

Heavy Impurity Transport in the Core of JET Plasmas

M Valisa

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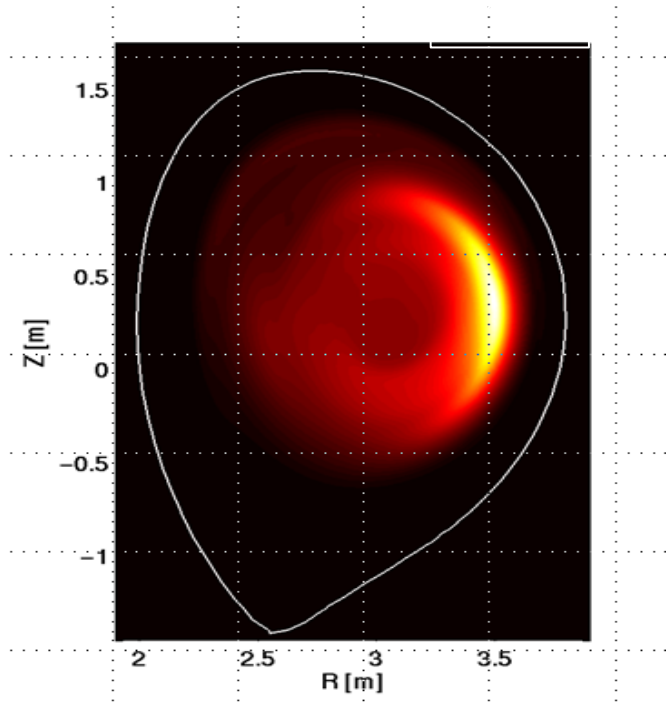
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**See the Appendix of F. Romanelli et al., Proceedings of this conference*

- **Introduction**
- **The analysis tools**
- **Results**
 - In both standard H-mode and hybrid scenarios, the path towards W accumulation is determined by the inward neoclassical convection due to density peaking of the main plasma.
 - ICRH helps hampering W accumulation in the core of standard H-mode plasmas.
- **Summary and conclusion**

- JET is studying the impact of a ITER-like wall on the plasma:
Be wall and W divertor.
- (W: $Z=74$, 193 amu; the W cooling rate remains high over a large range of T_e *T Putterich et al Nucl. Fusion 50 (2010) 025012*
- W concentration in a reactor must be kept around 10^{-5} ,
its production minimized and core accumulation avoided.

W density distribution is often highly asymmetric
as observed for heavy impurities in many experiments



This sets requirements on the modelling tools, which must include:

- 2 dimensional description for both neoclassical and turbulent transport.
- Description of the poloidal structure of the equilibrium electric potential in presence of centrifugal forces and auxiliary heating.

*L C Ingesson, H Chen, P Helander, et al. PPCF42, 161 (2000).
M L Reinke, I H Hutchinson, J E Rice, et al.. PPCF54, 045004 (2012).*

Analysis tools

Integrating the parallel force balance equation:

$$n(r, \theta) = n_0(r) \exp \left\{ -\frac{Ze\Phi(r, \theta)}{T(r)} + \frac{m\Omega^2(r)}{2T(r)} \left(R(r, \theta)^2 - R_0(r)^2 \right) \right\}$$

Poloidal angle
Toroidal rotation frequency
Major radius

the electrostatic potential

must include all possible mechanisms affecting it: in our case centrifugal effects and anisotropy heating of minority species with ICRH

Bilato Maj Angioni, NF 54, 072003 (2014)

- Goal of modelling is to compute the flux surface averaged particle fluxes

$$\Gamma = -D \frac{dn}{dr} + Vn,$$

- Different time scales \rightarrow compute turb. and neocl. coefficients separately

$$\frac{R\Gamma_W}{n_W} = -(D_{W\text{NEO}} + D_{W\text{GKW}}) \frac{R}{L_{nW}} + (RV_{W\text{NEO}} + RV_{W\text{GKW}}), \quad \text{at equilibrium}$$

- Reduce sensitivity of turb. transport to gradients using ratios between particle and heat transport channels.

Normalize turbulent transport to empirical turbulent component of the power balance heat conductivity

$$\frac{R}{L_{nW}} = - \frac{\frac{\chi_{ian}}{\chi_{i\text{NEO}}} \frac{RV_{W\text{GKW}}}{\chi_{i\text{GKW}}} + \frac{RV_{W\text{NEO}}}{\chi_{i\text{NEO}}}}{\frac{\chi_{ian}}{\chi_{i\text{NEO}}} \frac{D_{W\text{GKW}}}{\chi_{i\text{GKW}}} + \frac{D_{W\text{NEO}}}{\chi_{i\text{NEO}}}} \quad \chi_{ian} = \chi_{i\text{PB}} - \chi_{i\text{NEO}}$$

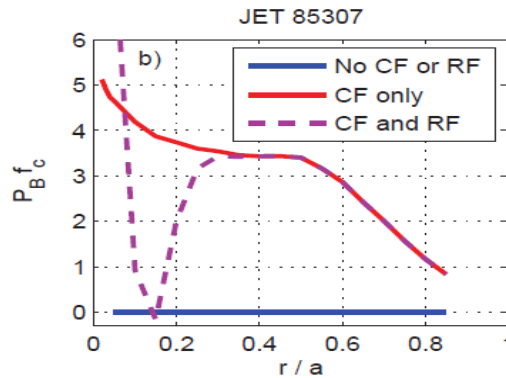
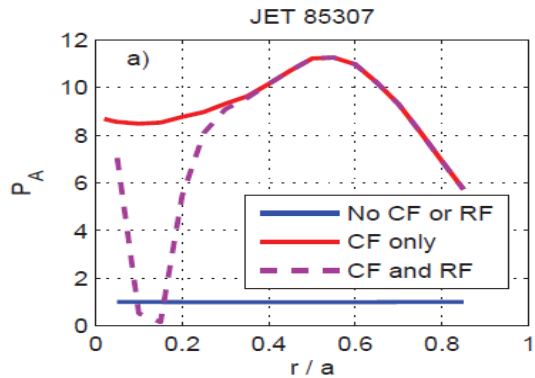
stationary, no impurity source

C. Angioni et al Nuclear Fusion 2014

Asymmetries in the electrostatic potential can strongly affect neoclassical transport

$$R \langle \Gamma_z^{\text{neo}} \cdot \nabla r \rangle \propto n_i T_i \nu_{ii} Z \left[P_A \left(-\frac{R}{L_{n_i}} + \frac{1}{2} \frac{R}{L_{T_i}} + \frac{1}{Z} \frac{R}{L_{n_z}} \right) - 0.33 P_B f_c \frac{R}{L_{T_i}} \right]$$

radial transport



fraction of passing particles

Wong PF 87;

M. Romanelli Ottaviani PPCF 98 ;

Angioni and Helander, PPCF 2014

Casson et al tbp on PPCF, <http://arxiv.org/abs/1407.1191>

Fulop Helander PoP 99;

Belli et al PPCF 2014 F ;

Theory

- Neoclassical transport: NEO *Belli PPCF 2008 and 2012*
- Turbulent transport: GKW *Peeters CPC 09, Casson PoP 10*

From the normalized density gradients the impurity densities to be compared with the experiments are derived.

Experiment

- W density recovered from SXR tomography, deconvolving W contribution from Bremsstrahlung due to hydrogen-like particles
T. Putterich et al 2012 IAEA FEC., San Diego, EX/P3-15
- JETTO/SANCO transport code to provide empirical W transport coefficients, and W densities. Based on best matching between synthetic data produced by JETTO and experimental SXR tomography and bolometry.

] Lauro Taroni L et al 1994 21st EPS Conf Montpellier, 1, (1994) 102.

Results

Electron density, initially hollow, evolves towards peaked profiles due to NBI core fuelling and Ware pinch.

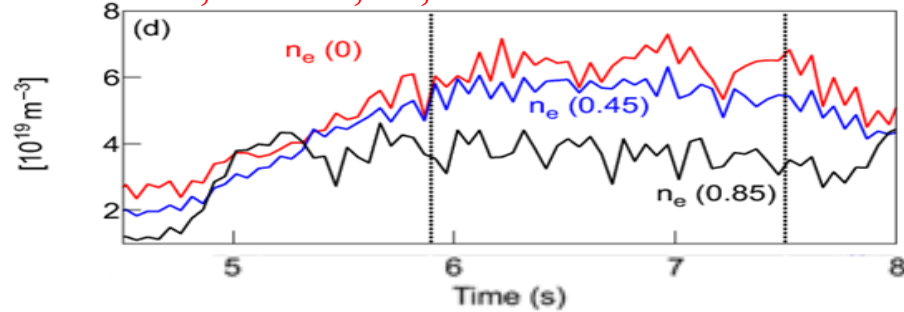
Hybrid**

#82722, 1.7 MA, 2T, 16MW NBI

P Mantica et al 40th EPS Conf., Helsinki 2013

C Giroud et al 41st EPS Conf, Berlin 2014

Loarte 2013 Nucl. Fusion 53 083031



*ne time evolution
@ three radii:
0, 0.45, 0.85 r/a*

Electron density, initially hollow, evolves towards peaked profiles due. NBI core fuelling and Ware pinch.

Hybrid**

#82722, 1.7 MA, 2T, 16MW NBI

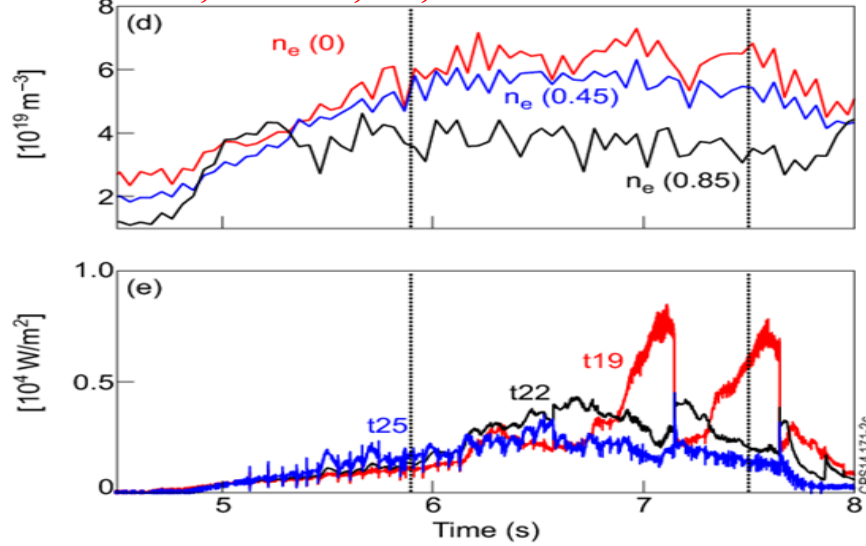
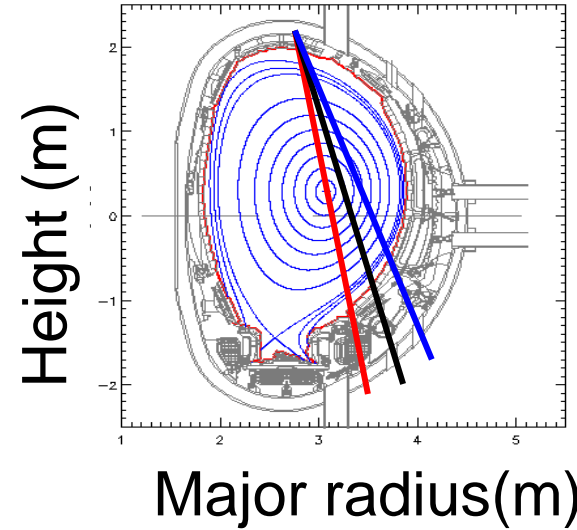
P Mantica et al 40th EPS Conf., Helsinki 2013

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*SXR LOS
Impact parameters
0, 0.2, 0.35 r/a*



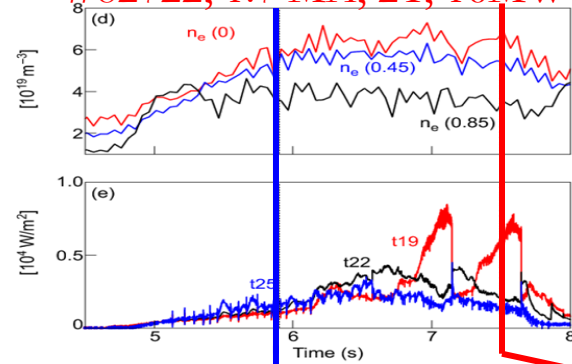
Electron density, initially hollow, evolves towards peaked profiles due. NBI core fuelling and Ware pinch.

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P Mantica et al 40th EPS Conf., Helsinki 2013

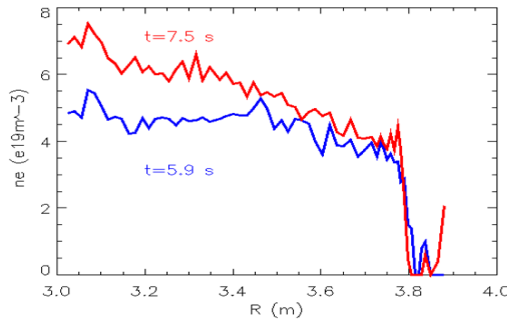
*C Giroud et al 41st EPS Conf, Berlin 2014
Loarte 2013 Nucl. Fusion **53** 083031*



*ne time evolution
@ three radii:
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*SXR LOS
Impact parameters
0, 0.2, 0.35 r/a*

*ne profiles
at selected times*



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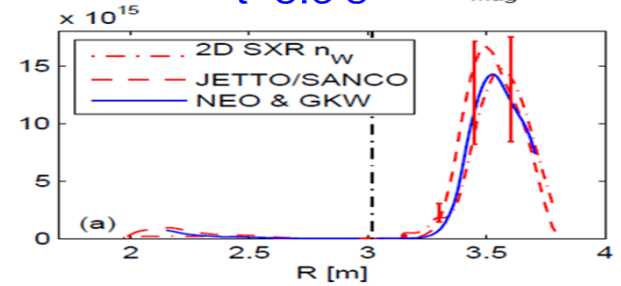
Loarte 2013 Nucl. Fusion 53 083031

$t=5.9\text{ s}$, $Z = Z_{\text{mag}}\text{ cut}$

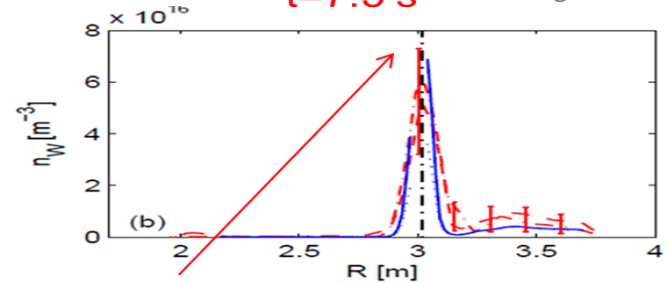
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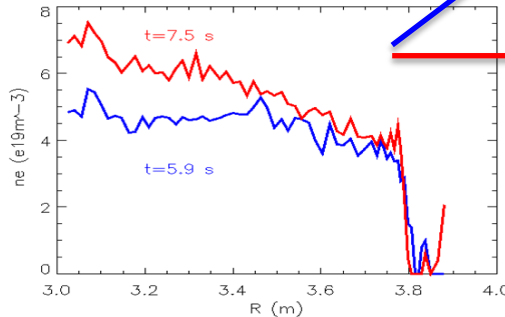
$n_W [\text{m}^{-3}]$



$t=7.5\text{ s}$, $Z = Z_{\text{mag}}\text{ cut}$



W concentration ~ 10⁻³



*ne profiles
at selected times*

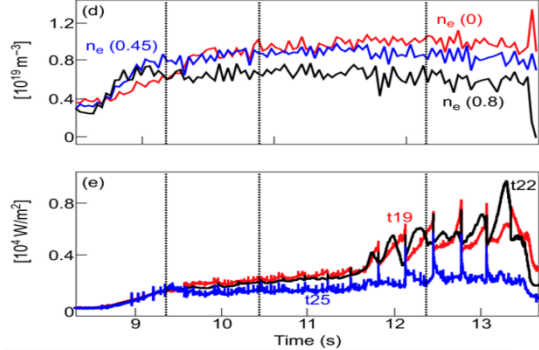
C. Angioni et al Nuclear Fusion 2014

Very similar situation for the standard Hmode.

More frequent sawteeth keep the W dynamics lower

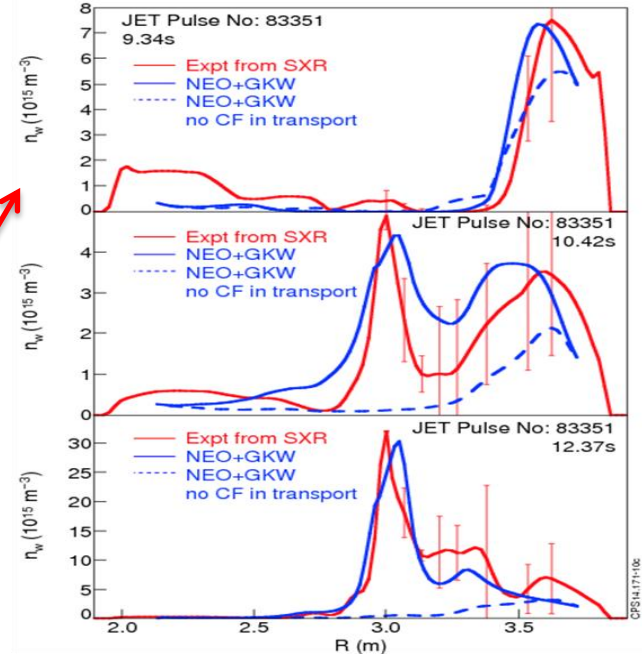
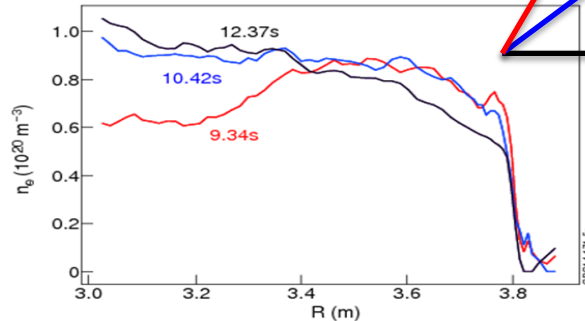
Standard H-mode

#83351, 2.75 MA, 2.6 T, 17.5 MW NBI,



ne time evolution
 @ three radii :
 0, 0.45, 0.8 r/a
SXR LOS impact parameters
 0, 0.2, 0.35 r/a

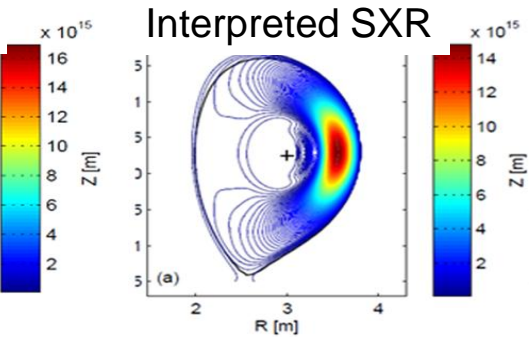
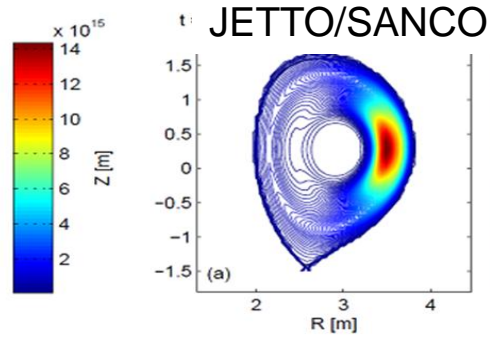
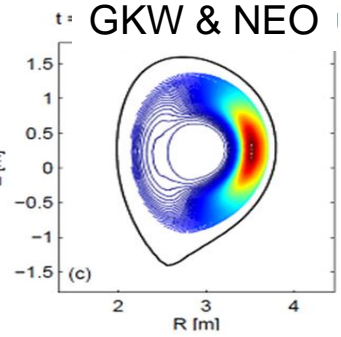
ne profiles at selected times



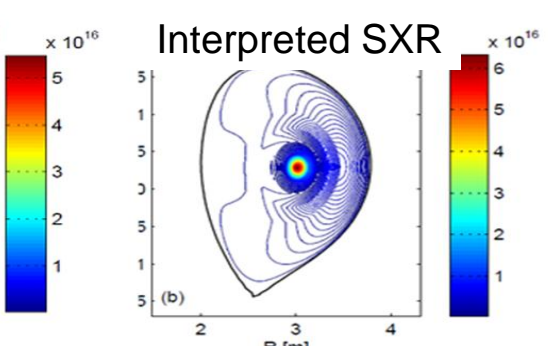
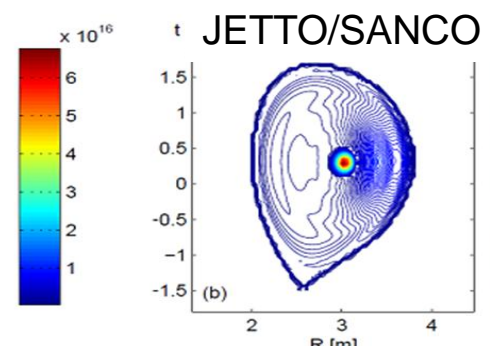
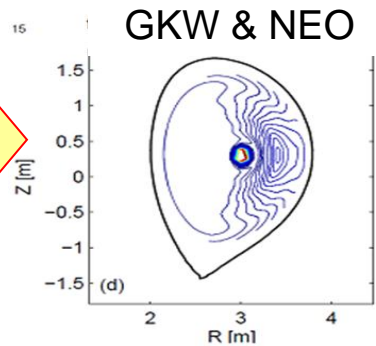
P Mantica et al 41st EPS Conf, 2014 Berlin

82722 Hybrid

Time slice
@ 5.9 s



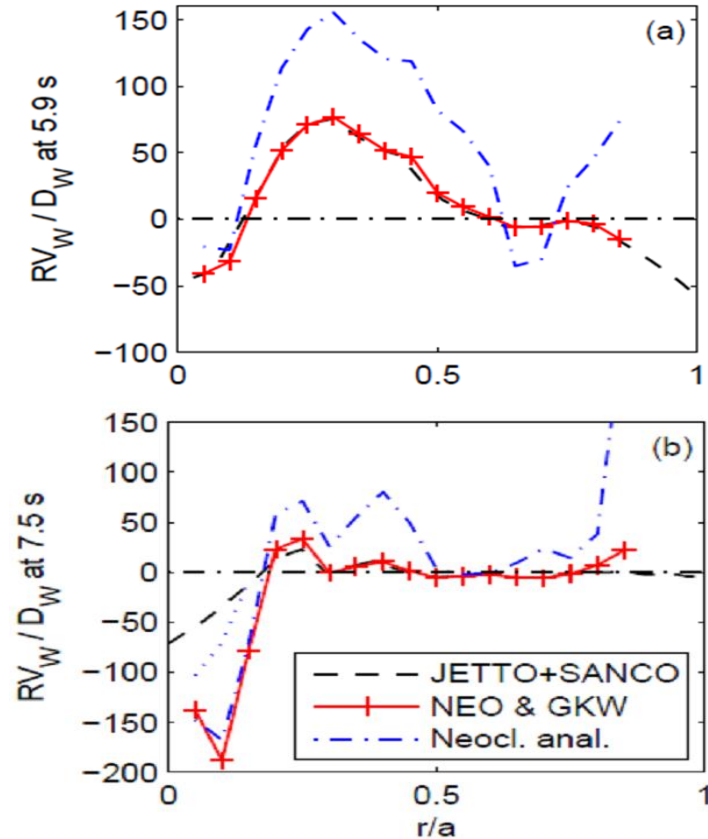
Time slice
@ 7.5 s
Center
accumulation



Time slice
@ 5.9 s

Convection to diffusion ratios
for W as computed by NEO
+ GKW and by
JETTO/SANCO

Time slice
@ 7.5 s



- MHD modes have complex interplay with W as they affect also the background kinetic profiles and thus the neoclassical transport drive.
- Sawtooth crashes clearly help flushing W out of the core.
- In presence of hollow W densities and peaked main plasma density the onset of an NTM accelerates the accumulation process. They facilitate the drift of W into inner regions where neoclassical inward pinch is particularly strong

C. Angioni et al Nuclear Fusion 2014

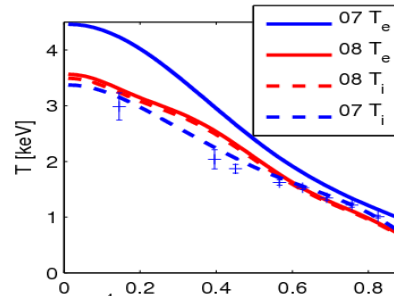
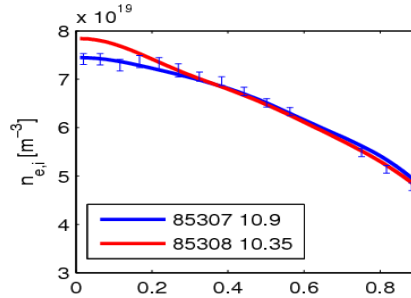
W transport and ICRH in standard H-mode

- Effects on background profiles and indirect impact on neoclassical transport of W
- Direct effects on W transport

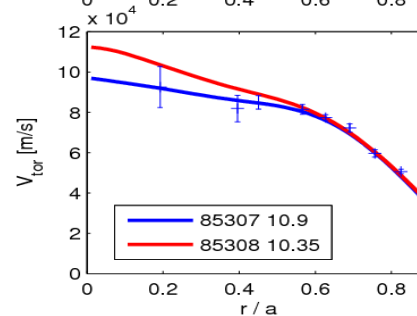
85308: 2.5 MA, 2.7 T , 19MW NBI ONLY

85307: 2.5 MA, 2.7 T , 14.7 MW NBI + 4.5 MW ICRH (H minority)**

Flatter ne



Higher Te
Similar Ti



Lower rotation

r/a

r/a

r/a

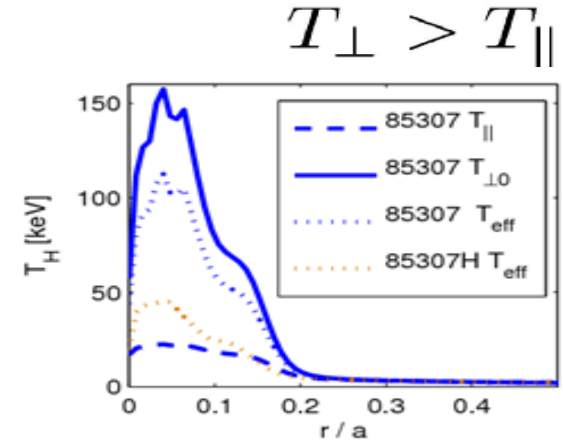
In order to match the experiment it is important to add the following mechanisms:

- Thermal screening due to minority species temperature gradients
- Anisotropy heating of minority species

$$n(\theta) = n_{R0} \frac{T_{\perp}(\theta)}{T_{\perp R0}} \exp\left(-\frac{eZ\Phi(\theta)}{T_{\parallel}} + \frac{m\Omega^2(R(\theta)^2 - R_0^2)}{2T_{\parallel}}\right)$$

$$\frac{T_{\perp}(\theta)}{T_{\perp R0}} = \left[\frac{T_{\perp R0}}{T_{\parallel}} + \left(1 - \frac{T_{\perp R0}}{T_{\parallel}}\right) \frac{B_{R0}}{B(\theta)} \right]^{-1}$$

Bilato Maj Angioni, NF 54, 072003 (2014)



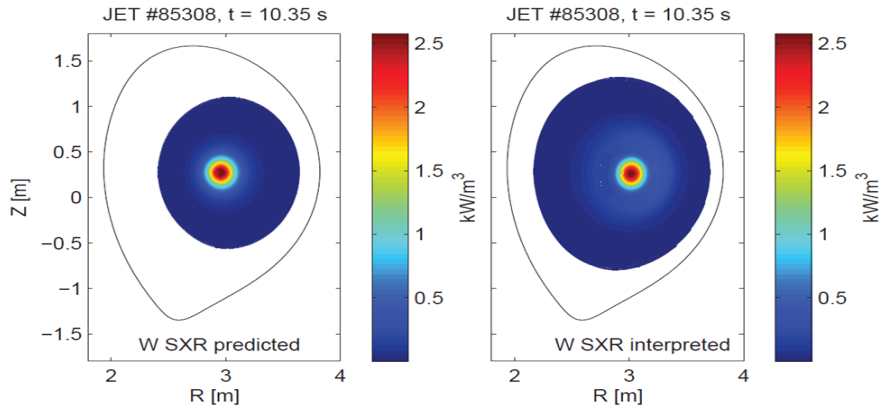
F Casson et al tbp on PPCF

from TORIC & SSPQL

R. Bilato, M. Brambilla, O. Maj, et al., Nucl. Fusion 51, 103034 (2011).

Again successful match between theory-based model and expt

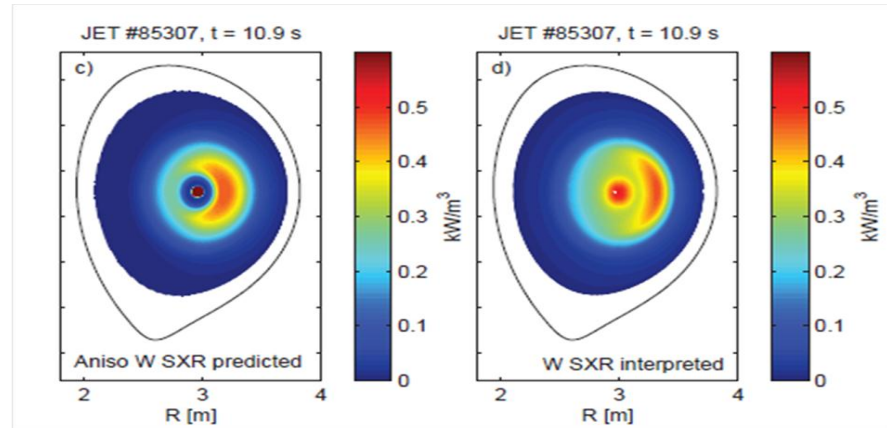
85308: NBI ONLY



Model

Experiment

85307: NBI + ICRH



Model

Experiment

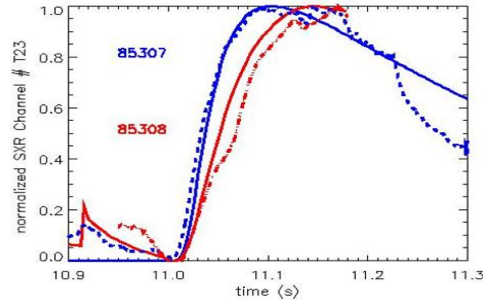
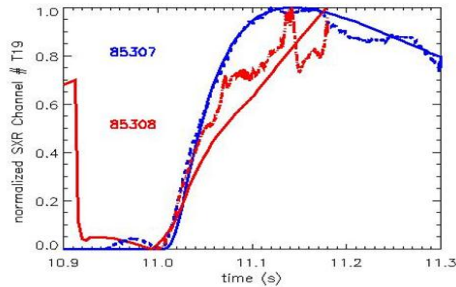
Includes anisotropy
heating of and thermal
screen by minority species

F Casson et al tbp in PPCF, <http://arxiv.org/abs/1407.1191>

R Bilato M Brambilla, O. Maj et al., Nucl. Fusion 51, 103034 (2011).

- Analysis of ICRH effects on W confirmed by LBO injections of Mo
- Simulation of Mo LBO with theory-based model coefficients fits well experiment in the two cases with and without ICRH .

85307 (with ICRH) **85308 (NBI only)**

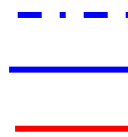


Simulation of two SXR vertical Lines of Sights From central (left) towards the LFS.

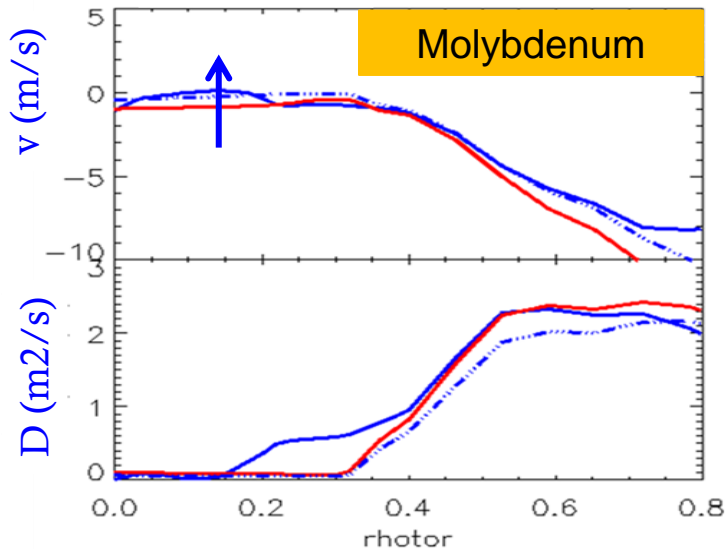
Model-based transport coefficients used in JETTO/SANCO to simulate Mo transient behavior (LBO)

85307 (with ICRH)

85308 (NBI only)



Centrifugal Effects only
CF and fast ion effects



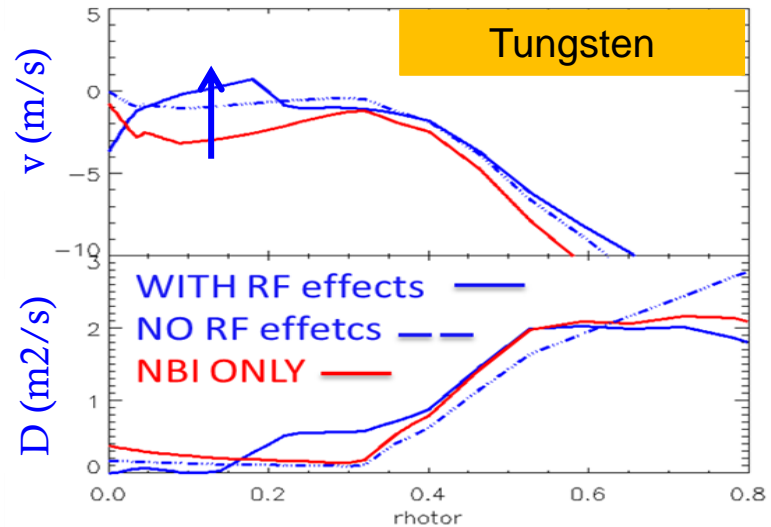
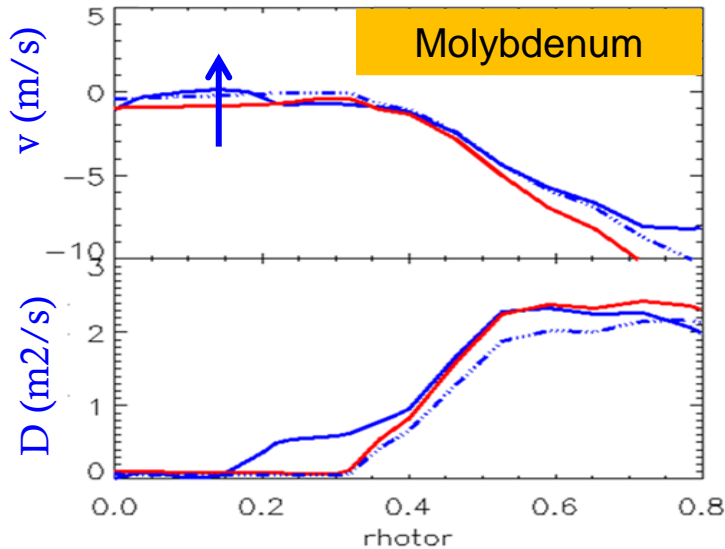
Model-based transport coefficients used in JETTO/SANCO to simulate Mo transient behavior (LBO)

85307 (with ICRH)

85308 (NBI only)



Centrifugal Effects only
CF + fast ion effects



- With advanced theory-based two dimensional transport models the complex behavior of W in the core of JET standard H-mode and hybrid discharges has been understood.
- The sensitivity to neoclassical transport of W is the main reason for its accumulation in JET discharges characterized by peaked density profiles.
- Central ICRH hampers W accumulation affecting the main kinetic profiles and the related neoclassical drive but also modifying directly W transport through thermal screening and anisotropy of heated minority species.