



This work has been carried out within the framework of the EUROfusion Consortium and has received funding from the Euratom research and training programme 2014-2018 under grant agreement No 633053. The views and opinions expressed herein do not necessarily reflect those of the European Commission.



EUROfusion



# Recent progress towards a quantitative description of filamentary SOL transport

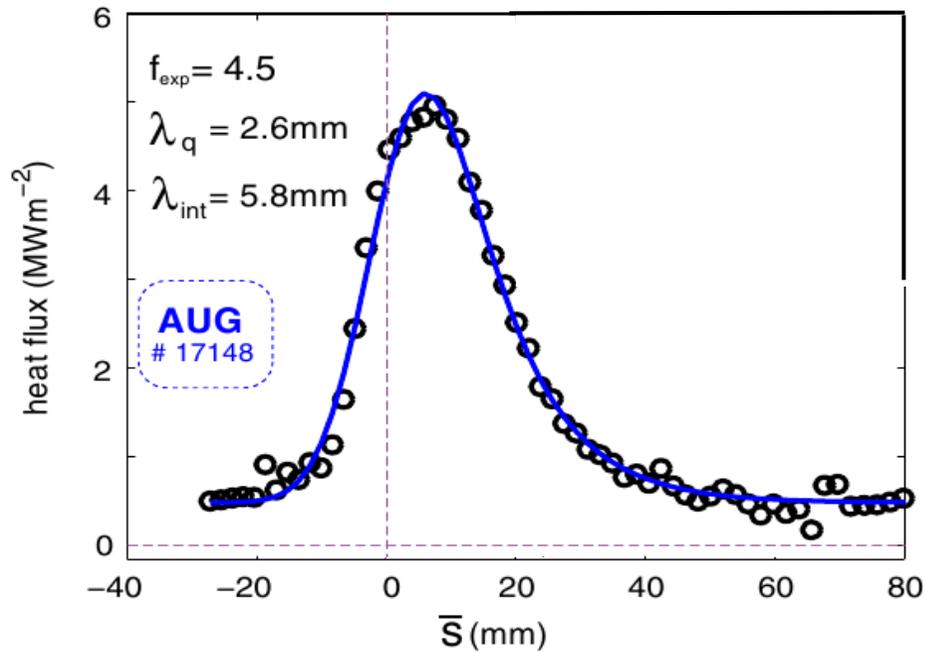
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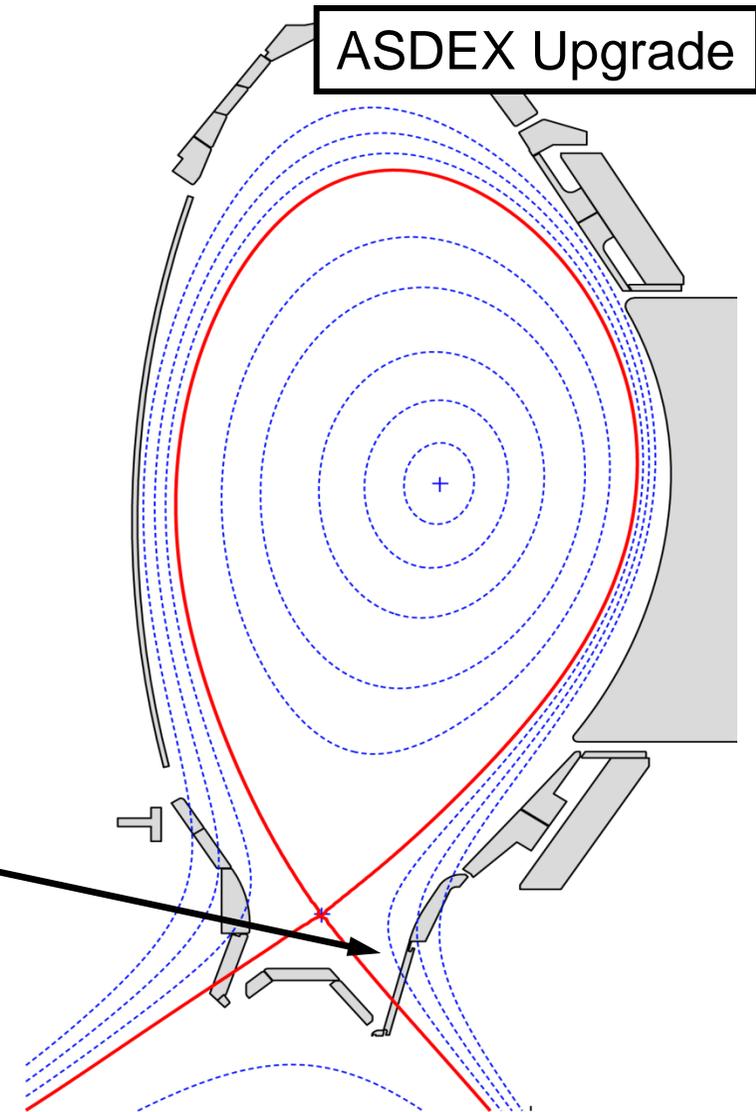
26<sup>th</sup> International Atomic Energy Agency Fusion Energy Conference, Kyoto, Japan  
October 19<sup>th</sup>, 2016

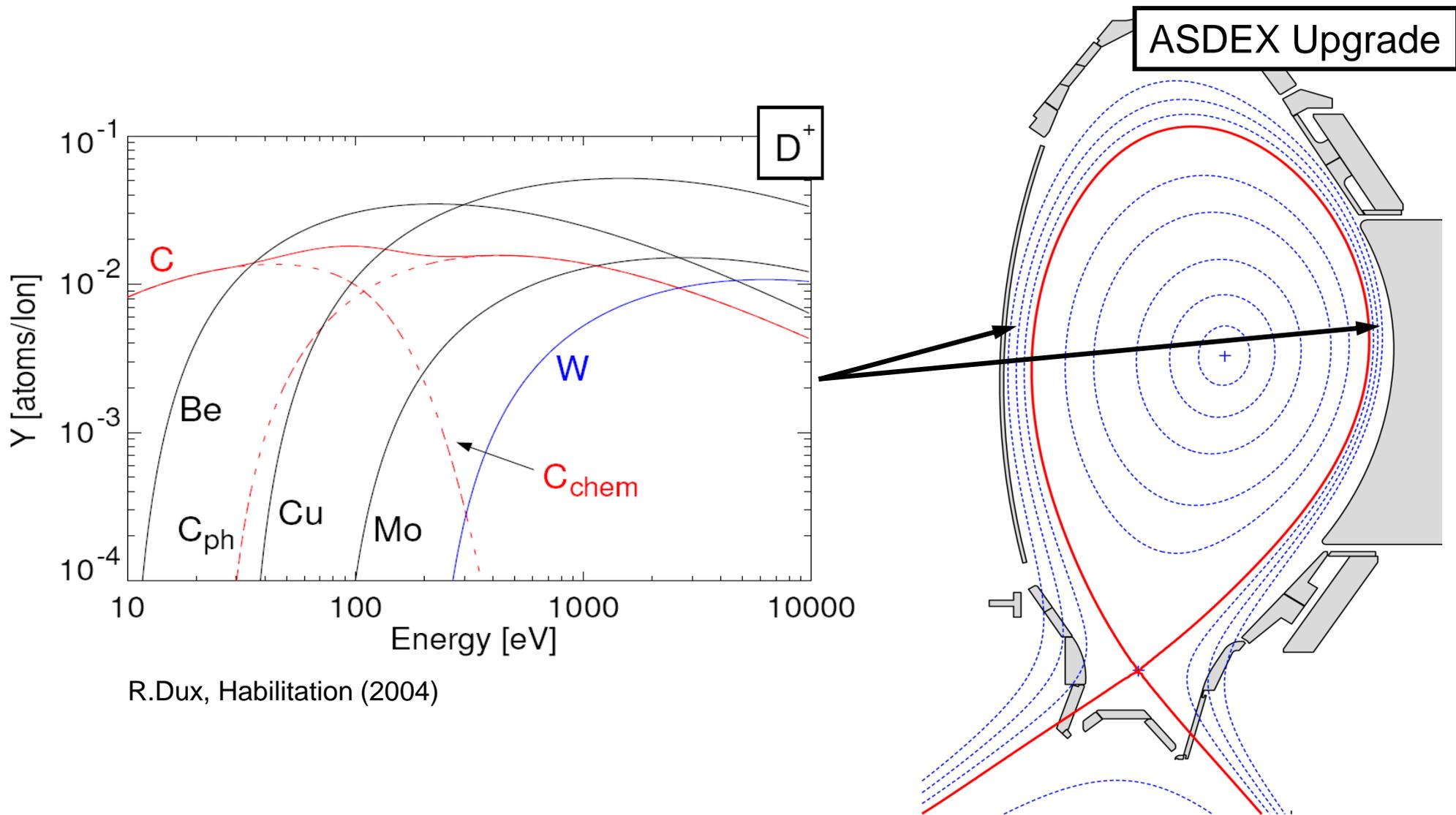
**Wetted area ITER :  $<1\text{m}^2$**   
**Total surfaces :  $680\text{ m}^2$**

## Inter-ELM, attached power deposition profile



*T.Eich et al. PRL 107, (2011)*

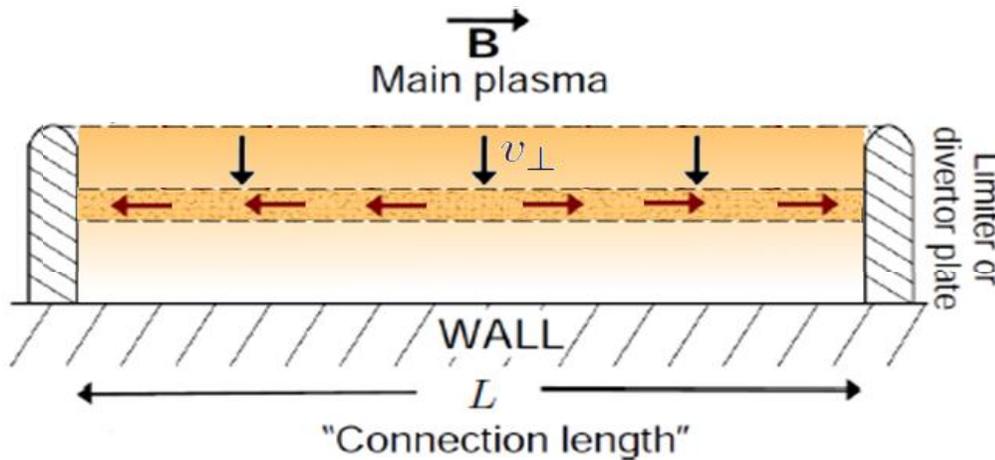




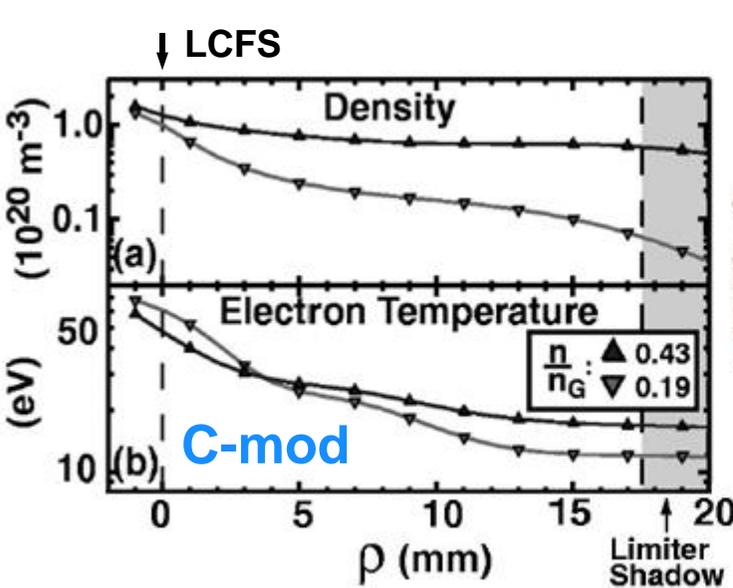
A key question for ITER & DEMO:

How are heat and particle fluxes distributed over the various plasma facing components?

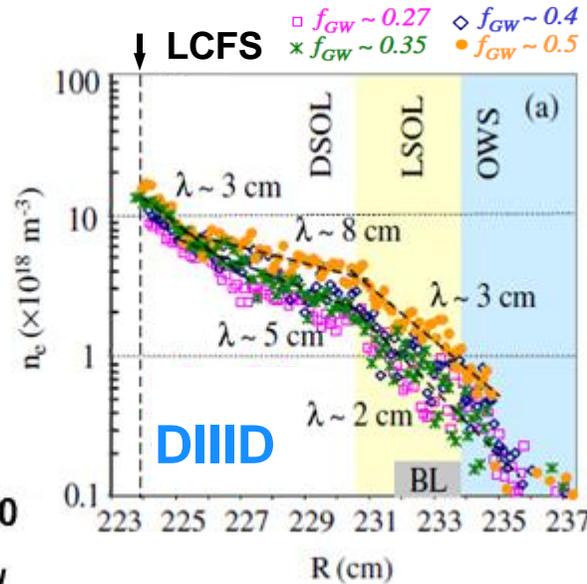
In the SOL of a tokamak, this is strongly affected by the competition between **parallel conduction** and **perpendicular convection**.



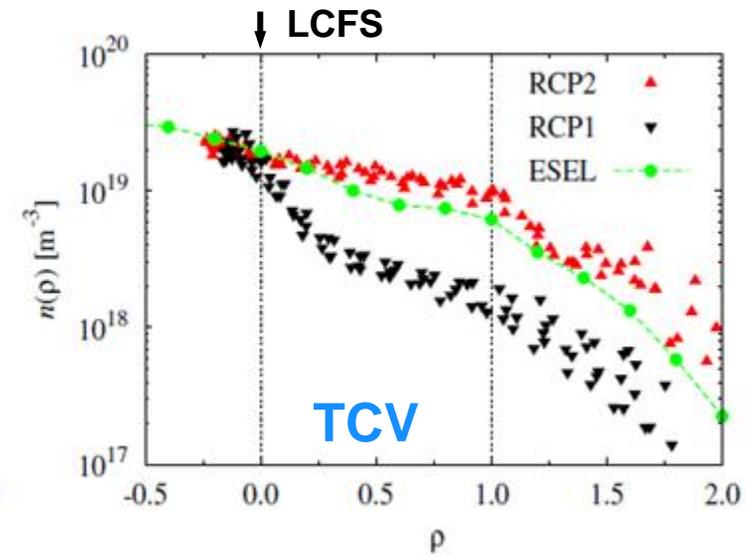
$$\tau_{\parallel} = L/c_s$$
$$\lambda_{SOL} = \frac{v_{\perp}L}{c_s}$$



B. LaBombard, PoP, (2001)



D.L. Rudakov, NF, (2005)

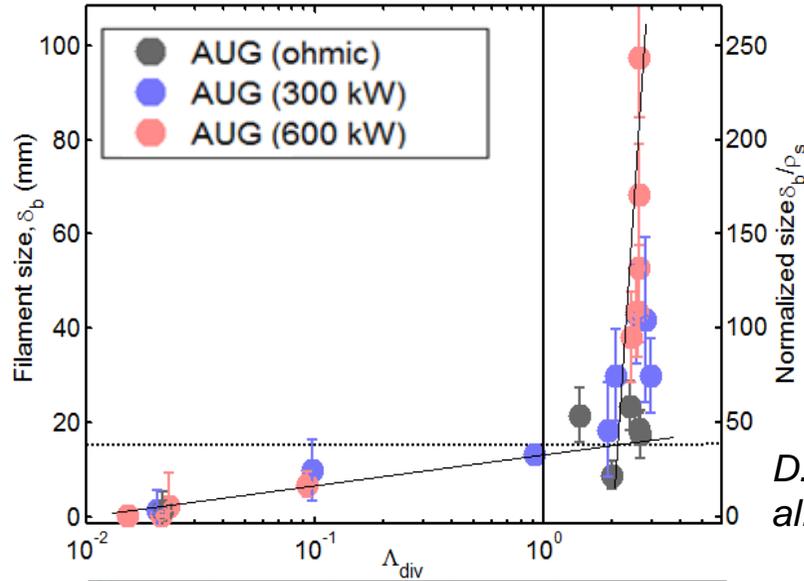


O. E. Garcia et al, JNM, (2007)

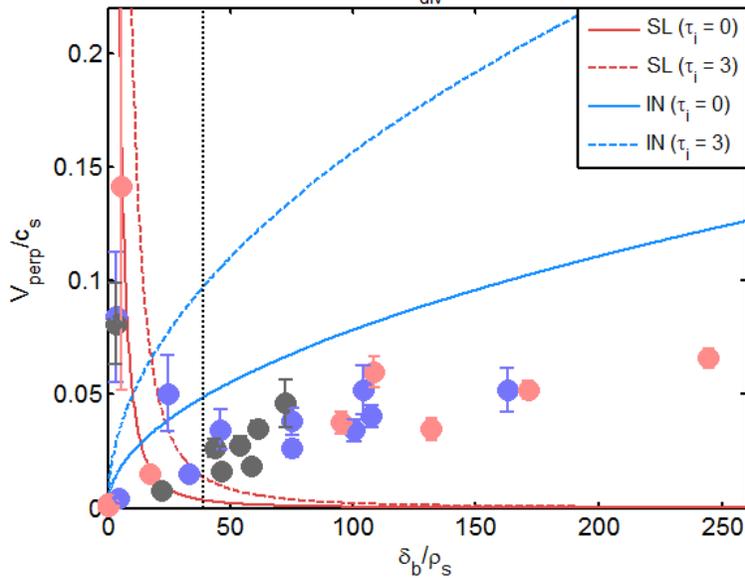
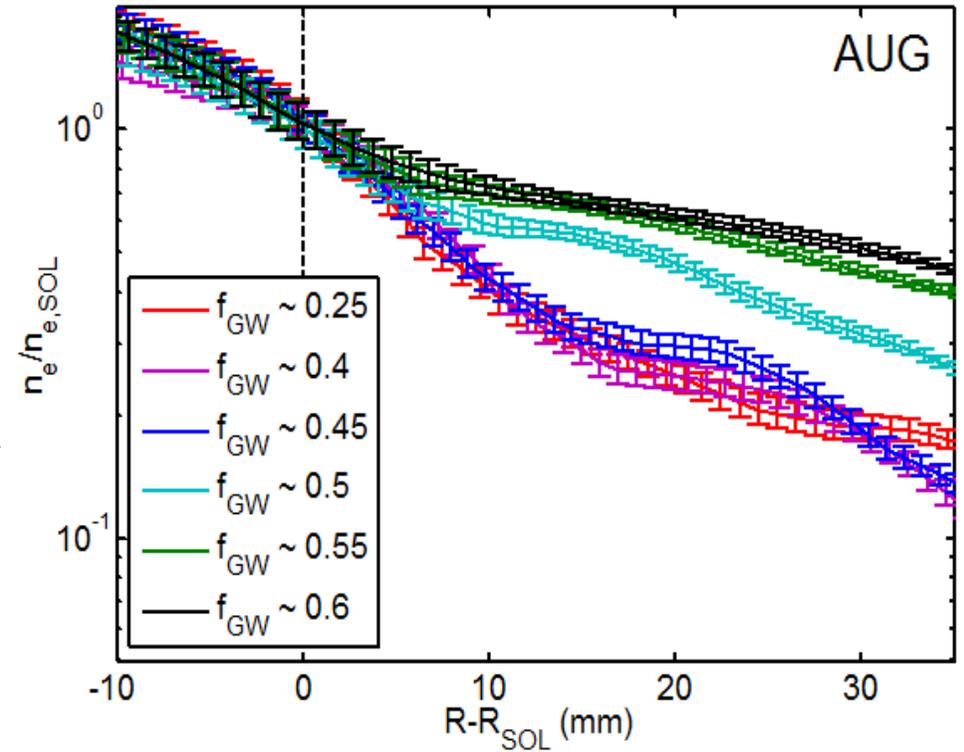
A well known experimental fact in tokamaks:

**In L-mode, SOL density profiles flatten over a certain density threshold.**

An increase of  $\Gamma_r$  associated to a **filament transition** has been proposed as the explanation in the literature.

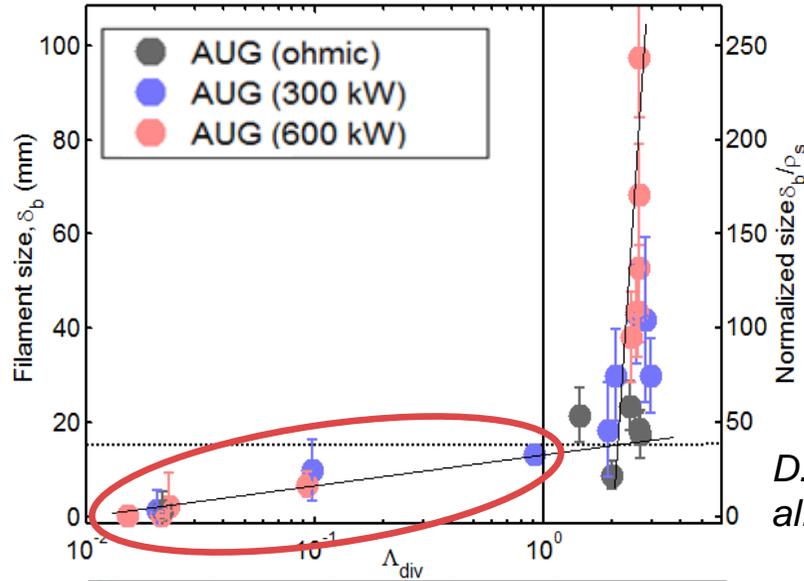


*D. Carralero et al., PRL, (2015)*

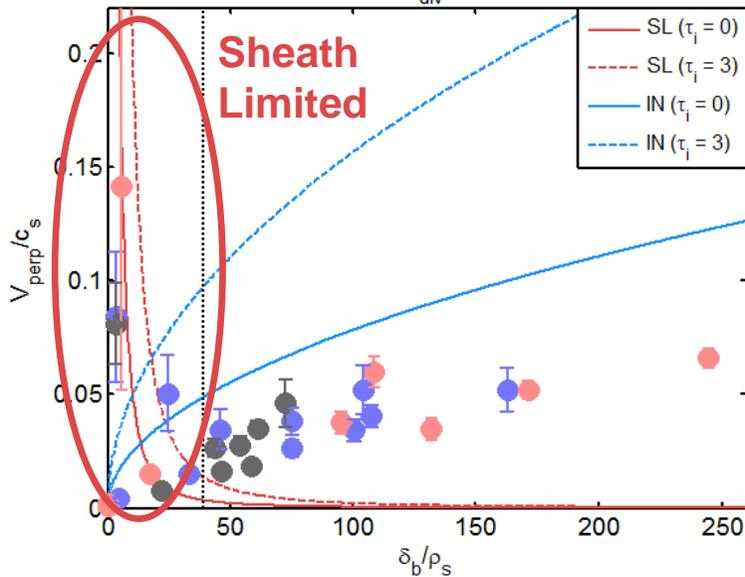
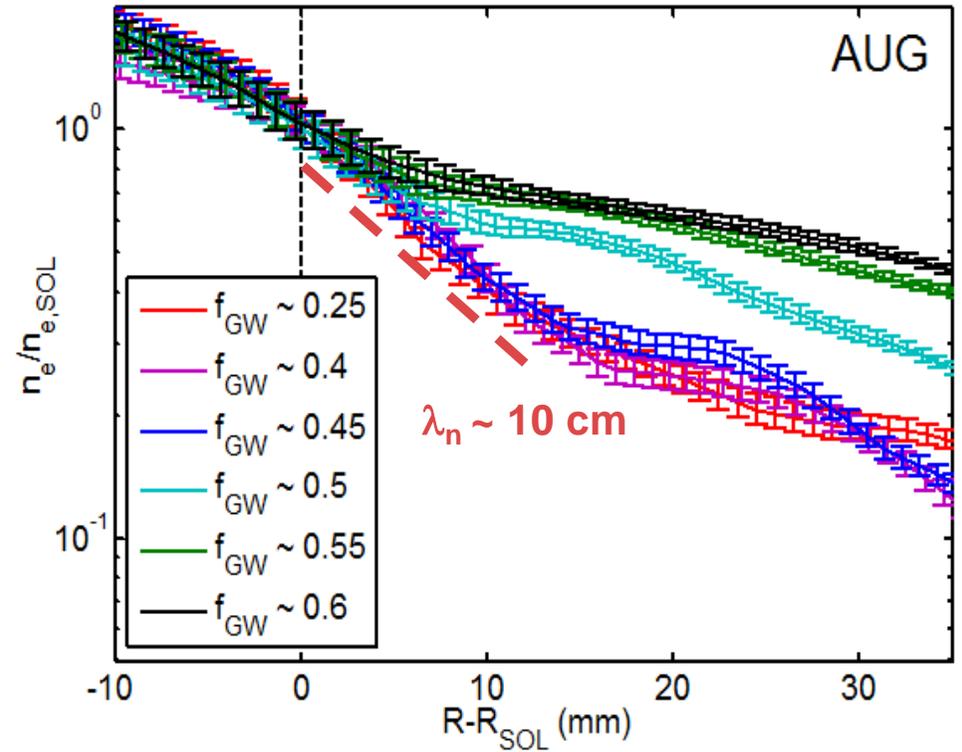


$$\Lambda = \frac{L_{\parallel} / c_s}{1 / v_{ei}} \frac{\Omega_i}{\Omega_e} \simeq \frac{\tau_{\parallel}}{\tau_{ei}}$$

*J.R. Myra et al., PoP, (2006)*



D. Carralero et al., PRL, (2015)

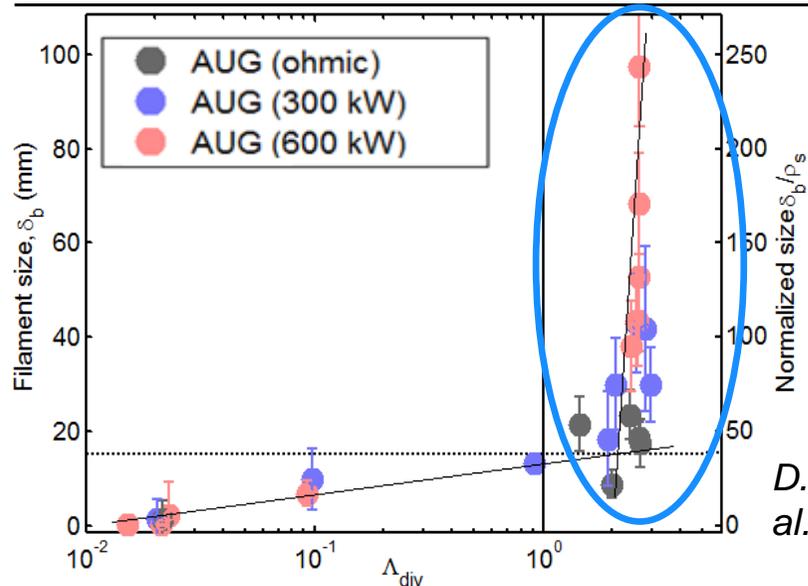


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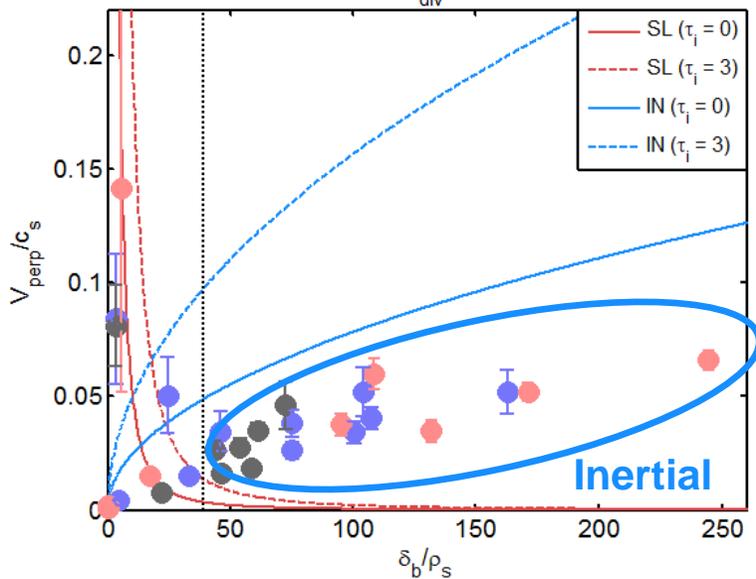
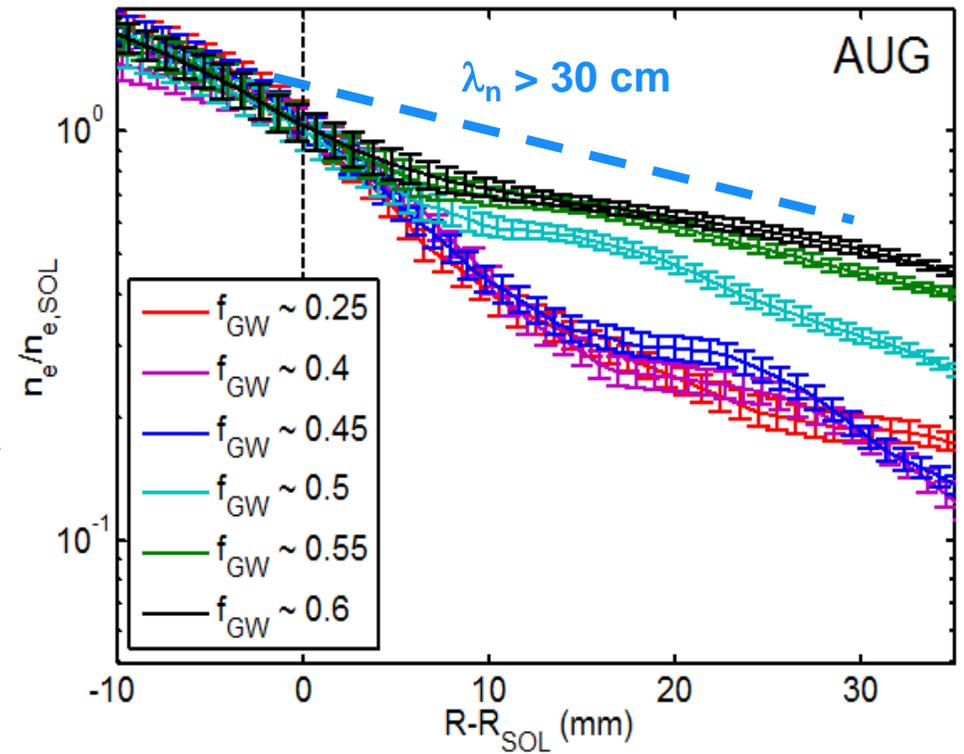
J.R. Myra et al., PoP, (2006)

$$\Lambda < 1$$

$$v_b^{SL} = c_s (1 + \tau_i) \frac{L_{\parallel}}{R} \frac{\tilde{n}}{\bar{n} + \tilde{n}} \left( \frac{\rho_s}{\delta_b} \right)^2$$



*D. Carralero et al., PRL, (2015)*

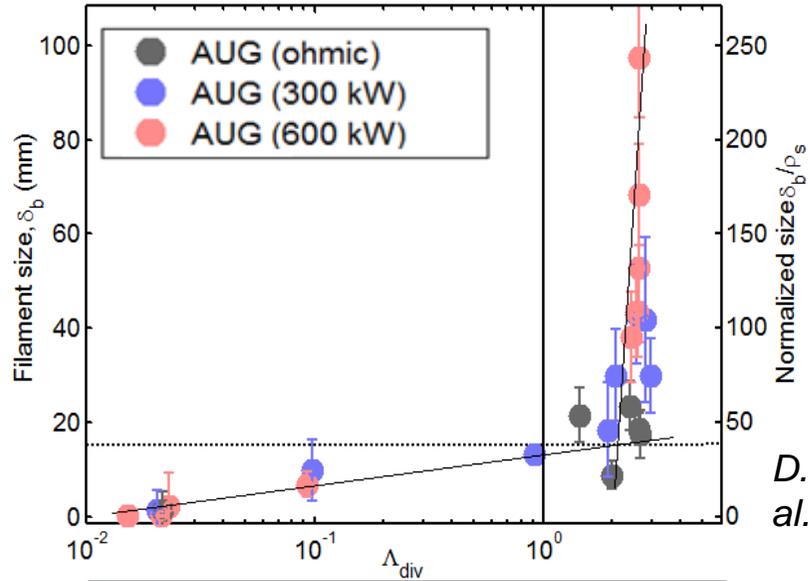


$$\Lambda = \frac{L_{\parallel} / c_s}{1 / v_{ei}} \frac{\Omega_i}{\Omega_e} \simeq \frac{\tau_{\parallel}}{\tau_{ei}}$$

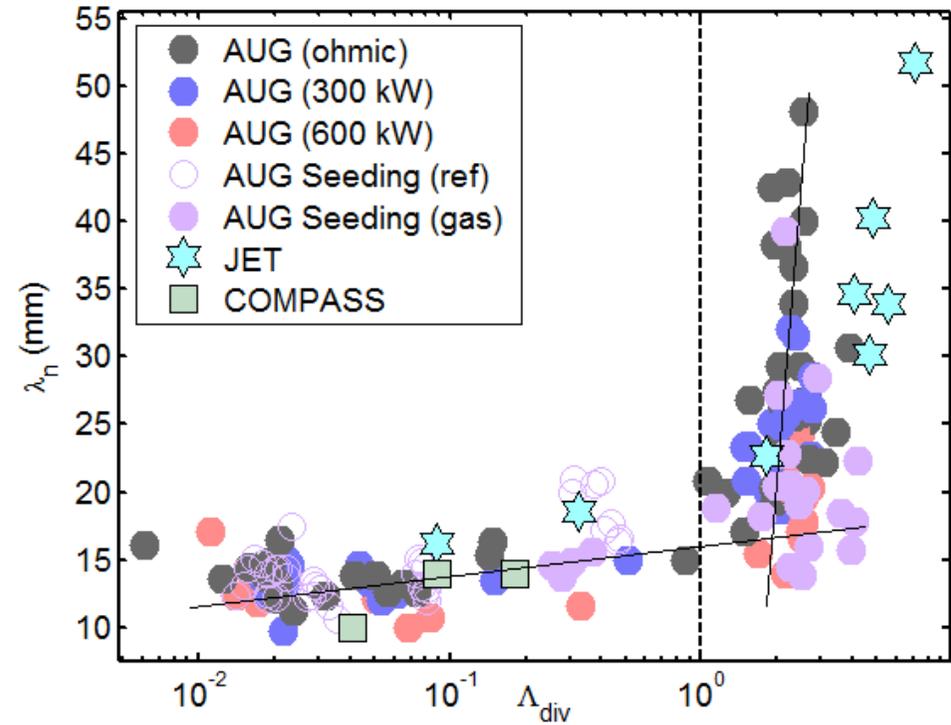
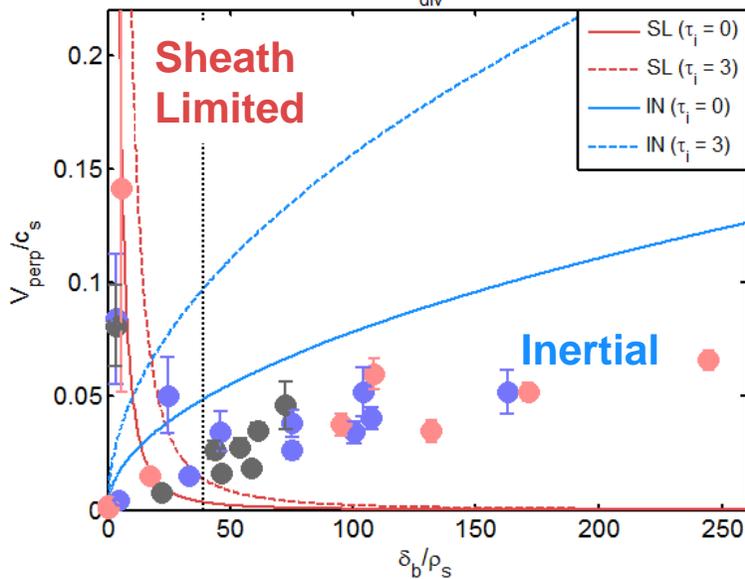
*J.R. Myra et al., PoP, (2006)*

$$\Lambda > 1$$

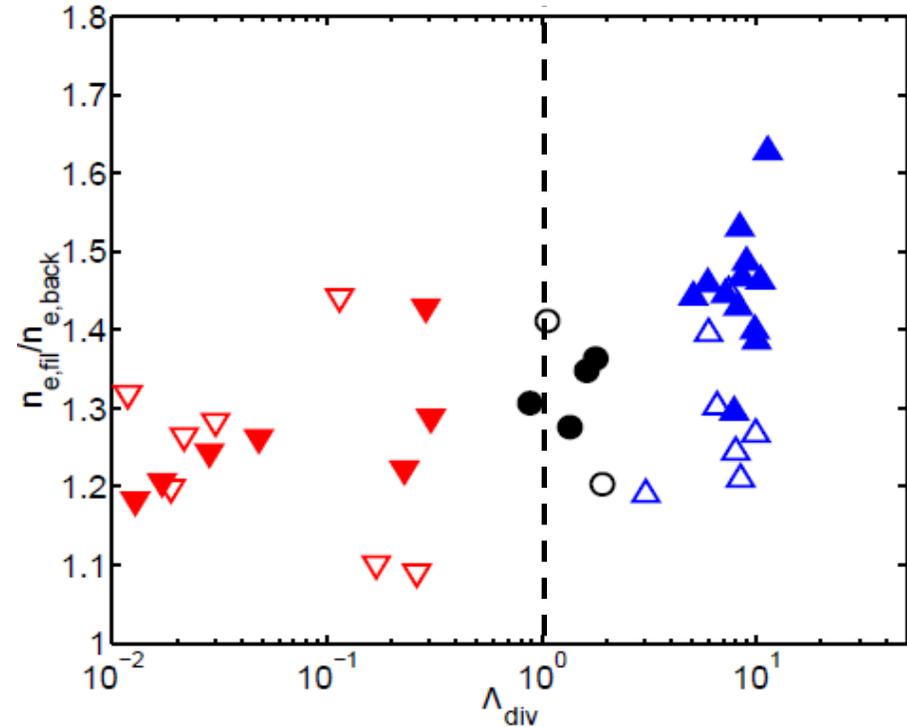
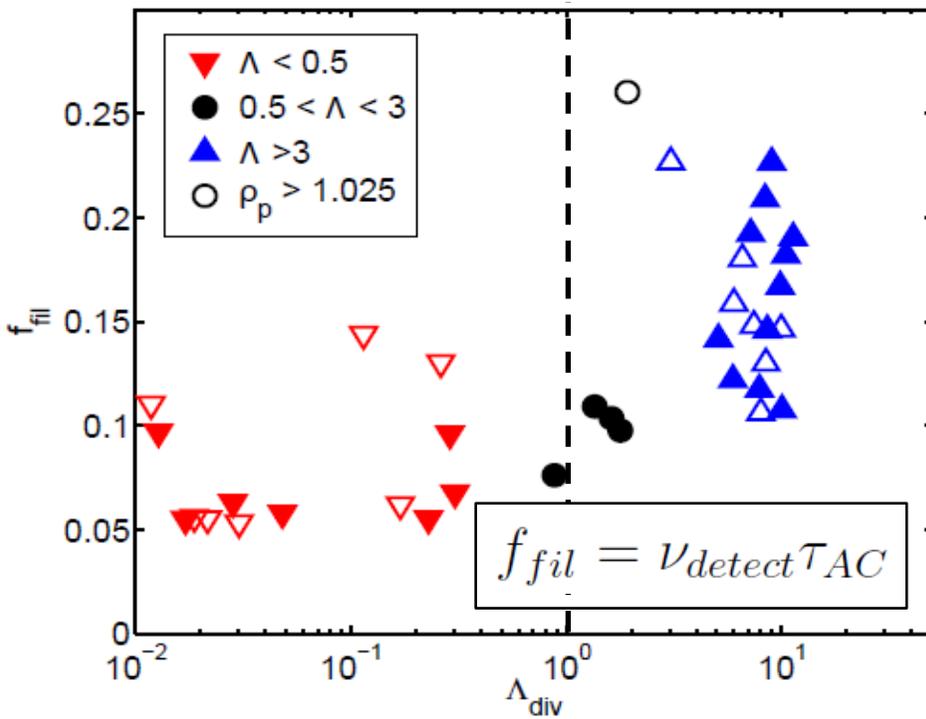
$$v_b^{IN} = c_s \sqrt{(1 + \tau_i) \frac{\tilde{n}}{\bar{n} + \tilde{n}} \frac{\delta_b}{R}}$$



*D. Carralero et al., PRL, (2015)*

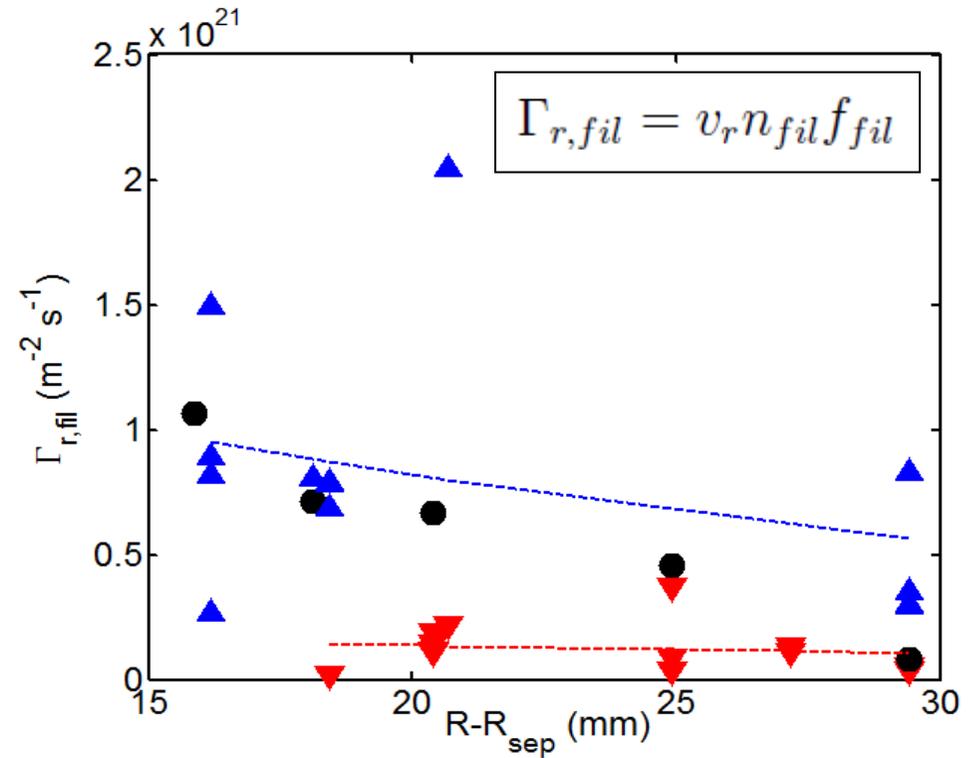
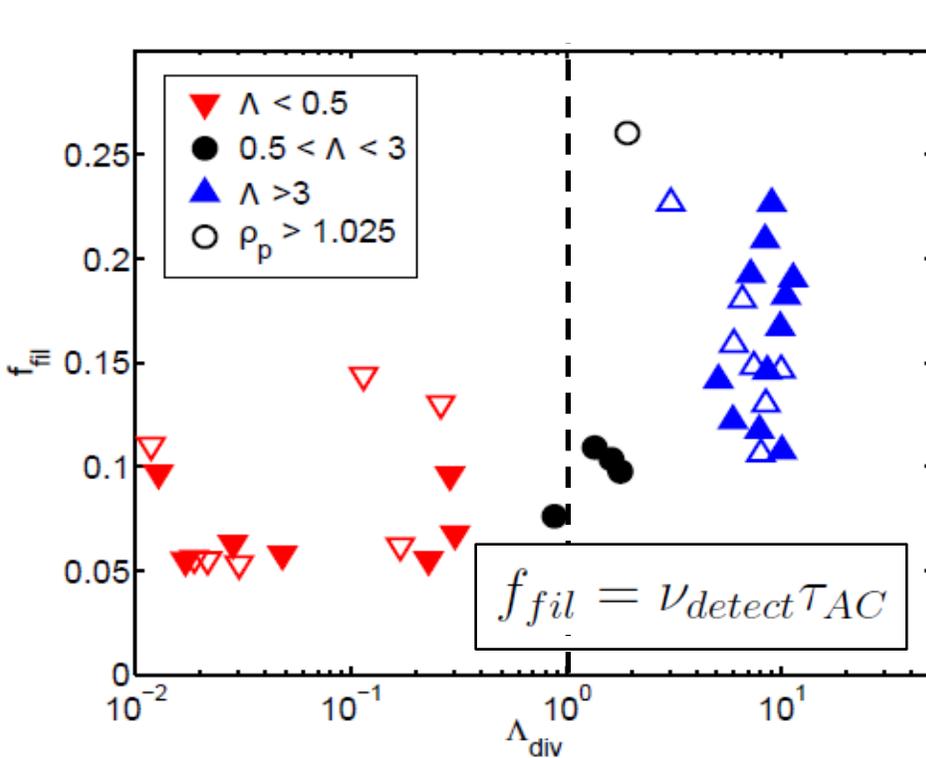


$\Lambda_{div}$  scaling of the shoulder formation has been demonstrated on the COMPASS-AUG-JET ITER Stepladder



The filament transition has clear effects on filamentary transport:

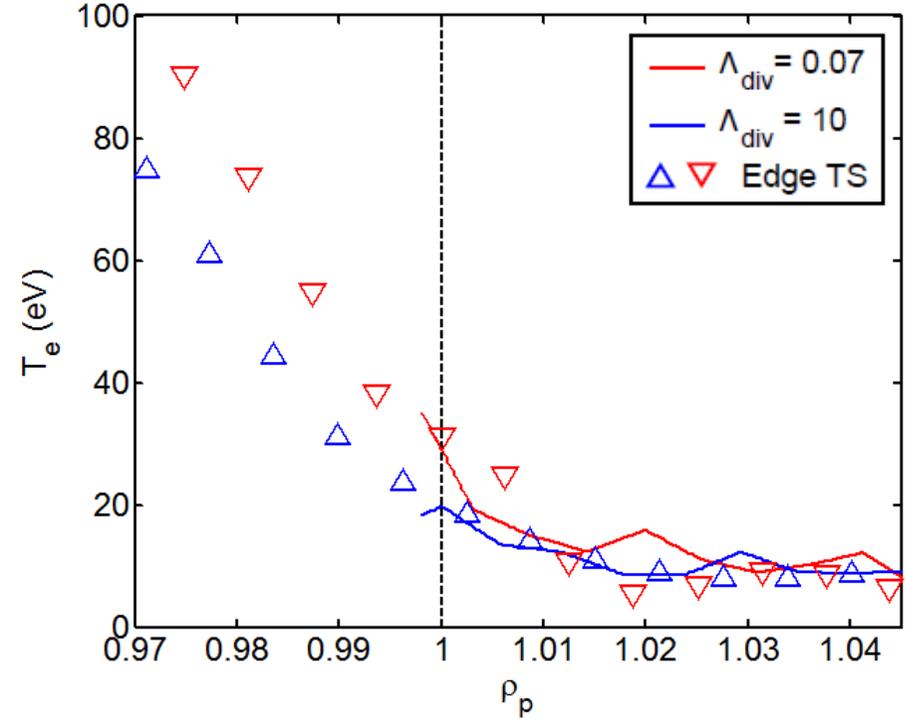
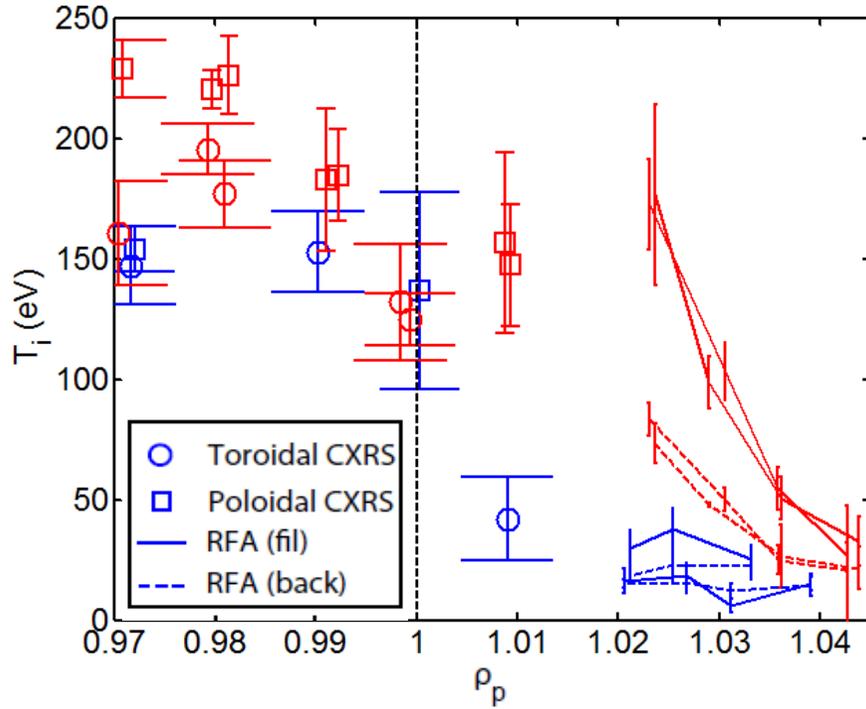
- ▶ Packing fraction,  $f_{fil}$  and amplitude of fluctuations  $n_{fil}/n_{back}$  **substantially increased**



The filament transition has clear effects on filamentary transport:

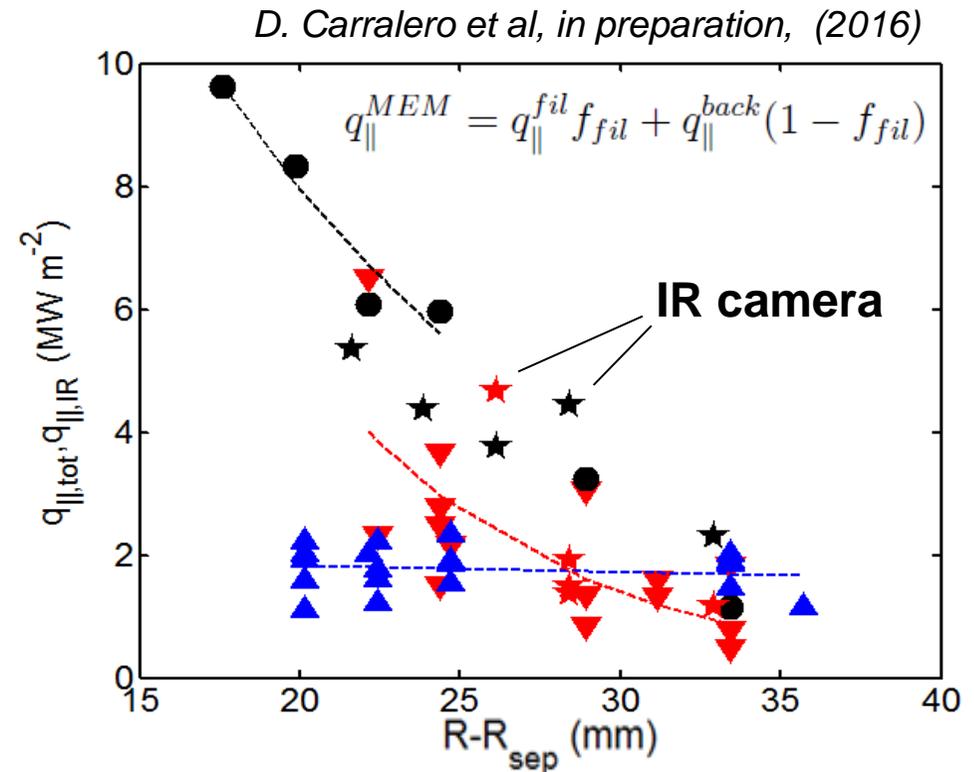
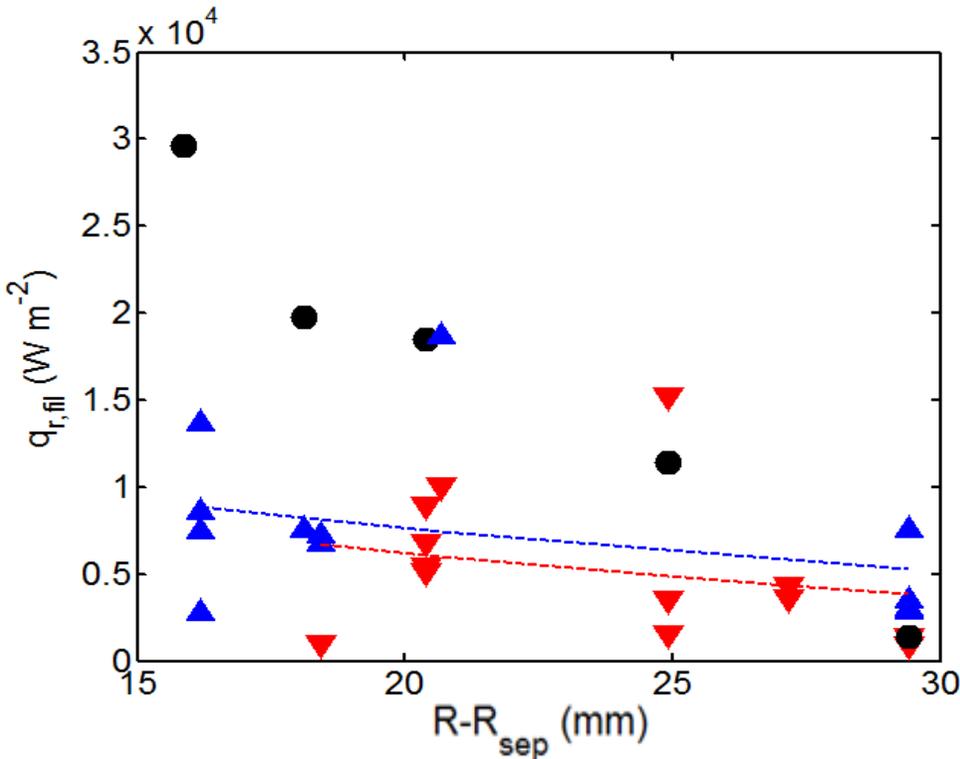
- ▶ Packing fraction,  $f_{fil}$  and amplitude of fluctuations  $n_{fil}/n_{back}$  **substantially increased**
- ▶  $\Gamma_r$  across the separatrix **is increased by a factor 3-4**
- ▶ Similar results reported from JET (Guillemaut PSI 2016)

*D. Carralero et al, EPS, (2015)*



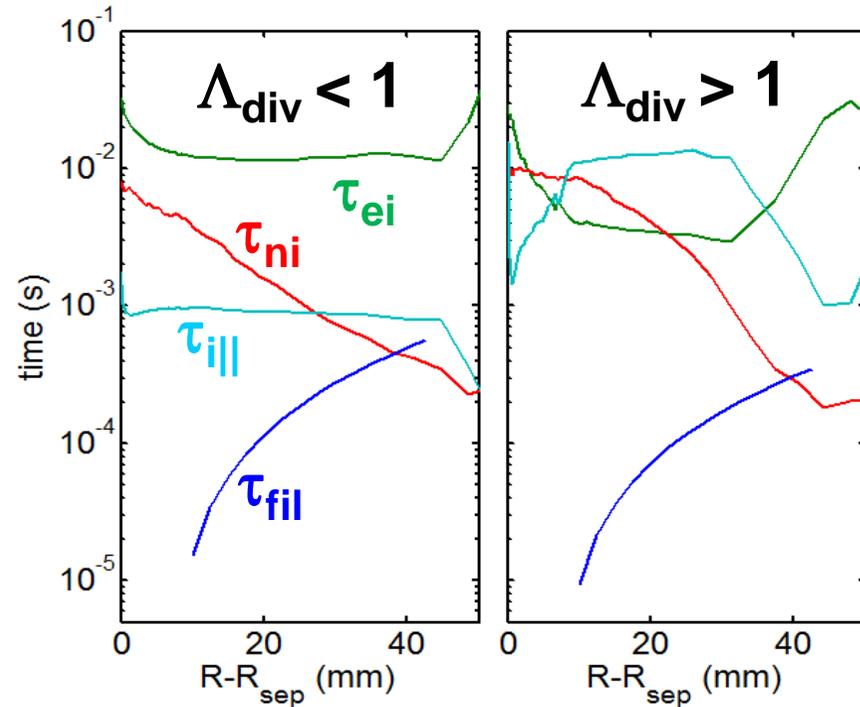
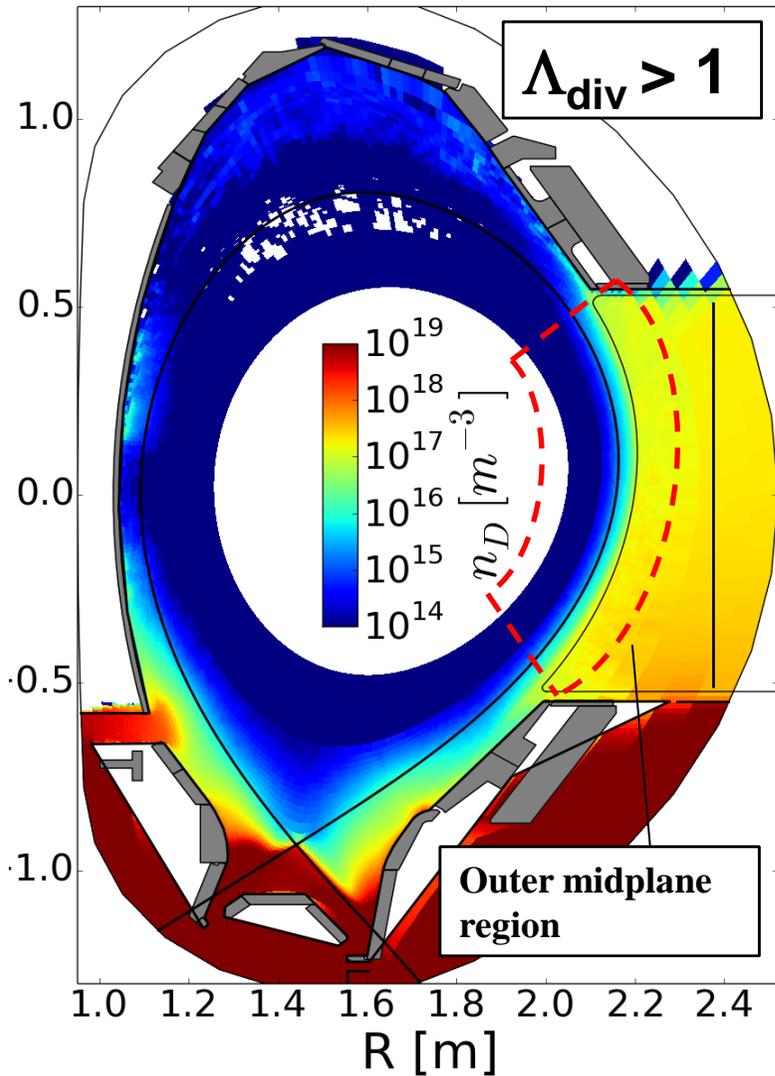
**Electrons and ions show very different behaviour** through the transition:

- ▶  $T_{e,fil} \sim 1.2 T_{e,back}$ ,  $T_e$  roughly constant across the SOL.
- ▶  $T_{i,fil} > T_{i,back}$  for  $\Lambda_{div} < 1$ . Slow radial decay,  $\lambda_{Ti} \sim 40$  mm.
- ▶  $T_{i,fil} \sim T_{i,back} \sim 25$  eV for high  $\Lambda_{div} > 1$ . Fast radial decay,  $\lambda_{Ti} \sim 10$  mm.



## Heat transport is not affected by the transition

- ▶  $q_{r,fil}$  near the wall are similar. Agreement with JET (Guillemaut, PSI 2016)
- ▶ A maximum in  $q_{r,fil}$  is reached around the transition.
- ▶ Good agreement with IR measurements of the  $q_{||}$  at the manipulator.



$$\tau_{ie}^{-1} \simeq \frac{m_e \nu_{ie}}{m_i} \frac{T_i - T_e}{T_i}$$

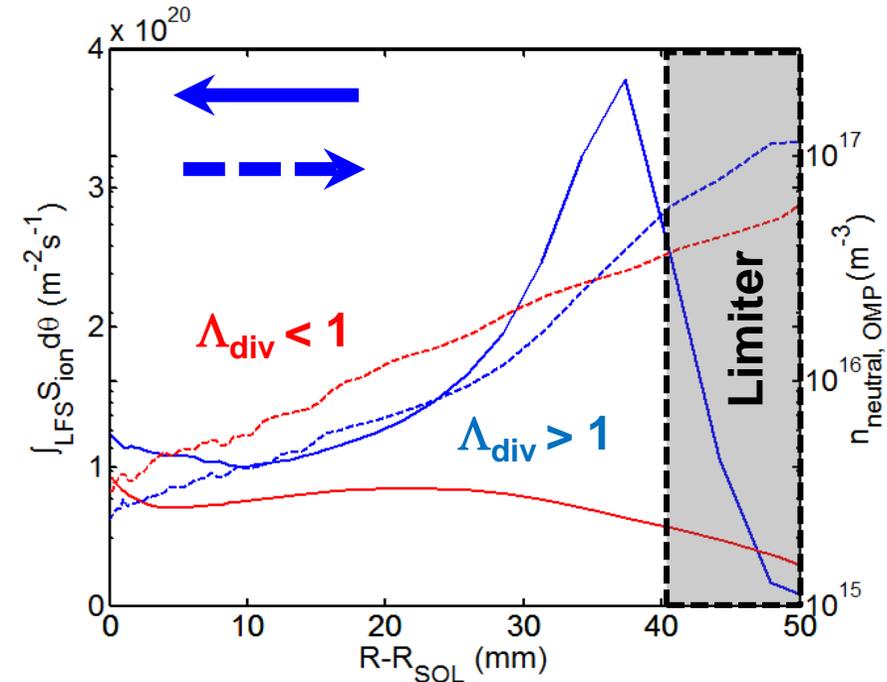
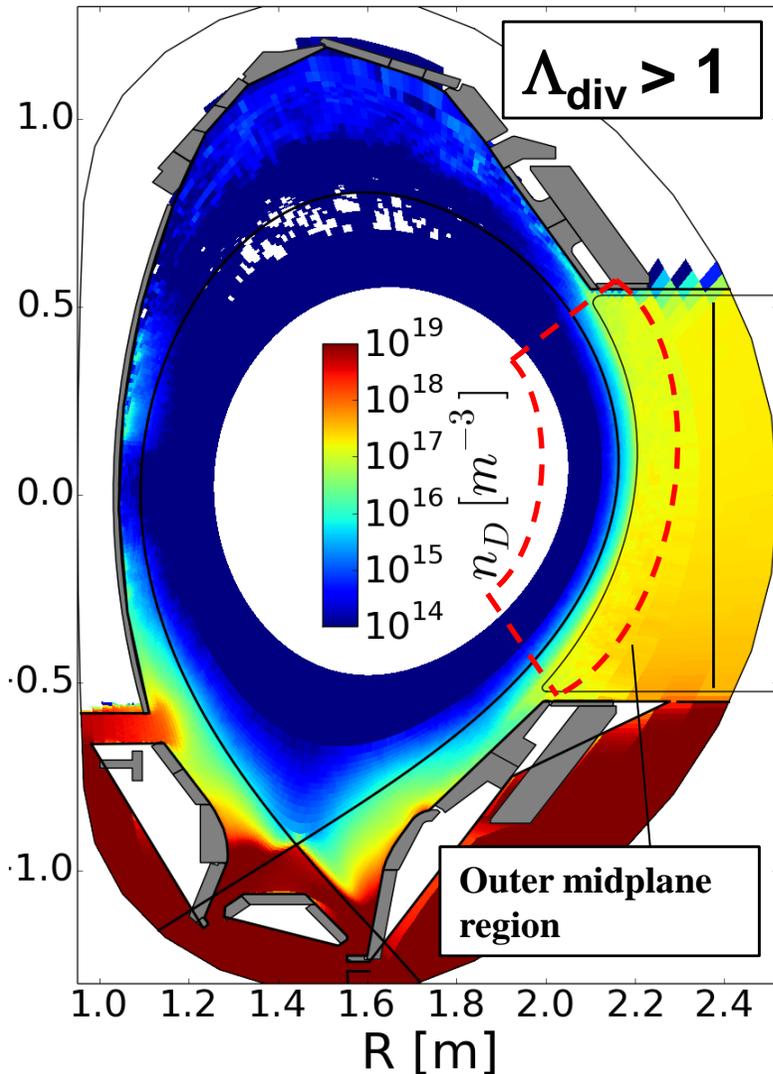
$$\tau_{in}^{-1} \simeq \nu_{in} \frac{T_i - T_n}{T_n}$$

$$\tau_{i||}^{-1} \simeq \frac{T_i}{m_i \nu_{ii}} \frac{1}{L_{||}^2}$$

$$\tau_{fil}^{-1} \simeq \frac{R - R_{SOL}}{v_{\perp, fil}}$$

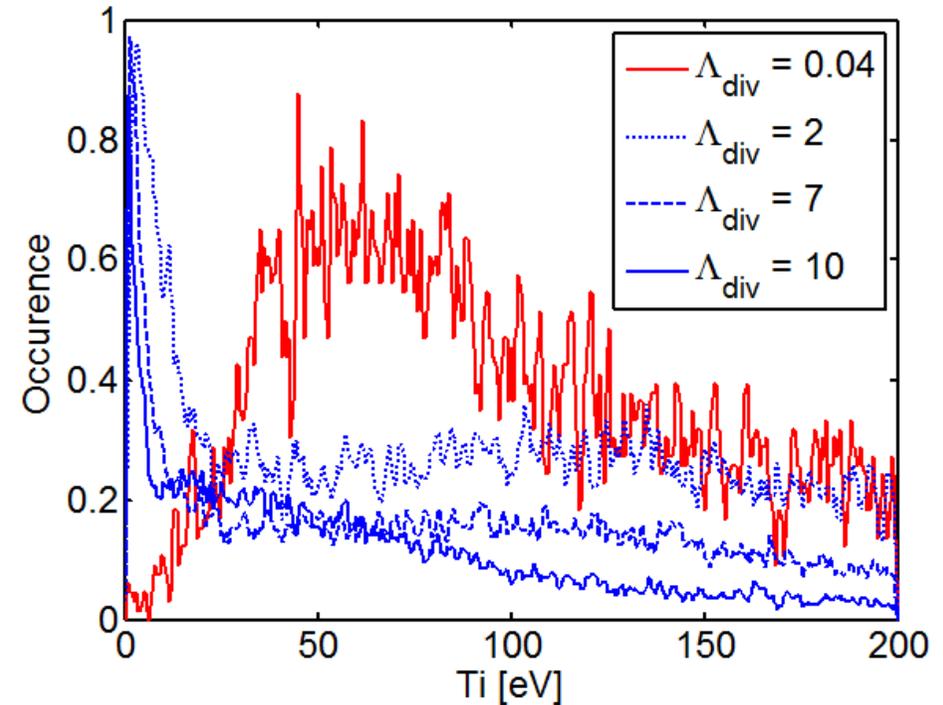
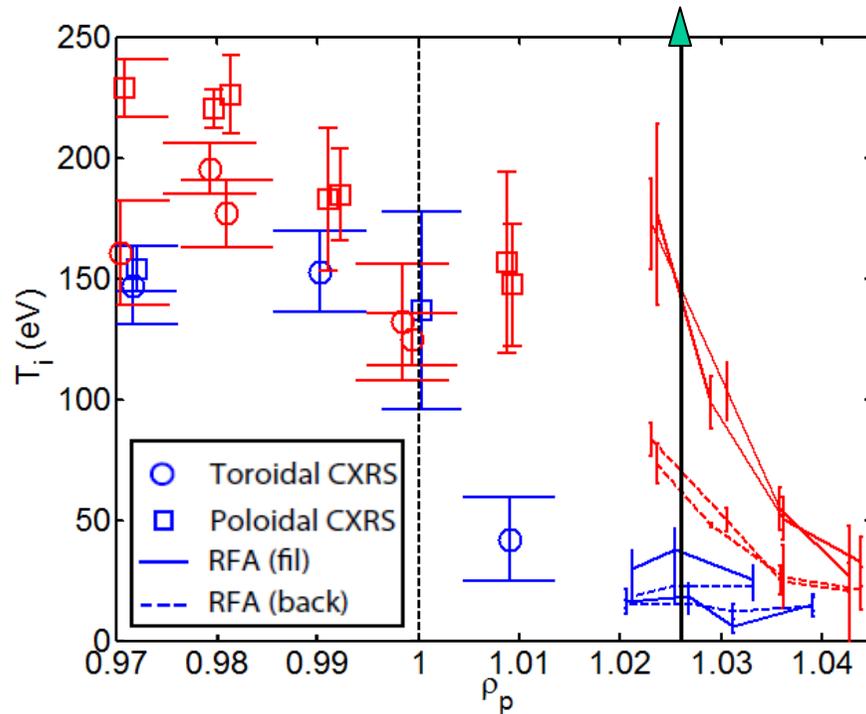
EMC3-EIRENE code is used to simulate  $\Lambda_{div} > 1$  and  $\Lambda_{div} < 1$  scenarios

- ▶ **No thermalization** mechanism can reduce  $T_i$  as observed in the experiment.



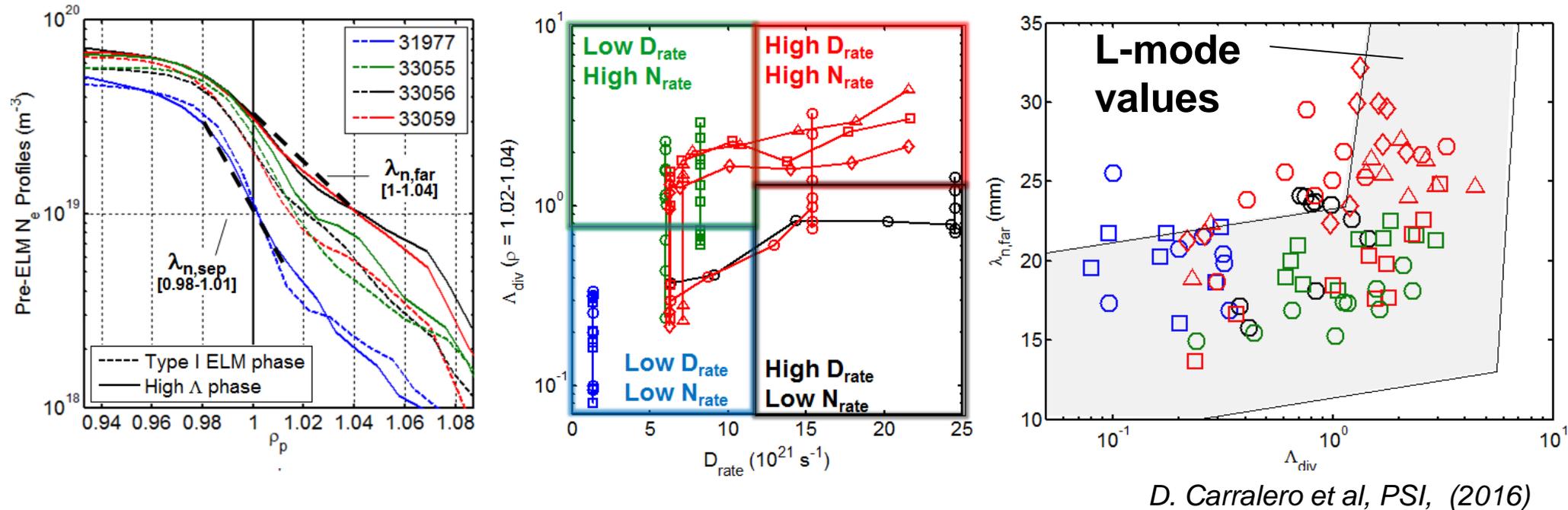
EMC3-EIRENE code is used to simulate  $\Lambda_{\text{div}} > 1$  and  $\Lambda_{\text{div}} < 1$  scenarios

- ▶ **No thermalization** mechanism can reduce  $T_i$  as observed in the experiment.
- ▶ An **ionization front** builds **in front of the limiter** shadow in the  $\Lambda_{\text{div}} > 1$  case.



ExB analyzer experiments are consistent with RFA measurements :

- ▶  $\Lambda_{\text{div}} < 1 \rightarrow$  **Monoenergetic distribution** with a positive tail, consistent with  $T_{i,\text{back}}$  and  $T_{i,\text{fil}}$ .
- ▶  $\Lambda_{\text{div}} > 1 \rightarrow$  **Two-energies distribution** with cold ions around the F.C.  $T_{i,\text{fil}}$  not necessarily greater than  $T_{i,\text{back}}$ .



Equivalent experiments have analyzed the shoulder formation on inter-ELM H-mode plasma:

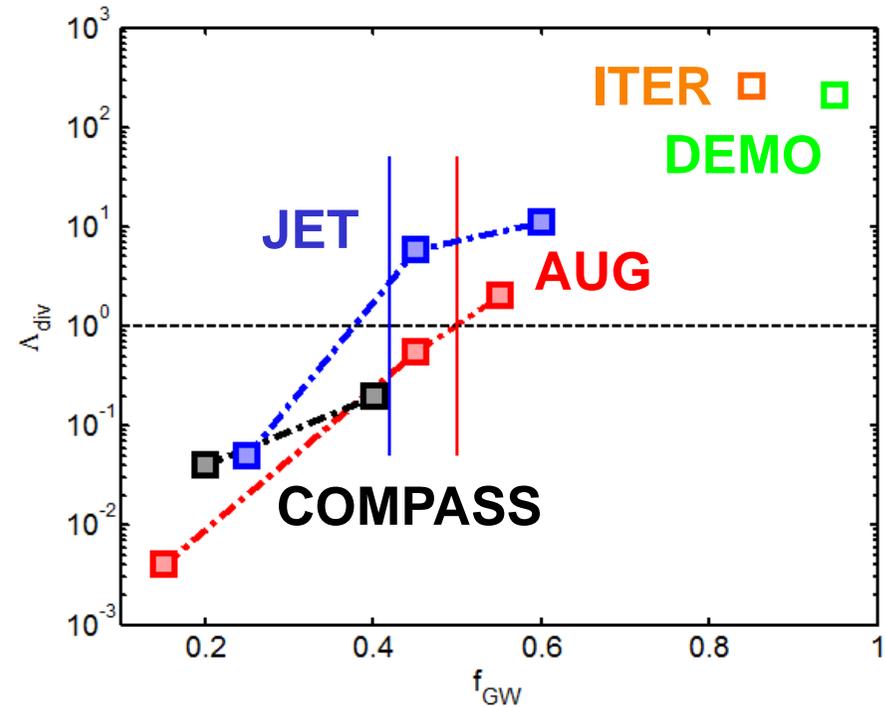
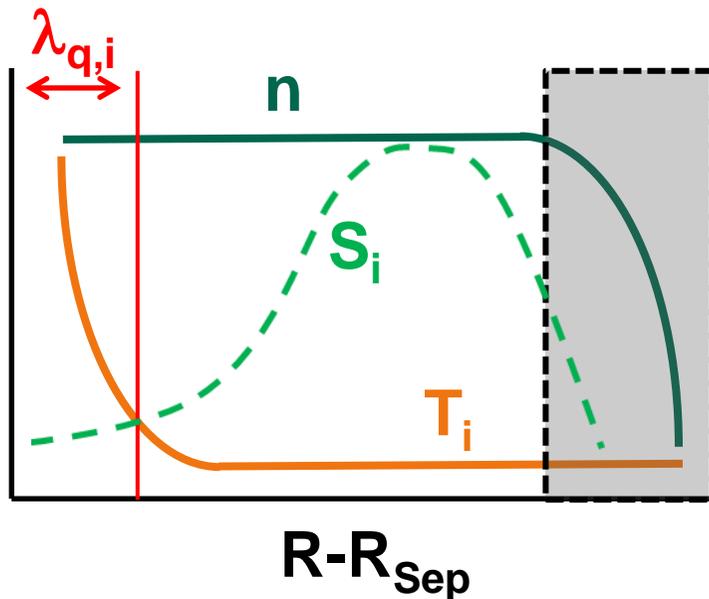
- ▶ **A similar shoulder** has been observed.
- ▶ An **equivalent filament transition** takes place for  $\Lambda_{div} > 1$
- ▶  $\Lambda_{div} > 1$  is a necessary but not sufficient condition. **A minimum level of D fueling is also required.** This is consistent with L-mode experiments in TCV [N. Vianello, this conference, EX/P8-26].

Several scenarios can be proposed for DEMO:

**Standard** (Extrapolation of current results):

$$\lambda_{Ti} \sim 10 \text{ mm}, \lambda_n \sim 40 \text{ mm};$$

$$1/\lambda_{q,i} \sim 1/\lambda_{Ti} + 1/\lambda_n \sim 10 \text{ mm}$$



$\Lambda_{div} > 1$  ▶ Shoulder formation

$\lambda_i \sim R_w - R_{sep}$  ▶  $\lambda_{Ti} \ll R_w - R_{sep}$

$\lambda_{q,i} \ll R_w - R_{sep}$  ▶ High divertor load

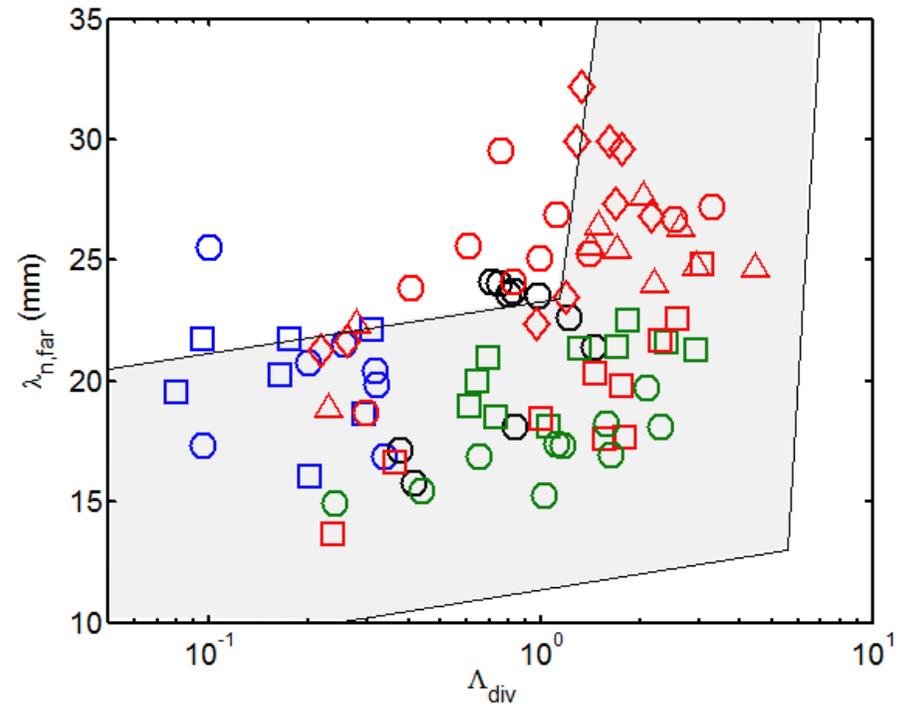
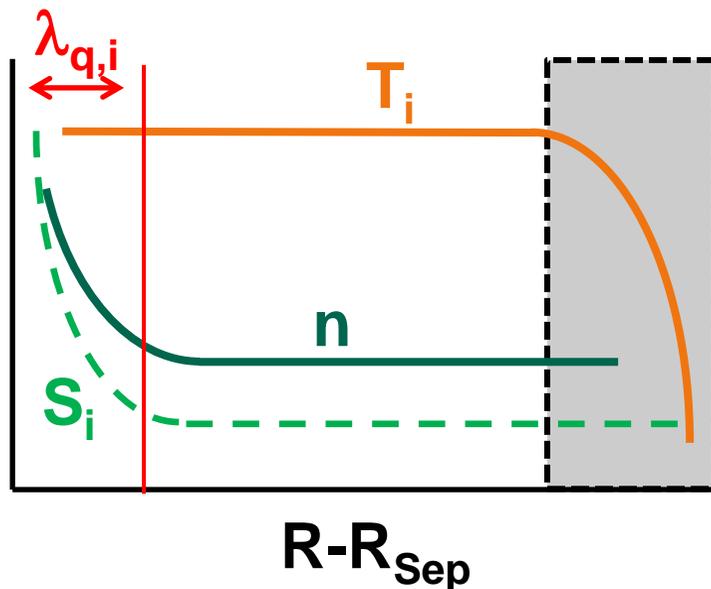
$T_{i,wall} \sim 10\text{-}20 \text{ eV}$  ▶ Low wall sputtering

Several scenarios can be proposed for DEMO:

**Worst case** (No shoulder formation):

$$\lambda_{Ti} \sim 40 \text{ mm}, \lambda_n \sim 10 \text{ mm};$$

$$1/\lambda_{q,i} \sim 1/\lambda_{Ti} + 1/\lambda_n \sim 10 \text{ mm}$$



Insufficient  $n_n$  in the SOL

$$\lambda_i \gg R_w - R_{Sep} \rightarrow \lambda_{Ti} \sim R_w - R_{Sep}$$

$$\lambda_{q,i} \ll R_w - R_{Sep} \rightarrow \text{High divertor load}$$

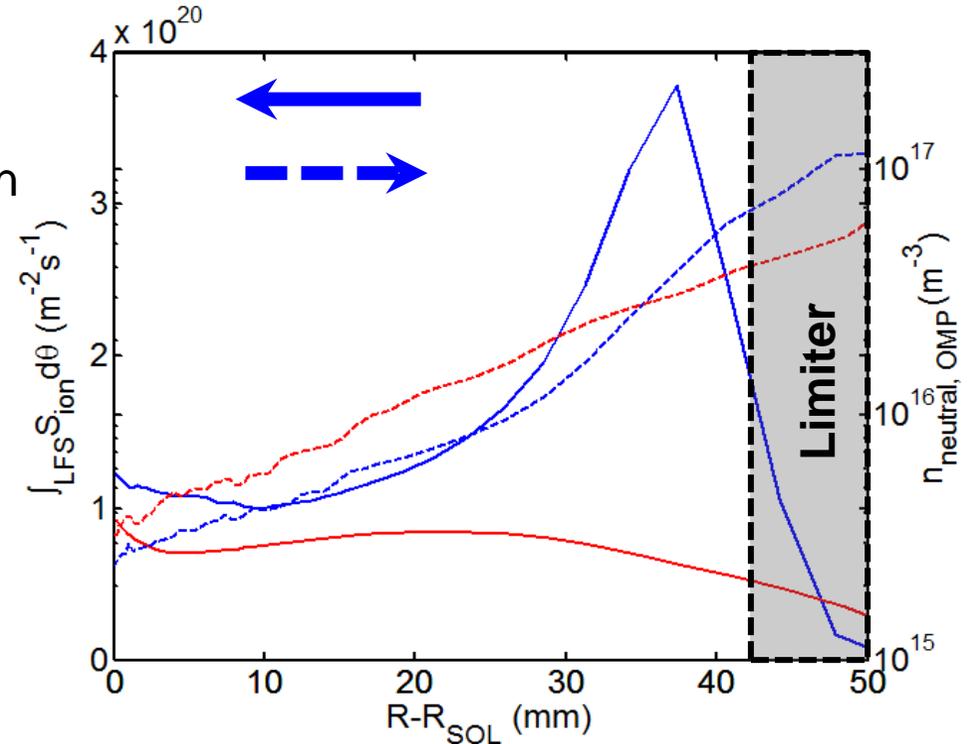
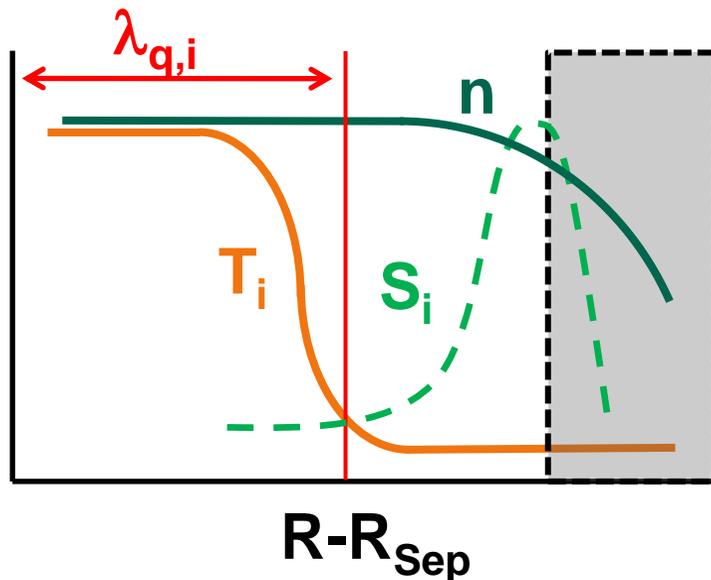
$$T_{i,wall} \gg 10\text{-}20 \text{ eV} \rightarrow \text{High wall sputtering}$$

Several scenarios can be proposed for DEMO:

**Best case** (independent shoulder formation/ion cooling, widened near-SOL):

$$\lambda_{Ti} \sim 40 \text{ mm}, \lambda_n \sim 40 \text{ mm};$$

$$1/\lambda_{q,i} \sim 1/\lambda_{Ti} + 1/\lambda_n \sim 40 \text{ mm}$$



$\Lambda_{div} > 1$  ▶ Shoulder formation

$\lambda_i < R_w - R_{Sep}$  ▶  $\lambda_{Ti} < R_w - R_{Sep}$

$\lambda_{q,i} \sim R_w - R_{Sep}$  ▶ Low divertor load

$T_{i,wall} \sim 10-20 \text{ eV}$  ▶ Low wall sputtering

L-mode experiments on AUG, JET and COMPASS have shown the **relation between a filament transition and the shoulder formation**. This process is triggered by the  $\Lambda_{\text{div}} > 1$  threshold.

The transition increases  $\Gamma_{\text{perp}}$  **by a factor 3** after the transition, while  $q_{\text{perp}}$  **remains roughly constant** due to the drop in  $T_i$ .

Simulations indicate that **ion cooling is not the result of thermalization**, but of the ionization of cold, recycled neutrals in front of the limiter. This is consistent with the **observation of a cold ion population for  $\Lambda_{\text{div}} > 1$** .

H-mode experiments indicate **that a shoulder can form also between ELMs**.  $\Lambda_{\text{div}} > 1$  is a necessary condition, but **sufficient D fueling level is also required**.

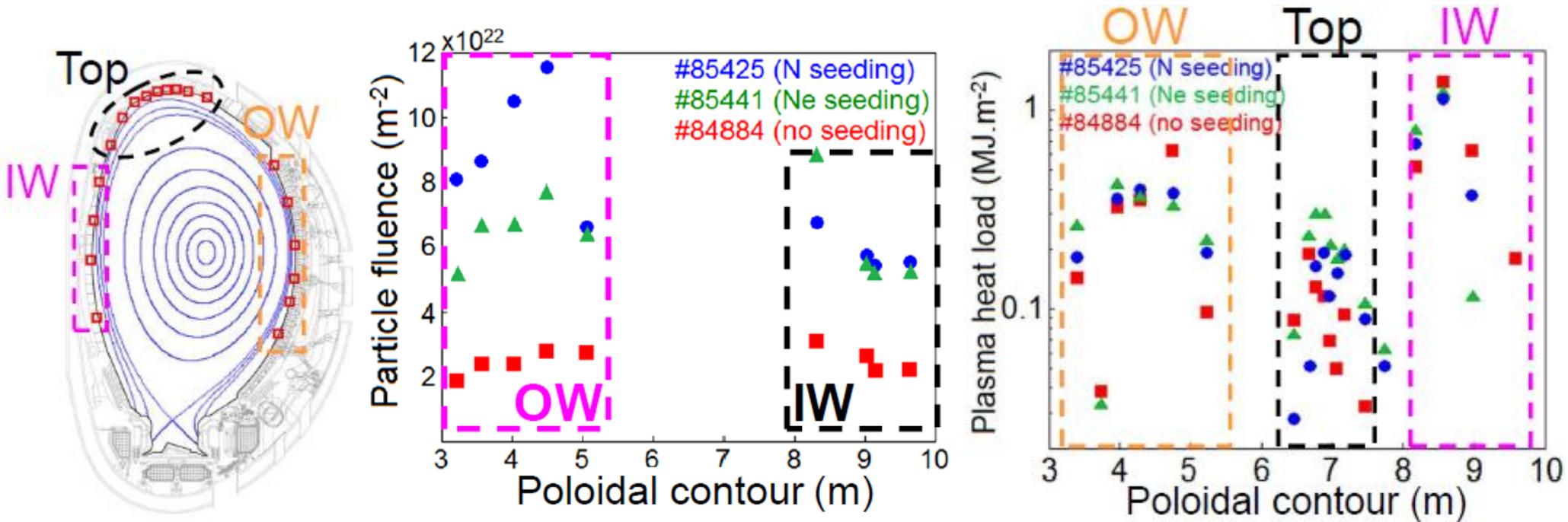
**The shoulder is probably the result of a feedback loop** between increased  $\Gamma_{\text{perp}}$ , wall recycling and ionization of reflected neutrals, leading to ever increased transport.

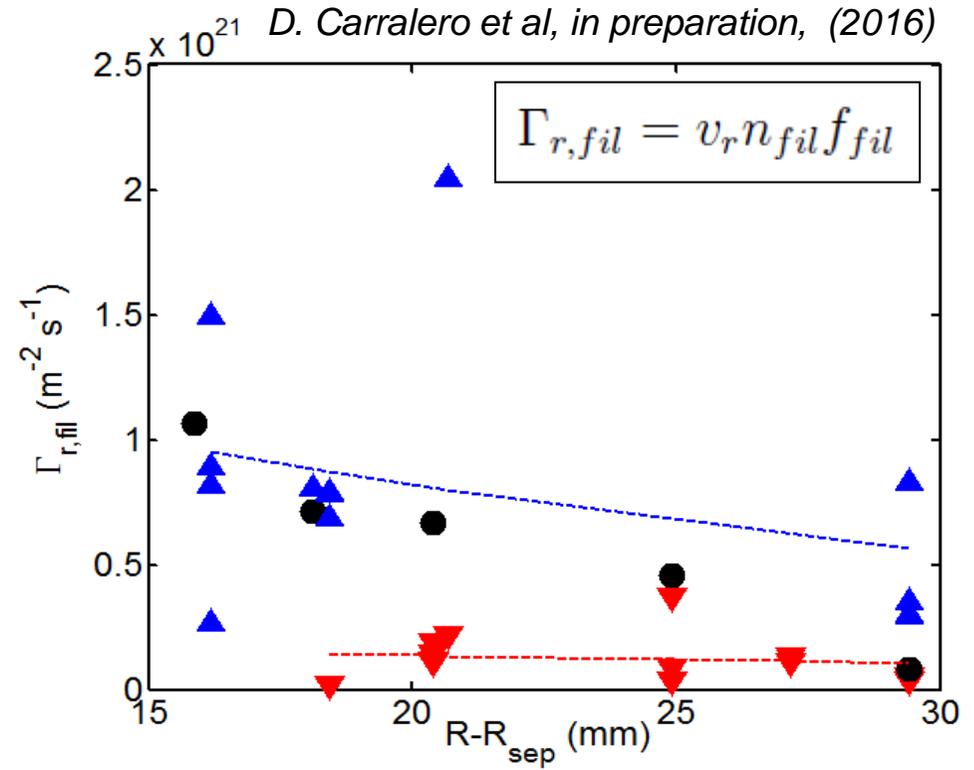
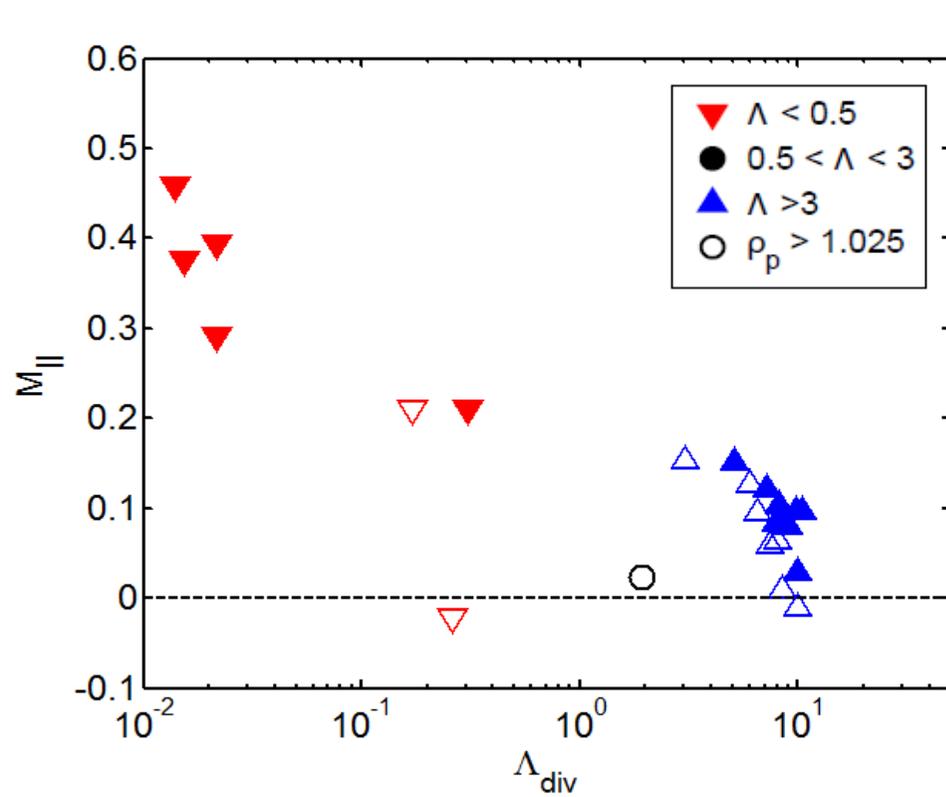
**$\lambda_{T_e}$  and  $\lambda_{T_i}$  are decoupled**, which could lead to enhanced  $\lambda_q$  under certain conditions in DEMO. Also, there is a risk of **high  $T_i$   $D^+$  ions arriving to the first wall**.

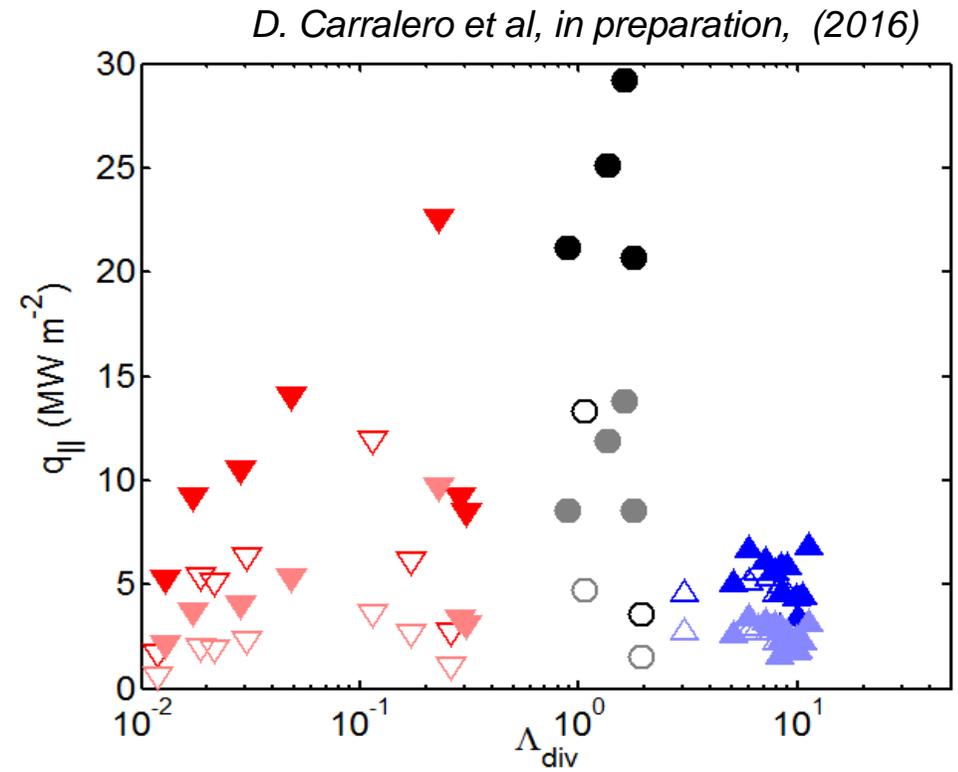
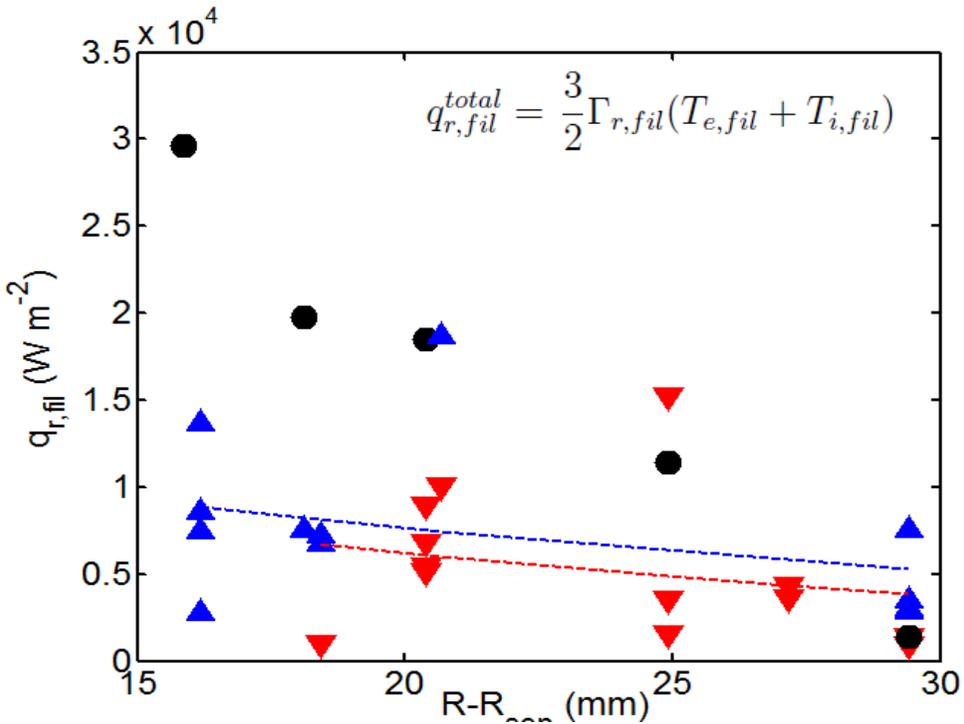
## Additional Slides



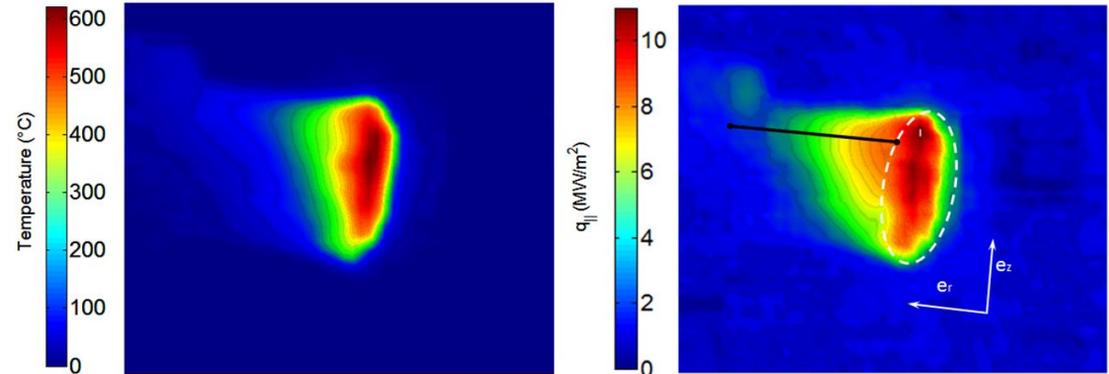
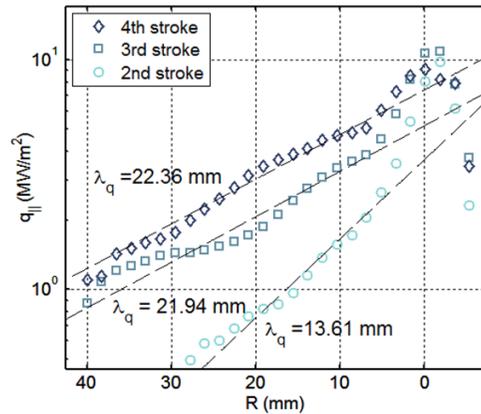
C. Guillemaut, et al., PSI Rome, (2016)





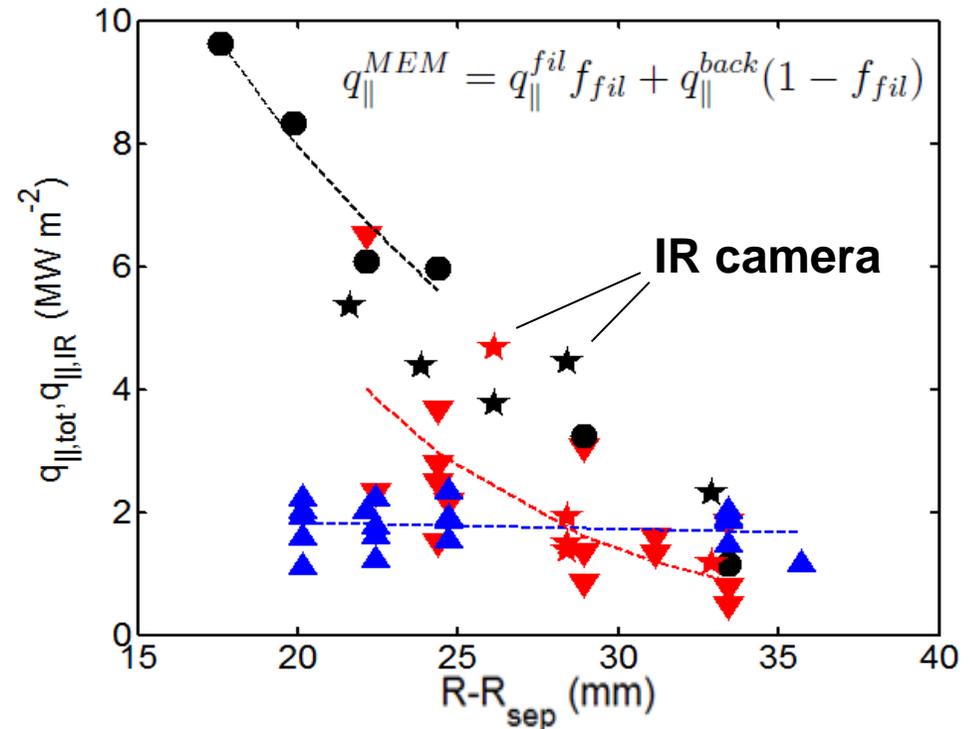


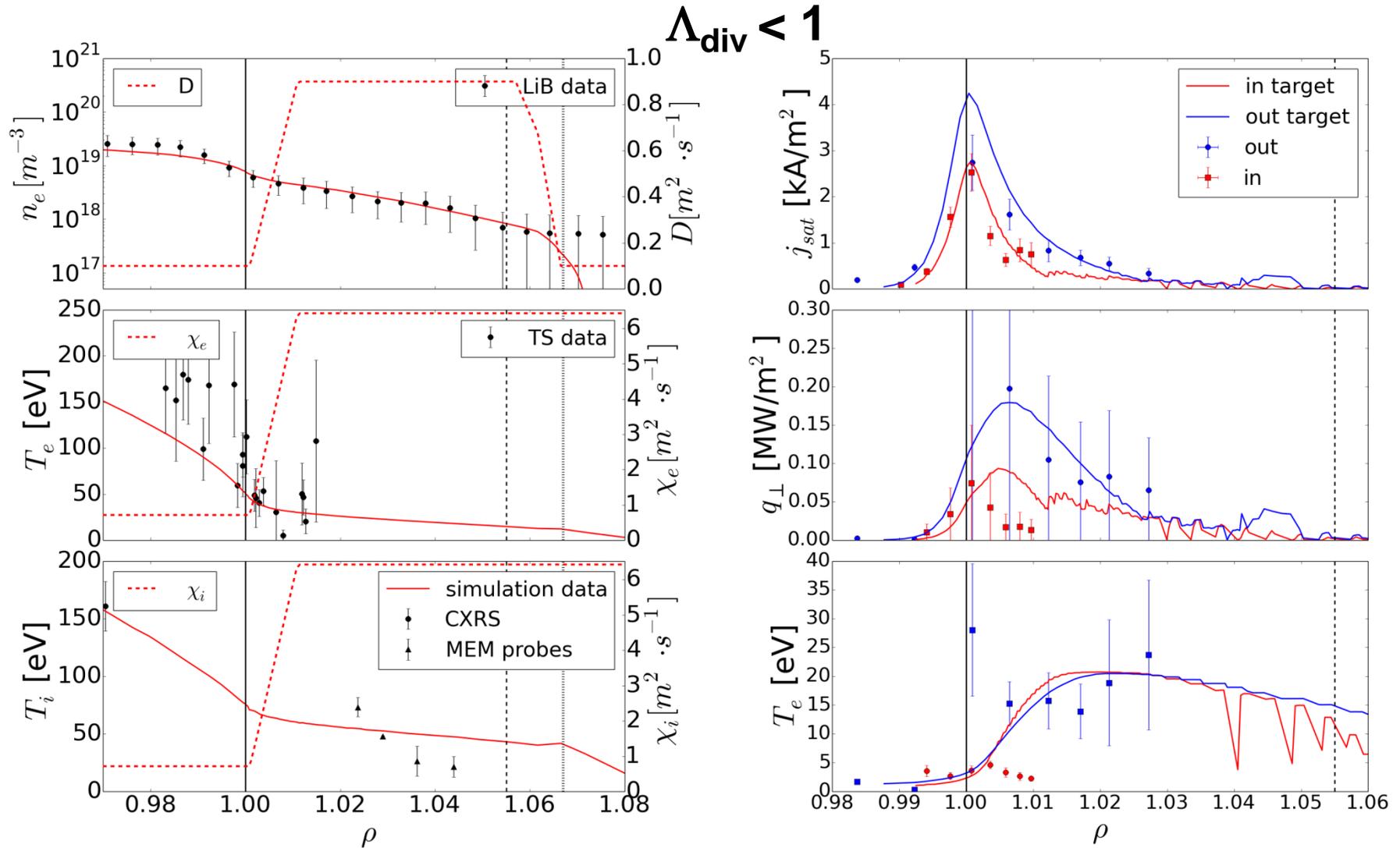
*Validated by IR and JET*

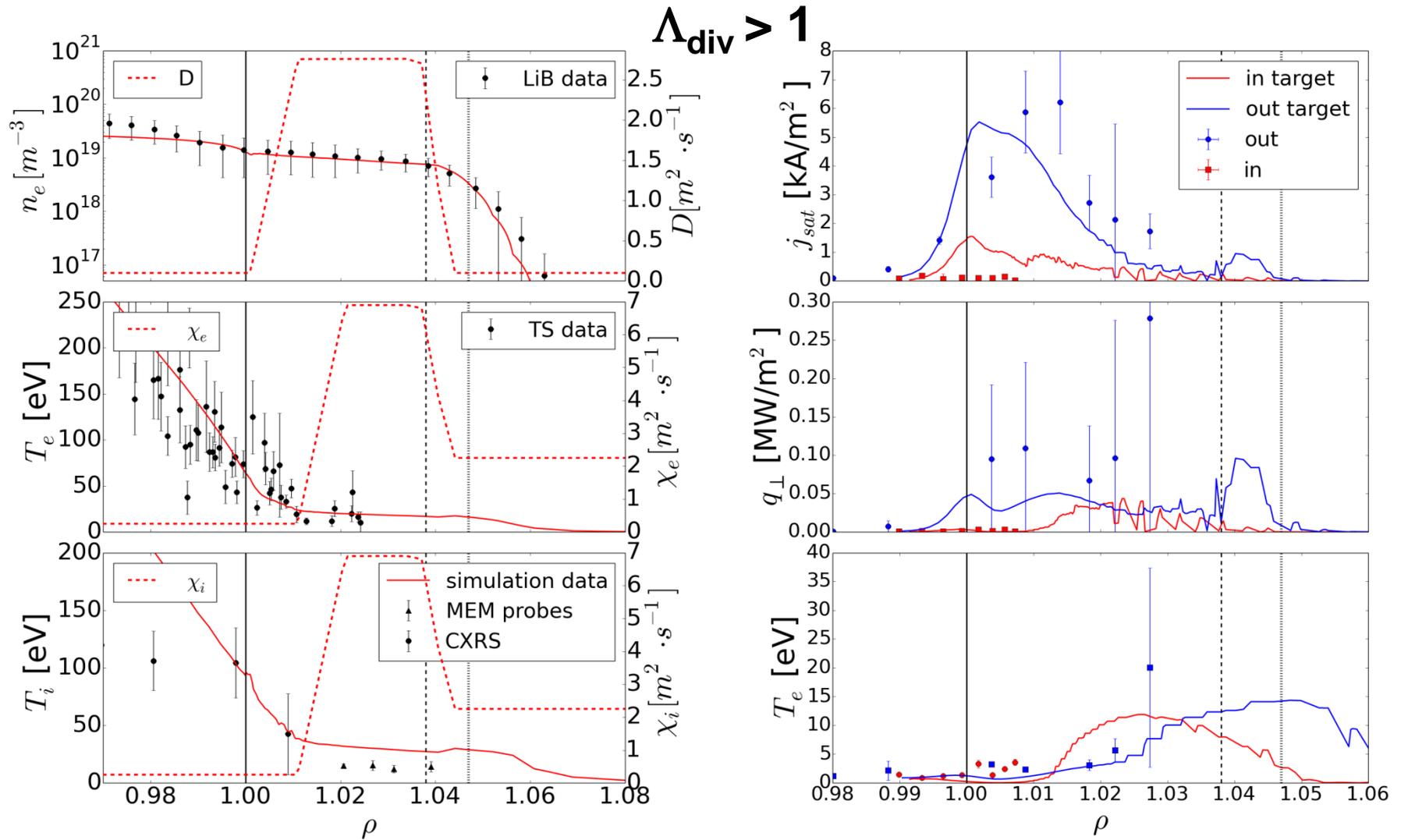


Total  $q_{||}^{MEM}$  can be calculated and compared to IR measurements:

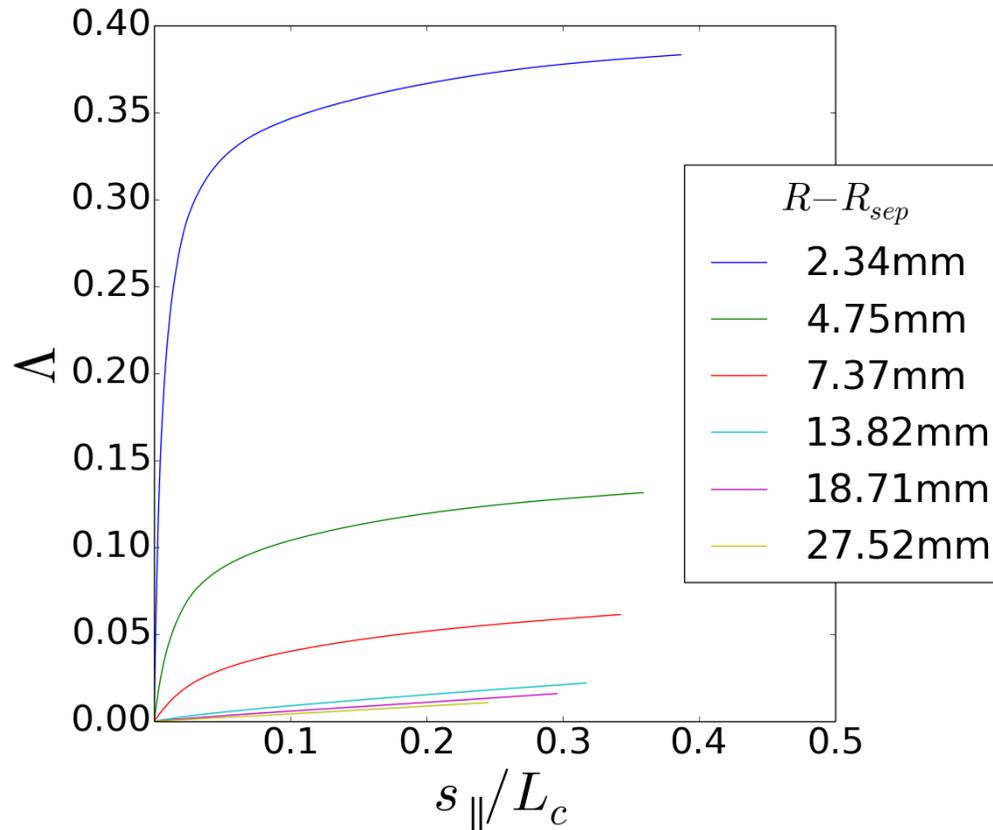
$$q_{||}^{MEM} = q_{||}^{fil} f_{fil} + q_{||}^{back} (1 - f_{fil})$$



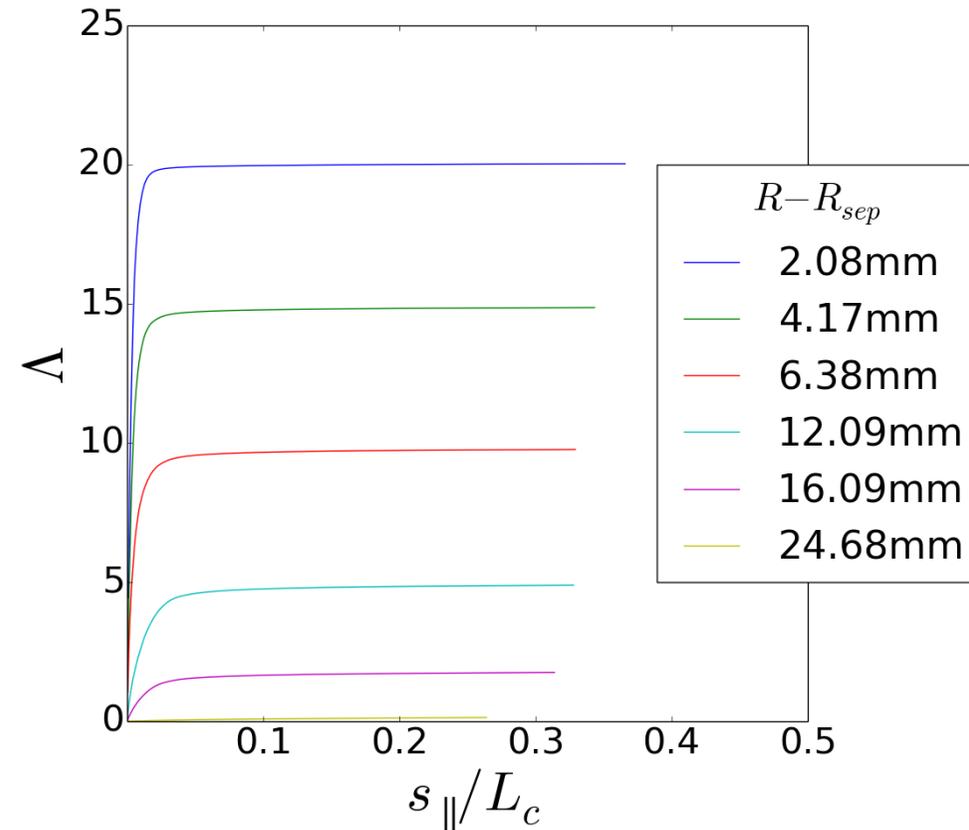




## Low density

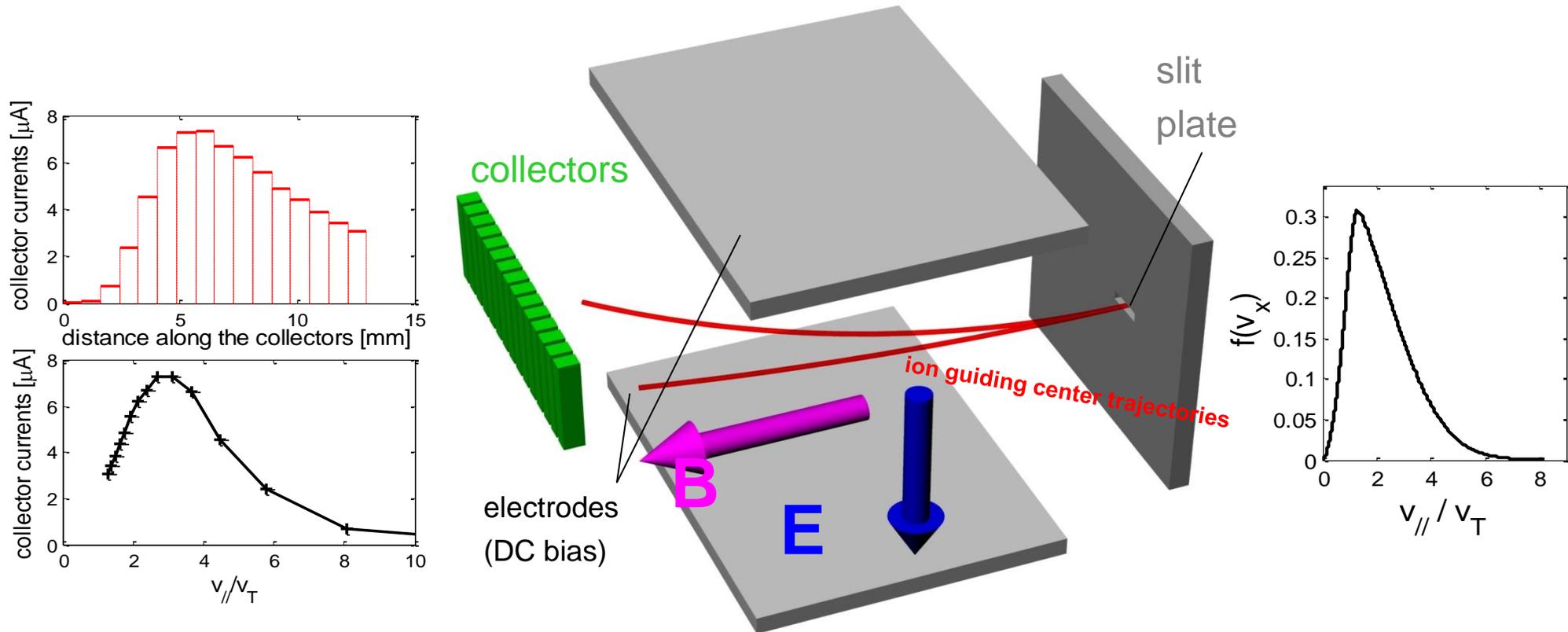


## High density

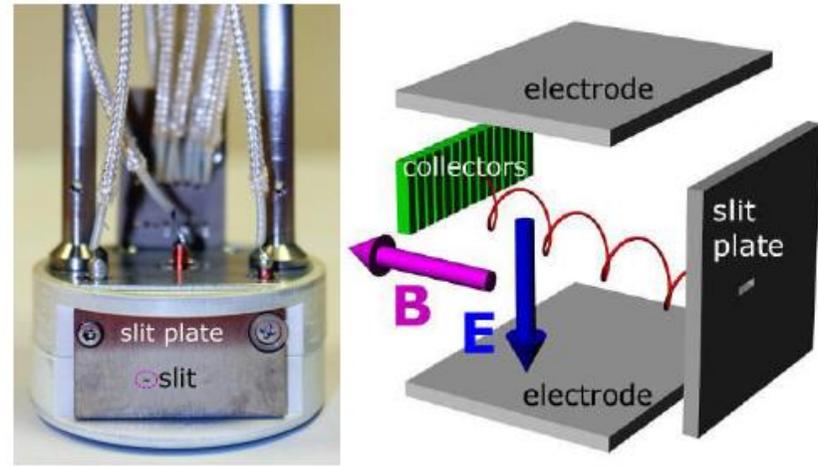
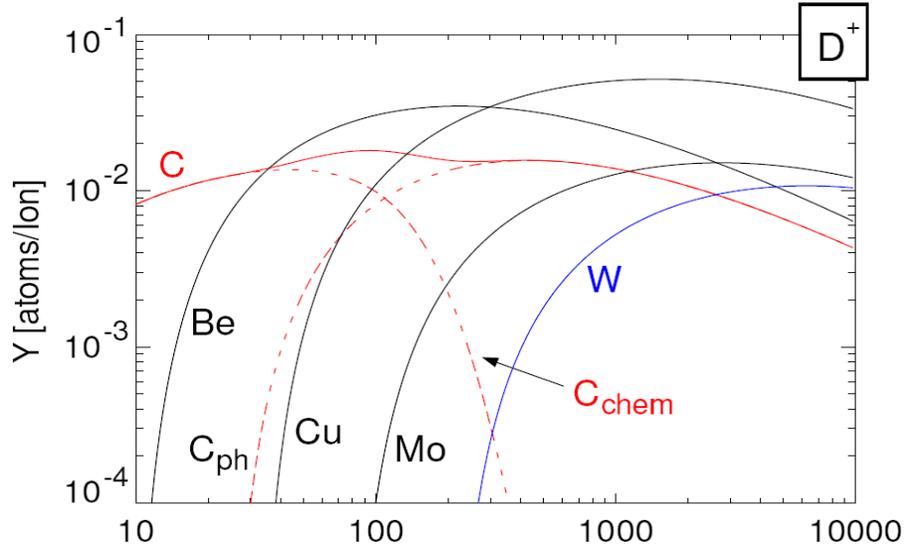


Probe head designed to measure ion energy distribution in real time.

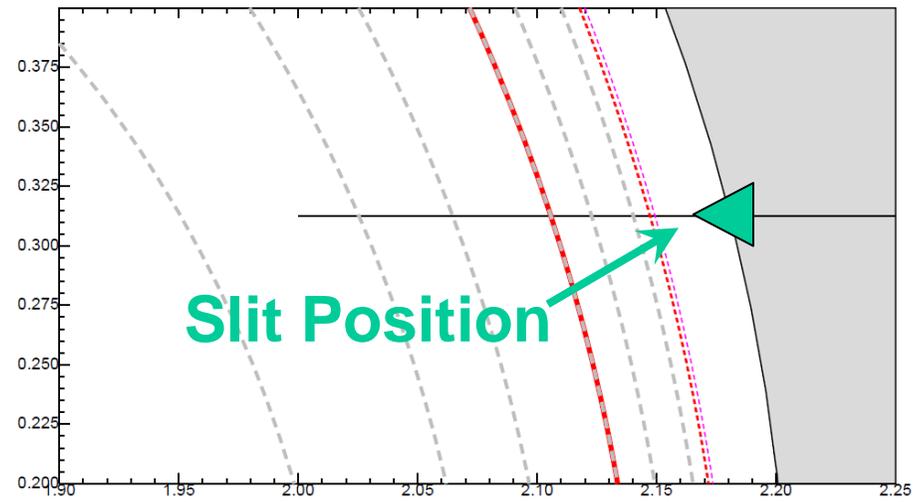
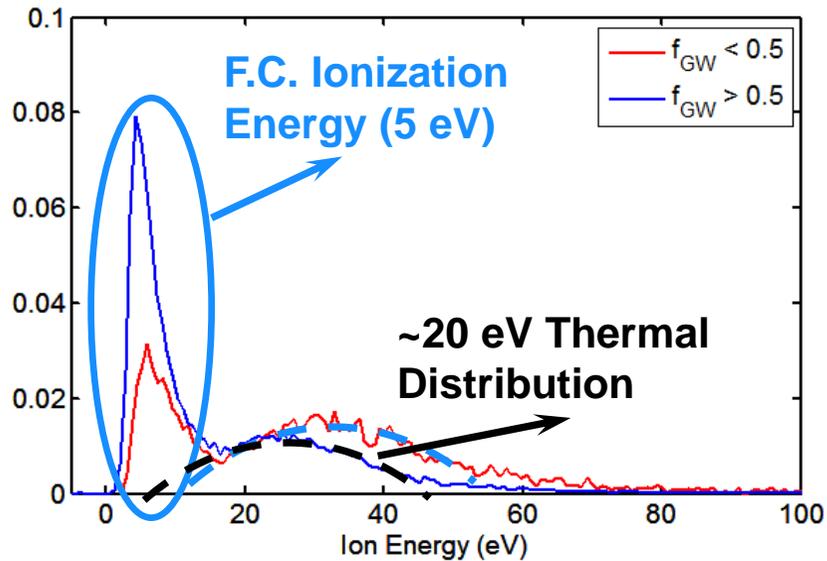
Collaboration between AUG and COMPASS teams.



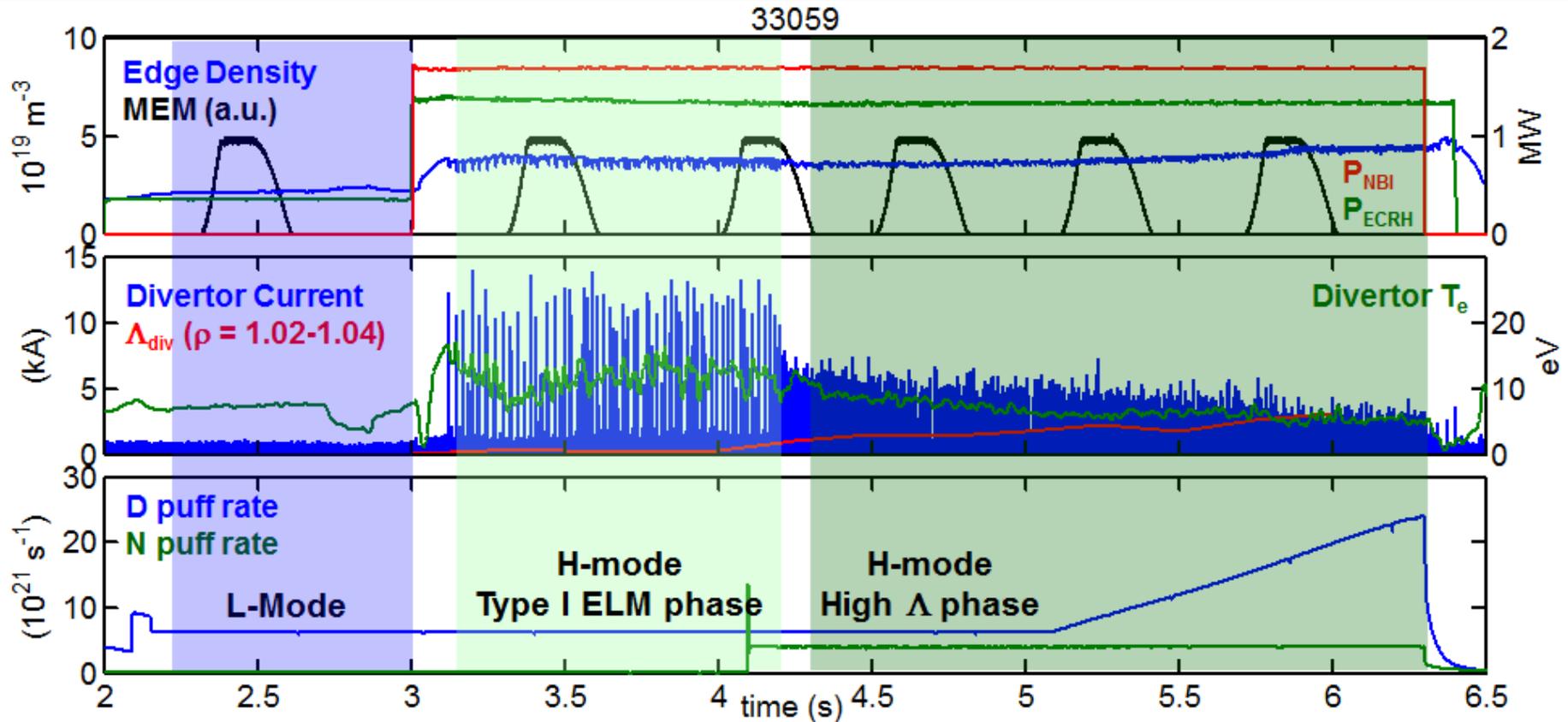
Time resolution determined by DAQ system (2 MHz)



M. Komm, et al., 40<sup>th</sup> EPS, (2013)

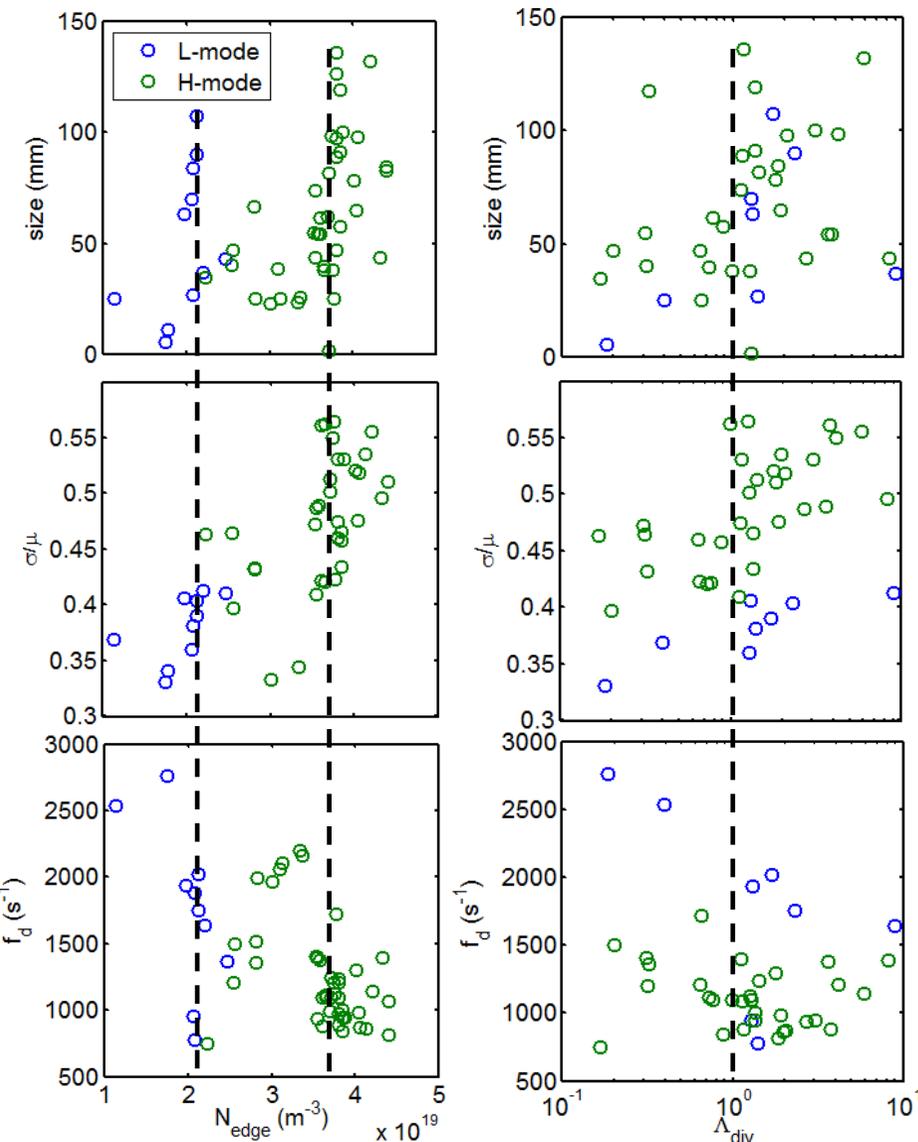


D. Carralero et al., 21<sup>st</sup> PSI, (2014)



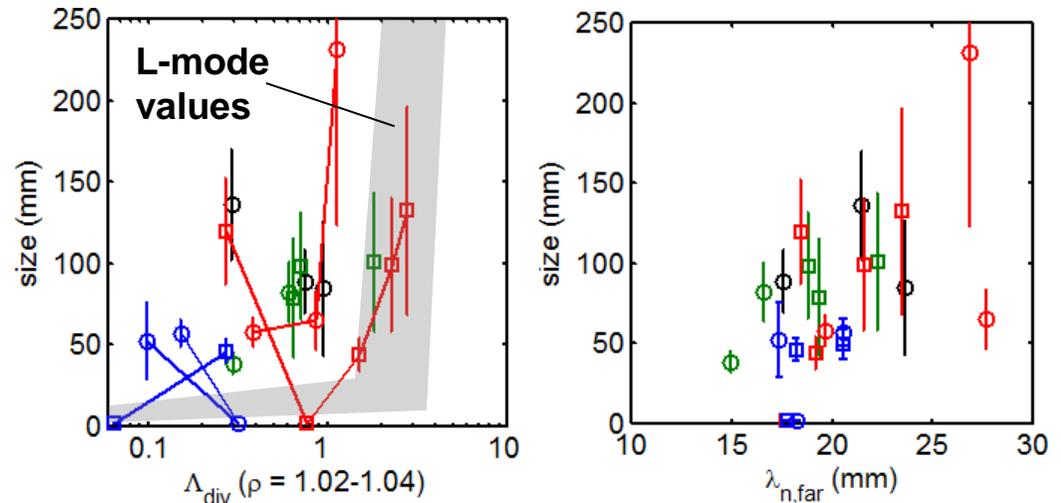
Series of  $\Delta_{\text{div}}$  sweeps with the same magnetic configuration (LSND Edge optimized) and parameters ( $B_T = 2.5 \text{ T}$ ,  $I_p = 800 \text{ kA}$ ,  $q_{95} = 4.85$ ) as L-mode.

Enough  $P_{\text{heat}}$  to access H-mode, but avoid damage to the manipulator.  $\Delta_{\text{div}} > 1$ , is achieved using N seeding and/or large D fueling.

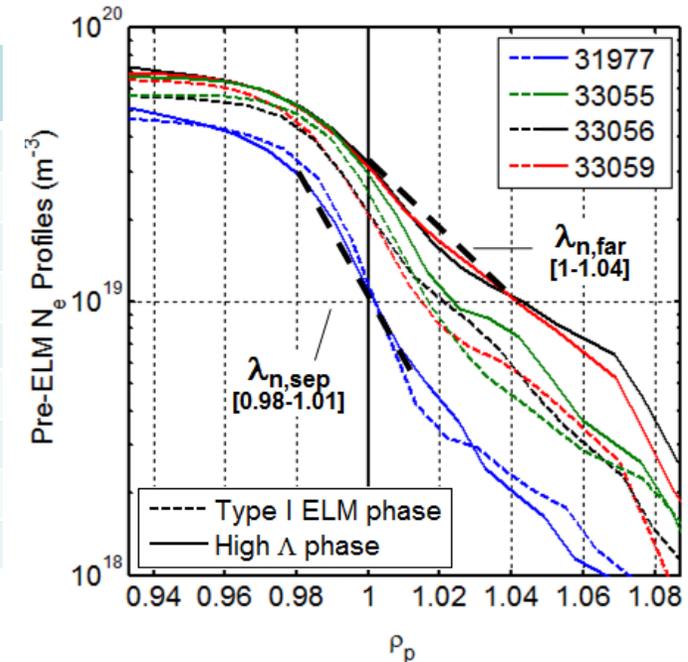


Perpendicular filament **size transitions are observed** for different  $N_{\text{edge}}$  values at L-mode and H-mode. However, **both happen around  $\Lambda_{\text{div}} = 1$** , in agreement with the SL-IN regime transition model.

Individual discharges show limited data. However, there is **general correlation between  $\delta_b$  and  $\lambda_{n, \text{far}}$** .



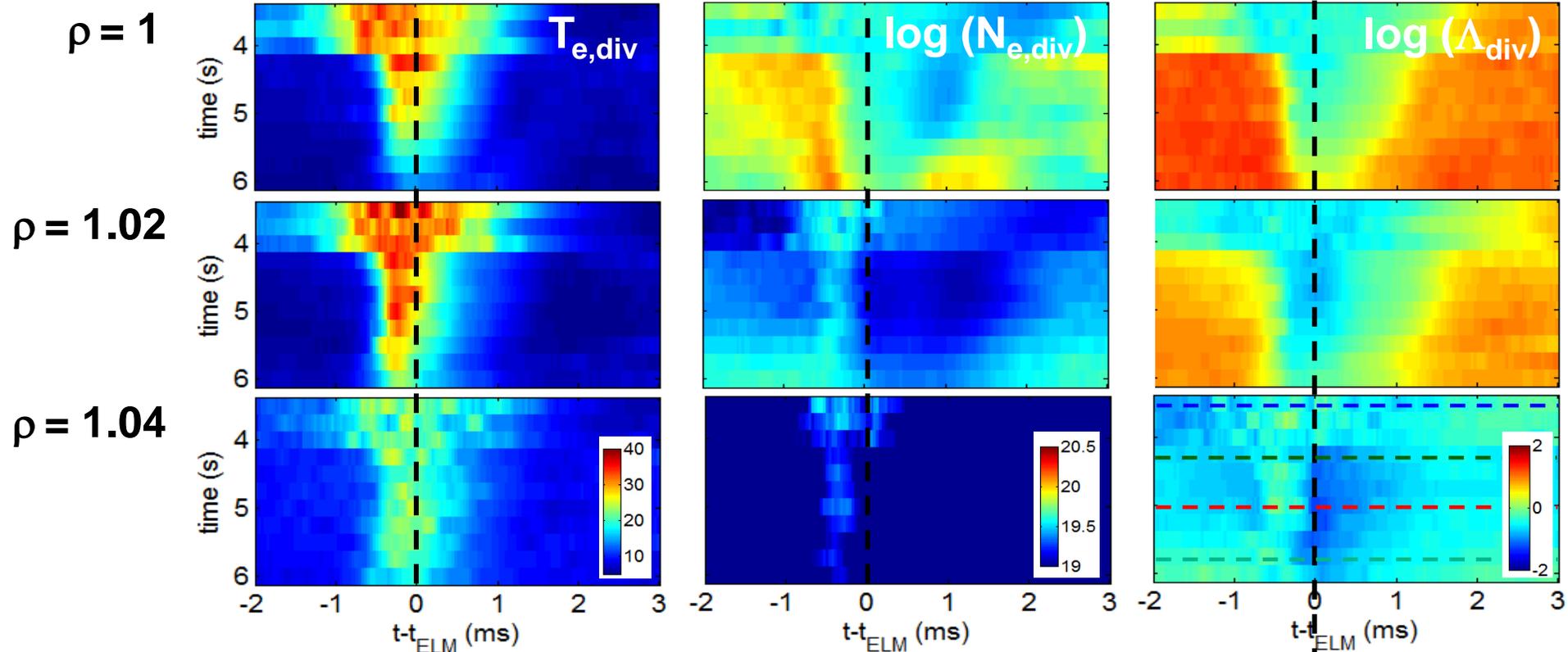
| Shot      | Scenario       | $P_{ECH}$<br>(MW) | $P_{NBI}$<br>(MW) | $D_{rate\ t=3.5\ s}$<br>( $10^{21}\ s^{-1}$ ) | $N_{rate\ t=4.5\ s}$<br>( $10^{21}\ s^{-1}$ ) | $D_{rate\ max}$<br>( $10^{21}\ s^{-1}$ ) | $N_{rate\ max}$<br>( $10^{21}\ s^{-1}$ ) | $T_{e,div\ min}$<br>(eV) |
|-----------|----------------|-------------------|-------------------|---|---|--|--|--------------------------|
| 31974 (○) | Low N, Low D   | 1.3               | -                 | 1.3   | 2.8   | 1.3                                      | 3.4                                      | 1                        |
| 31977 (□) |                | 1.9               | -                 | 1.3   | 2.8   | 1.3                                      | 3.8                                      | 0                        |
| 33055 (○) | High N, Low D  | 1.3               | 1.7               | 6   | 1.8   | 6  | 8  | 1.5                      |
| 33057 (□) |                | 1.3               | 1.7               | 8.2   | 1.8   | 8.2                                      | 8  | 3                        |
| 33056 (○) | High D, no N   | 1.3               | 1.7               | 6.3   | -   | 24.5                                     | -  | 10                       |
| 33058 (○) | High D, High N | 1.3               | 1.7               | 6.3   | 2   | 15.4                                     | 8  | 5                        |
| 33059 (□) |                | 1.3               | 1.7               | 6.3   | 4   | 24                                       | 4  | 5                        |
| 33475 (◇) |                | 1.4               | 1.7               | 6.2   | 5   | 24.4                                     | 5  | 7.5                      |
| 33478 (△) |                | 1.4               | 2.4               | 7   | 9.2   | 24.4                                     | 9.2                                      | 1.5                      |



To achieve different levels of  $\Lambda_{div}$ , and disentangle the contributions of  $D_{rate}$  and  $N_{rate}$ , **4 different scenarios are defined**, based on  $P_{heat}$  and the relative amplitude of D fueling and N seeding.

To quantify the degree of shoulder formation,  $\lambda_{n,far}$  is defined in the far SOL.

**Some scenarios develop a clear shoulder, some others do not.**



During H-mode phase, ELM ejection causes a major, intermittent perturbation of SOL conditions.

The **beginning of the pre-ELM** phase is considered to be a good approximation of inter-ELM conditions

