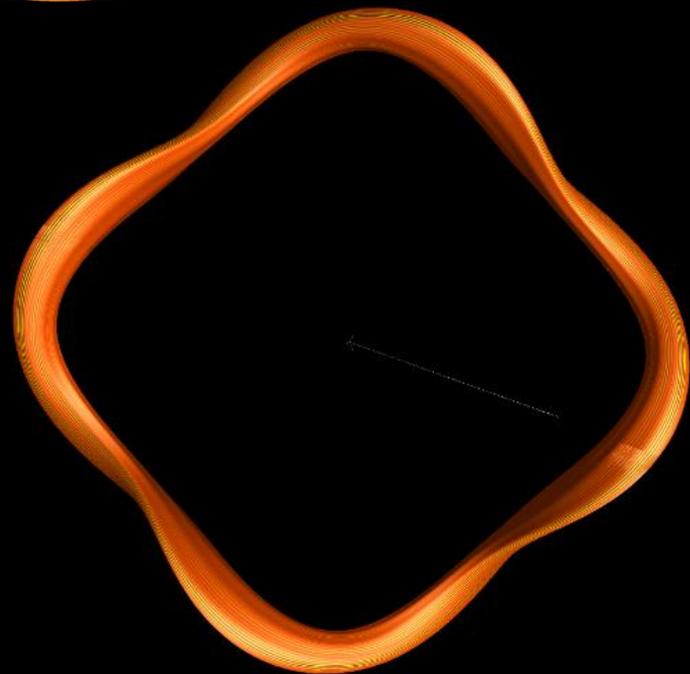
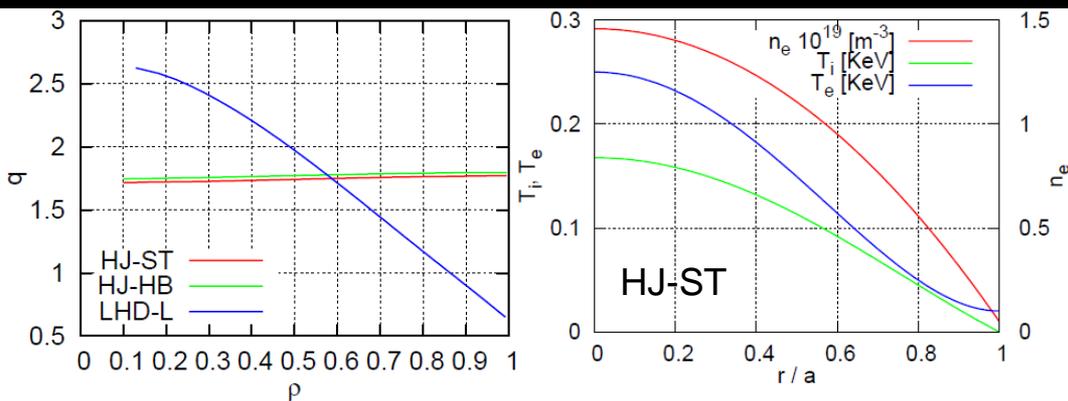
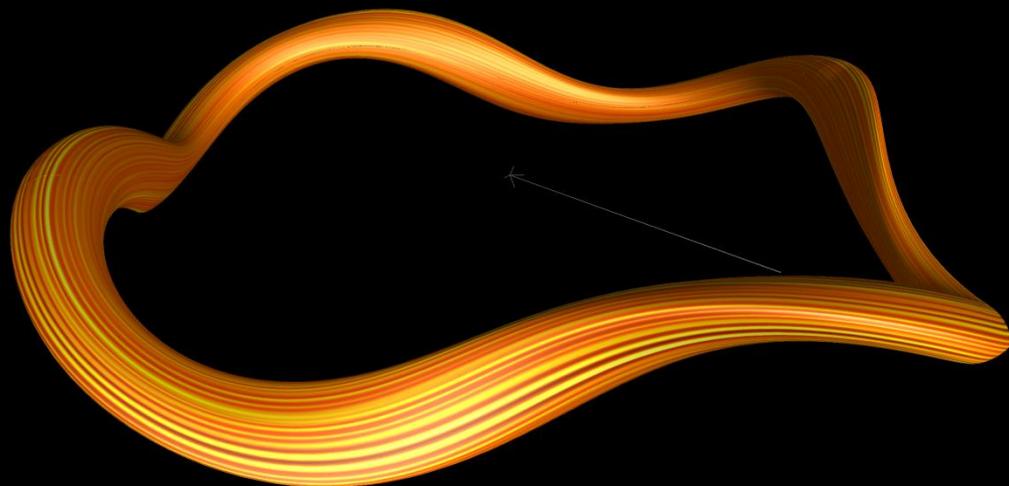
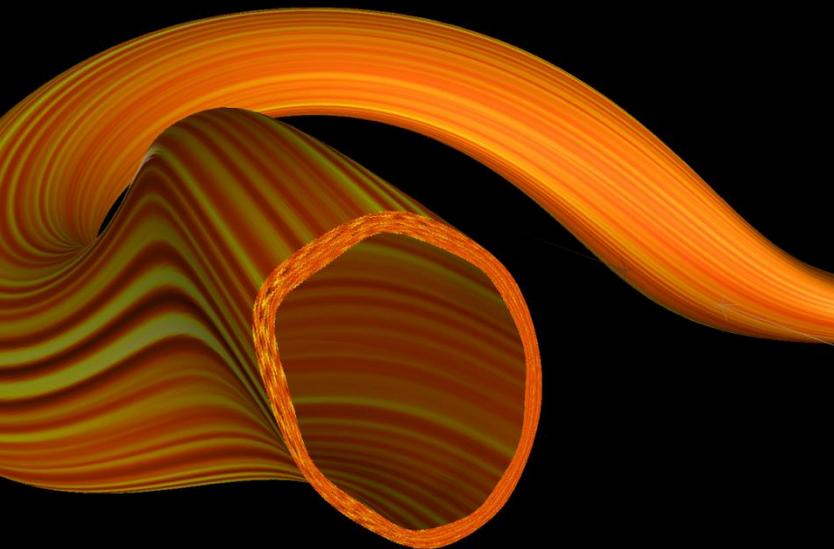
# Importance of kinetic electrons in validation

Ion and electron energy fluxes due to turbulence,  $Q_i$  and  $Q_e$ .



- The energy fluxes are in good agreement with experimental results at  $\rho > 0.7$  in the low-Ti phase (LHD-L,  $T_i = 1.6 \text{ keV}$ ).
- The electron energy flux is in good agreement with experimental results in the high-Ti phase (LHD-H,  $T_i = 3.9 \text{ keV}$ ).
- Prediction of temperature gradient length by flux matching has 20% error.
- There is no short-fall problem, which suffers GK analysis of some tokamaks.

# Turbulence in Heliotron J

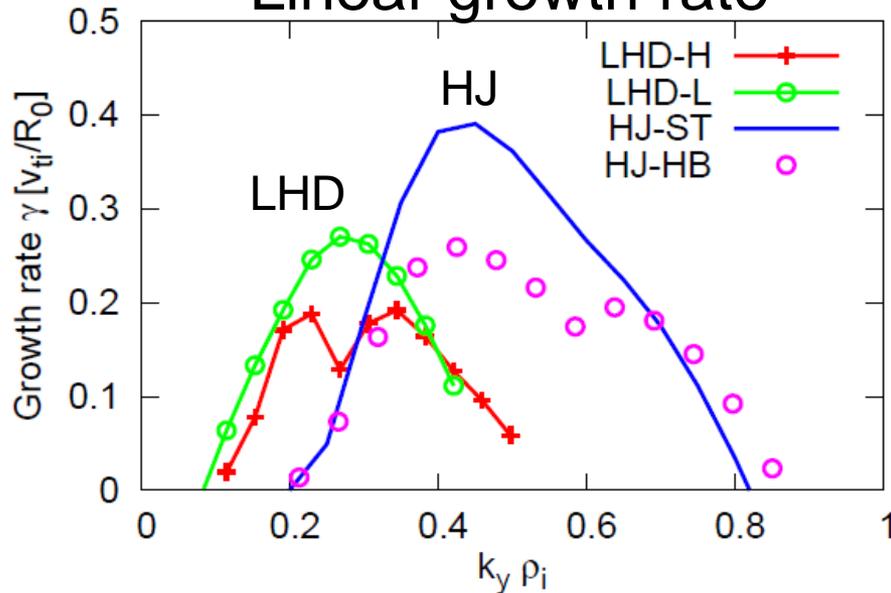


HJ-ST: standard configuration

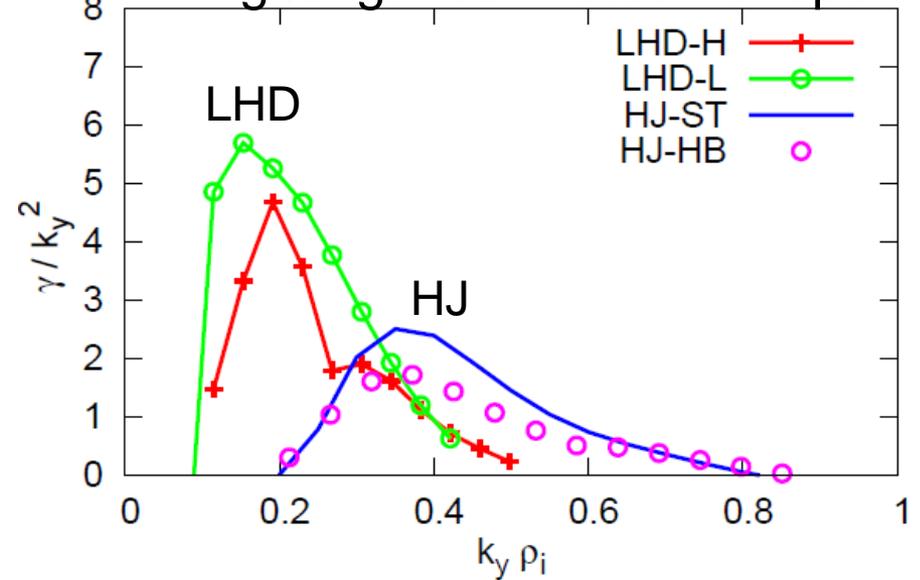
HJ-HB: High-bumpiness (high-mirror ratio)

# HJ plasmas are unstable against ITG mode

Linear growth rate



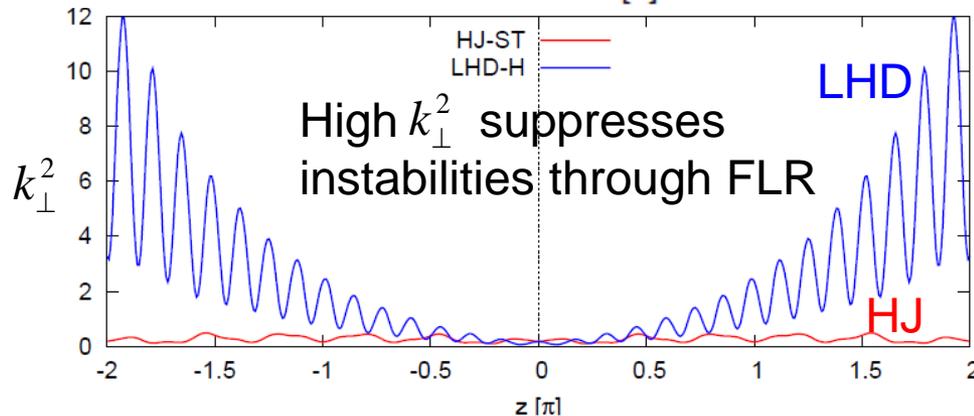
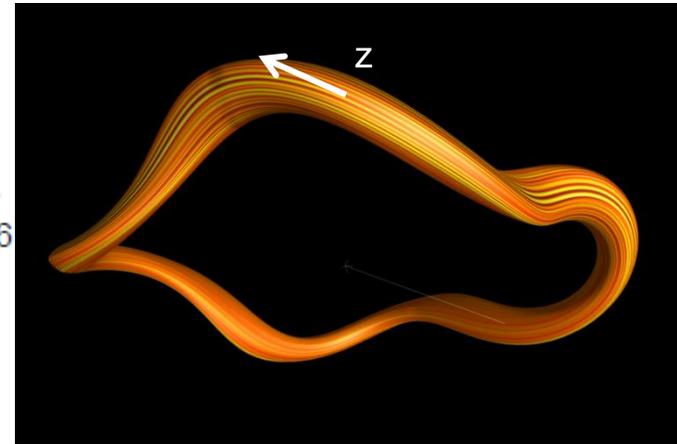
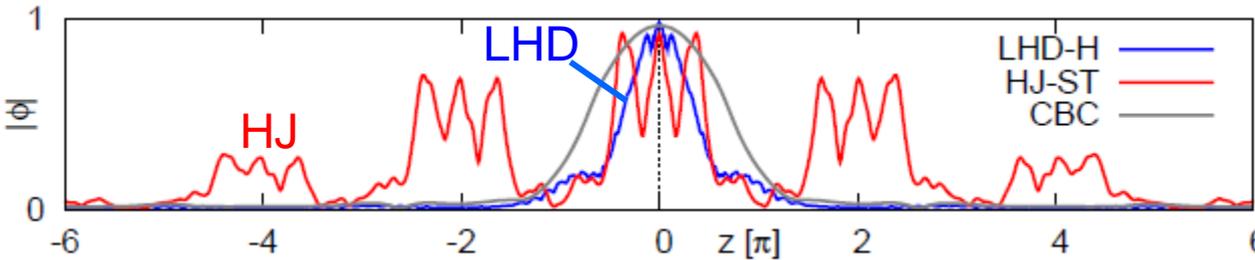
Mixing length estimated transport



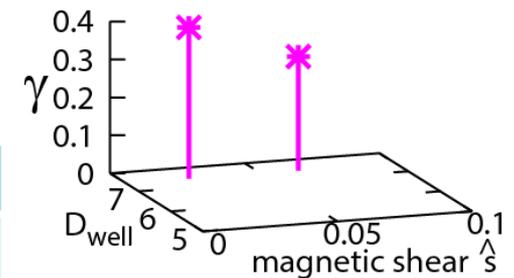
- HJ plasma is unstable at higher wavenumber regime than LHD.
- Mixing length estimation predicts lower transport in HJ than LHD.

# Elongated mode structure in HJ

- The mode structure is elongated along the field line in HJ.
- That is due to the weak shear and clearly seen in the profile of  $k_{\perp}^2$
- The stabilizing effect of shear is confirmed by the reduction of growth rate by increasing the shear.



	HJ standard	HJ with shear	LHD
Magnetic shear $\hat{s}$	0.023	0.064	1.2
Growth rate $\gamma$	0.4	0.3	0.26
$\gamma / k_{\perp}^2$	2.5 [GB]		5.7 [GB]

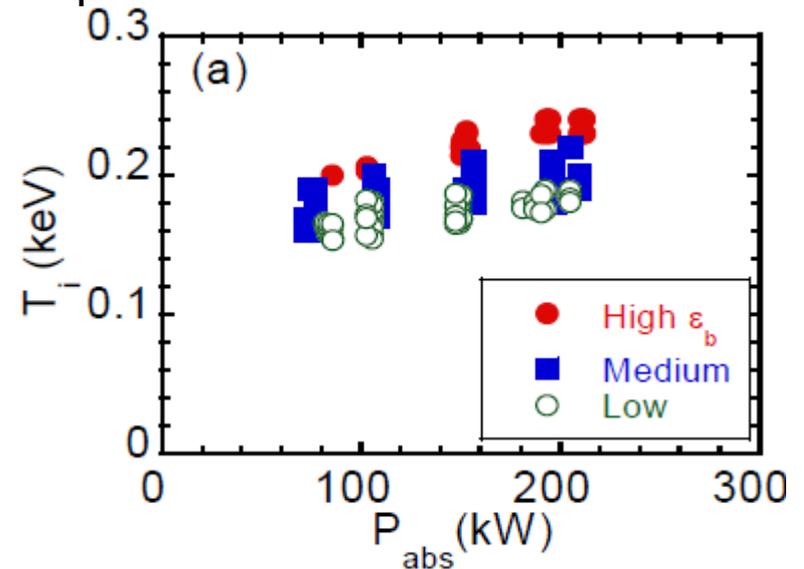


# High mirror-ratio reduces turbulent transport

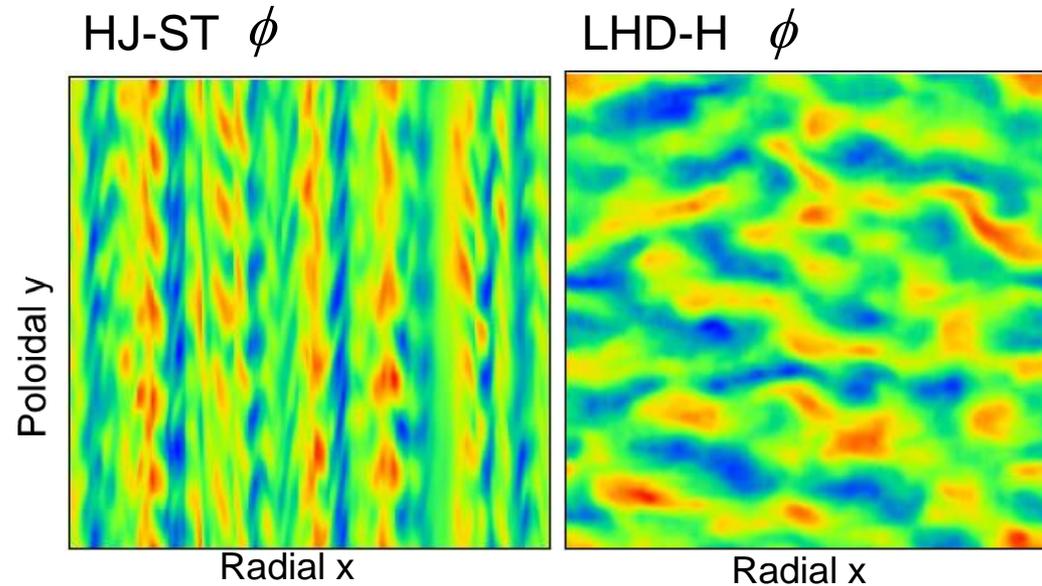
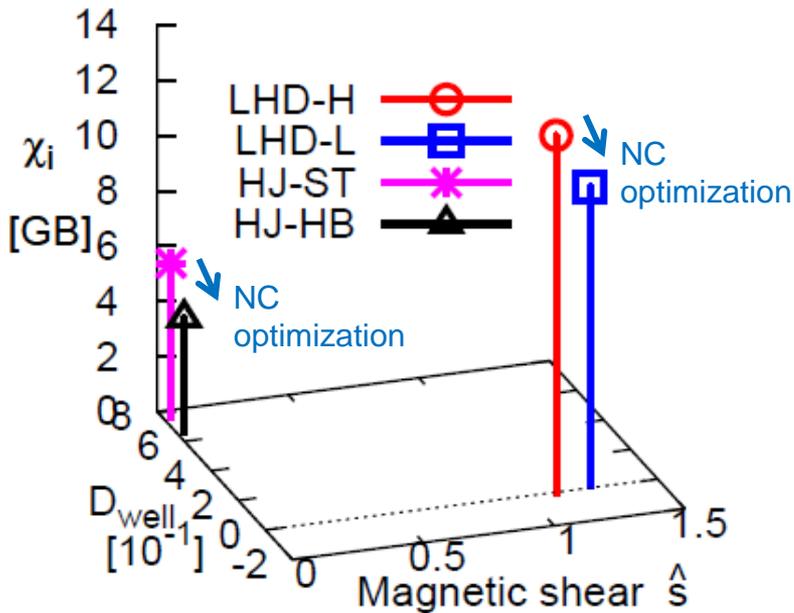
Heliotron J	Standard (HJ-ST)	High mirror-ratio (HJ-HB)
Growth rate $\gamma$	0.40	0.26
Heat transport $\chi_i, \chi_e$	5.9, 2.4 [GB]	4.2, 1.7 [GB]
$\gamma/k_{\perp}^2$	2.5 [GB]	1.7 [GB]

- High mirror-ratio (HJ-HB) reduces ITG mode, and thus suppresses turbulent transport.
- Qualitatively consistent with the experimental observation.

Experimental observation in Heliotron J



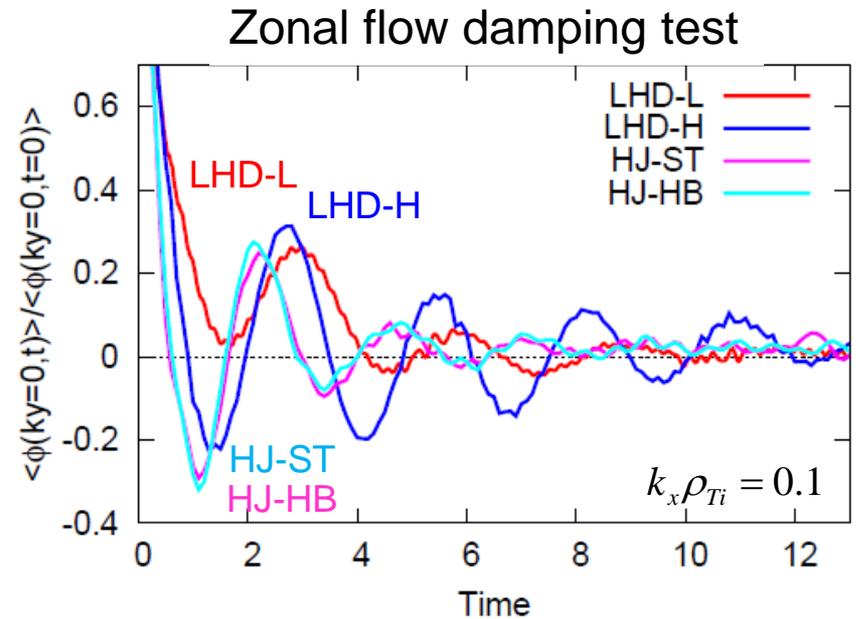
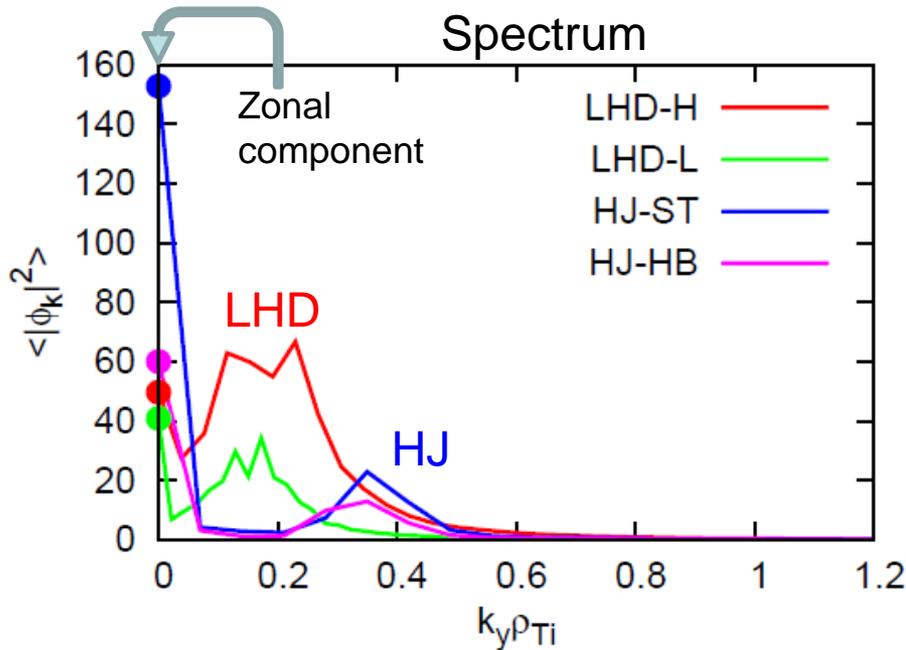
# Turbulent transport and zonal flows



- The neoclassical optimization (high-bumpiness) improves turbulent transport in Heliotron J.
- Weak magnetic shear of Heliotron J does not lead to high turbulent transport.
- Zonal flow in HJ is stronger than LHD.

	LHD-L	HJ-ST
Instability	ITG	ITG
$\gamma [v_{Ti}/R_0]$	0.27	0.4
$R_0/L_T - R_0/L_{Tcrit}$	2.6	5.2
$\chi_i [v_{Ti}\rho_{Ti}^2/R_0]$	11.	5.9
$\chi_e [v_{Ti}\rho_{Ti}^2/R_0]$	4.8	2.4

# Strong zonal flow in Heliotron J



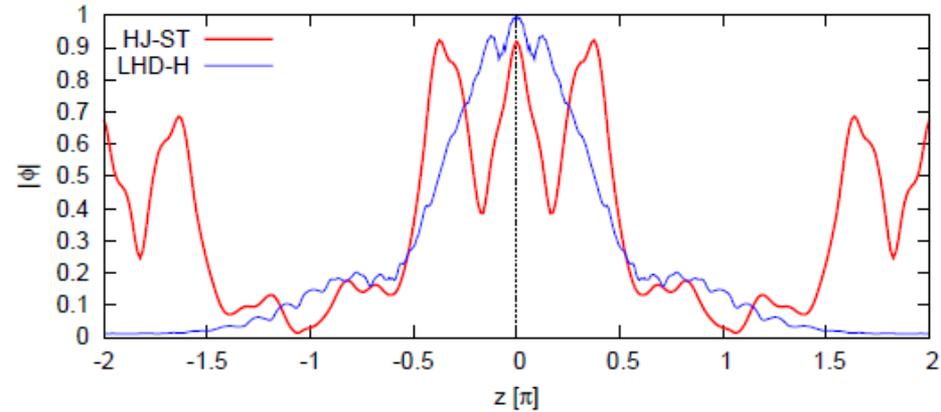
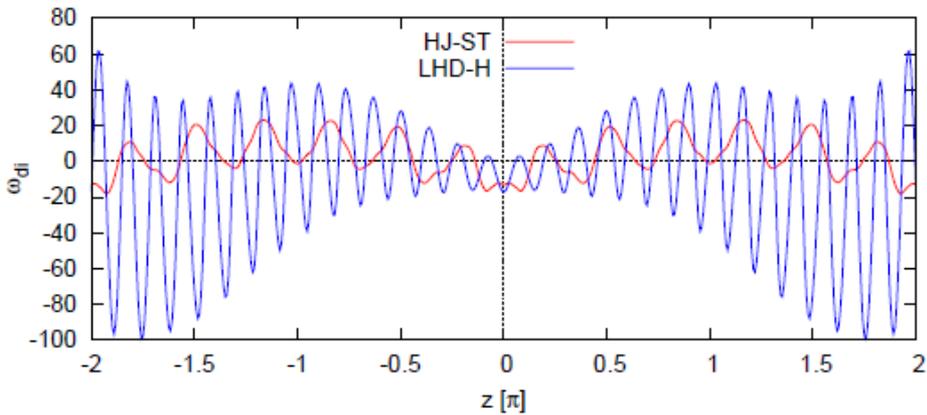
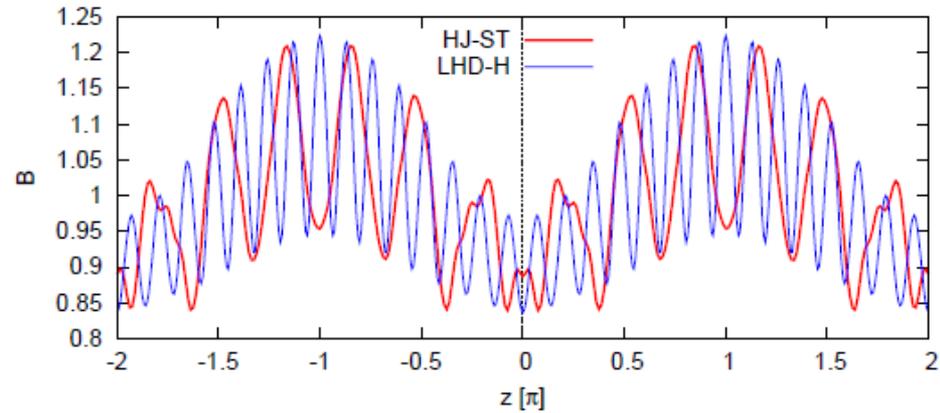
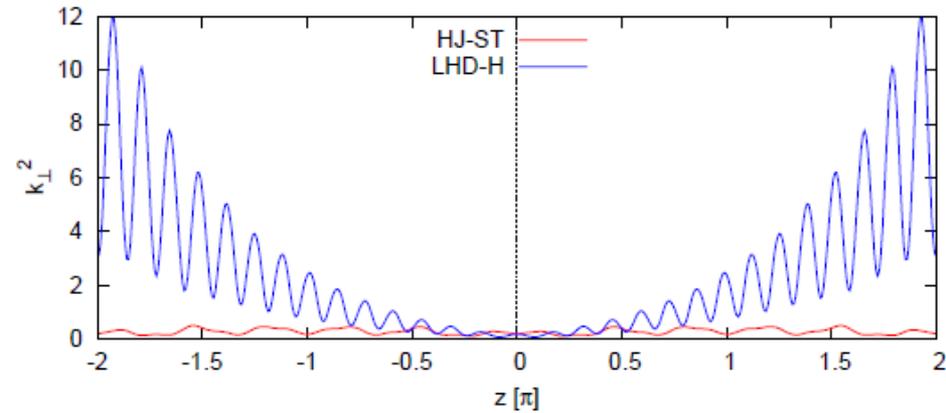
- LHD
  - $\chi_{i,e}$  in LHD-L is smaller than LHD-H, while  $\gamma / k_{\perp}^2$  in LHD-L is larger than LHD-H. This is explained by more inward shifted axis and larger  $\tau_{ZF}$  in LHD-L.
- Heliotron J
  - $\chi_{i,e}$  in HJ-HB is smaller than HJ-ST, which is explained by smaller  $\gamma / k_{\perp}^2$  in HJ-HB.
- Zonal flow relaxation time and the residual level do not explain the strong zonal flow in HJ.
- The strong zonal flow is expected to be produced by nonlinear interaction.

# Summary

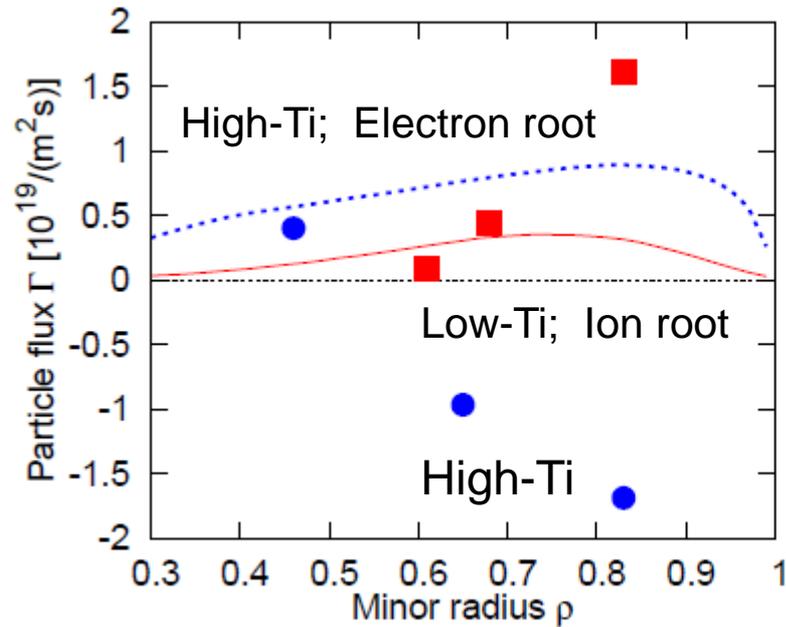
- We have investigated turbulent transport in helical systems, LHD and Heliotron J, by gyrokinetic simulations.
- LHD (Validation)
  - Including kinetic electrons is crucial for the validation.
  - There is no short-fall problem near the edge.
  - The inward-shifted magnetic axis of LHD-L leads to smaller turbulent transport than LHD-H in GB unit because of longer zonal flow relaxation time.
- Heliotron J
  - High mirror ratio (neoclassical optimization) reduces turbulent heat transport.
- In this comparison of those typical cases, turbulent transport in HJ is lower than that of LHD in GB unit. Lower mixing-length estimated transport and higher amplitude of zonal flows can be the mechanism.
- Strong zonal flow in HJ can be produced by elongated mode structure of ITG due to weak shear.

# Additional slides

# Profiles along the field line



# Turbulent particle flux



$$\Gamma = \Gamma_{neo-classical} + \Gamma_{turbulence}$$