



# Synergetic effects of collisions, turbulence and sawtooth crashes on impurity transport.

X. Garbet<sup>1</sup>, J.H. Ahn<sup>1</sup>, S. Breton<sup>1</sup>, P. Donnel<sup>1</sup>, D. Esteve<sup>1</sup>, R.Guirlet<sup>1</sup>, H.Lütjens<sup>2</sup>, T. Nicolas<sup>3</sup>, Y. Sarazin<sup>1</sup>, C. Bourdelle<sup>1</sup>, O. Février<sup>1</sup>, G. Dif-Pradalier<sup>1</sup>, P. Ghendrih<sup>1</sup>, V. Grandgirard<sup>1</sup>, G.Latu<sup>1</sup>, J.F. Luciani<sup>2</sup>, P.Maget<sup>1</sup>, A. Marx<sup>2</sup>, A. Smolyakov<sup>4</sup>

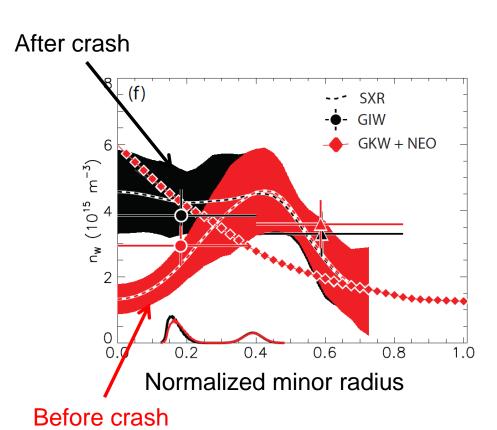
- 1) CEA, IRFM, F-13108 Saint Paul-lez-Durance, France
- 2) Centre de Physique Théorique, Ecole Polytechnique, CNRS, Palaiseau, France
- 3) Swiss Plasma Center, Switzerland
- 4) University of Saskatchewan, Saskatoon, Canada



#### **Motivation: impurity transport**



- Tungsten plasma facing components → impurity accumulation in the core ?
- Neoclassical and turbulence transport processes compete Casson 13, Angioni 14
- Interplay with MHD events:
   tearing modes, ELMs,
   sawtooth crashes Hender 16,
   Sertoli 15



Asdex Upgrade – tungsten density- Sertoli 15



#### **Outline**



- Interaction between turbulent and neoclassical transport.
- Interaction between sawtooth cycles and neoclassical transport.

Punchline: interplay between different contributions to impurity transport are mediated by large scale flows



#### Impurity transport modelled with gyrokinetics or MHD with closure



- Gyrokinetic description (GYSELA code): d<sub>t</sub>F=C(F) + Poisson equation→ neoclassical and turbulent transport Grandgirard 16
- MHD equations (XTOR code) + impurity density and momentum equations Lütjens 10

  pinch velocity

$$\frac{\partial N}{\partial t} + \nabla \cdot (N\mathbf{V}) = \nabla \cdot (D\nabla N - V)$$
 collisional friction force  $\rightarrow$  neoclassical flux 
$$Nm\left(\frac{\partial}{\partial t} + \mathbf{V} \cdot \nabla\right)\mathbf{V} = Ne\left(\mathbf{E} + \mathbf{V} \times \mathbf{B}\right) - \nabla \cdot \mathbf{\Pi} + \mathbf{R}$$

- → Pfirsch-Schlüter transport included in the fluid dynamics
- → neoclassical and MHD transport



#### Impurity neoclassical flux is related to parallel friction force

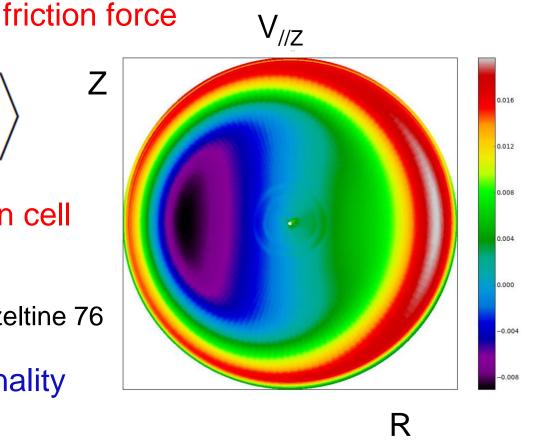
// collisional



Neoclassical flux

$$\Gamma^{\psi} = -\frac{B_T R}{Ze} \left\langle \begin{array}{c} R_{\parallel} \\ B \end{array} \right\rangle$$

- Pfirsch-Schlüter convection cell due to perpendicular compressibility Hinton & Hazeltine 76
- controls Γ<sup>Ψ</sup> at high collisionality
   v\*<sub>7</sub>>1





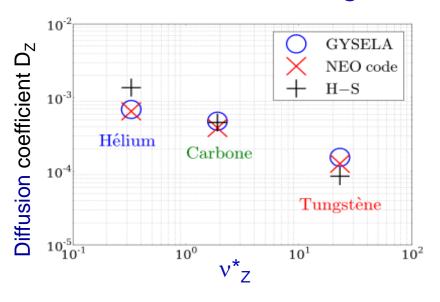
#### Neoclassical thermal screening works against accumulation

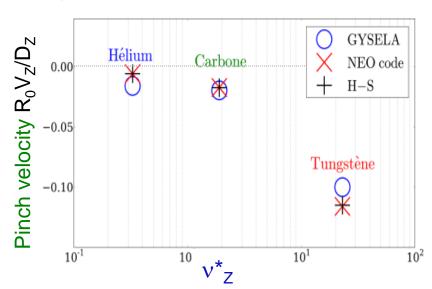


General form of the impurity flux Hirshman & Sigmar 81

$$\frac{\Gamma_{Z\psi}}{D_{neo}N_Z} = -\frac{\partial \ln N_Z}{\partial r} + Z\frac{\partial \ln N_i}{\partial r} + HZ\frac{\partial \ln T_i}{\partial r}$$
 Accumulation Thermal screening

- Impurity collisional, ions weakly collisional → H = -1/2 Hirshman 76
- GYSELA benchmarked against theory and NEO code Belli 08



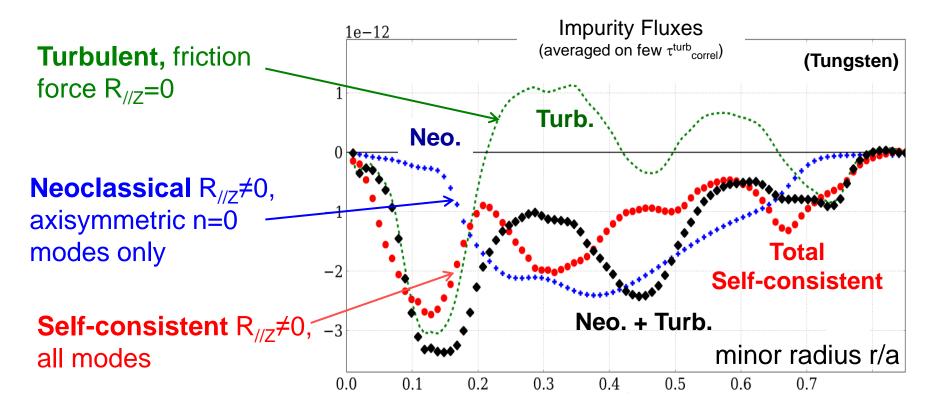




## Neoclassical and turbulent transport processes are synergetic



- Neoclassical and turbulent contributions isolated by playing with collisionality and symmetries
- Total flux ≠ neoclassical + turbulent

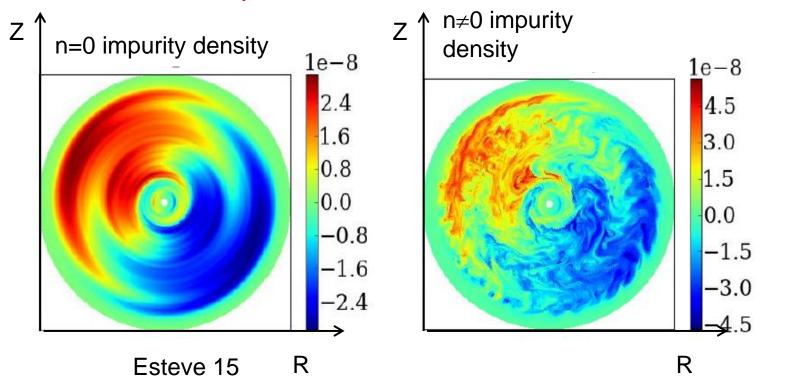




### Interplay is mediated by poloidal convective cells



- Turbulent Reynolds stress → poloidal convective cells
- Poloidal asymmetries → change neoclassical impurity flux
- // momentum transport, turbulence self-regulation Diamond 05

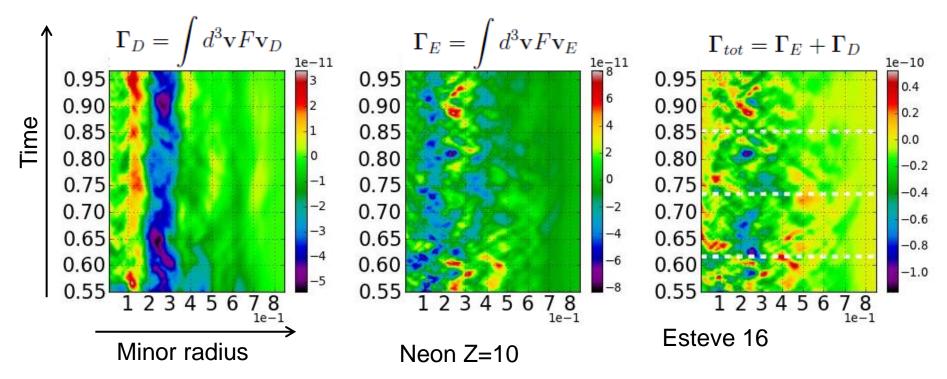




#### Curvature and ExB fluxes are anticorrelated



- Anti-correlation due to poloidal convective cells
- Thermal screening factor H>-1/2: consequence of static density poloidal asymmetries? Romanelli 98, Fülöp 99, Angioni 14, Breton 16



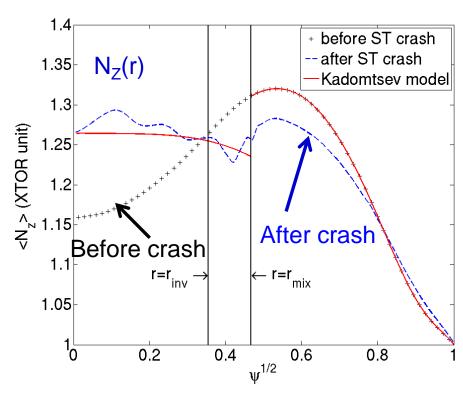


### Fast relaxation of the impurity density profile during a sawtooth crash



- $\nabla N_i = 0$ ,  $\nabla T_i \neq 0 \rightarrow$  screening
- Crash time << collision time → neoclassical transport processes inefficient during crash
- Post-crash profile
   consistent with Kadomtsev
   model Kadomtsev 75, Porcelli 96,
   Nicolas 15



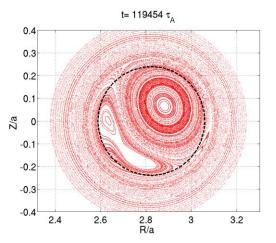


Normalized minor radius



### ExB drift is the main cause of impurity transport during a sawtooth crash





t= 120198 τ<sub>A</sub>

0.4

0.3

0.2

0.1

-0.1

-0.2

-0.3

-0.4

2.4

2.6

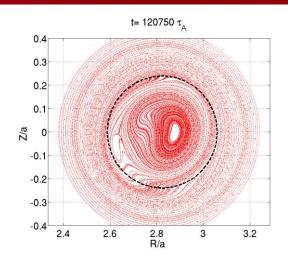
2.8

8/a

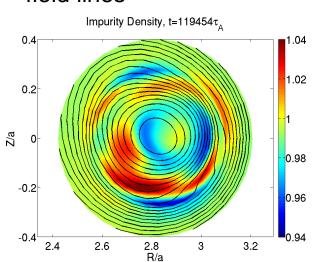
3

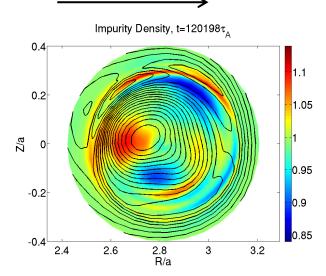
3.2

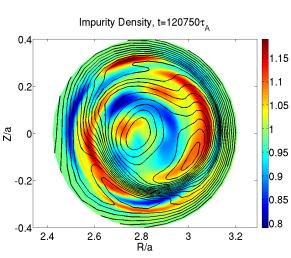
time



Poincaré map of magnetic field lines







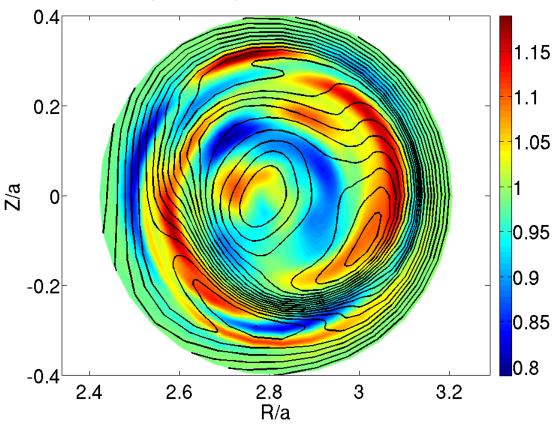
Impurity density and stream function



## ExB drift is the main cause of impurity transport during a sawtooth crash (cont.)







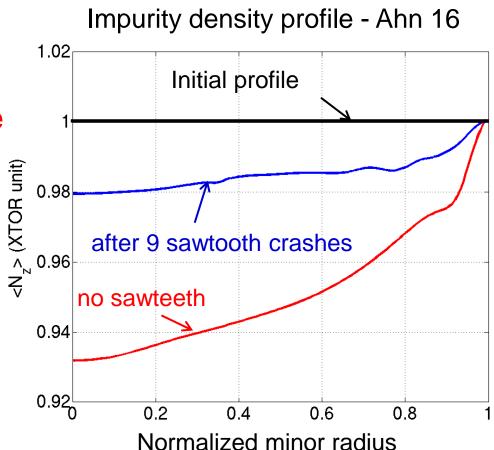
- ExB impurity flux ~10 flux due to magnetic flutter
- Consistent with SXR measurements on TFTR Nagayama 91



### Sawteeth change the impurity profile on long time scales



- Neoclassical transport
   dominant during recovery
   phase, but ion temperature
   gradient is lower
- → weaker thermal screening effect
- Overall temperature profile flatter with sawteeth





#### Conclusion



- Interplay between turbulent and neoclassical transport processes:
- poloidal convective cells generated by turbulence → poloidal asymmetries – total flux ≠ turbulent + neoclassical calculated separately. Should play at low rotation speed (e.g. EAST, WEST, ITER)
- thermal screening gets weaker
- Sawteeth cycles affect neoclassical transport
- crashes flatten impurity density profile + lower main ion temperature gradient → thermal screening less efficient