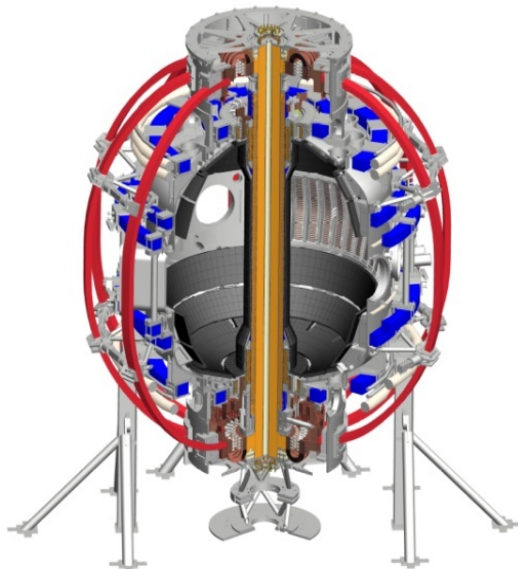


# Effects of MHD instabilities on Neutral Beam current drive

M. Podestà, M. Gorelenkova, D. S. Darrow,  
E. D. Fredrickson, S. P. Gerhardt,  
W. W. Heidbrink, R. B. White  
*and the NSTX-U Research Team*

**25<sup>th</sup> IAEA Fusion Energy Conference**  
**St Petersburg, Russian Federation**  
**13-18 October, 2014**



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*CEA, Cadarache*  
*IPP, Jülich*  
*IPP, Garching*  
*ASCR, Czech Rep*

# Reliable, quantitative predictions of Energetic Particle (EP) dynamics are crucial for burning plasmas

- EPs from Neutral Beam (NB) injection, alphas, RF tails drive instabilities,
    - e.g. Alfvénic modes – AEs
  - With instabilities, ‘classical’ EP predictions (e.g. for NB heating, current drive) can fail
- > *Predictive tools are being developed, validated for integrated modeling of these effects in present and future devices (ITER, Fusion Nuclear Science Facility – FNSF)*

# Outline

- NSTX discharges with strong MHD are used to test and validate EP transport models
- Modeling methods beyond ‘classical’ EP physics are developed to account for MHD effects
- New model captures MHD modifications of EP phase space leading to Neutral Beam current redistribution

# Outline

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# Alfvénic modes (AEs) and kink-like modes degrade fast ion confinement, plasma performance

## NSTX

Major radius 0.85 m

Aspect ratio 1.3

Plasma current  $\sim 1$  MA

Toroidal field  $< 0.55$  T

Pulse length  $< 2$  s

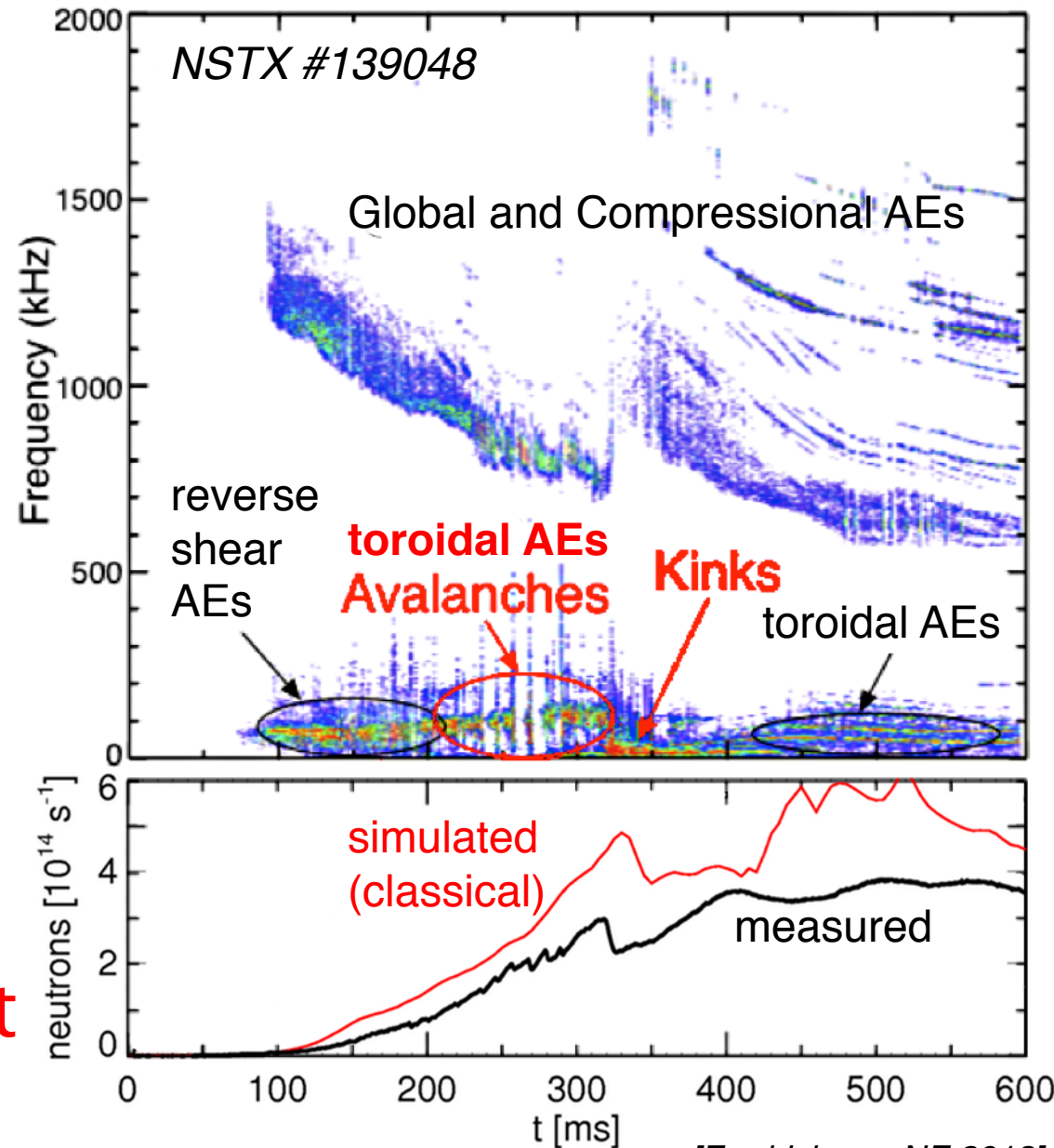
Neutral Beam sources:

$$P_{\text{NBI}} \leq 6 \text{ MW}$$

$$E_{\text{injection}} \leq 95 \text{ keV}$$

$$\Rightarrow 1 < v_{\text{fast}} / v_{\text{Alfvén}} < 5$$

Super-alfvénic ions,  
high  $\beta_{\text{fi}}$ : plethora of fast  
ion driven instabilities



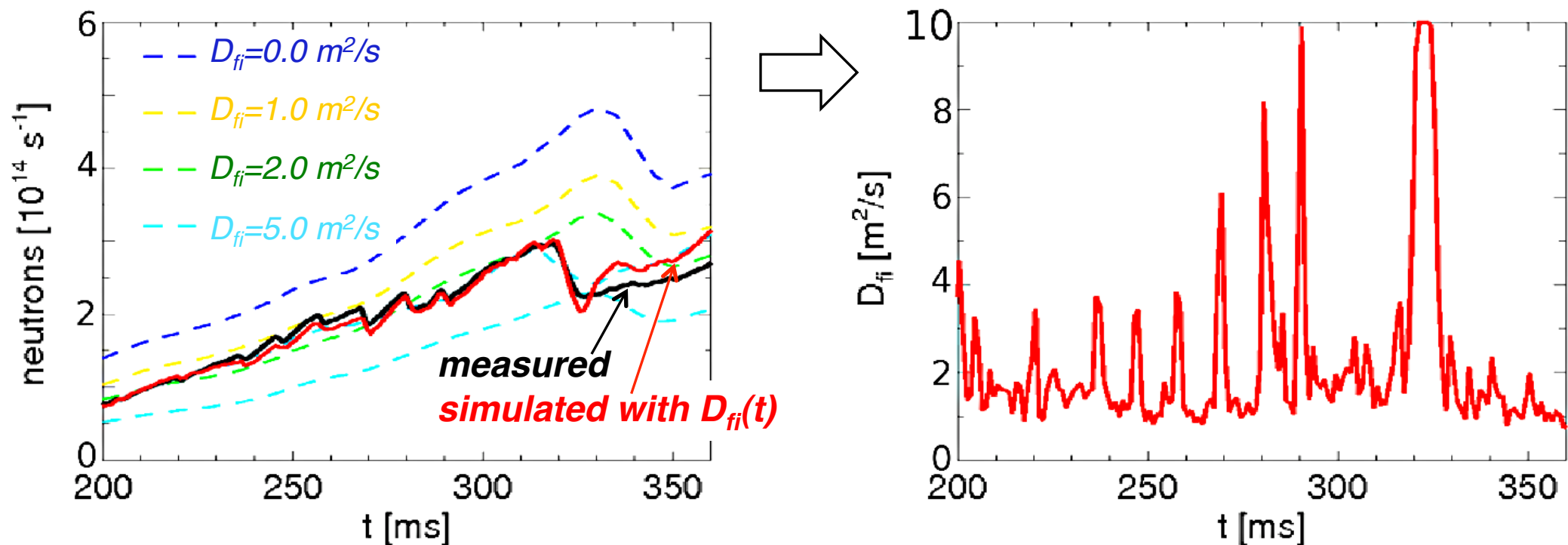
[Fredrickson, NF 2013]

# Outline

- NSTX discharges with strong MHD are used to test and validate EP transport models
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# Transport code TRANSP includes NUBEAM module for classical fast ion physics

- Additionally, *ad-hoc* diffusivity  $D_{fi}$  is used to mimic enhanced fast ion transport
  - Assumed uniform in radius, pitch, energy in this work
- Metric to set  $D_{fi}$ : match neutron rate,  $W_{mhd}$



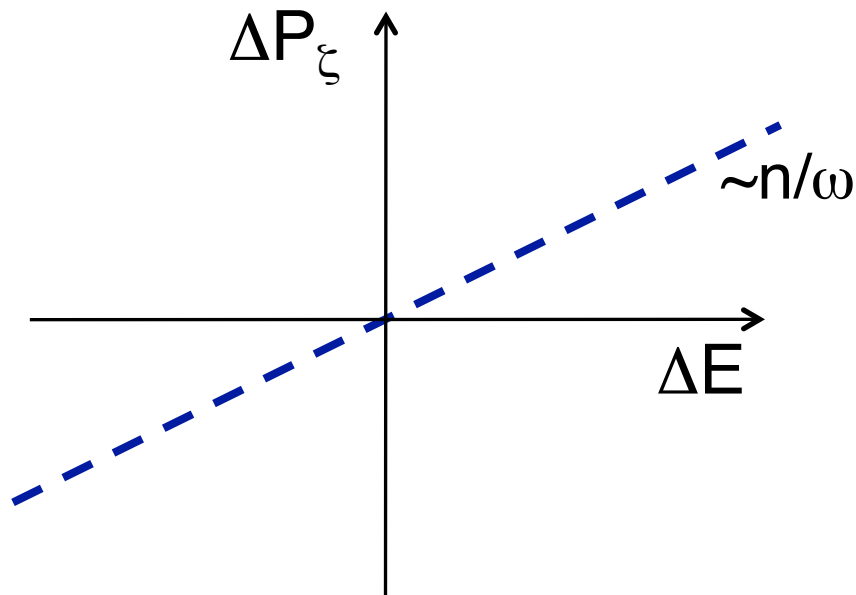
# However: instabilities introduce fundamental constraints on particle dynamics

From Hamiltonian formulation – single resonance:

$$\omega P_{\zeta} - nE = \text{const.} \implies \Delta P_{\zeta} / \Delta E = n / \omega$$

$\omega = 2\pi f$ , mode frequency

$n$ , toroidal mode number



$E$ , energy

$P_{\zeta} \sim mRv_{\text{par}} - \Psi$ , canonical angular momentum

$\mu \sim v_{\text{perp}}^2 / (2B)$ , magnetic moment

where  $\Psi$  : poloidal flux

$R$  : major radius

$m$  : mass



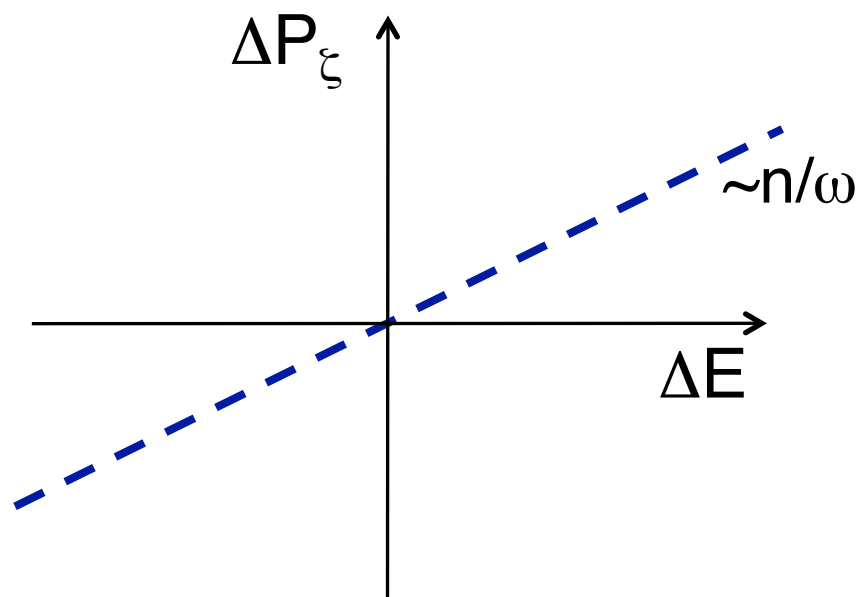
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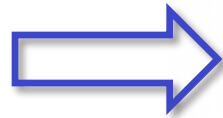
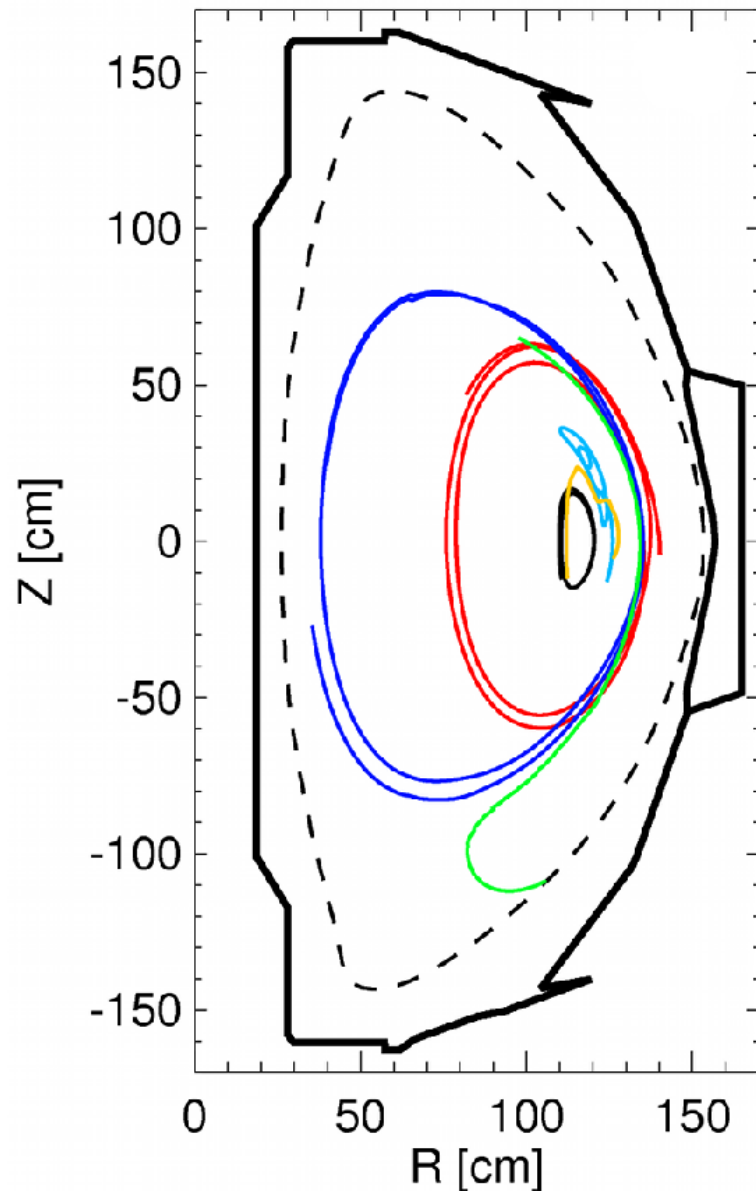
$m$  : mass

*These effects are not accounted for by ad-hoc  $D_{fi}$ .*

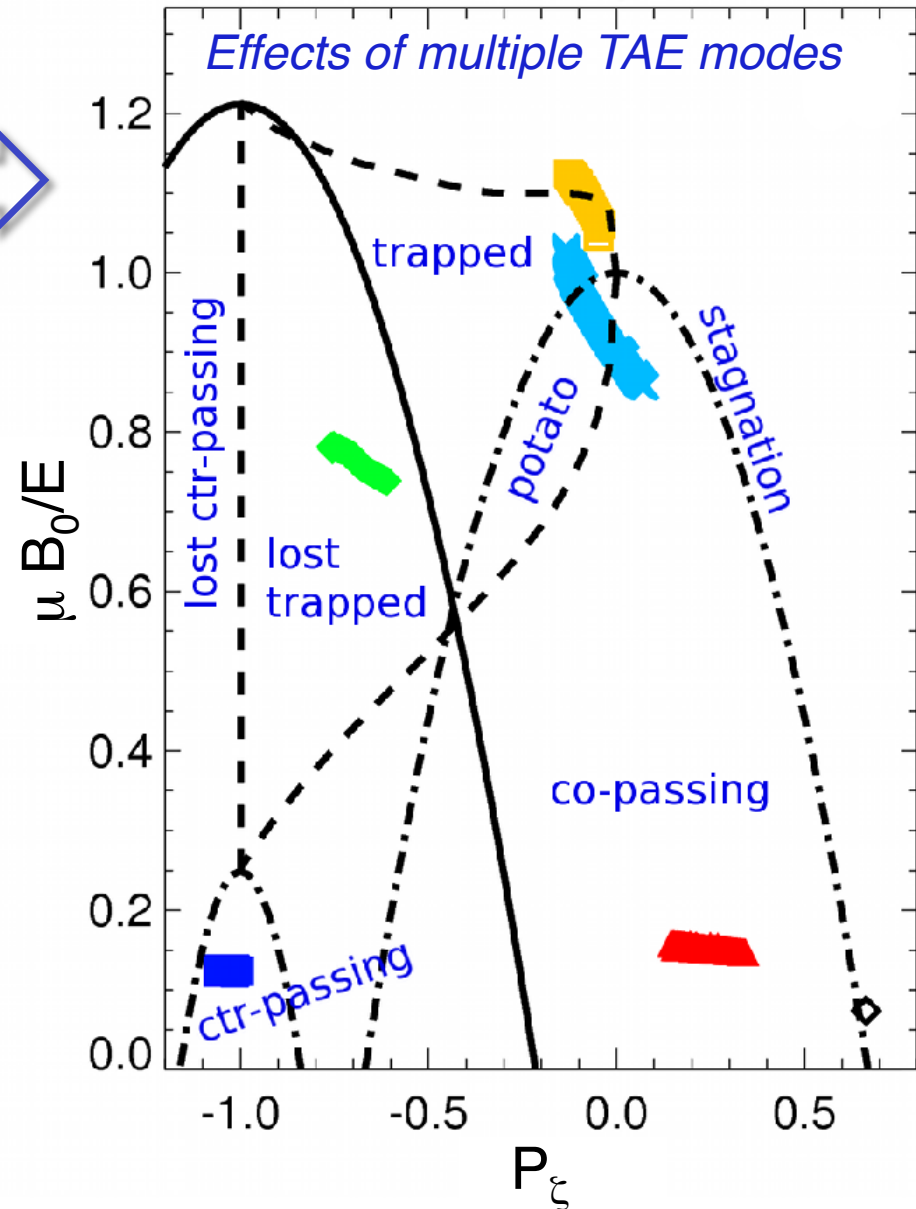
*A new method is needed to include them in integrated modeling.*

# Constants of motion ( $E, P_\xi, \mu$ ) are the natural variables to describe wave-particle interaction

NSTX poloidal section



Phase space,  $E_0=80.0\text{keV}$



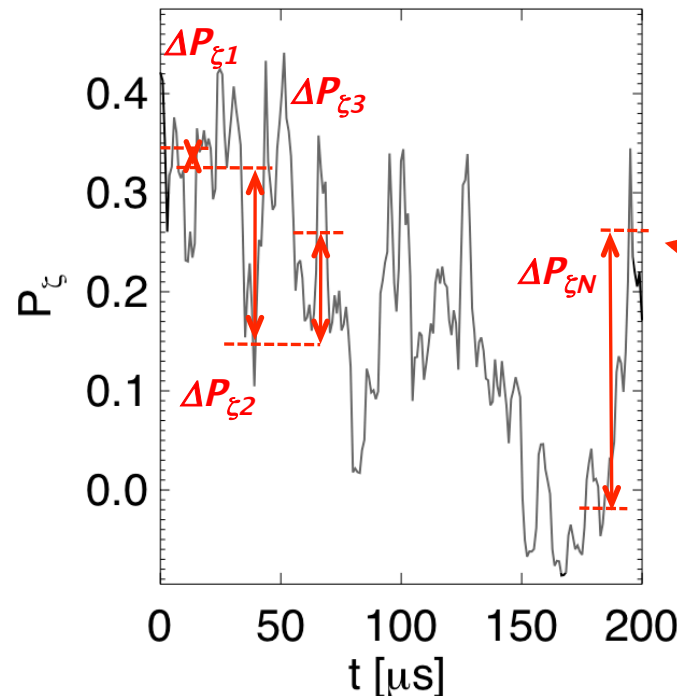
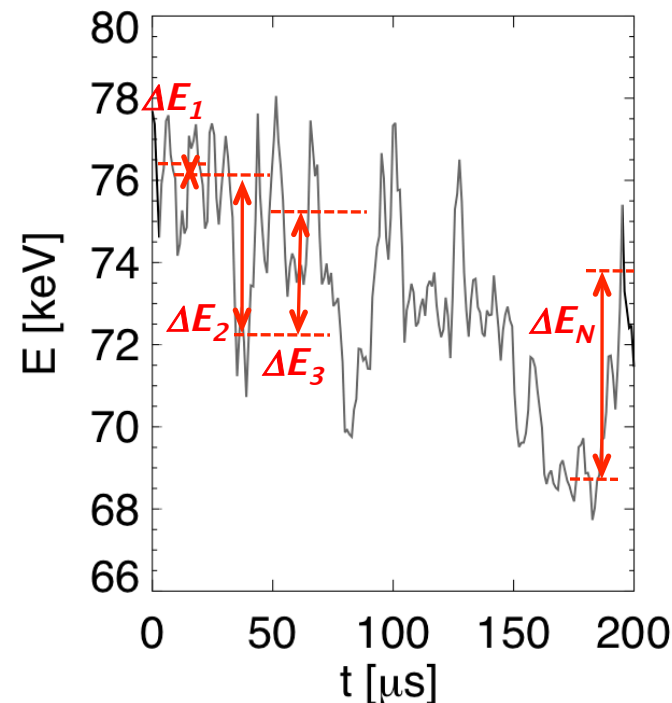
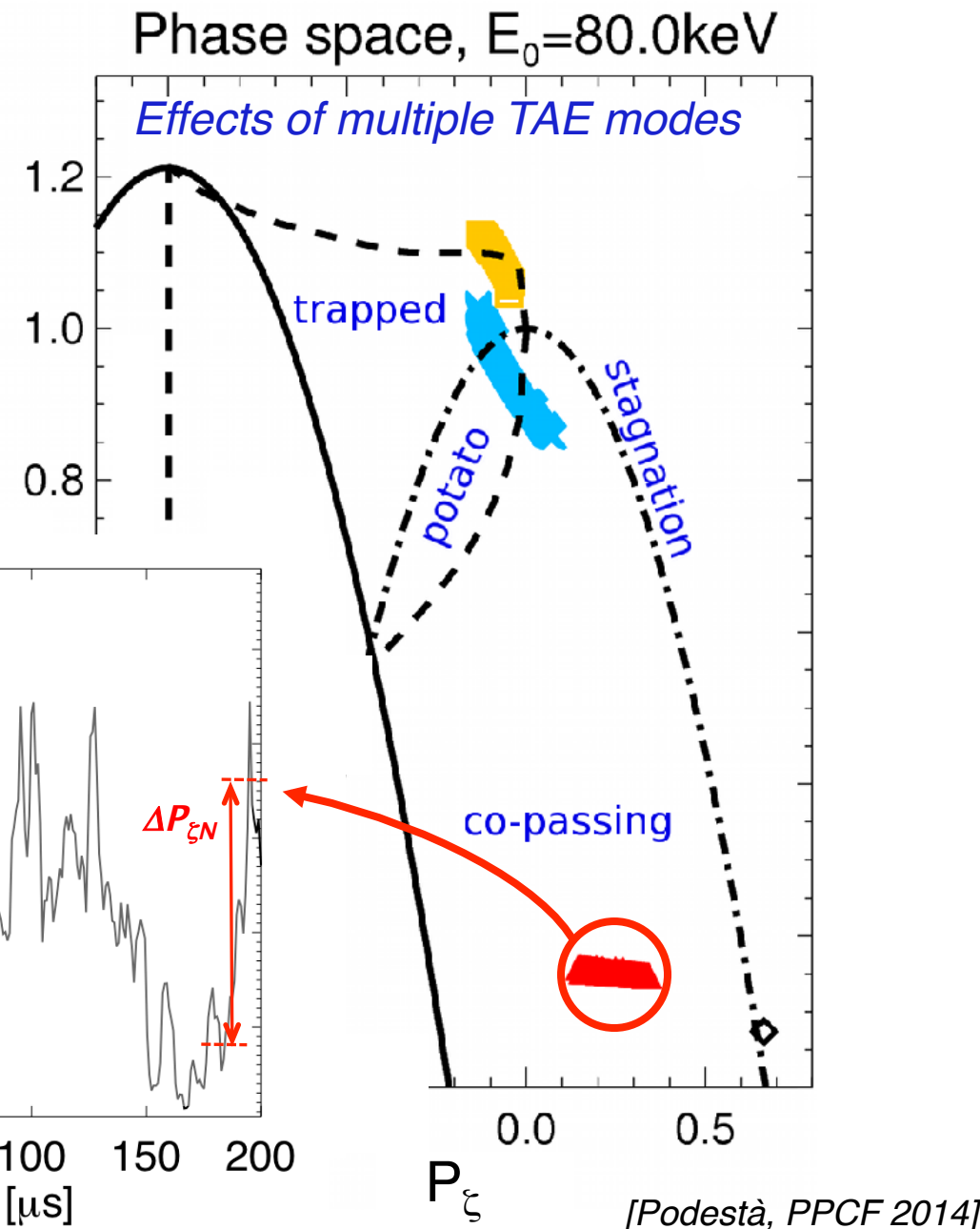
*R. B. White, Theory of toroidally confined plasmas,  
Imperial College Press (2014)*

# Particle-following codes are used to extract distribution of 'kicks' $\Delta E$ , $\Delta P_\xi$ for each *bin* ( $E, P_\xi, \mu$ )

- ORBIT code: record  $E, P_\xi, \mu$  vs. time for each particle
- Compute average kicks over multiple wave periods:

$$\underbrace{1/f_{\text{wave}}}_{\text{neglected}} < \underbrace{\tau_{\text{resonance}}}_{\text{relevant time scale}} < \underbrace{\tau_{\text{collisions}}}_{\text{classical}}$$

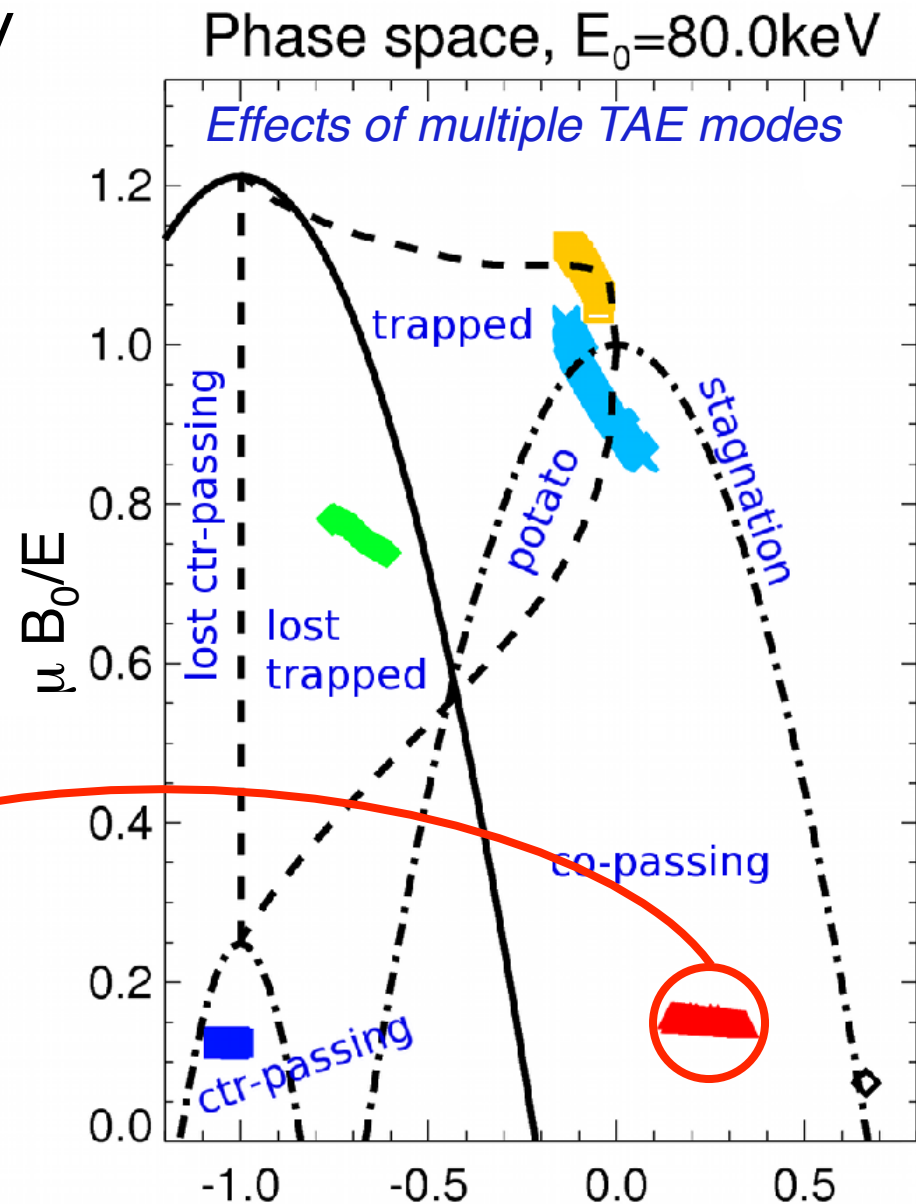
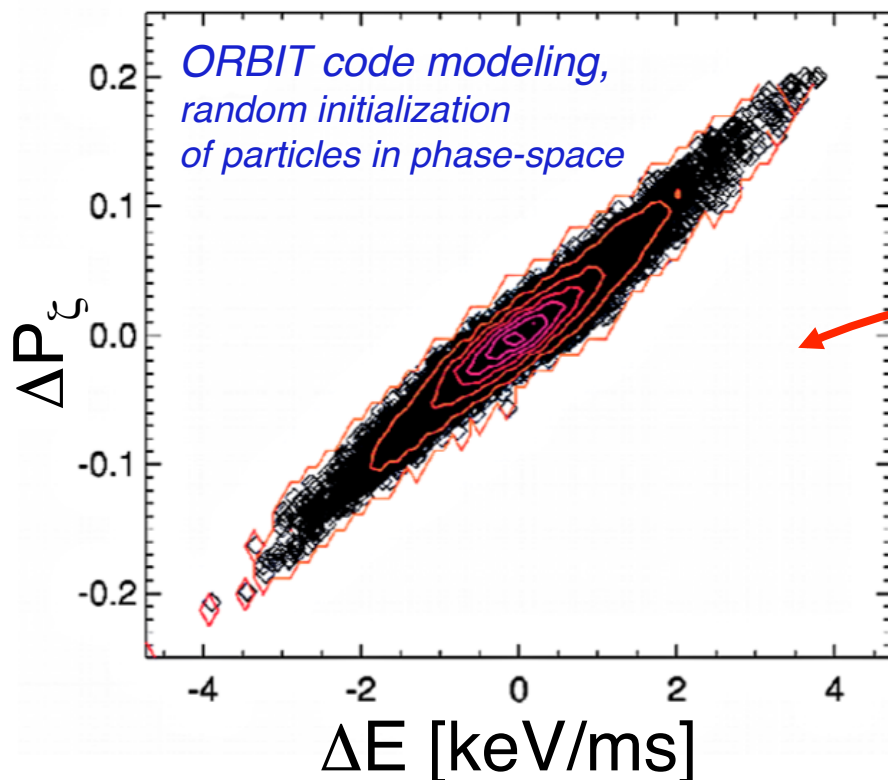
- Re-bin for each ( $E, P_\xi, \mu$ ) region



# New 'kick model' uses a *probability distribution function* for particle transport in $(E, P_\xi, \mu)$ space

Kicks  $\Delta E, \Delta P_\xi$  are described by  $p(\Delta E, \Delta P_\xi | P_\xi, E, \mu, A)$  which includes the effects of multiple modes, resonances.

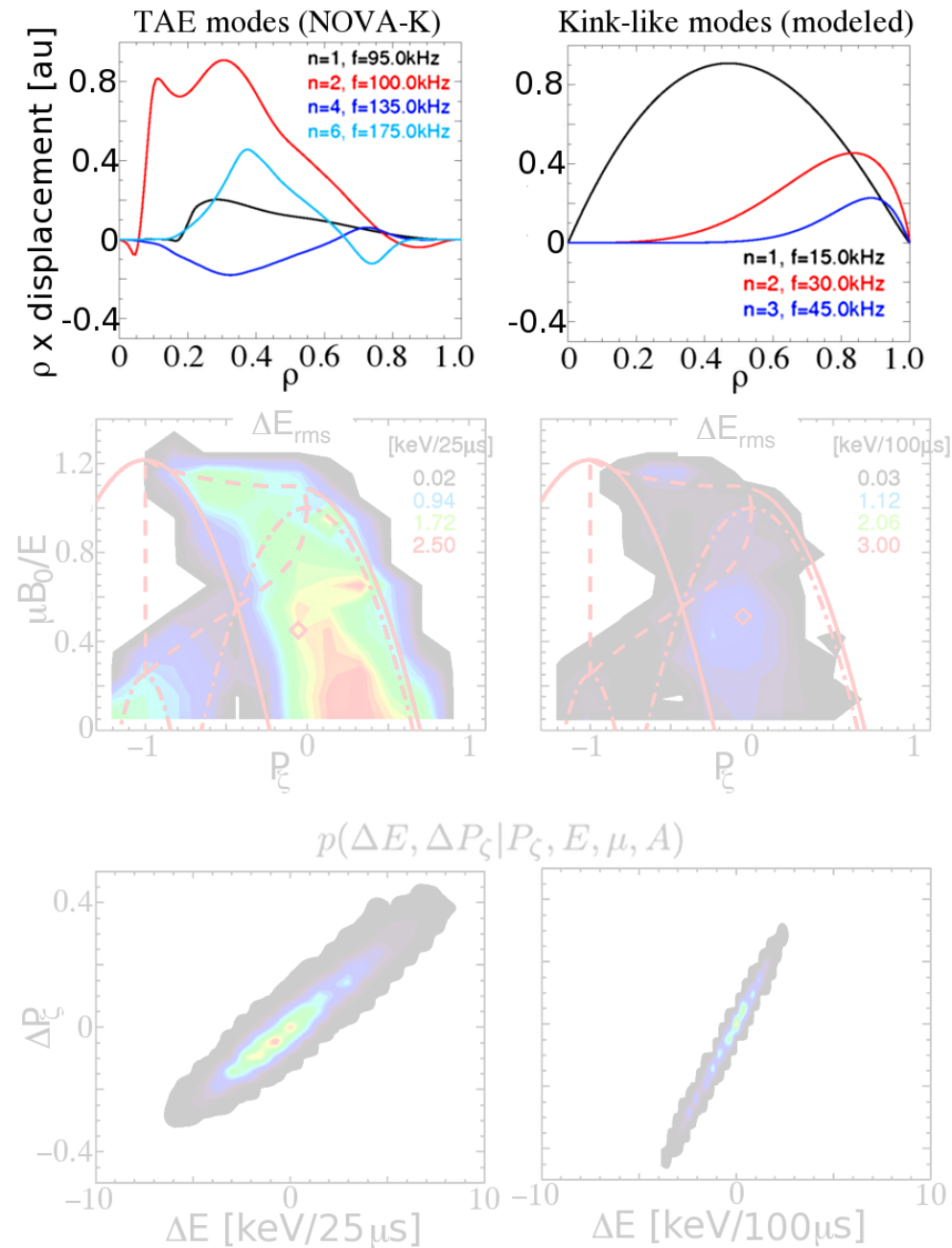
→ *correlated random walk in  $E, P_\xi$*



[Podestà, PPCF 2014]

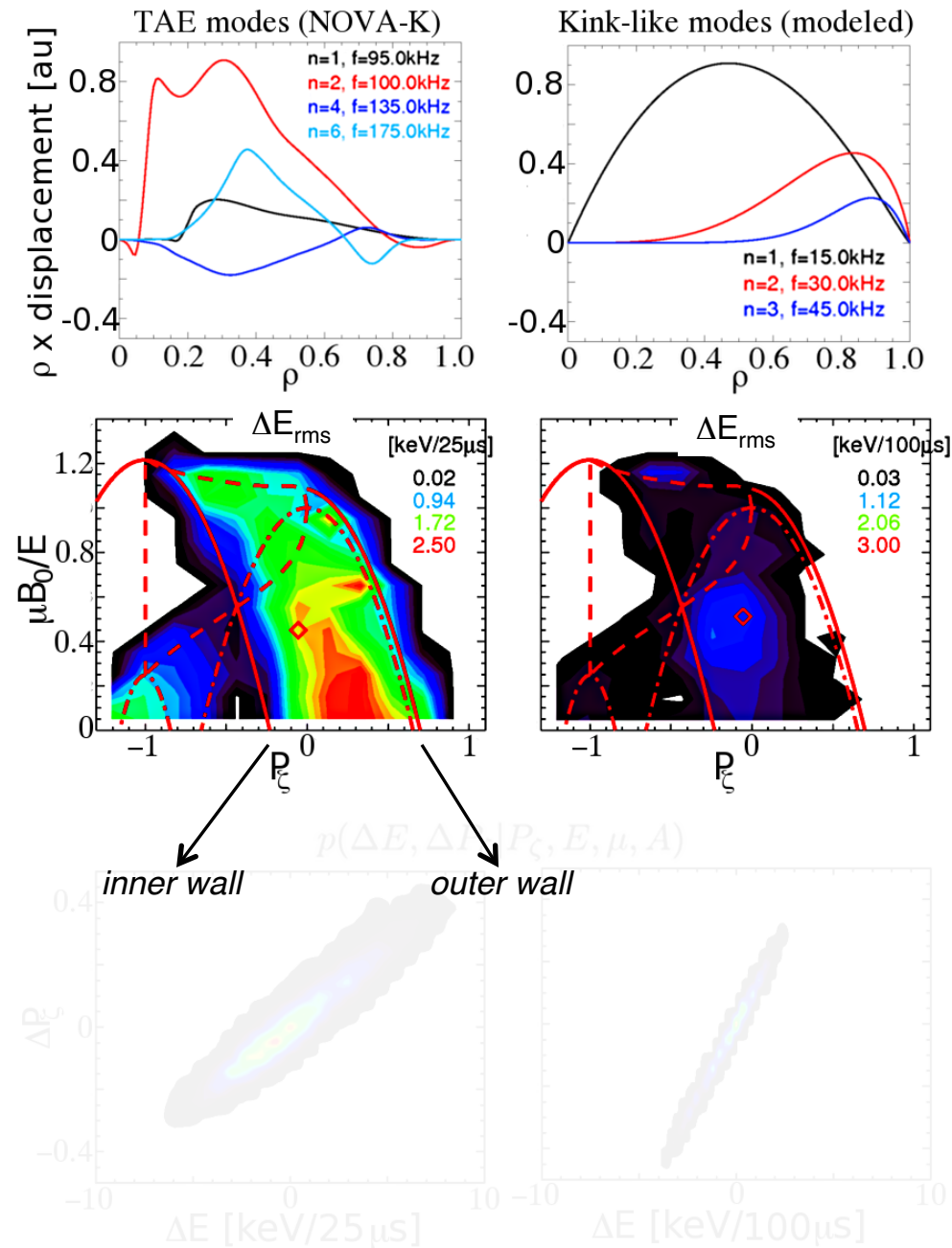
# $p(\Delta E, \Delta P_\zeta | P_\zeta, E, \mu)$ and a time-dependent ‘mode amplitude scaling factor’ enable multi-mode simulations

- Example: toroidal AEs (TAEs) and low-frequency kink
- $p(\Delta E, \Delta P_\zeta | P_\zeta, E, \mu)$  from particle-following code ORBIT
- Each type of mode has separate  $p(\Delta E, \Delta P_\zeta)$ ,  $A_{\text{mode}}(t)$
- TAEs and kinks act on different portions of phase space
- Amplitude vs. time can differ, too
- Effects on EPs differ
  - > TAEs: large  $\Delta E$ ,  $\Delta P_\zeta$
  - > kinks: small  $\Delta E$ , large  $\Delta P_\zeta$



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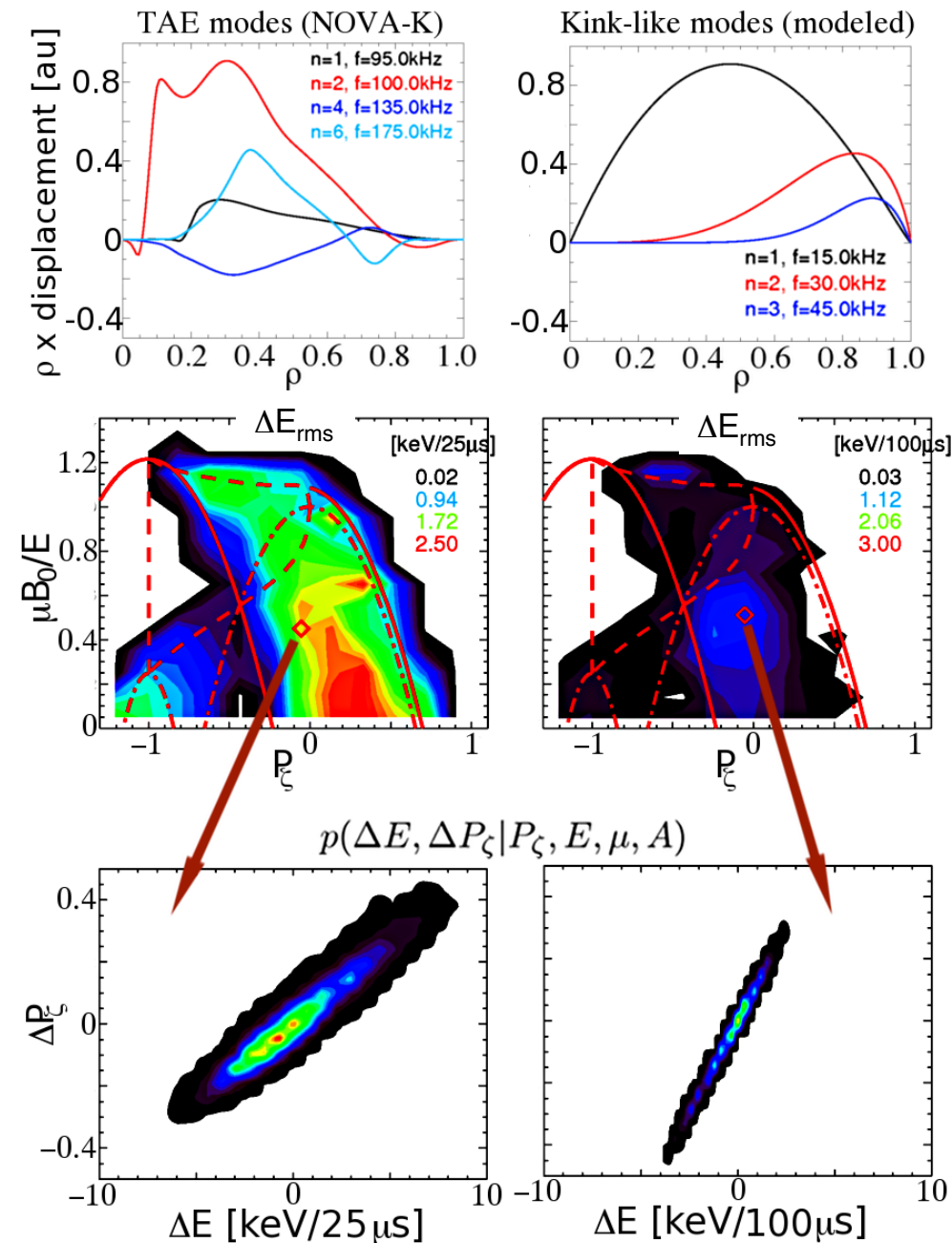
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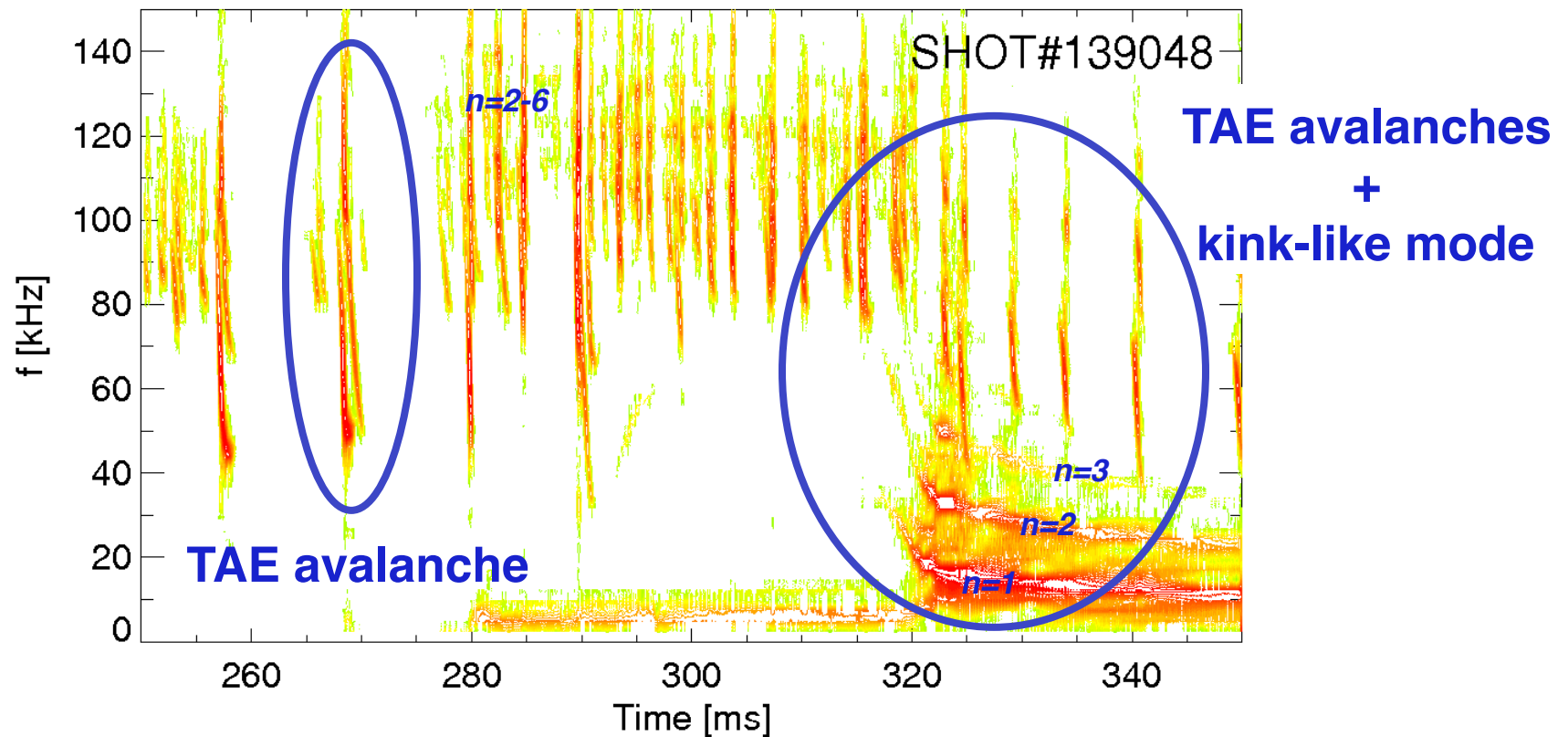


# Outline

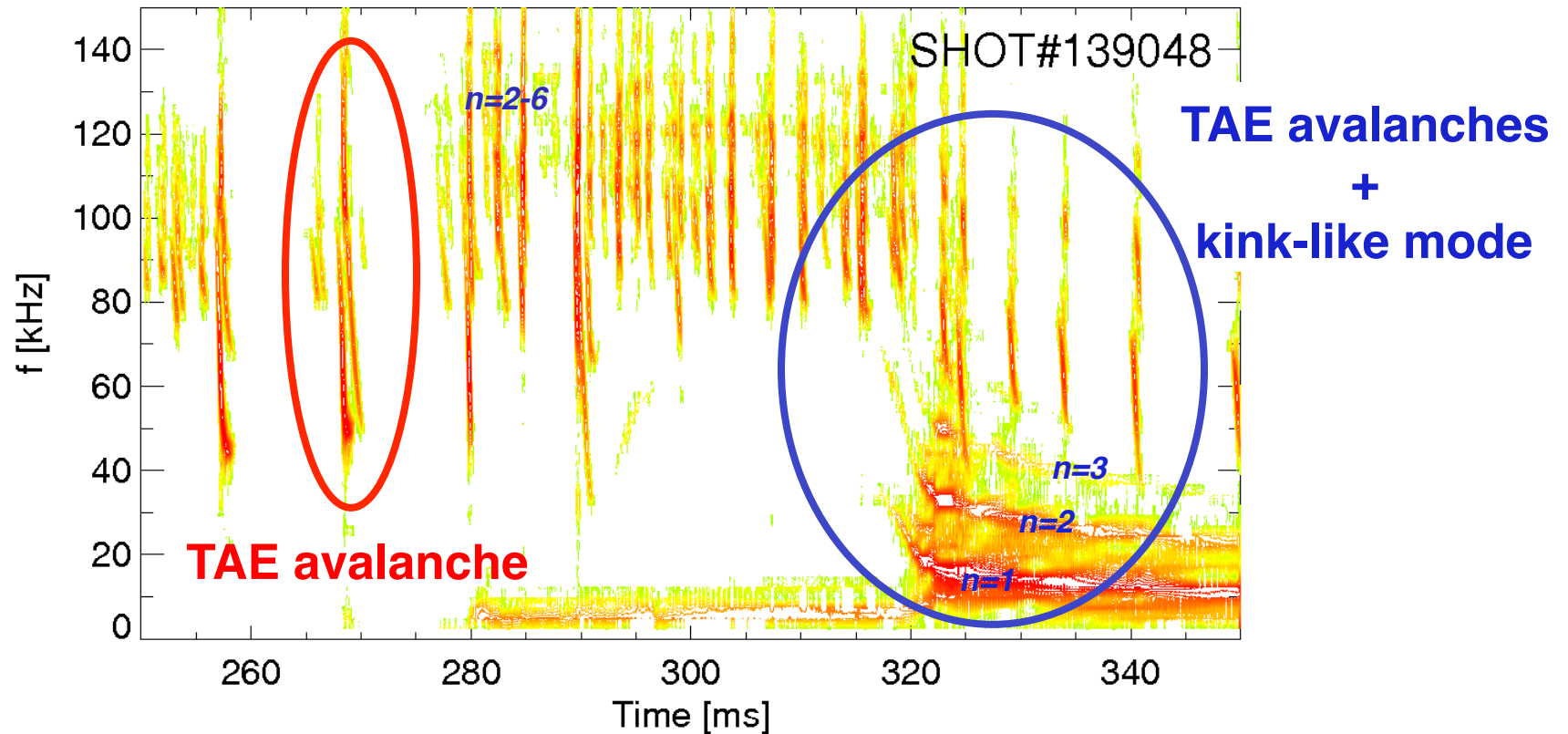
- NSTX discharges with strong MHD are used to test and validate EP transport models
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- **New model captures MHD modifications of EP phase space leading to Neutral Beam current redistribution**



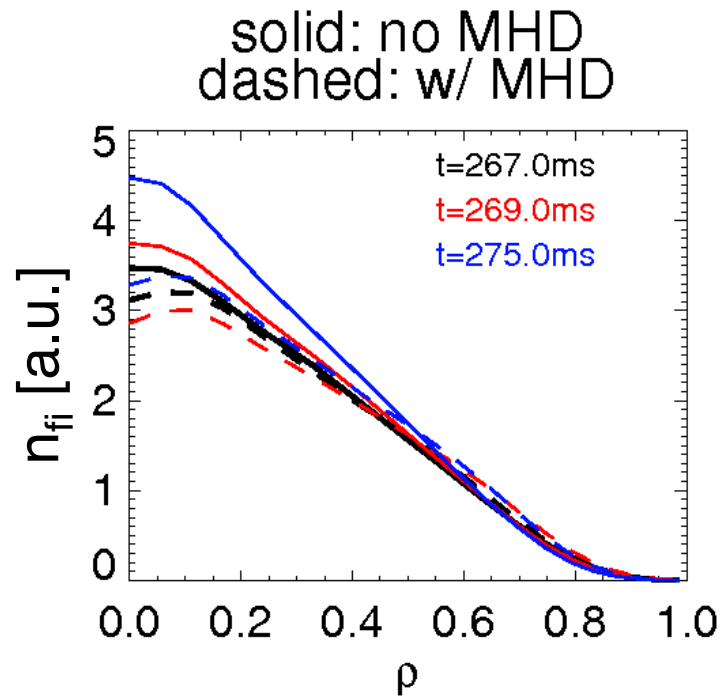
# Two NSTX cases are analyzed in detail: TAE avalanche and avalanche + kink-like mode (multi-mode scenario)



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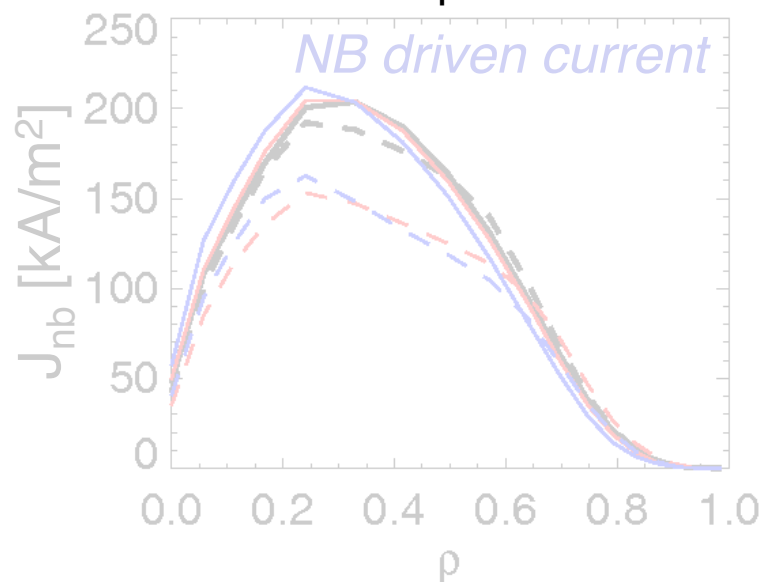


# TAE avalanches cause an abrupt drop in fast ions and up to ~40% reduction in local NB-driven current density



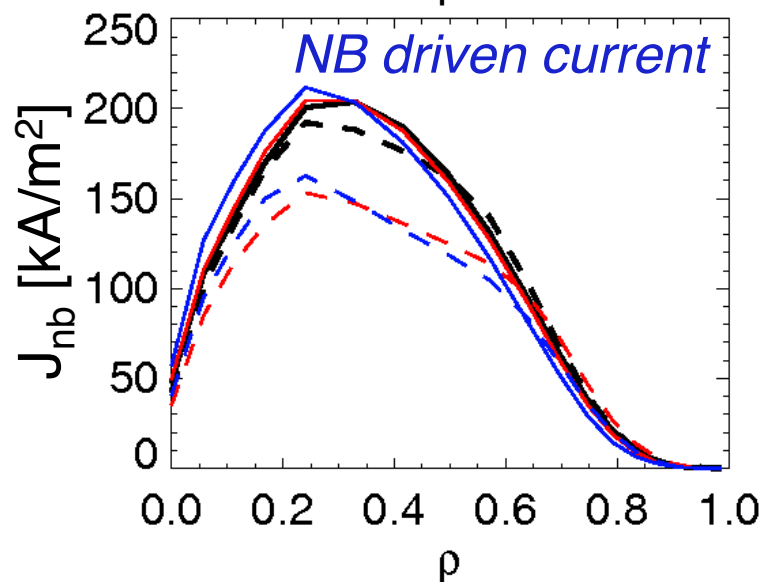
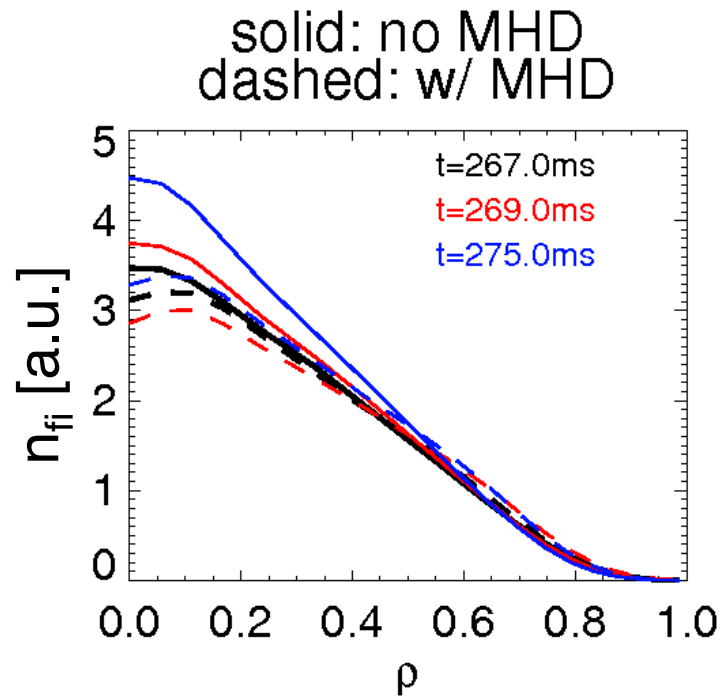
- Results from ‘kick model’
- Fast ions redistributed outward, lose energy
  - Consistent with constraints from resonant interaction:

$$\Delta P_{\zeta} / \Delta E = n / \omega$$



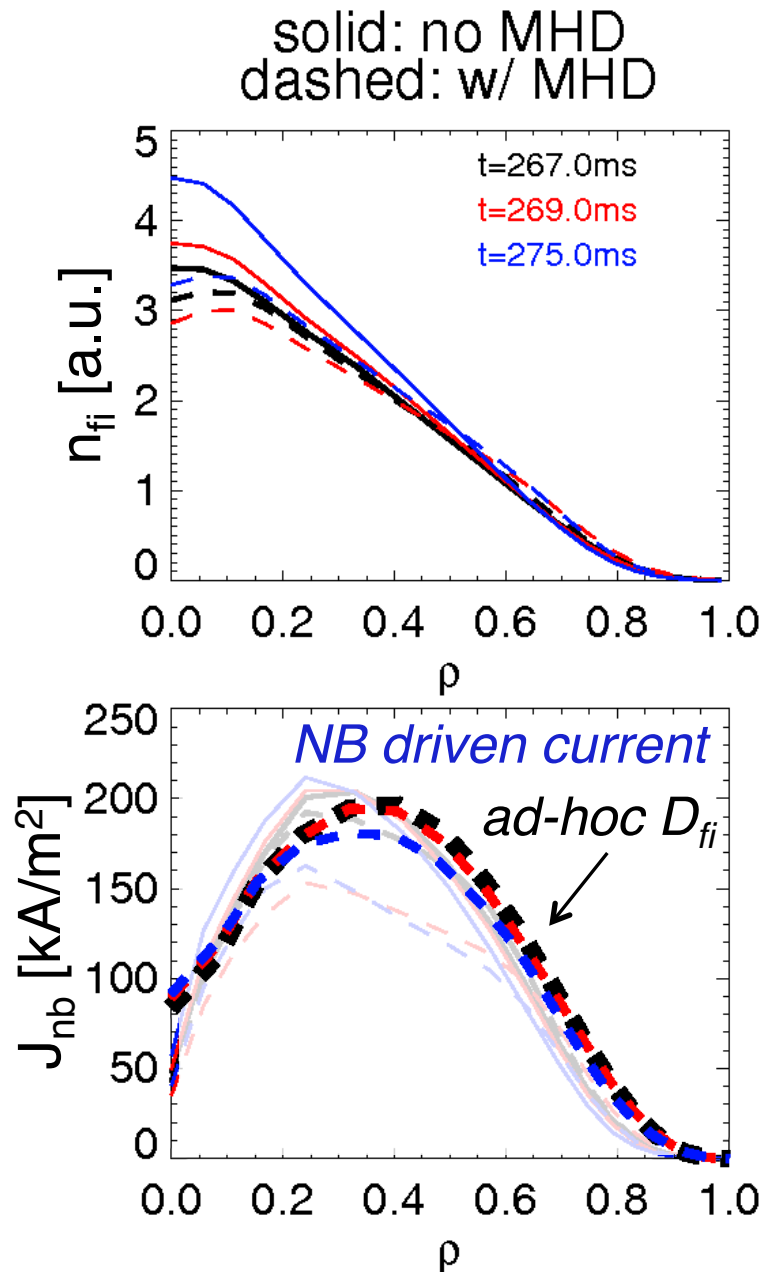
- NB-driven current  $J_{nb}$  is also redistributed out
- $J_{nb}(r)$  modification largely unpredicted by *ad-hoc*  $D_{fi}$  in this case

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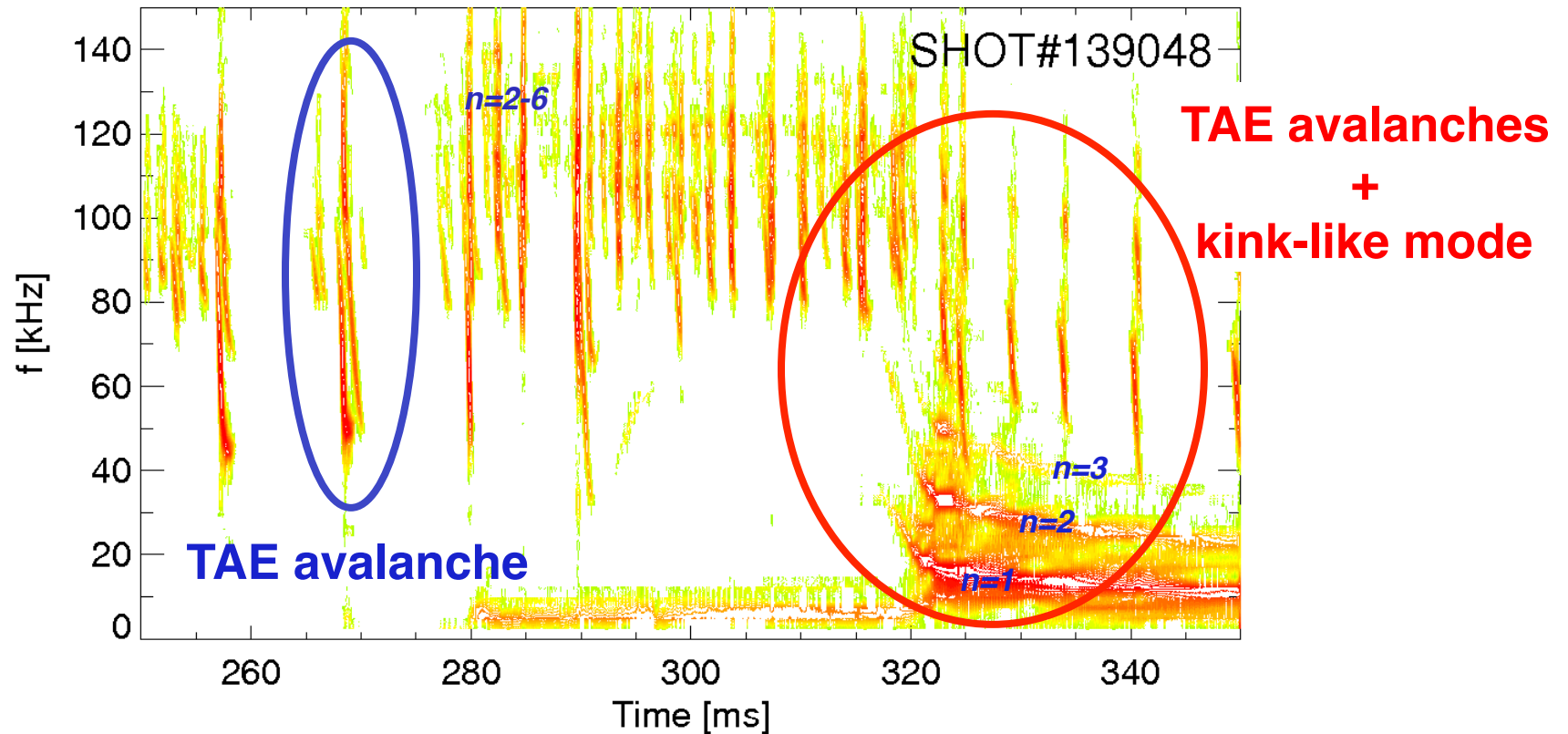
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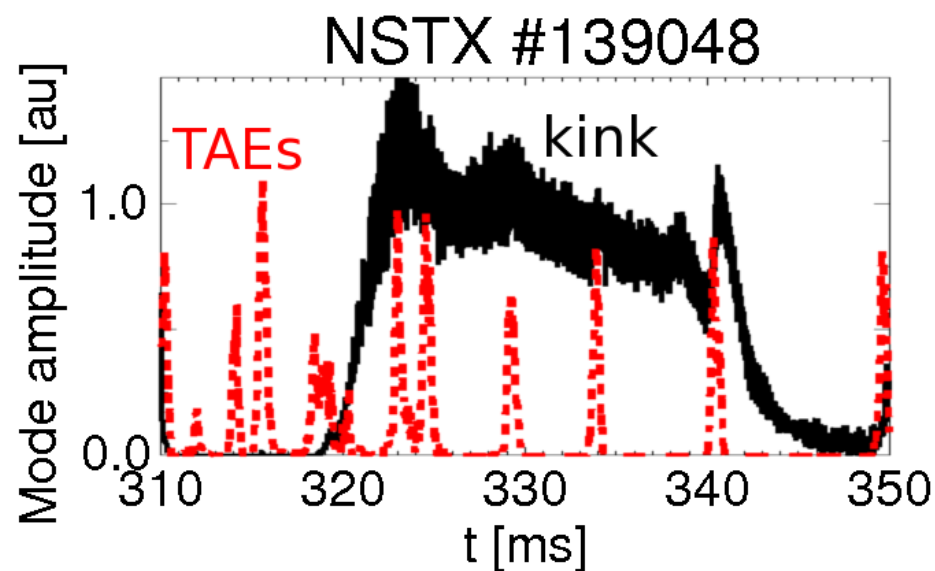


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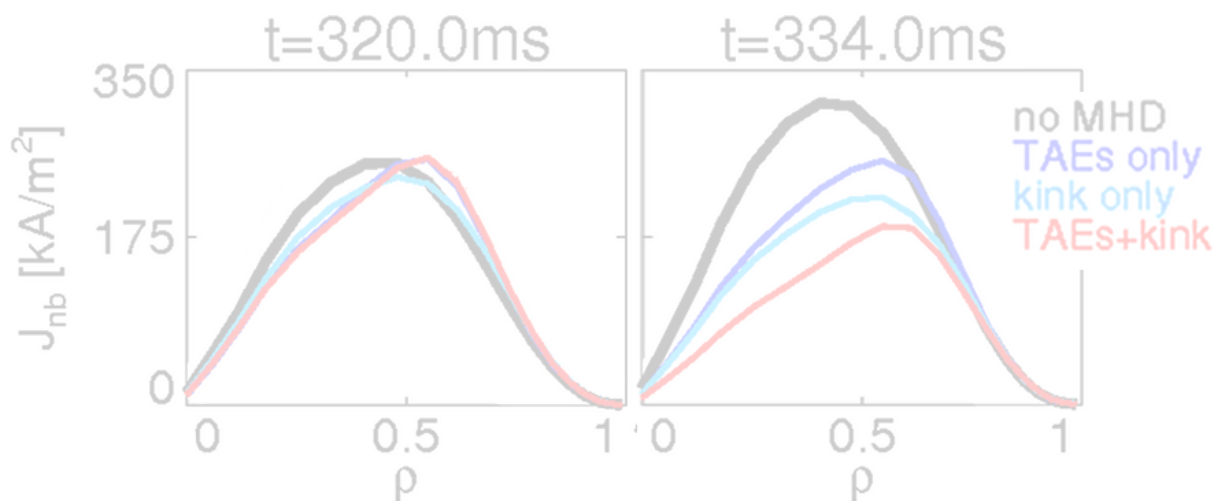
# Two NSTX cases are analyzed in detail: TAE avalanche and avalanche + kink-like mode (multi-mode scenario)



# Synergy between different classes of instabilities modifies MHD effects on $J_{nb}(r)$ – not captured by *ad-hoc* $D_{fi}$

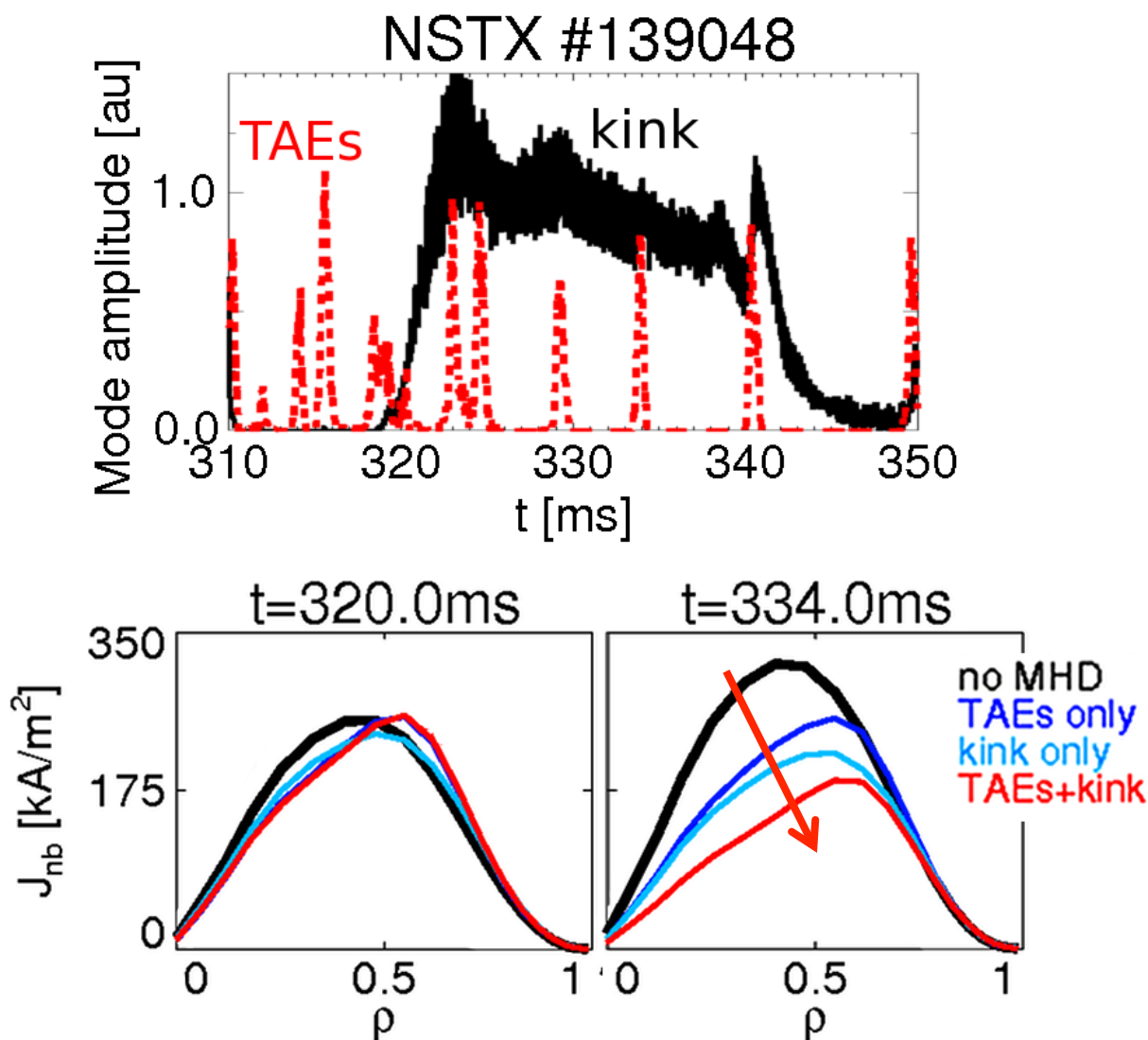


- Kinks have broad radial structure, connect core to boundary



> Synergy arises from mode overlap in phase space

# Synergy between different classes of instabilities modifies MHD effects on $J_{nb}(r)$ – not captured by *ad-hoc* $D_{fi}$

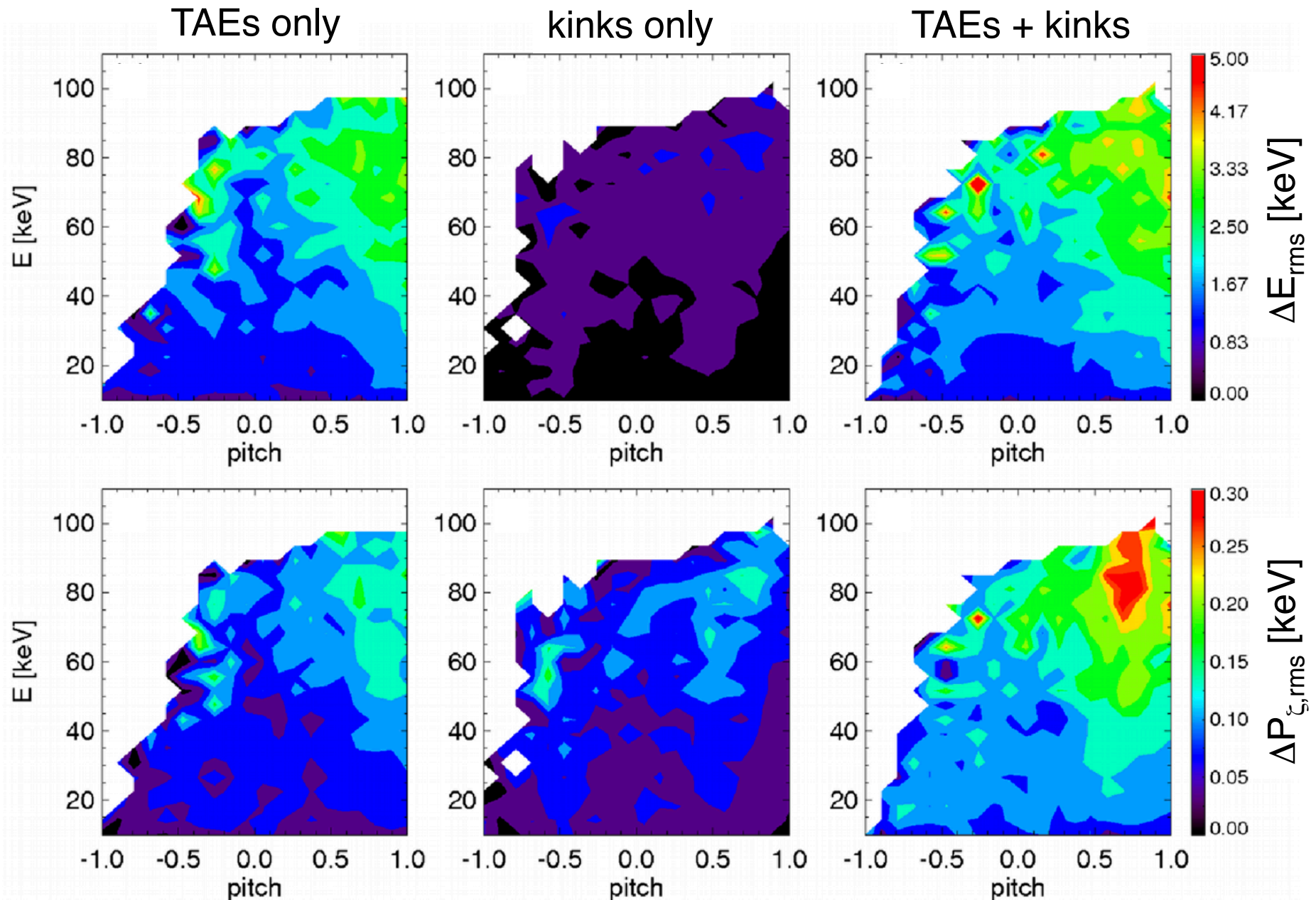


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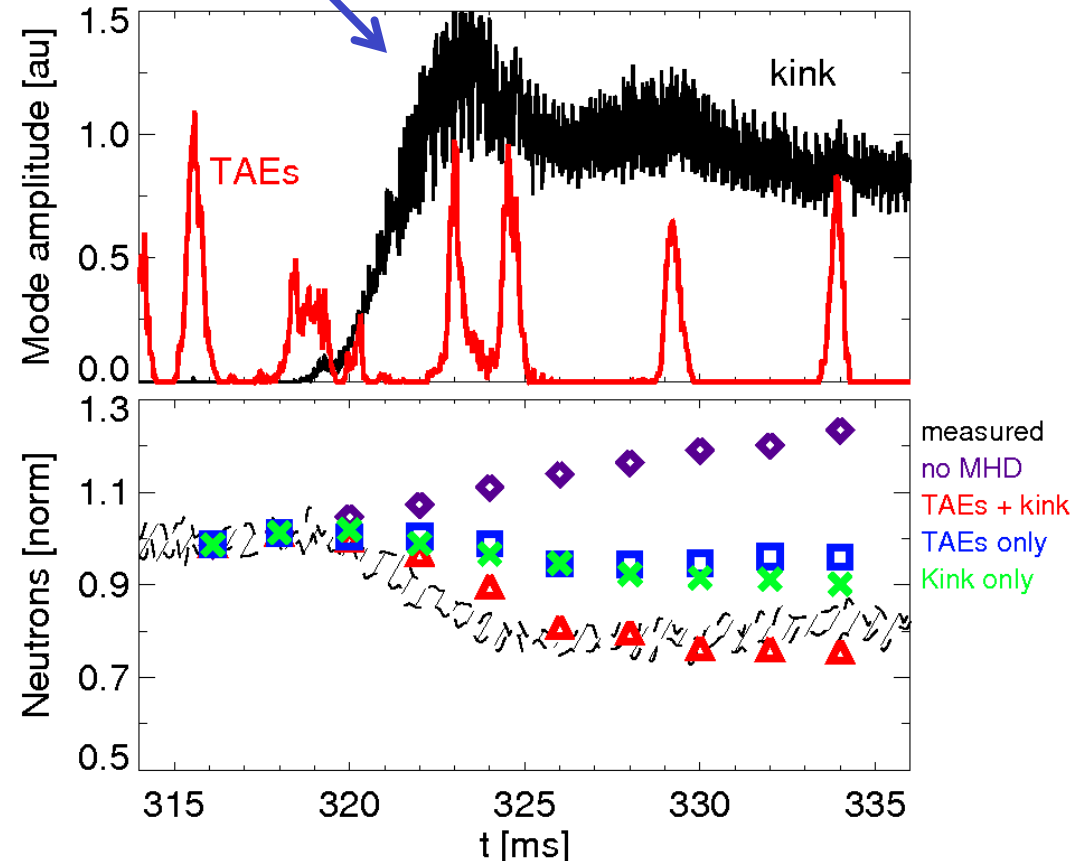
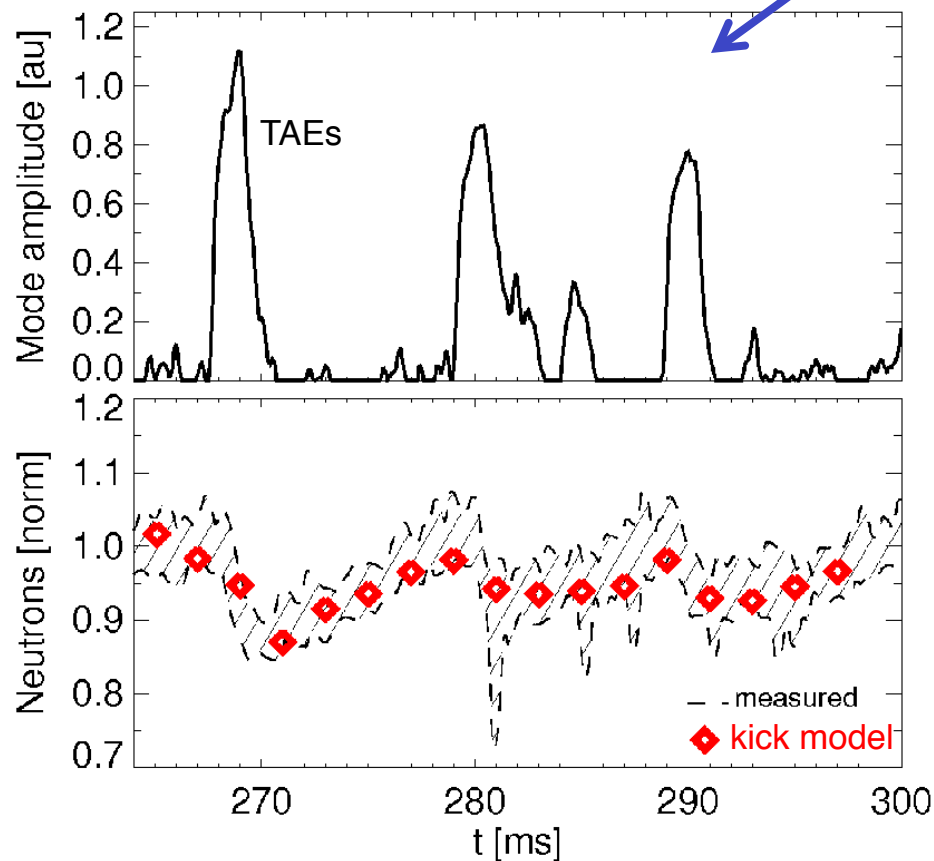
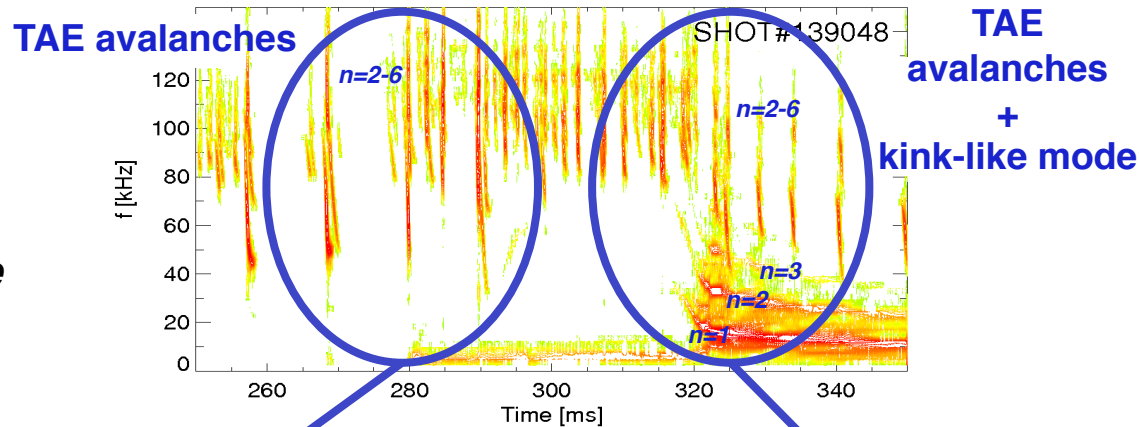


# Phase-space is *selectively* modified by instabilities: TAEs $\rightarrow \Delta P_\zeta / \Delta E = n/\omega$ , kinks $\rightarrow$ mostly $\Delta P_\zeta$



# Simulated neutron rate agrees with experiments for both TAE avalanches & multi-mode cases

Use 'kick model'  
coupled to stand-alone  
NUBEAM



# Summary

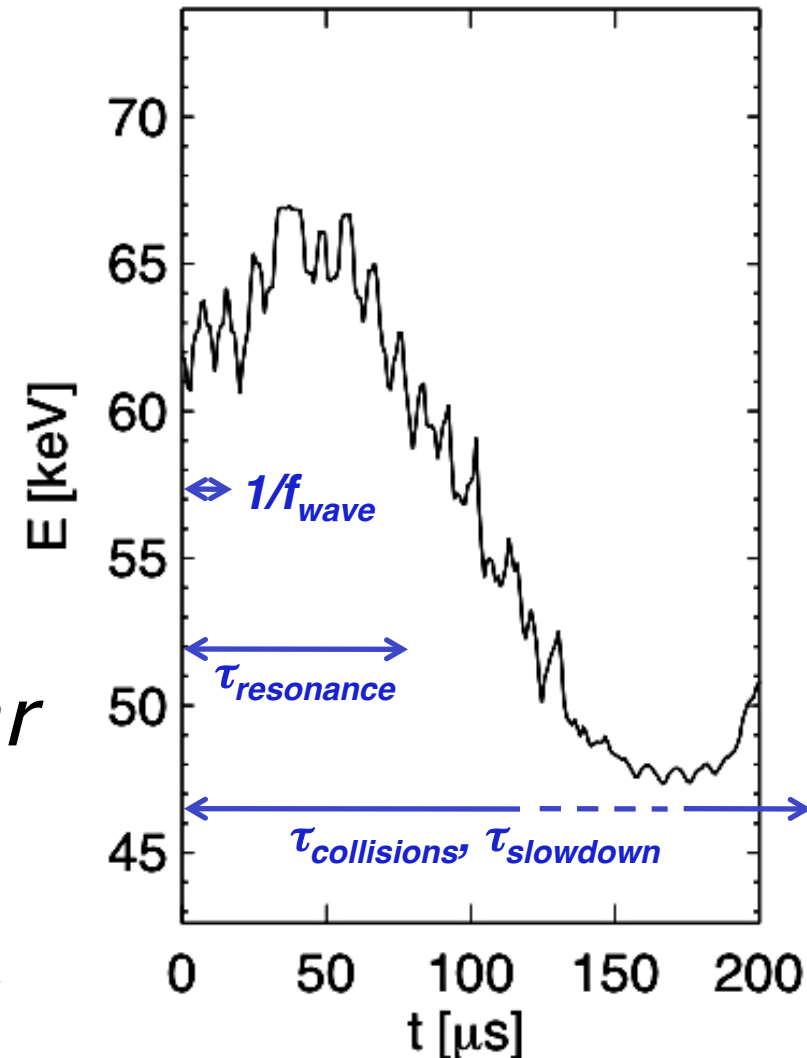
- NB-driven current profile can be strongly affected by MHD instabilities
  - Not all effects properly captured by classical EP physics
- A new model is implemented in TRANSP for EP simulations including phase-space details
  - Validation within TRANSP framework is in progress
- New tools will improve scenario development on NSTX Upgrade & future devices
  - NB current drive optimization
  - NB-driven current profile control for high- $q_{\min}$  steady state operations

# Backup slides

---

# 'Kick' model exploits separation of typical time scales between instabilities and collisional processes

- 3 time scales characterize particle motion in the presence of instabilities:
  - $1/f_{\text{wave}} \sim 10\text{'s } \mu\text{s}$
  - $\tau_{\text{resonance}} > 10 \times \tau_{\text{transit}} > 100\text{'s } \mu\text{s}$
  - $\tau_{\text{collisions}}, \tau_{\text{slowdown}} \gg 1 \text{ ms}$
- Relevant time scale for *secular*  $\Delta E, \Delta P_{\xi}$  by waves is  $\tau_{\text{resonance}}$
- Classical mechanisms already included in IM codes (TRANSP)
  - E.g. collisions, slowing down, atomic physics



## Reduced models offer advantages for Integrated Modeling (IM), plasma control over *first-principles* codes

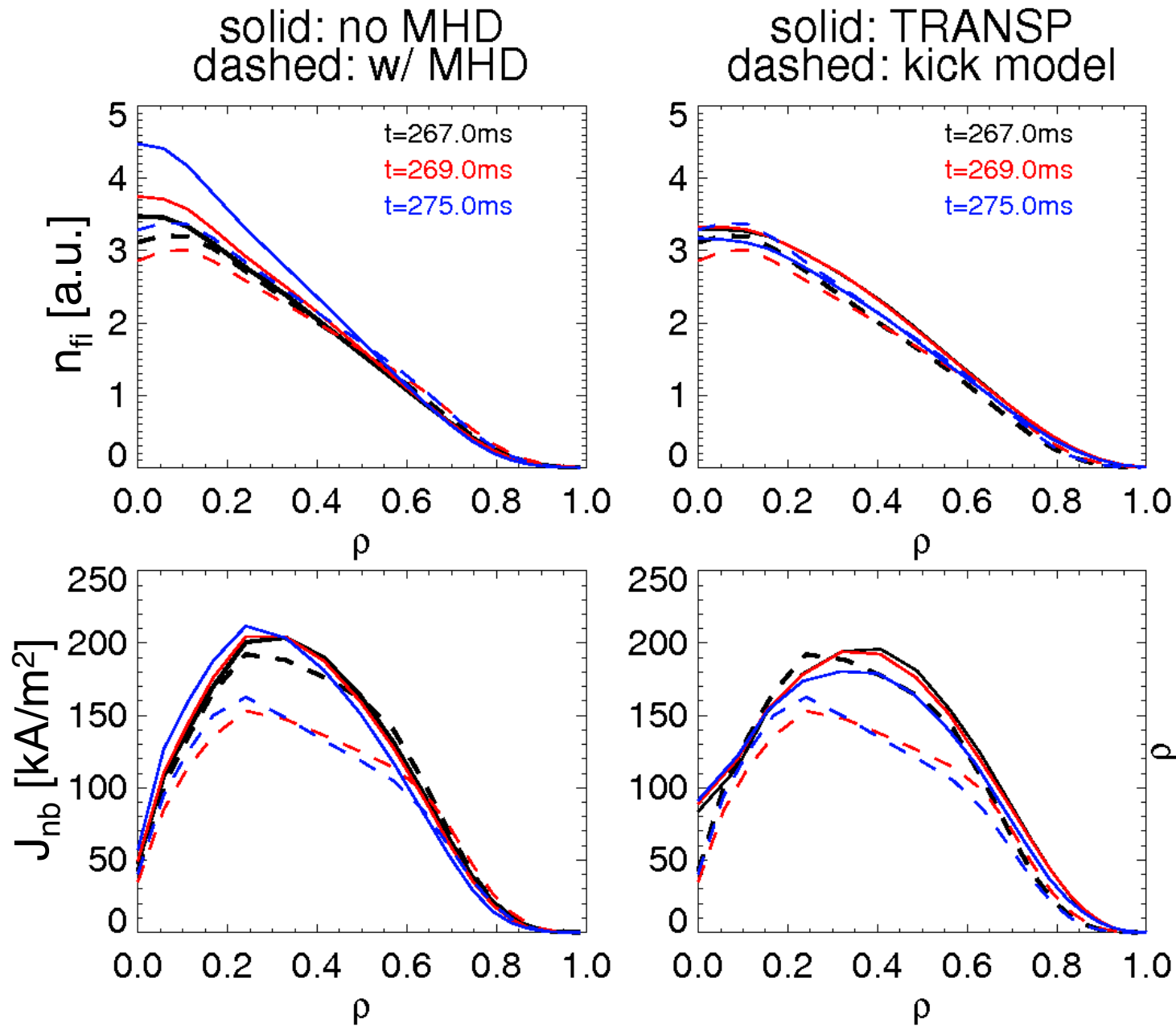
- *First-principles* codes not (yet) suitable for extensive ‘scans’ with multiple shots, long time-scale simulations
  - Inclusion in real-time control schemes also impractical
- IM codes (e.g. TRANSP) have accurate treatment of atomic physics, ‘classical’ mechanisms
  - Reduced models for EP transport are good complement
- IM codes have much broader scope than just EP physics
  - Physics-based reduced models improve accuracy of simulations, retaining ‘generality’ of IM codes

# Summary comparison of some reduced models used for EP transport

		<i>ad-hoc <math>D_{fi}</math></i>	<i>CGM model (*)</i>	<i>kick' model</i>
<i>physics-based</i>		no	yes	yes
<i>required input</i>		$D_{fi}(\rho, t)$	growth/damping rates	probability, mode amplitude
<i>applicability</i>				
	<i>multi-mode</i>	indirectly	multiple AEs	AEs, kinks, NTMs. Fishbones/EPs?
	<i>steady-state</i>	yes	yes	yes
	<i>transients</i>	yes	only for $\tau > \tau_{relax}$	yes
<i>phase-space selectivity</i>		modest	no	yes
<i>predictive runs</i>		requires guess $D_{fi}$	requires mode spectrum: growth/damping	requires mode spectrum, amplitude
<i>improvements</i>		none planned	extend to 2D in velocity space	remove $\mu$ conservation

(\*) CGM – Critical Gradient Model  
see Gorelenkov TH/P1-2, Heidbrink EX/10-1

# Simulations with *ad-hoc* $D_{fi}$ show similar fast ion drops, but largely underestimate $J_{nb}(r)$ modification

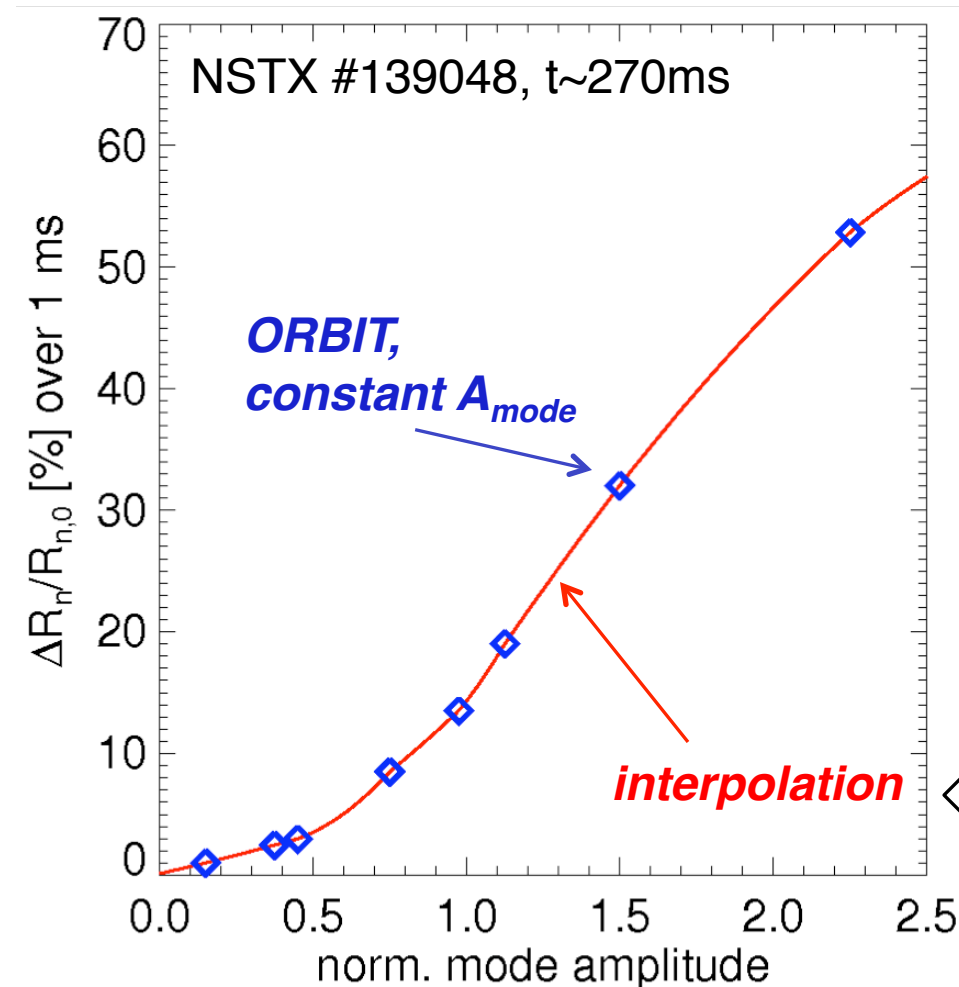


- Uniform  $D_{fi}$  acts in the same way on *all* particles at *all* radii
- No constraints from wave-particle interaction



# Scaling factor $A_{mode}(t)$ is obtained from measurements, or from other *observables* such as neutron rate + modeling

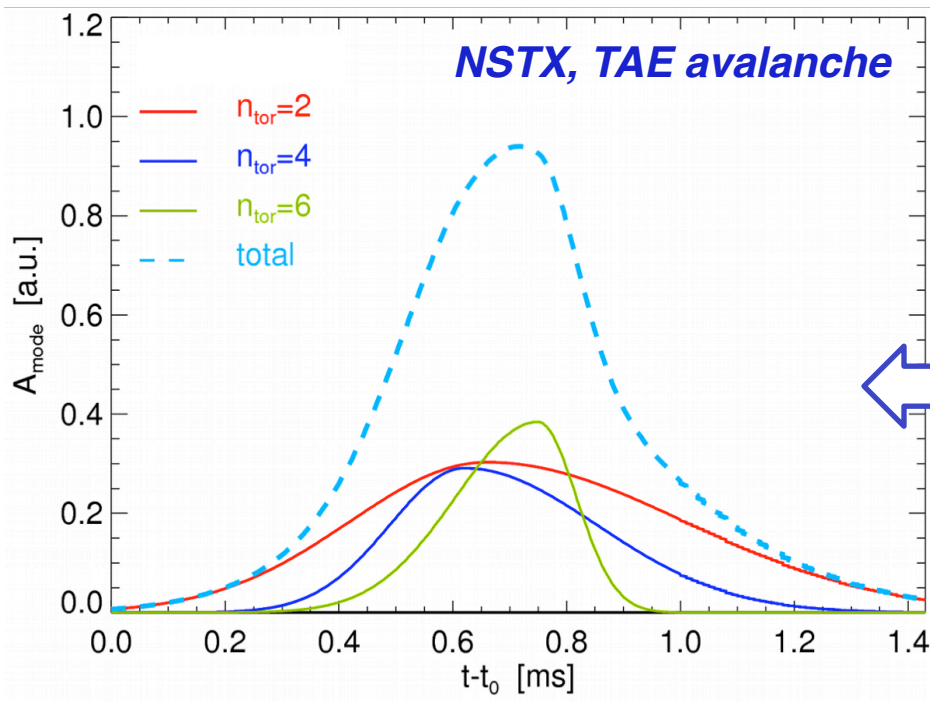
- If no mode data directly available,  $A_{mode}$  can be estimated based on other measured quantities



## Example: use measured neutron rate

- Compute ideal modes through NOVA
- Rescale relative amplitudes from NOVA according to reflectometers
- Rescale total amplitude based on computed neutron drop from ORBIT
- Scan mode amplitude w.r.t. experimental one,  $A_{mode}=1$ : get table
- Build  $A_{mode}(t)$  from neutrons vs. time, table look-up

# Mode amplitude can evolve on time-scales shorter than typical TRANSP/NUBEAM steps of $\sim 5\text{--}10\text{ ms}$

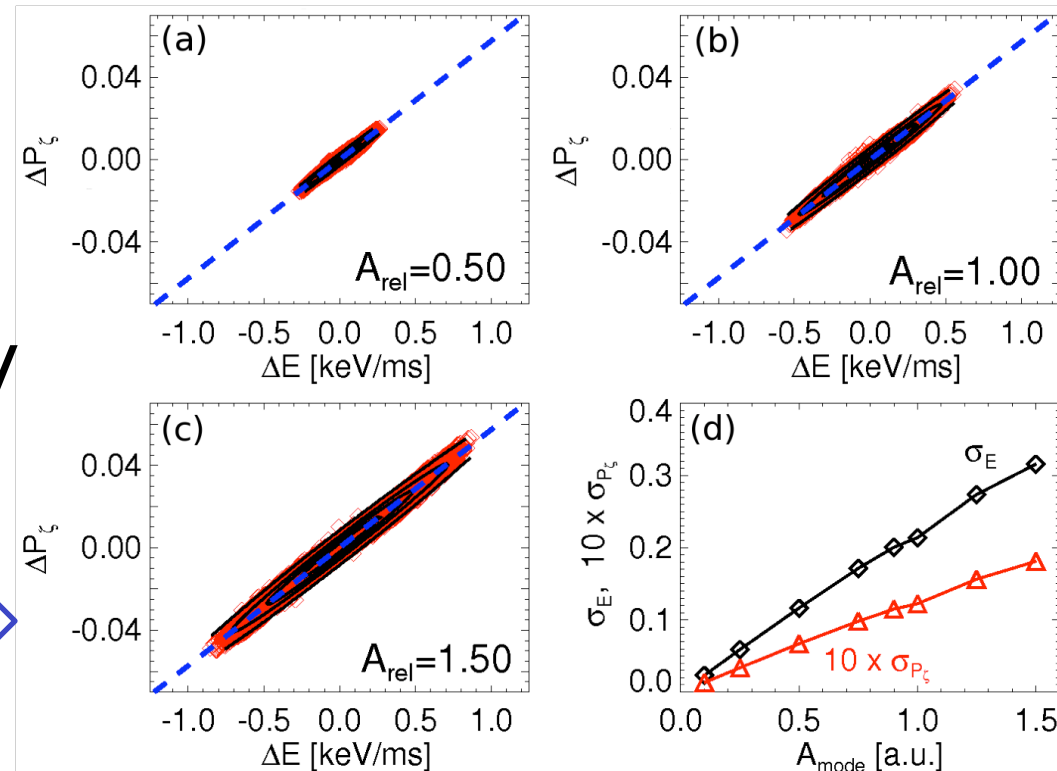


$F_{\text{nb}}$  evolution must be computed as a sequence of sub-steps

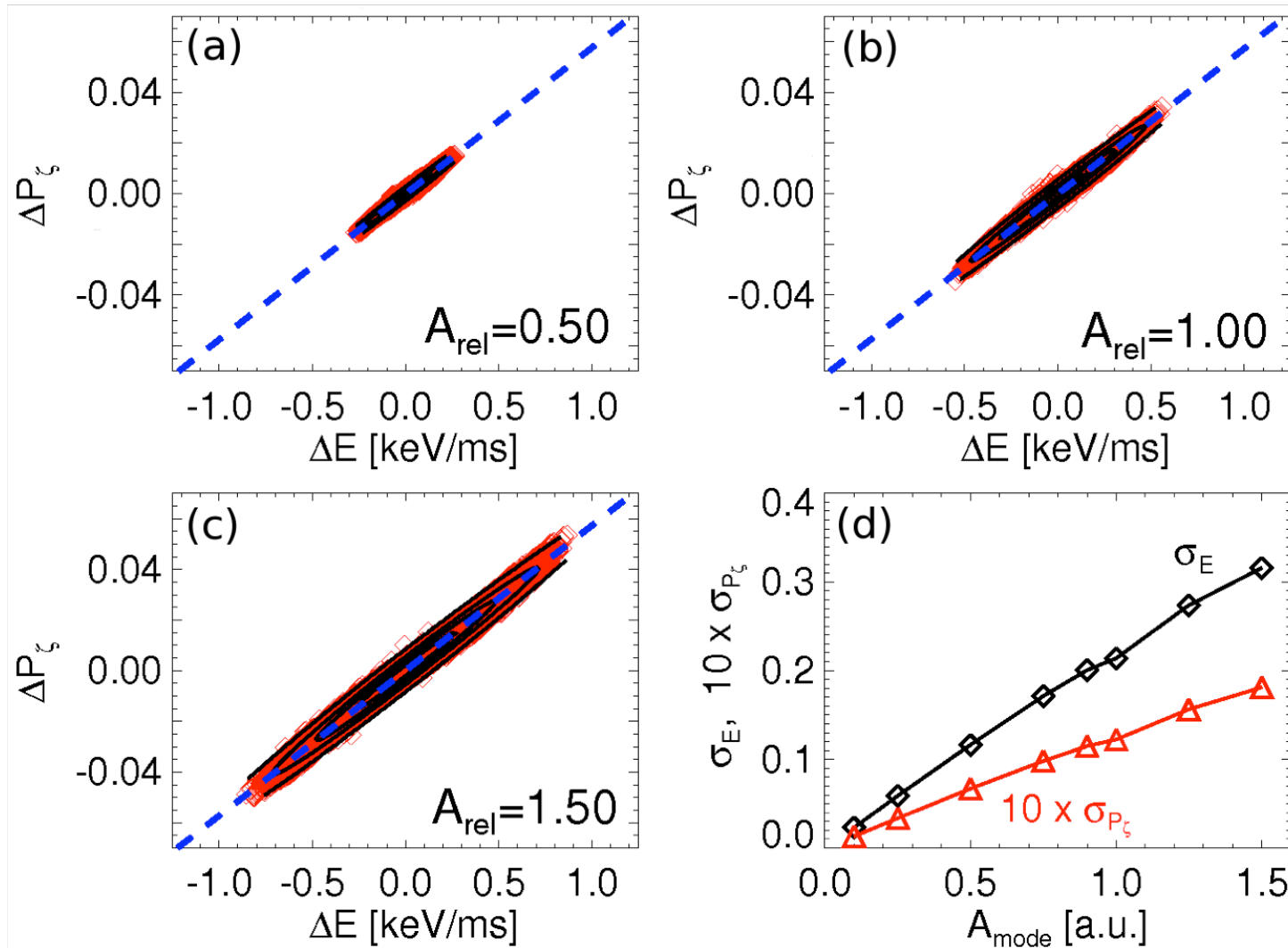
- Duration  $\delta t_{\text{step}}$  sufficiently shorter than time-scale of mode evolution
- Examples here have  $\delta t_{\text{step}} \sim 25\text{--}50\text{ }\mu\text{s}$

Energy and  $P_{\zeta}$  steps assumed to scale linearly with mode amplitude

- Roughly consistent with ORBIT simulations



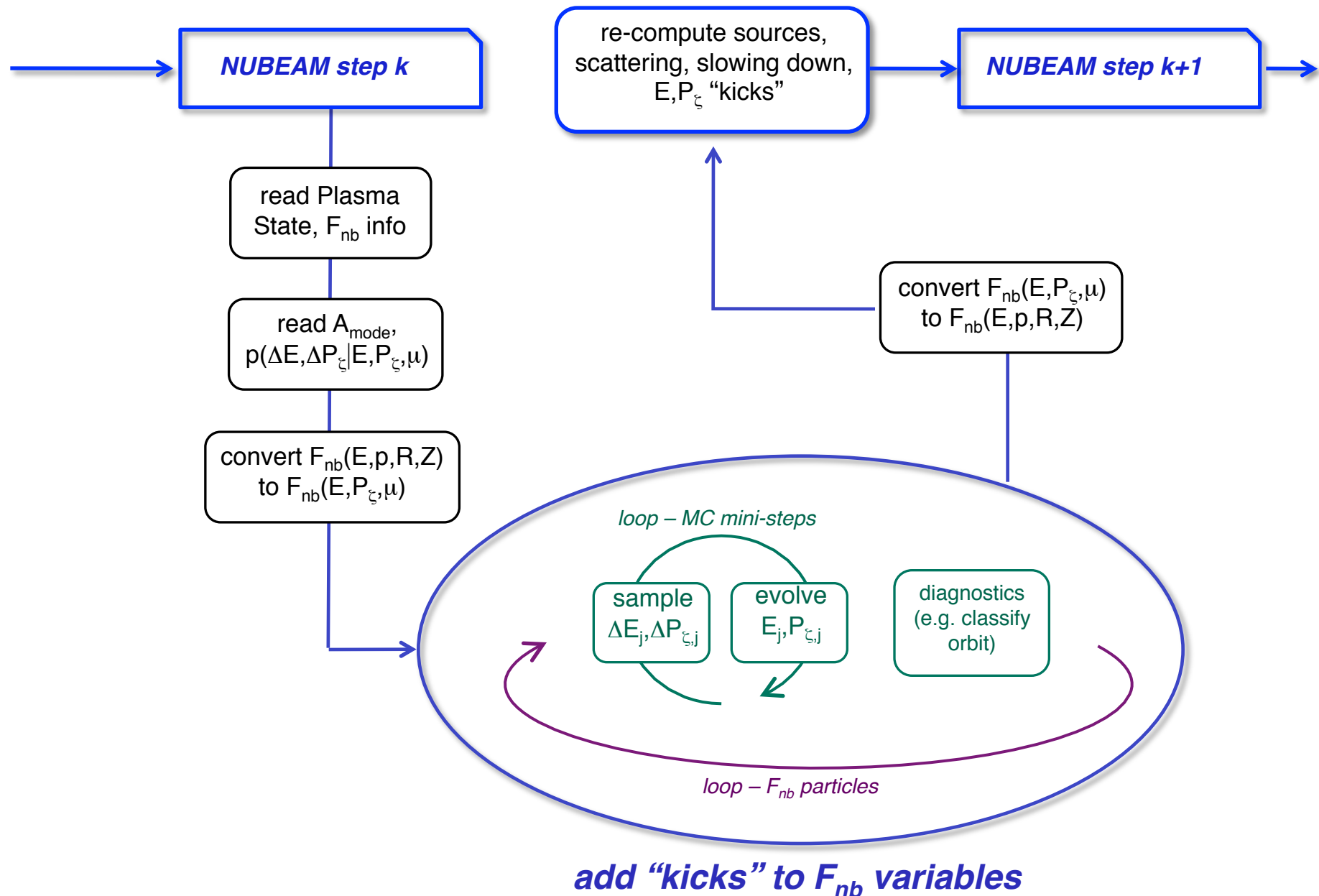
# Each type of mode is characterized by its own amplitude vs time (e.g. from experiments)



For each type of mode, energy and  $P_z$  steps assumed to scale linearly with mode amplitude

- Consistent with ORBIT simulations

# Scheme to advance fast ion variables according to transport probability in NUBEAM module of TRANSP

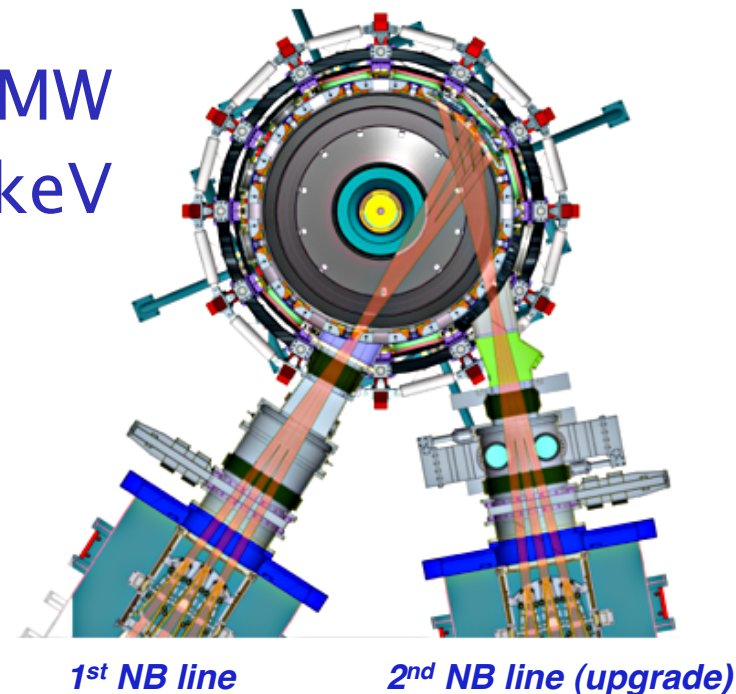


# Spherical torus NSTX is well suited for NB physics studies, model validation

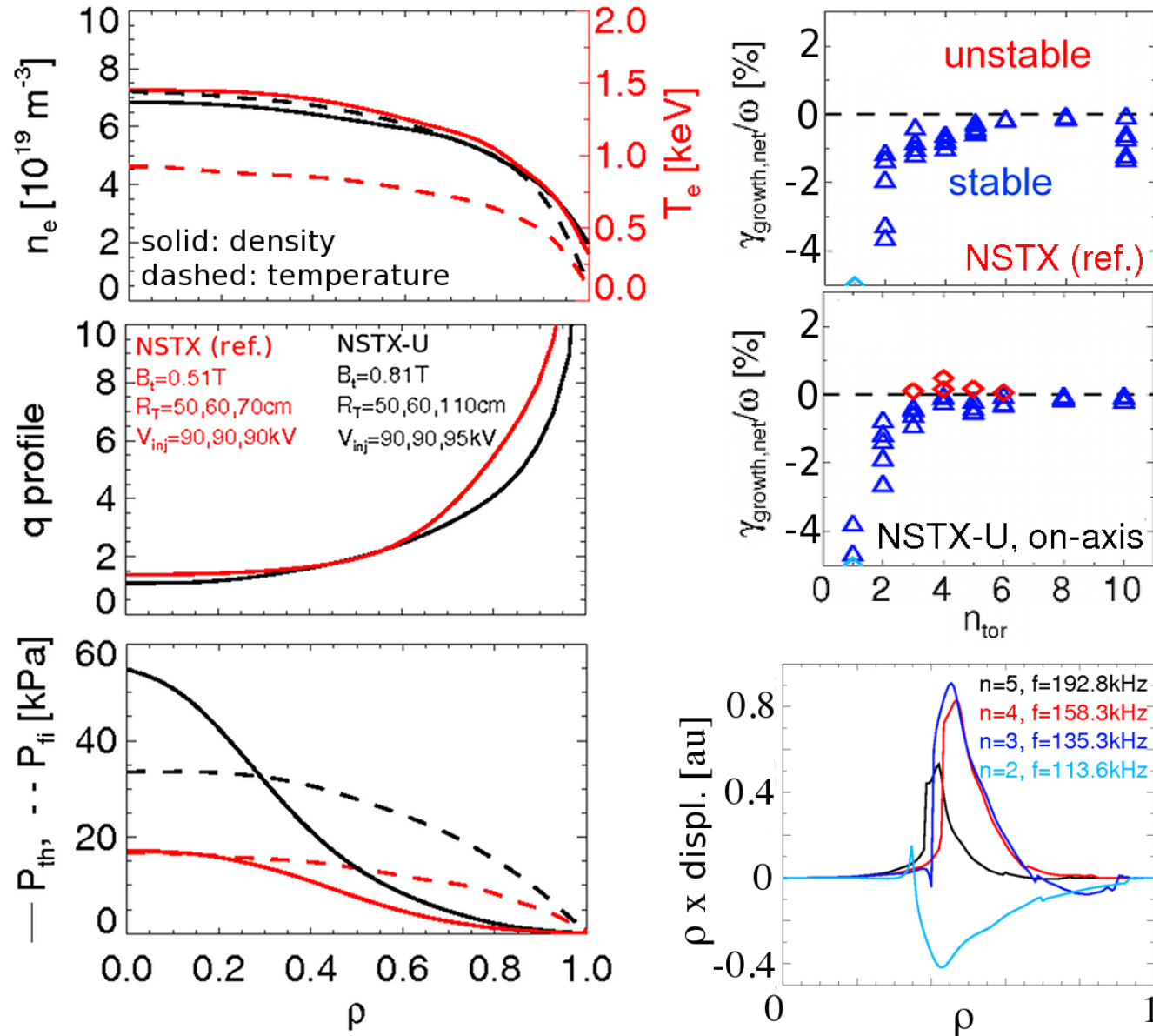
	NSTX	NSTX-U
Major radius	0.85 m	0.9 m
Aspect ratio	1.3	1.5
Plasma current	~1 MA	<2 MA
Toroidal field	<0.55 T	<1 T
Pulse length	<2 s	<5 s
Neutral Beam sources:		
$P_{\text{NBI}} \leq 6 \text{ MW}$		$\leq 12 \text{ MW}$
$E_{\text{injection}} \leq 95 \text{ keV}$		$\leq 95 \text{ keV}$

***New NBI set on NSTX-U will enable more flexible NB current drive***

[Menard, NF 2012]

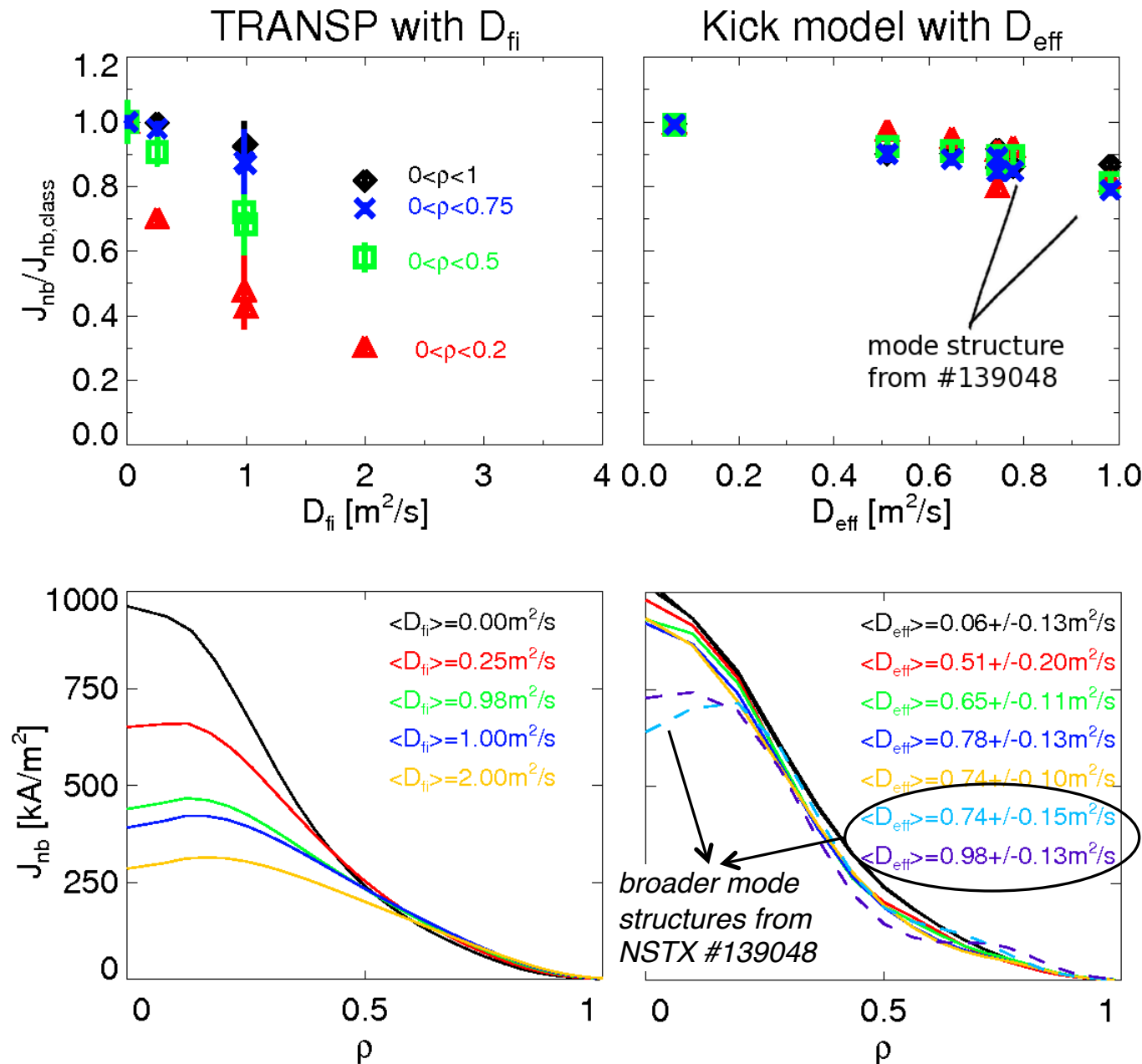


# Predicted NSTX-U scenario with strongly peaked fast ion pressure has unstable TAEs



- Fast ion pressure is  $>2$  times larger than in reference NSTX discharge
- NOVA-K finds spectrum of (linearly!) unstable TAEs with  $n=3-6$
- Predicted mode structure is narrower on NSTX-U than for typical NSTX

# ‘Kick’ and *ad-hoc* $D_{fi}$ models predict comparable reduction of total $J_{nb}$ – but profiles are very different



- Reduction in total  $J_{nb}$  is modest,  $< 20\%$
- Local  $J_{nb}(r)$  changes are much larger
- ‘Kick model’ predicts localized reduction of  $J_{nb}(r)$  because of narrow mode structures
- Non-linear physics may result in broader modes, though