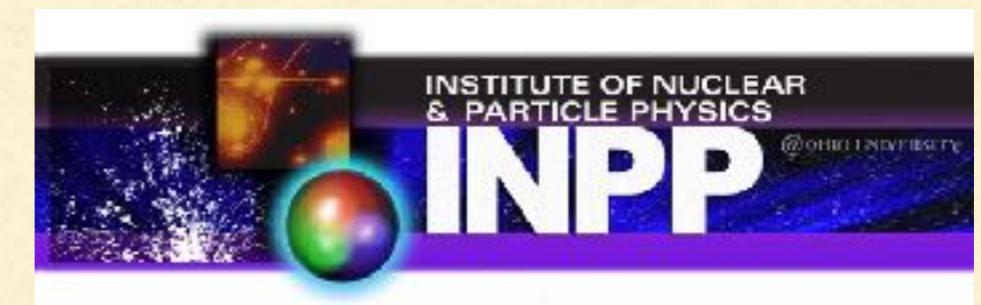


BRICK and Beyond: Bayesian analyses of low-energy ${}^3\text{He}$ - ${}^4\text{He}$ data using R-matrix and Effective Field Theory

Daniel Phillips
Ohio University

with Daniel Odell, Carl Brune,
James deBoer, Som Paneru,
Mahesh Poudel, & Andrius Burnelis



RESEARCH SUPPORTED BY THE DOE OFFICE OF SCIENCE, THE SSAP, AND THE NSF OAC

Outline

- Bayes' theorem in one slide
 - Bayesian R-matrix analysis of ${}^3\text{He} + {}^4\text{He} \rightarrow {}^7\text{Be} + \gamma$ and ${}^3\text{He} + {}^4\text{He}$ elastic scattering
 - Set up: data + model Odell, Brune, DP, deBoer, Paneru, Frontiers in Physics (2022)
Paneru, Brune Connolly, Odell, Poudel, DP, et al. Phys. Rev. C (2024)
 - Experimental imperfections
 - Why the full posterior?
 - Error propagation
 - EFT modeling of $\text{He} + {}^4\text{He} \rightarrow {}^7\text{Be} + \gamma$ and ${}^3\text{He} + {}^4\text{He}$ elastic scattering
 - What does the EFT amplitude for these reactions look like? Zhang, Nollett, DP, J. Phys. G (2020)
 - Low-energy S-factor data Poudel, DP, J. Phys. G. (2022)
Burnelis, DP
 - What about scattering data?
 - Summary and Future Work
-

The Bayes-ics

Thomas Bayes (1701?-1761)



$$\text{pr}(A|B, I) = \frac{\text{pr}(B|A, I)\text{pr}(A|I)}{\text{pr}(B|I)}$$

<http://www.bayesian-inference.com>

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Likelihood Prior
↓ ↓
↑ ↑
Posterior Model evidence

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Typically evaluated by MCMC sampling

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Probability as
degree of belief

Marginalization: $\text{pr}(x|\text{data}, I) = \int dy \text{pr}(x, y|\text{data}, I)$

Allows us to integrate out “nuisance” parameters, e.g., those associated with systematic uncertainties

$$\text{pr}(A|B, I) = \frac{\text{pr}(B|A, I)\text{pr}(A|I)}{\text{pr}(B|I)}$$

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Likelihood Prior
↓ ↓
Posterior Model evidence
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Typically evaluated by MCMC sampling

SONIK data set

- ^3He beam incident on SONIK apparatus

- Intensity of about 10^{12} s^{-1}

- Windowless gas target

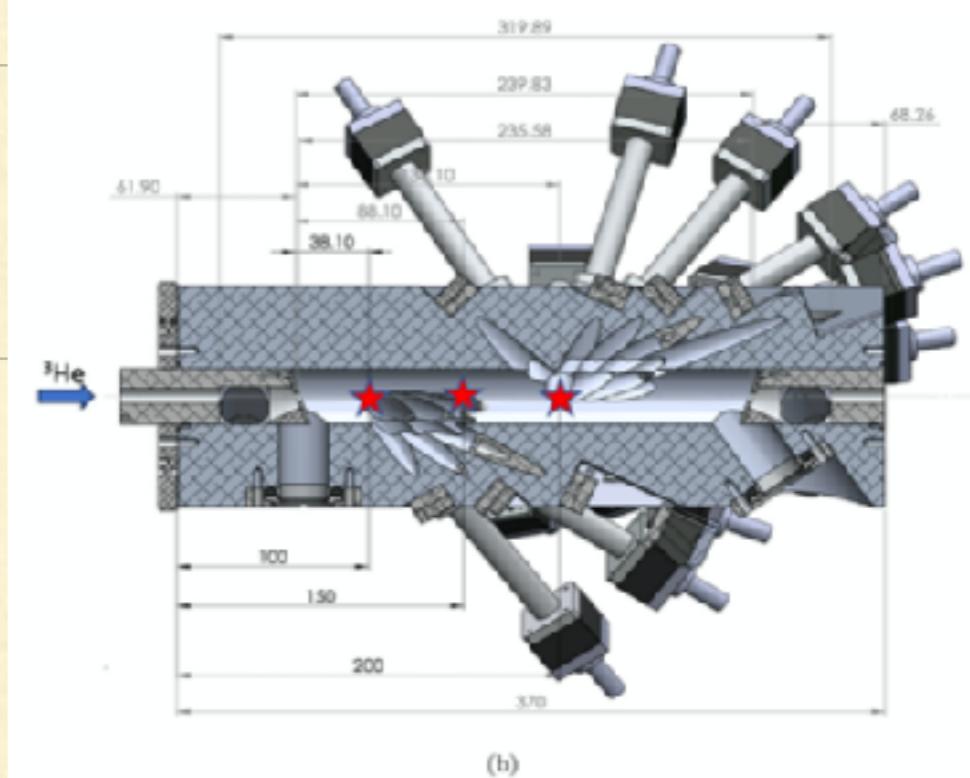
- Measured at nine beam energies $E[^3\text{He}] = 0.721 - 5.490 \text{ MeV}$

- Three different interaction regions, corresponding to slightly different c.m. energies of the collision

- Total systematic uncertainties that are independent of energy and angle estimated at 2%

- Telescopes have a 1.6% variation of aperture around mean value, so acceptance correction is angle dependent

- Consistent with older, Barnard data set at higher energies



Paneru et al., Phys. Rev. C (2024)

Full data set for model calibration

- 88 S-factor data
 - Seattle (S)
 - Weizmann
 - Luna (L)
 - Erna
 - Notre Dame
 - Atomki
- Scattering data
 - SONIK: 451 from 0.385 to 3.127 MeV
 - Barnard: 646 from 1.49 to 3.27 MeV
- Plus 34 branching-ratio data

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Two analyses:

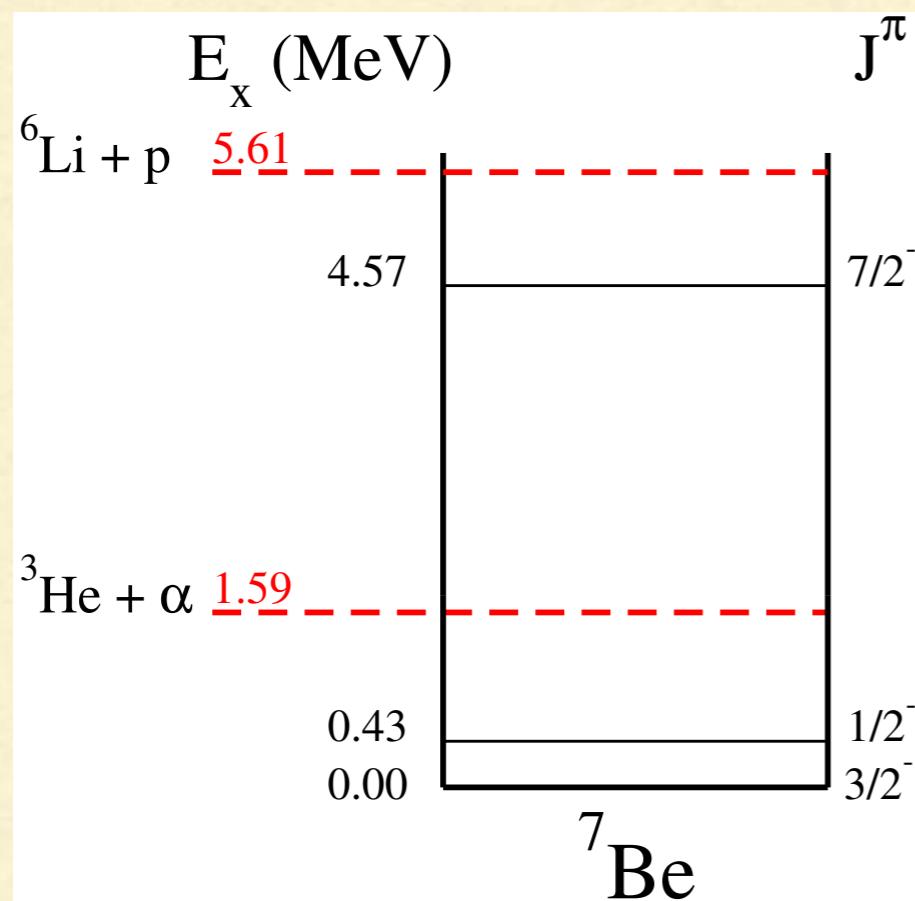
Capture + SONIK

Capture + SONIK + Barnard

R-matrix model

Odell, Brune, DP, deBoer, Paneru, Frontiers in Physics (2022)

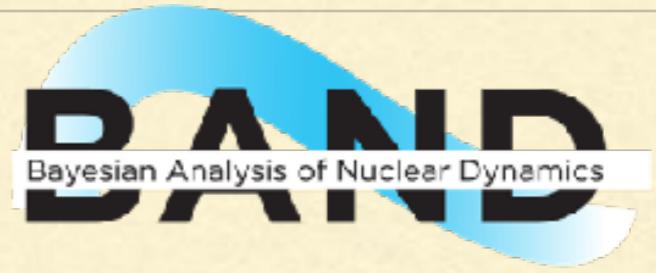
- Goal: describe scattering and capture data up to the $p^6\text{Li}$ threshold
- 3/2- and 1/2- bound states with prior ranges for ANCs from 1 to 5 MeV



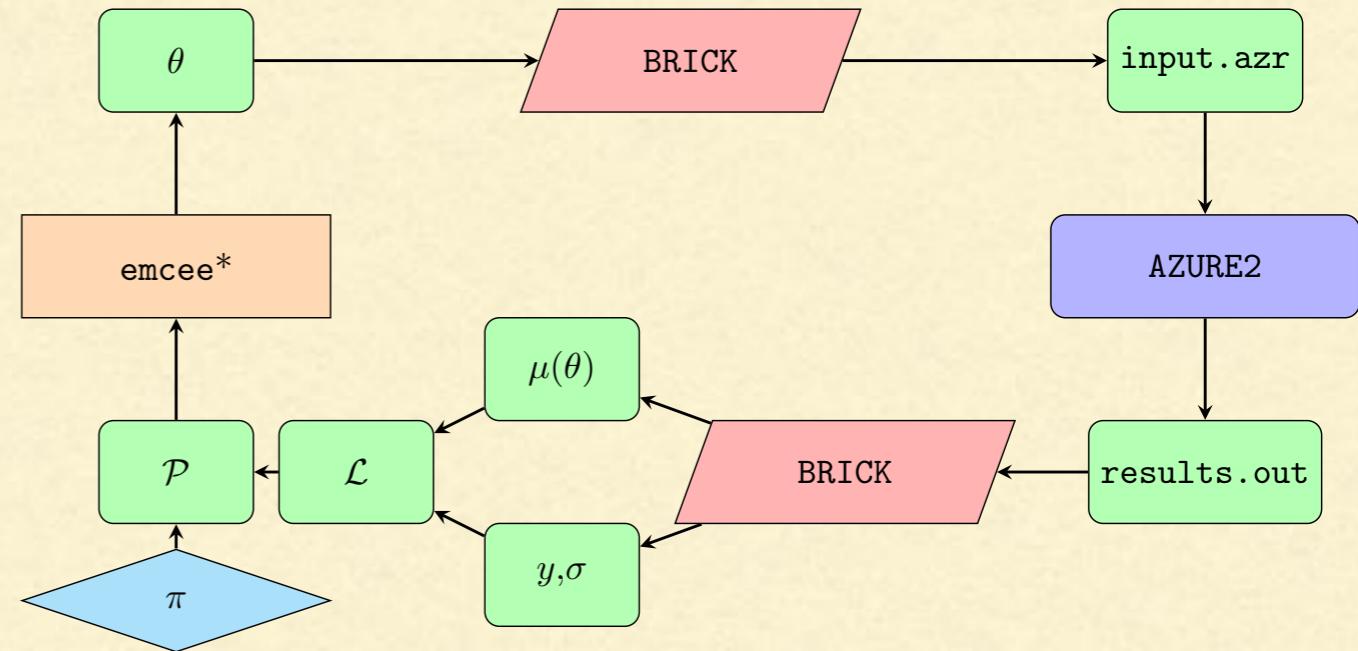
Background & resonance levels

	E (MeV)	Γ_a (MeV)
1/2-	21.6	[-200,200]
3/2-	21.6	[-100,100]
5/2-	7	[0,100]
7/2-	[2,10]	[0,10]
1/2+	14	[0,100]
3/2+	12	[0,100]
5/2+	12	[0,100]

Throw a BRICK at it



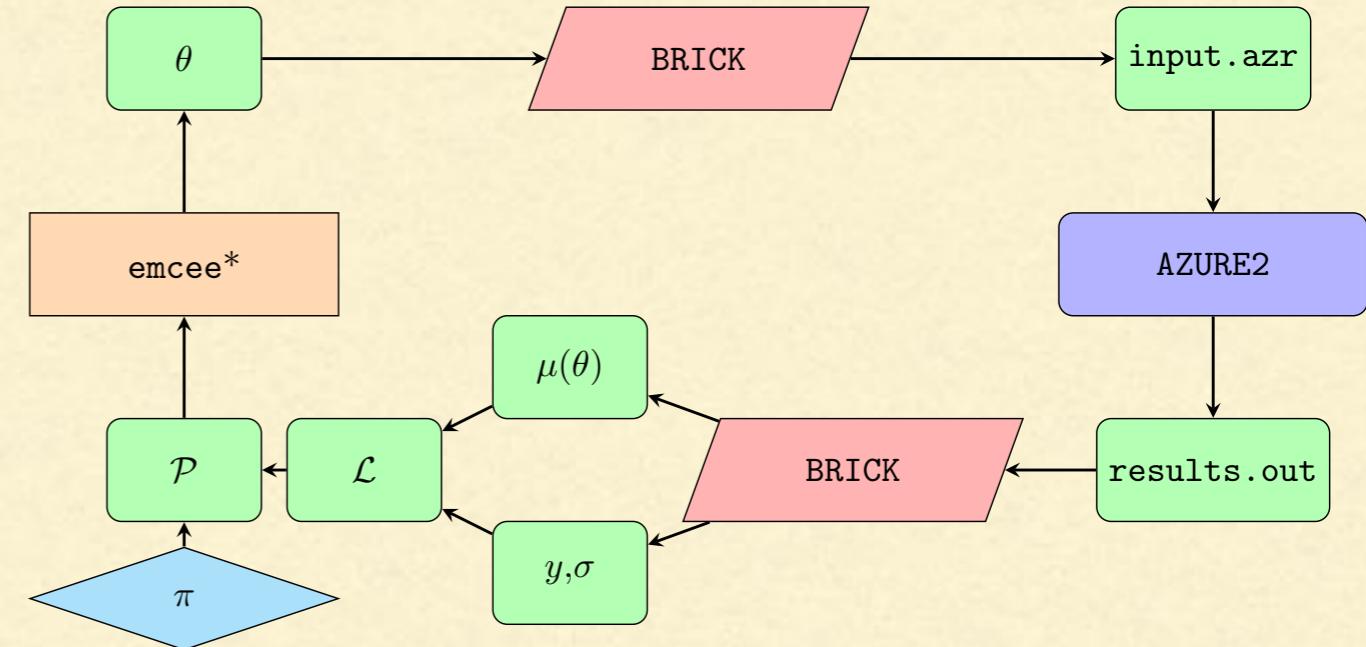
- Publicly available Python code
<https://github.com/odell/brick>
 - Available on PyPI
 - BAND Framework v0.4
bandframework.github.io
 - AZURE2 must be installed
 - User specifies R-matrix model & data set in AZURE2
 - Specify priors & likelihood



Throw a BRICK at it

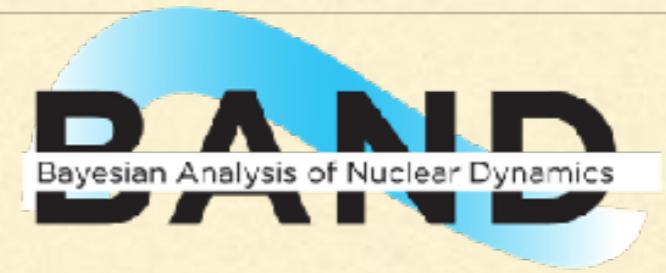


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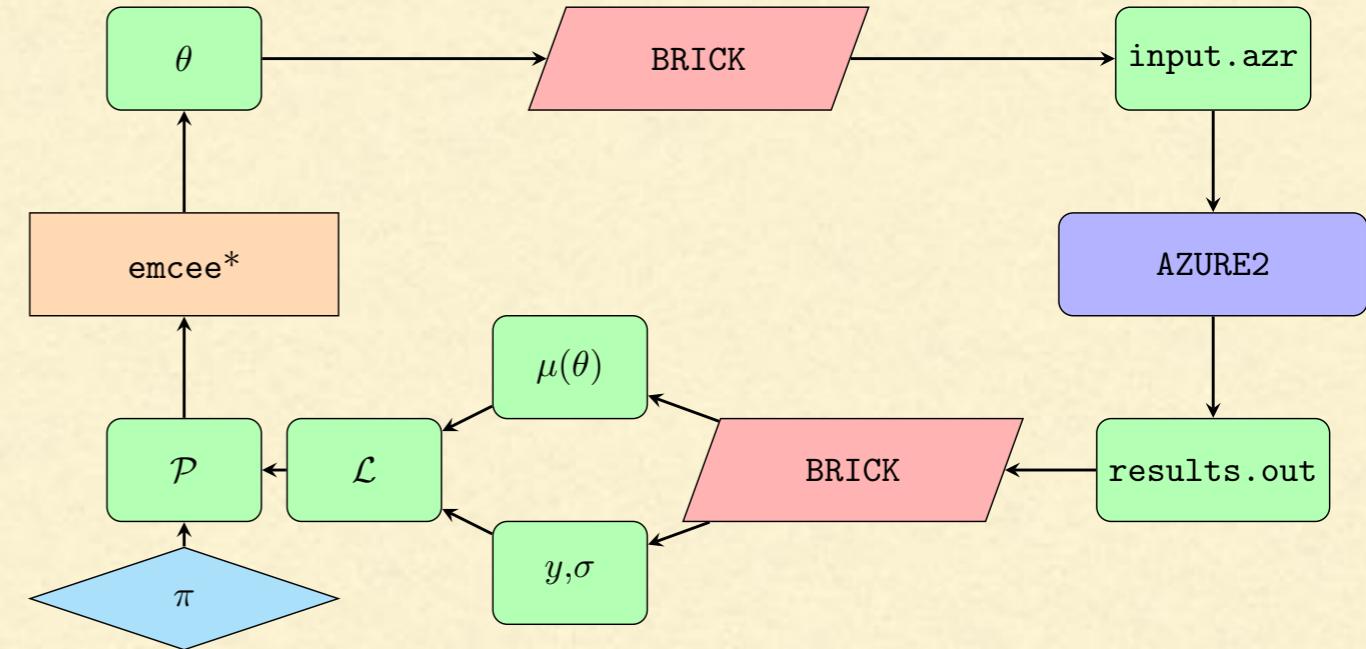


$$\mathcal{L} \propto \prod_{\alpha=1}^{N_{\text{sets}}} \prod_{j=1}^{N_{\alpha}} \exp \left(-\frac{(y_{j\alpha} - f_{\alpha} \mu(x_{j\alpha}; \theta_R))^2}{2\sigma_{j,\alpha}^2} \right)$$

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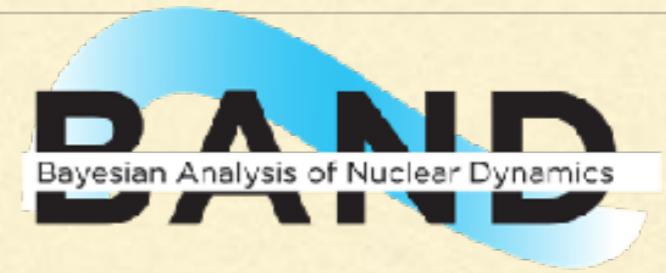


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Data

Two arrows point from the text "Data" to the terms $y_{j\alpha}$ and $\sigma_{j,\alpha}^2$ in the equation.

Throw a BRICK at it

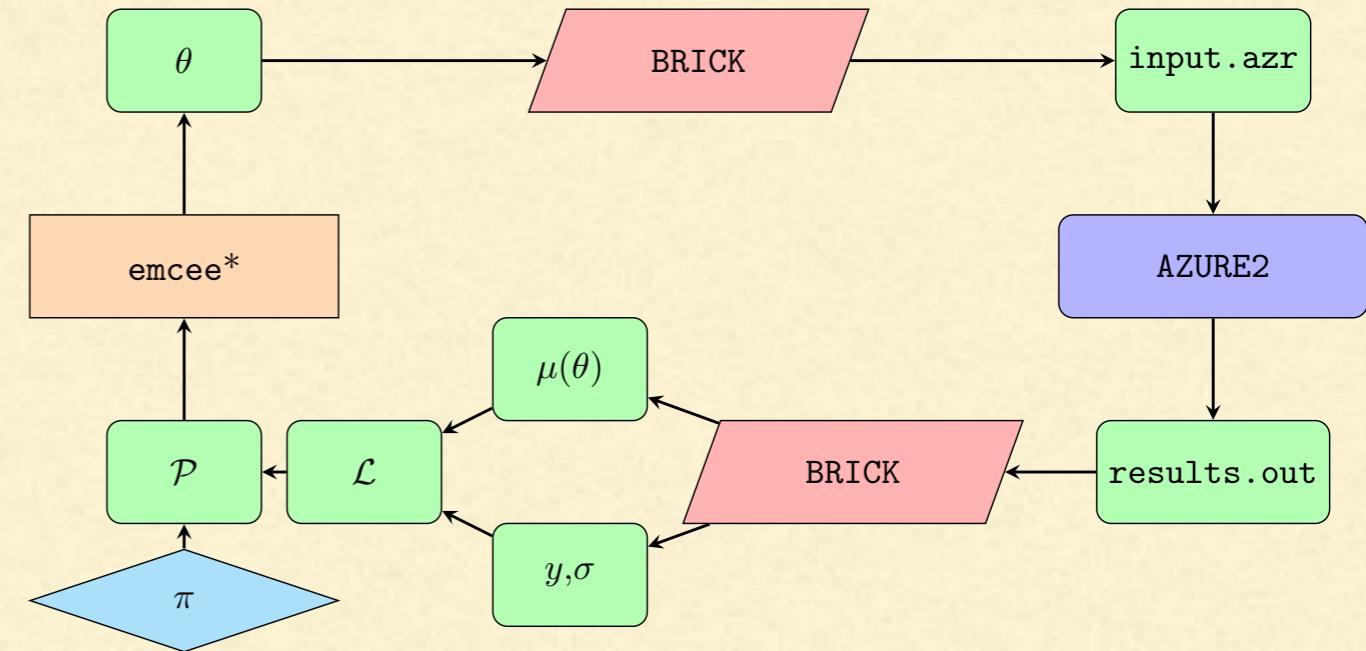


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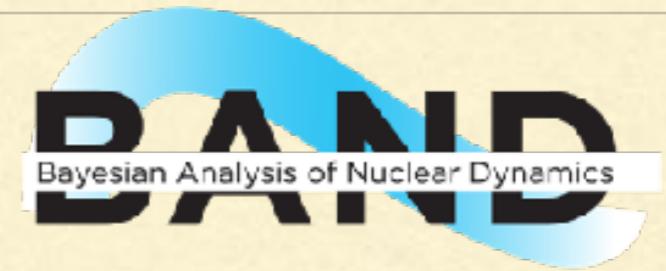
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Data

R-matrix number



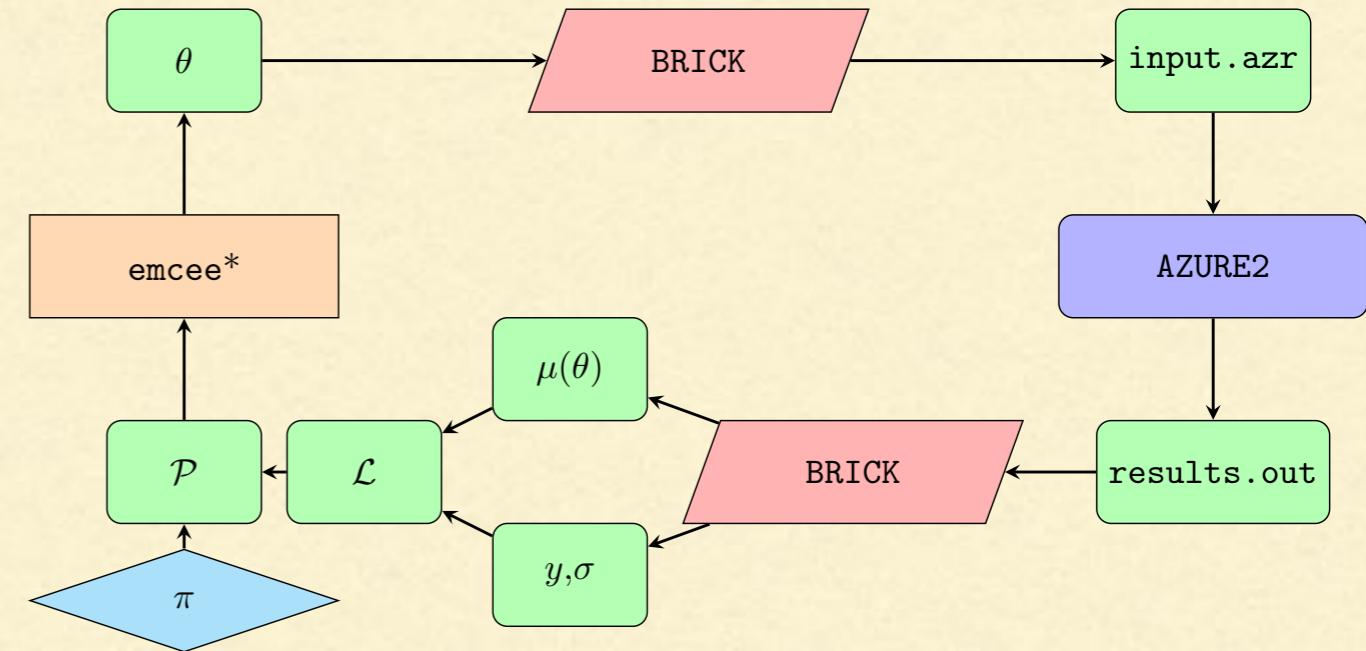
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Data R-matrix number



The benefits of Bayesian parameter estimation

The benefits of Bayesian parameter estimation

- Straightforward to introduce additional nuisance parameters to model experimental imperfections. Marginalizing over them includes impact of those imperfections on parameters and evaluated quantities
- Access to full multi-dimensional posterior for parameters, not just properties around a (local) minimum
- With samples of R-matrix parameters in hand, straightforward to evaluate any observable we want for all those samples \Rightarrow error propagation is a snap!

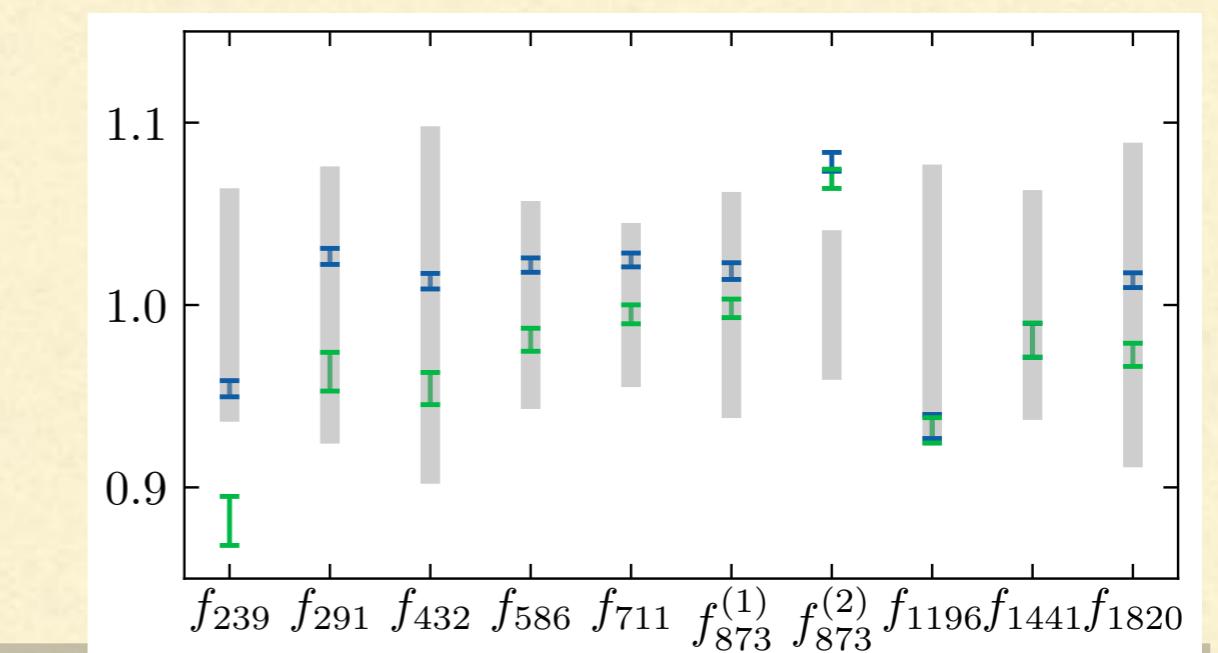
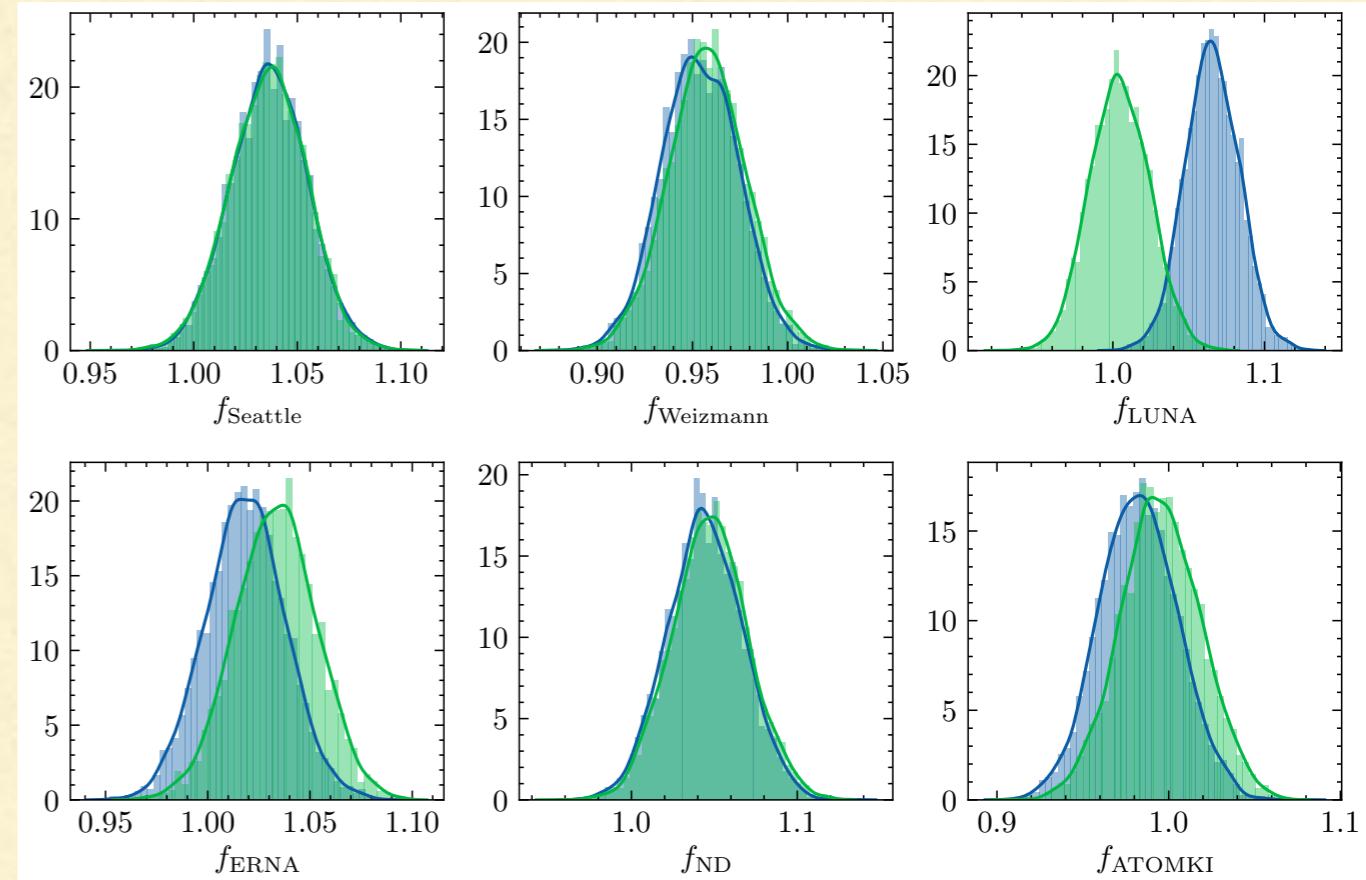
$$\text{pr}(S(E_0) | D, I) = \int d\vec{\theta} \delta(S(E_0) - S_{\text{R-matrix}}(E_0; \vec{\theta})) \text{pr}(\vec{\theta} | D, I)$$

- Not just experimental imperfections either! Theory imperfections can be accounted for too
-

Modeling of normalization uncertainties

- Analysis includes common-mode errors for all data sets, implemented by factor f_a to avoid d'Agostini bias
- For SONIK data set this normalization factor is assigned for each beam energy
- Almost all normalizations come out inside quoted CMEs, all are within $2 \times \text{CME}$, apart from LUNA in CSB analysis
- “Dialogue with the data”

CS
CSB



More sophisticated normalization modeling

Paneru, Brune, Connelly, Odell, ..., DP, et al., PRC (to appear)

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More sophisticated normalization modeling

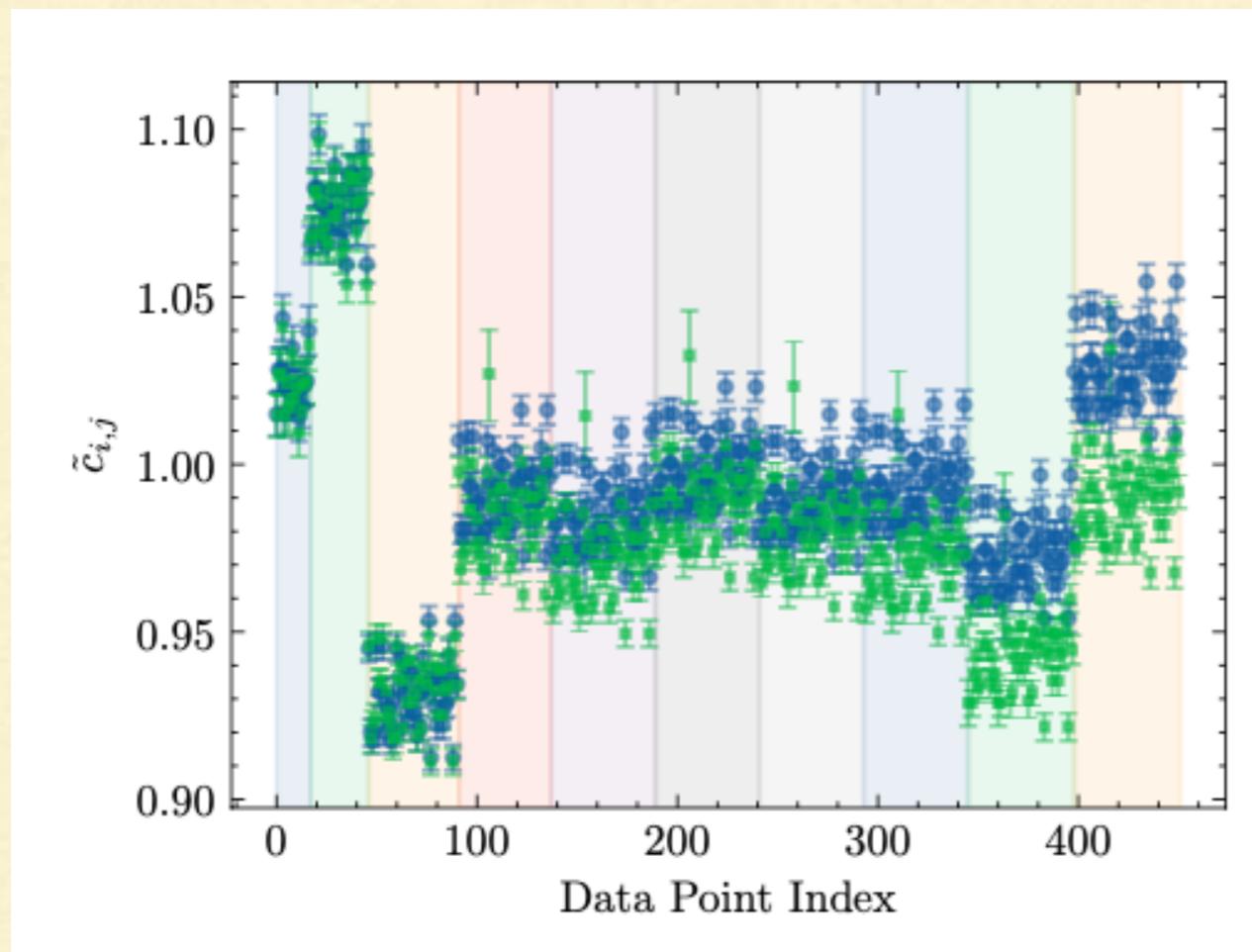
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- Consider not just beam normalization uncertainty, but also uncertainty due to acceptance (aperture variation) of each of 27 detectors

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$$y_{\text{exp}} = f_{\text{SONIK}} f_E f_{\text{det}} y_R + \delta y_{\text{exp}}$$

$$\tilde{c}_{i,j} = f_E f_{\text{det}}$$

Green: R-matrix
Blue: EFT

Beam energy shift

Beam energy shift

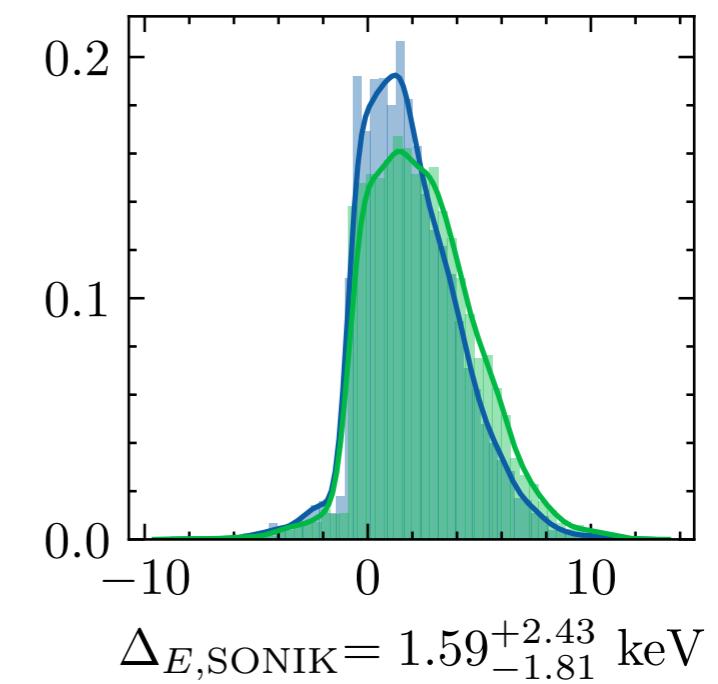
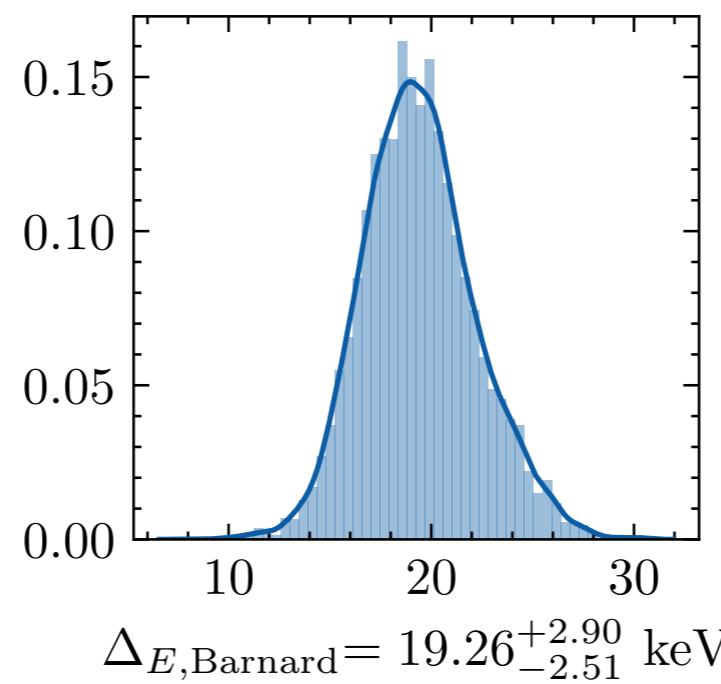
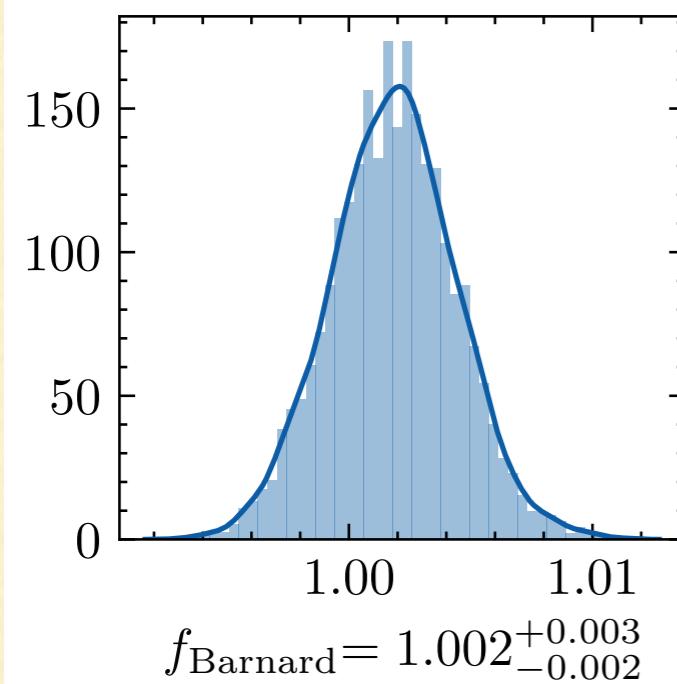
- Shift energy of Barnard data set by a constant to account for possible miscalibration of beam energy: $E \rightarrow E + \Delta$. Prior a Gaussian with standard deviation 40 keV ← information in paper

$$\mathcal{L} \propto \prod_{\alpha=1}^{N_{\text{sets}}} \prod_{j=1}^{N_{\alpha}} \exp \left(-\frac{(y_{j\alpha} - f_{\alpha}\mu(E_{j\alpha} + \Delta, \phi_{j\alpha}; \theta_R))^2}{2\sigma_{j,\alpha}^2} \right)$$

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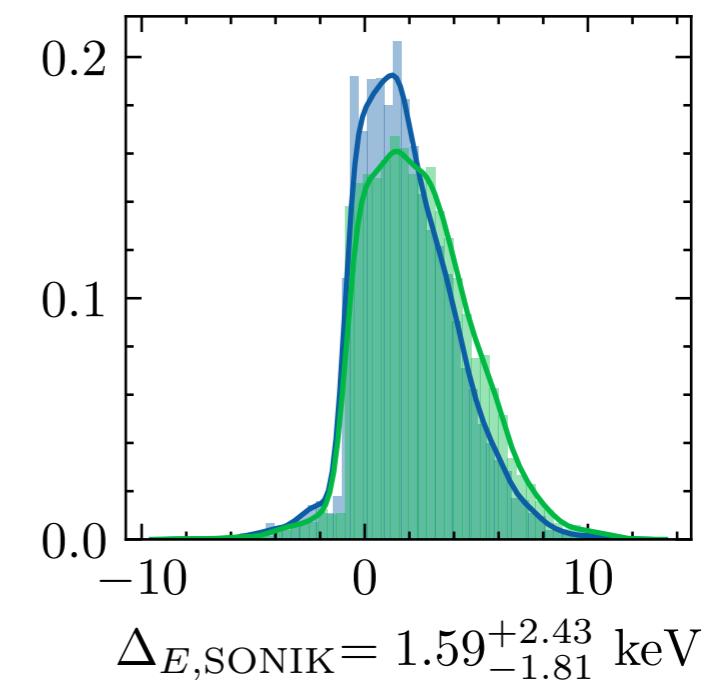
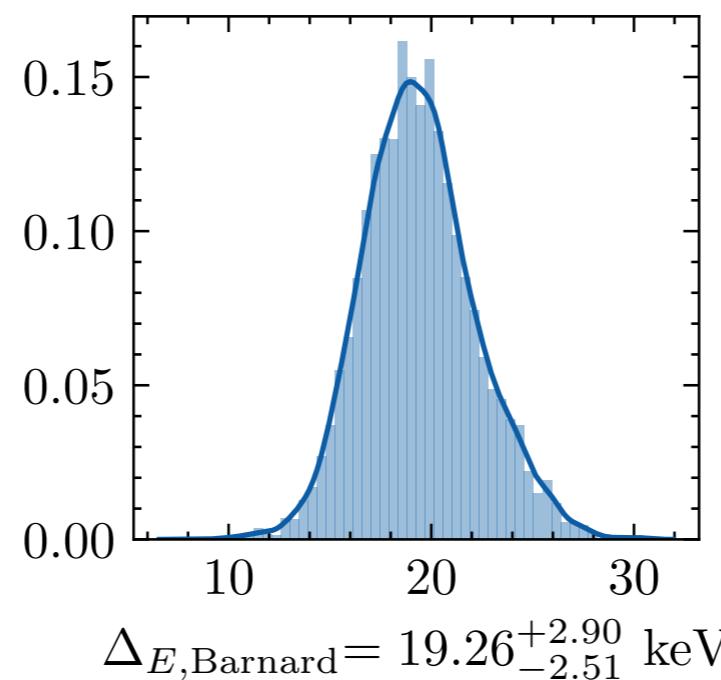
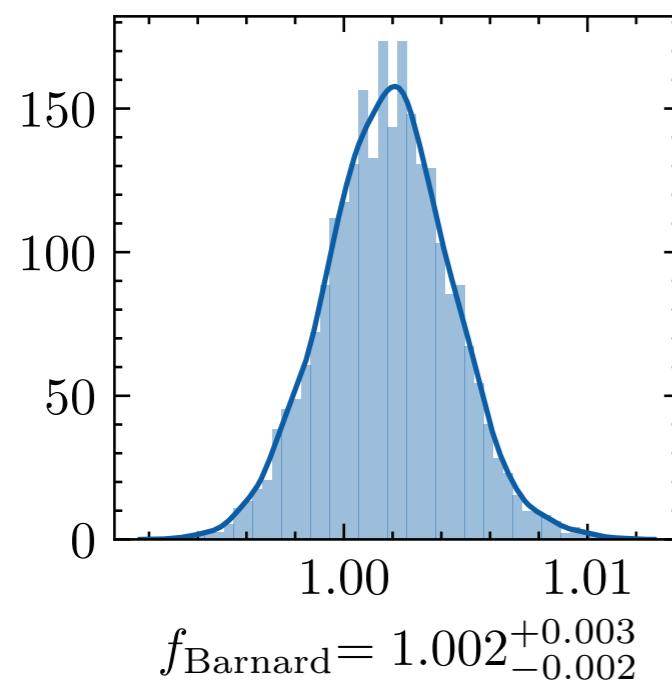


CS
CSB

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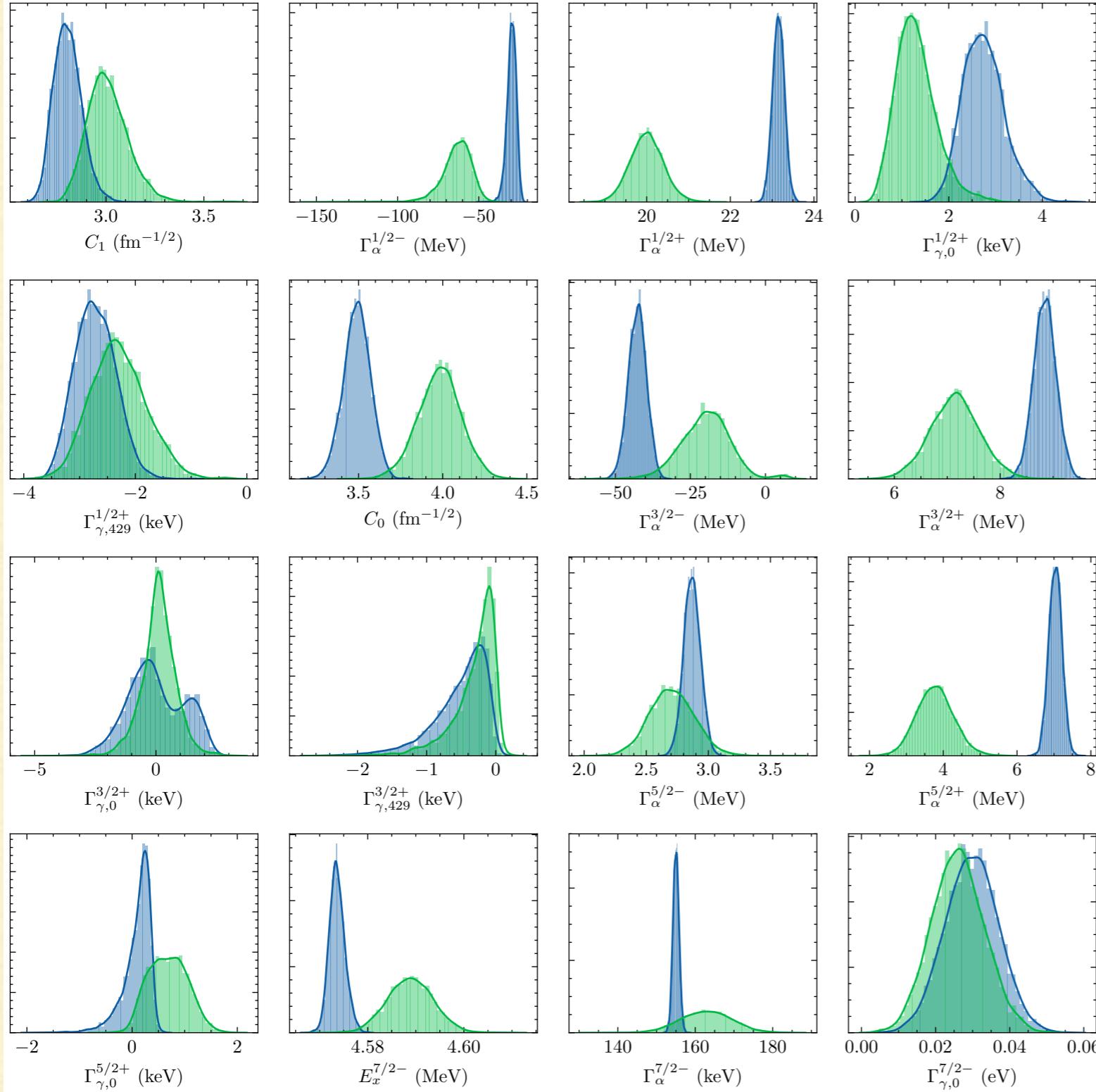
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CS
CSB

- No significant change in $\vec{\theta}_{\text{Barnard}}$ due to this though

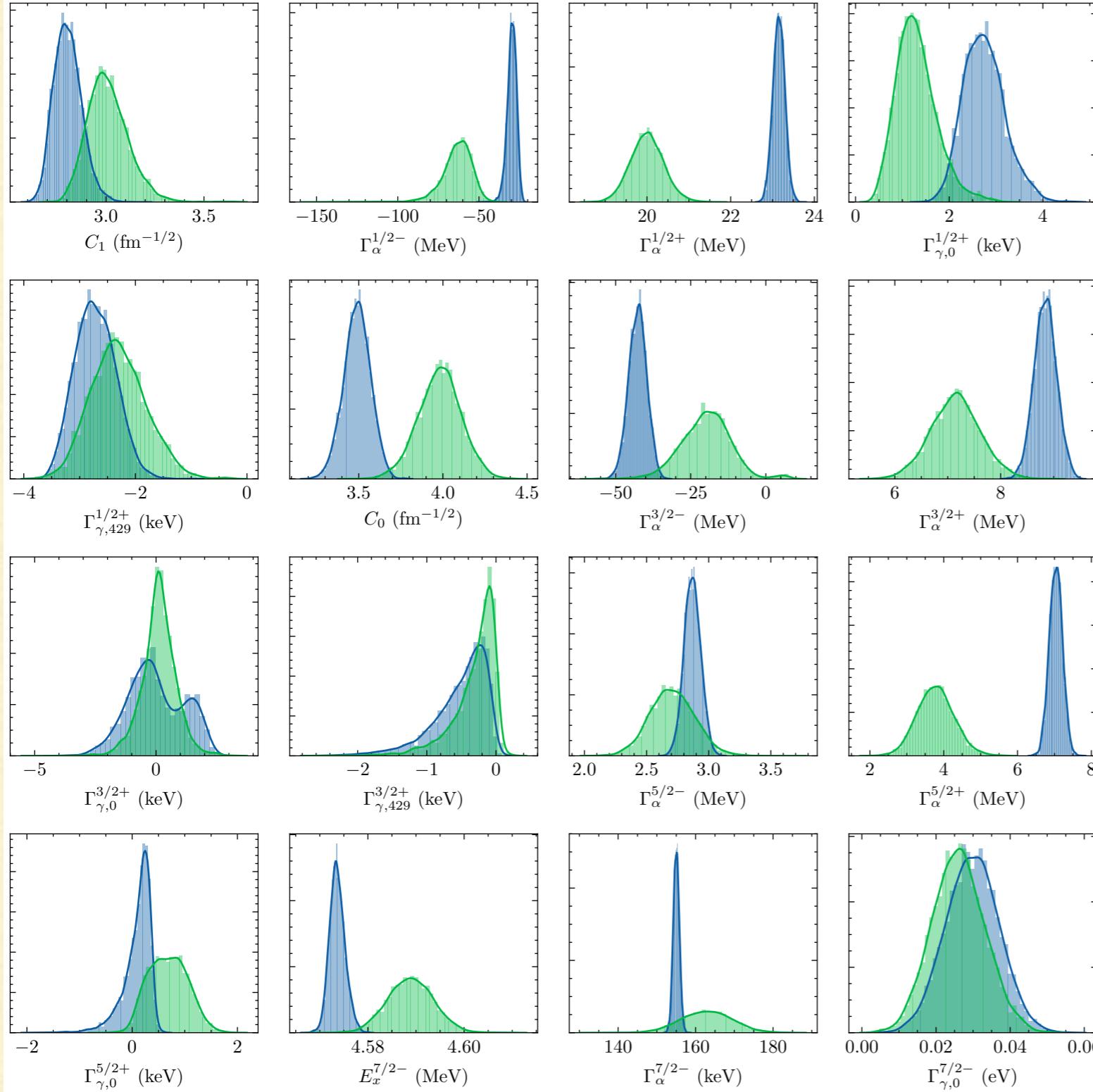
Posteriors for R-matrix parameters



Capture + SONIK

Capture + SONIK + Barnard

Posteriors for R-matrix parameters



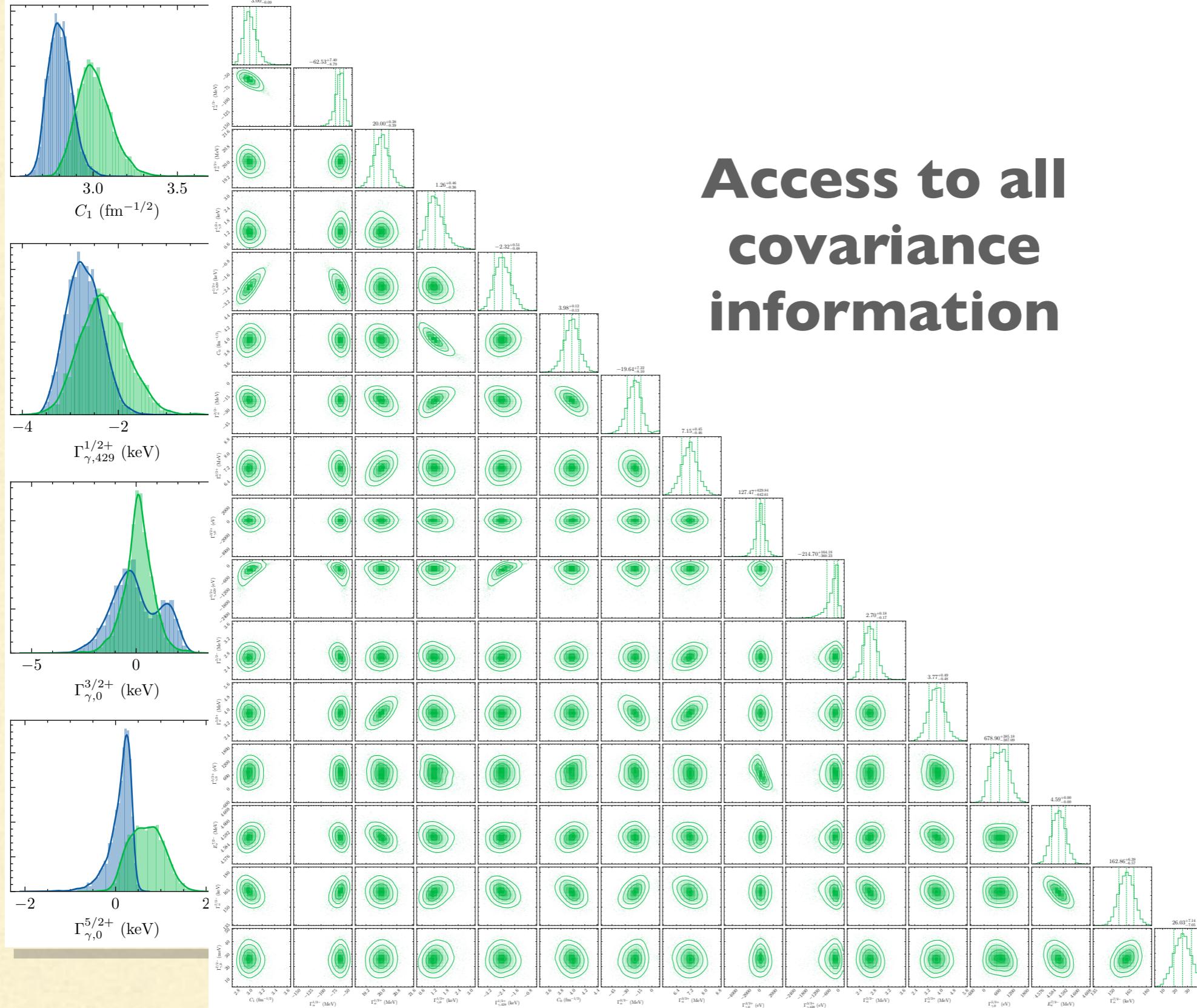
Capture + SONIK

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Notable points:

- ANCs
- $\Gamma_{\alpha}^{7/2-}$
- Non-Gaussianity

Posteriors for R-matrix parameters



re + SONIK

SONIK + Barnard

able points:

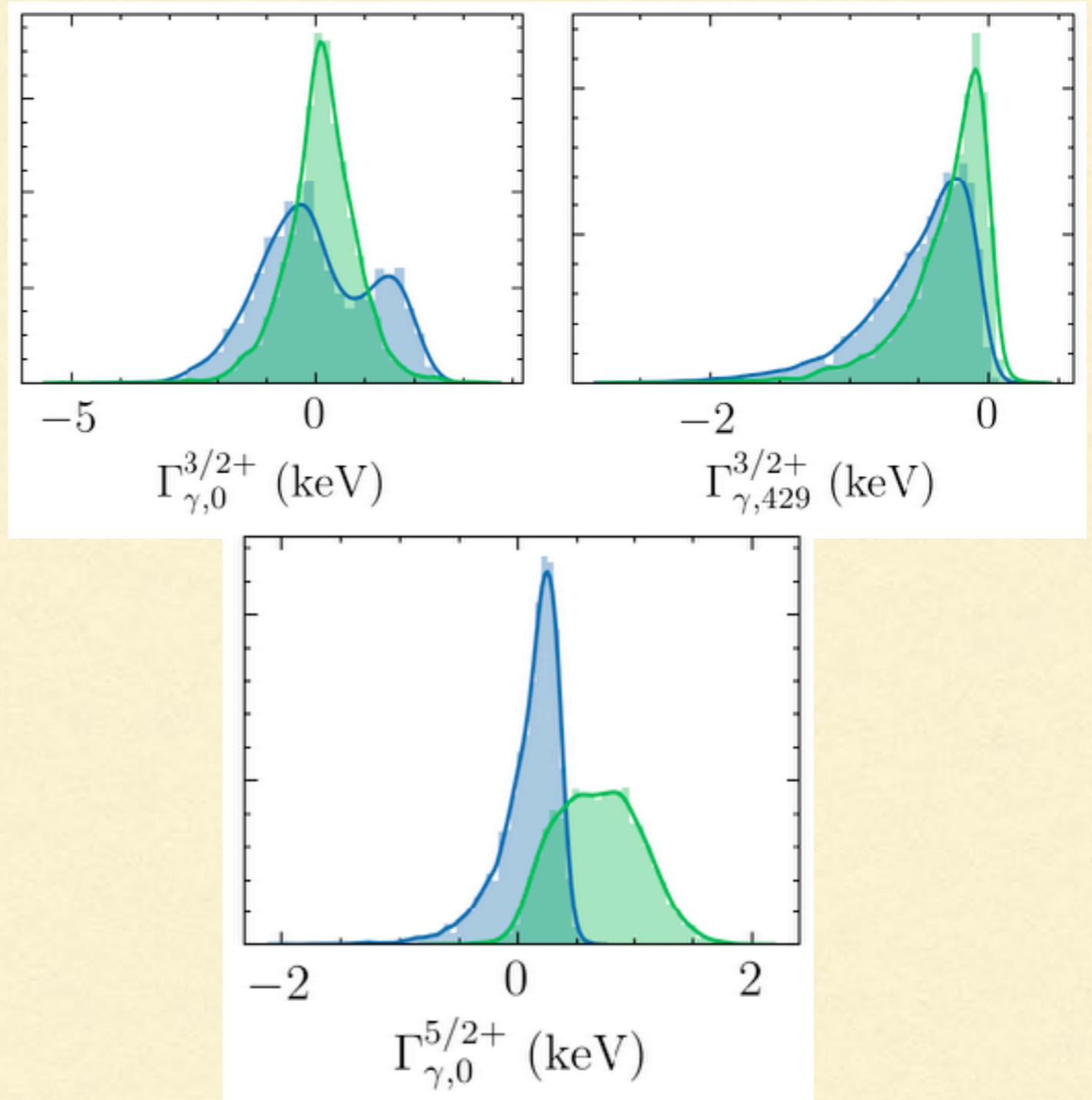
Cs

n-Gaussianity

But why didn't you just use MINUIT?

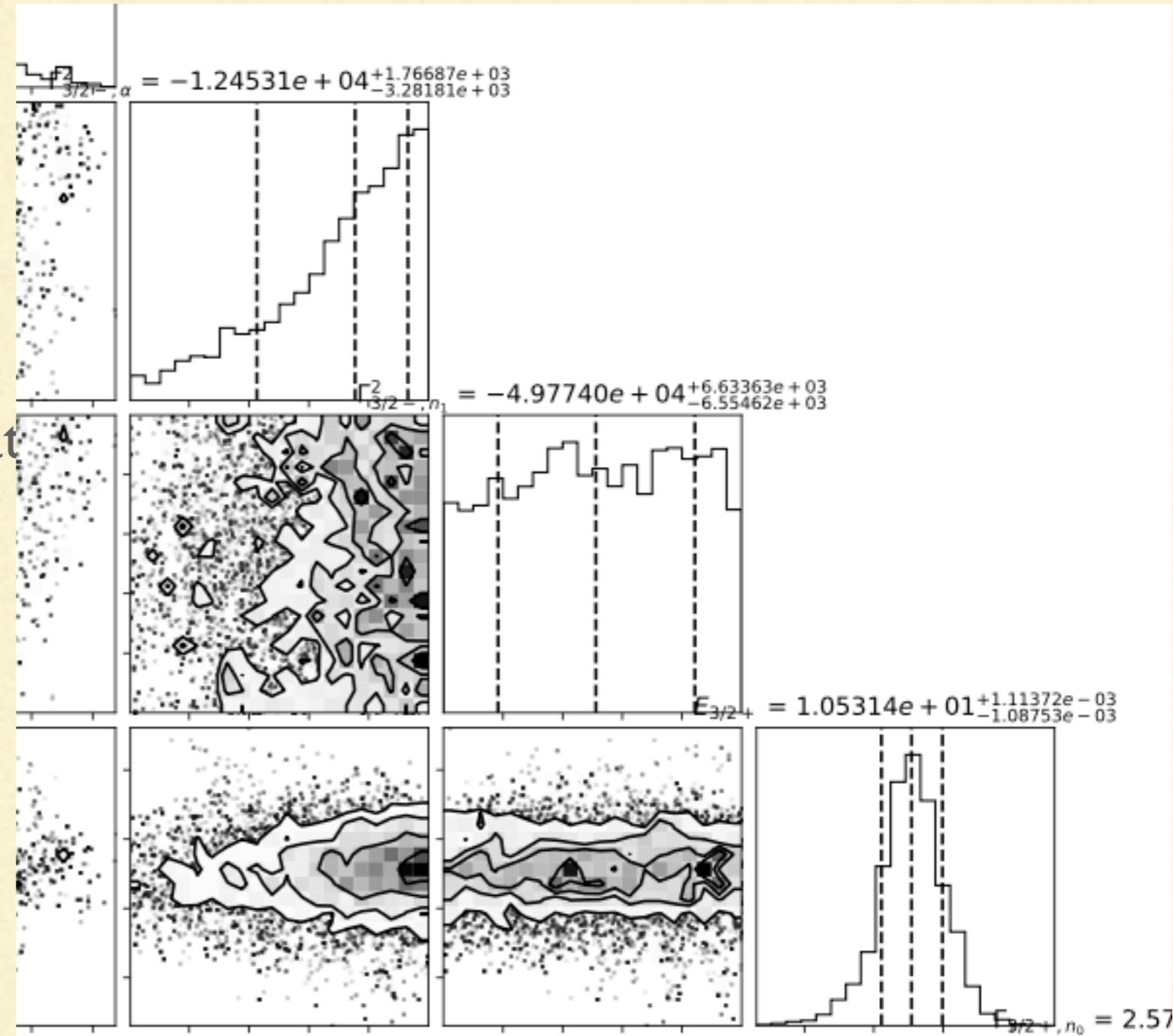
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- Diagnose non-Gaussianity



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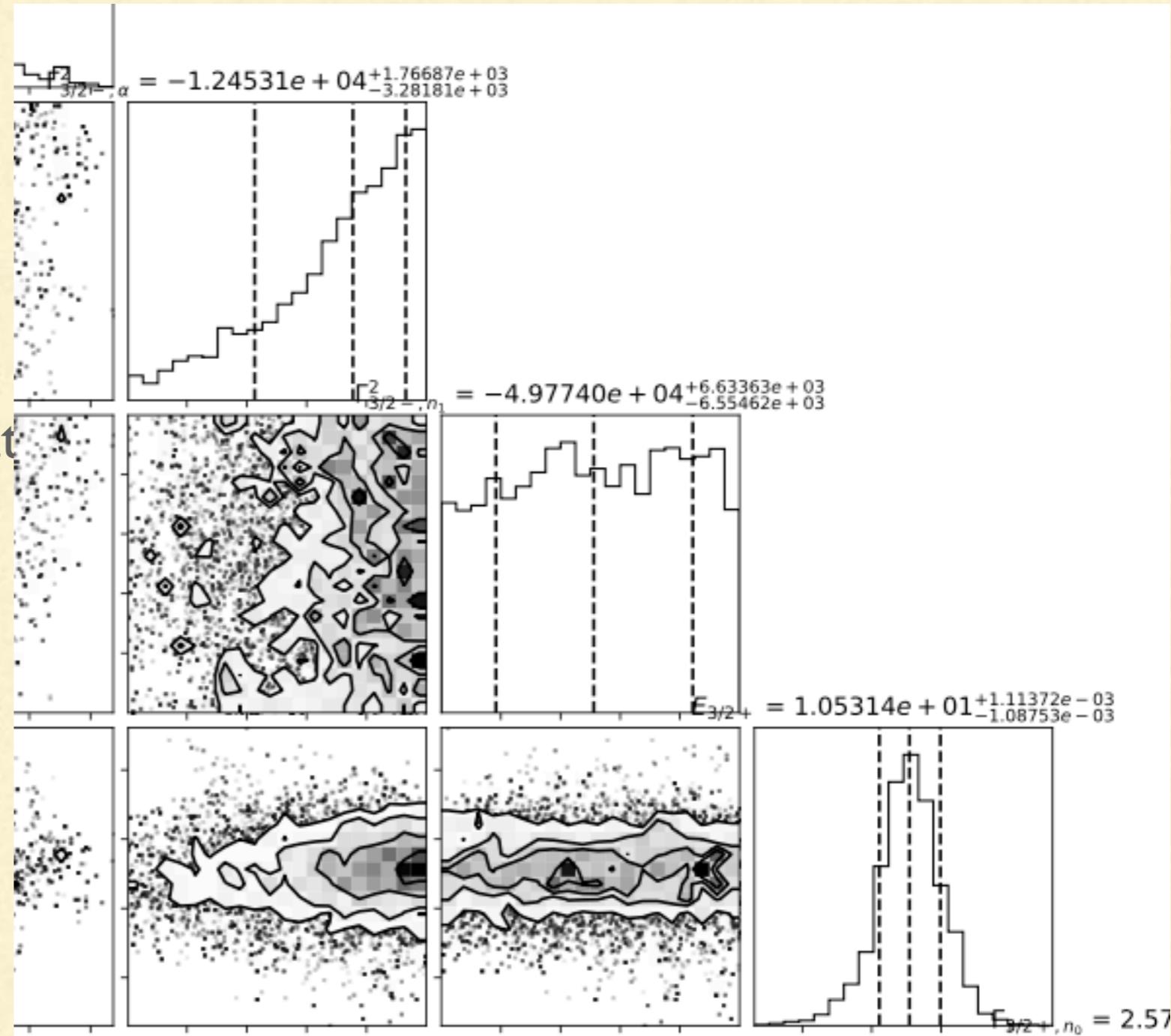
- Diagnose non-Gaussianity
- Very clear when posterior “returns the prior” for a particular parameter and so that parameter is not really needed for the fit.



$^{13}\text{C}(\alpha, n)$, courtesy James deBoer

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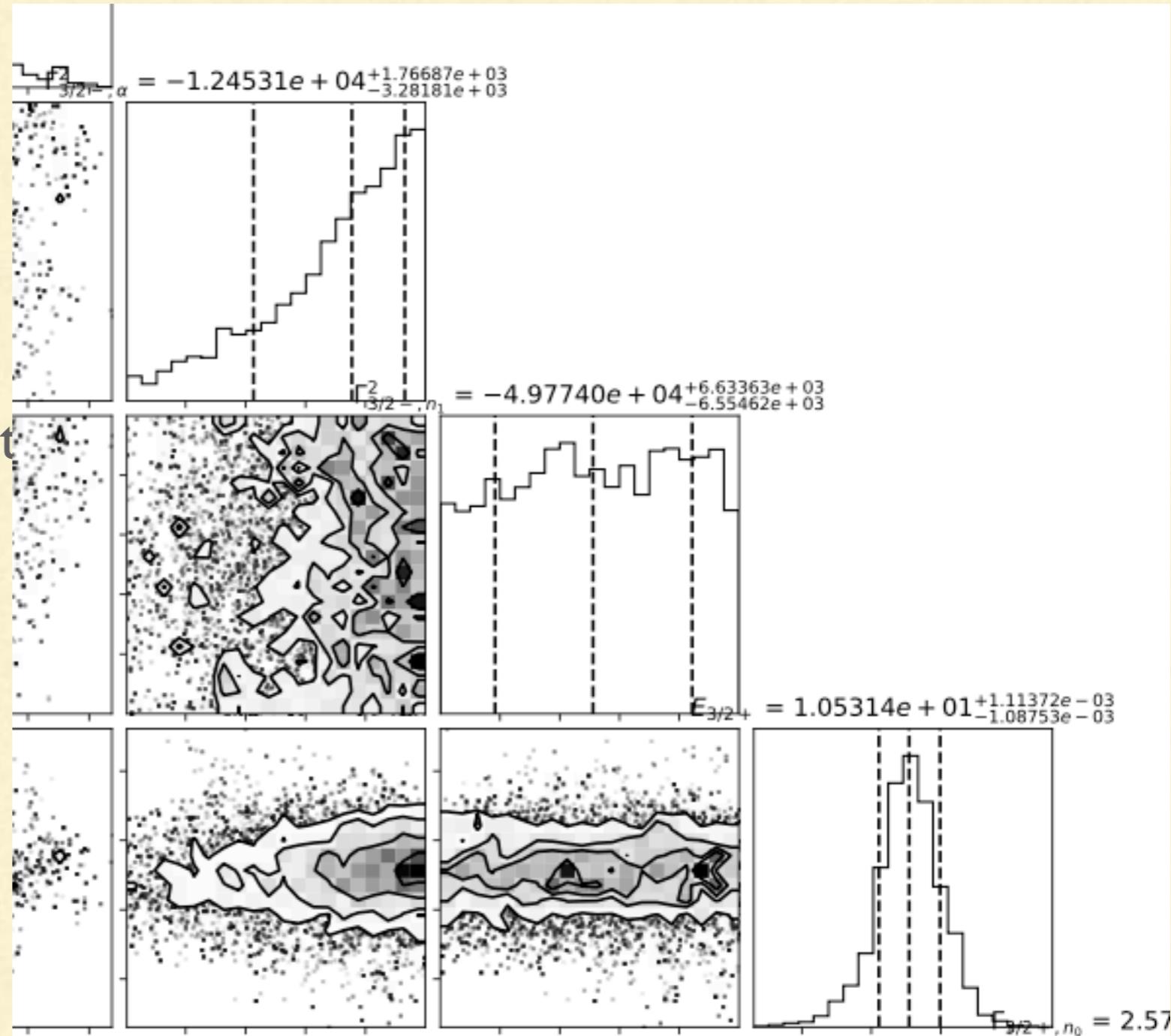
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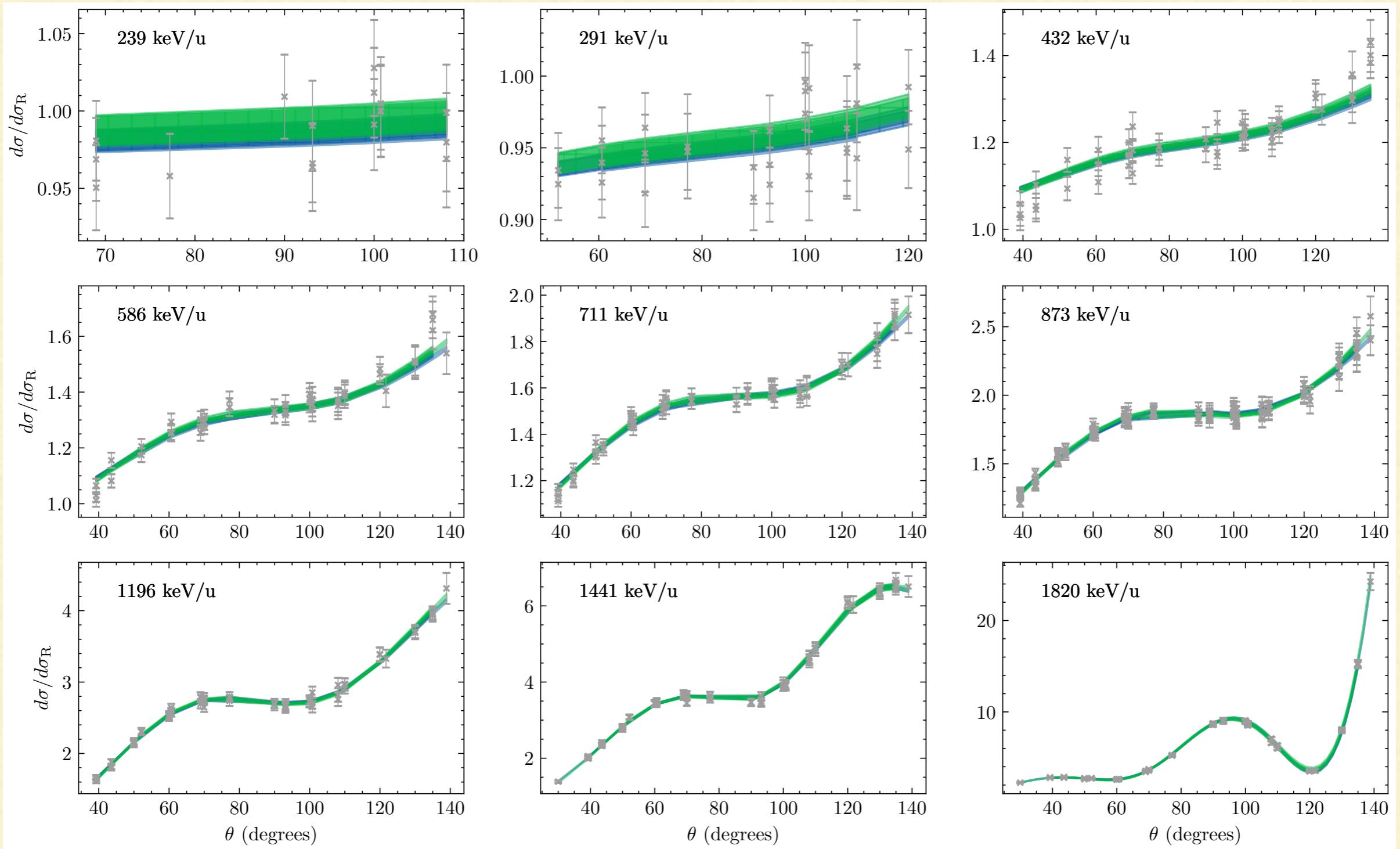
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- Diagnose non-Gaussianity
- Very clear when posterior “returns the prior” for a particular parameter and so that parameter is not really needed for the fit.
- (Note that it's also clear when prior is affecting shape of posterior.)
- Also, error propagation....

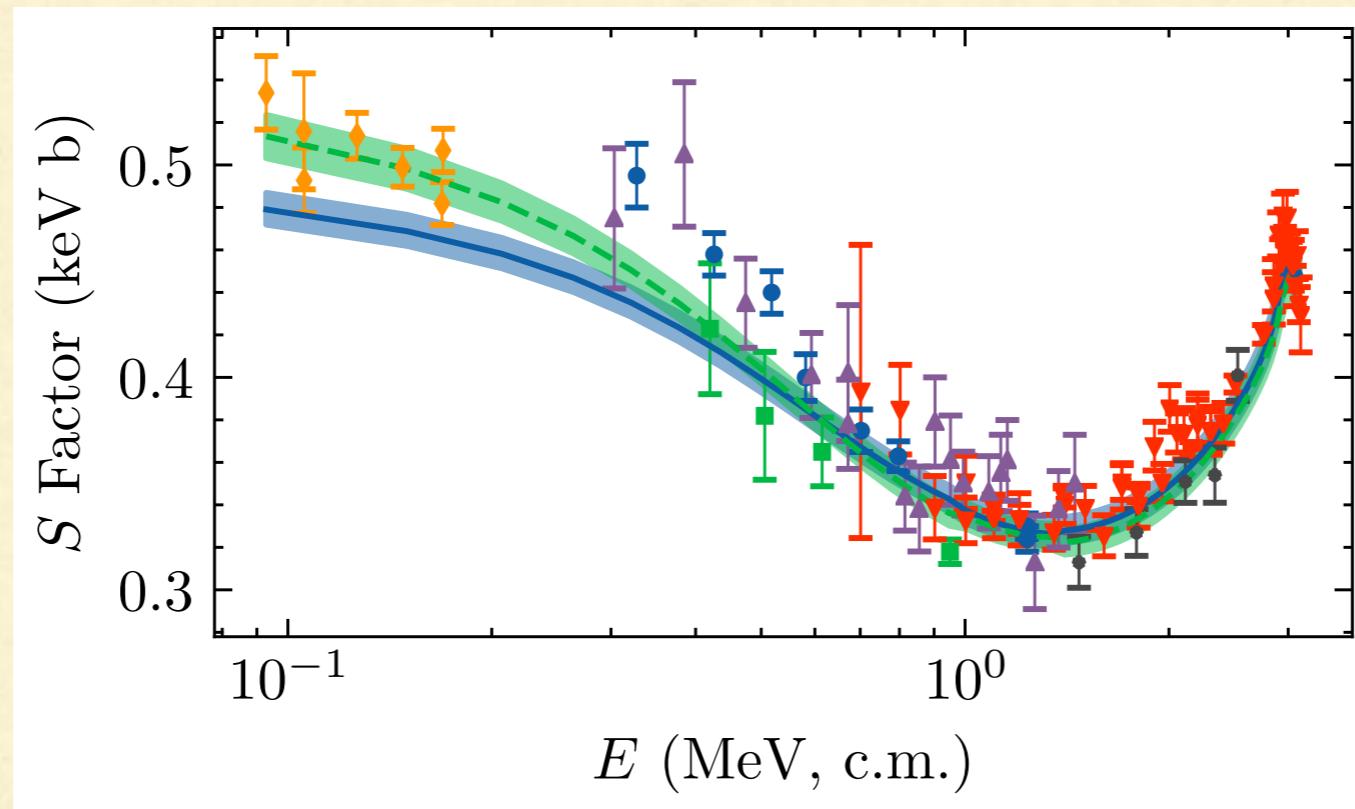


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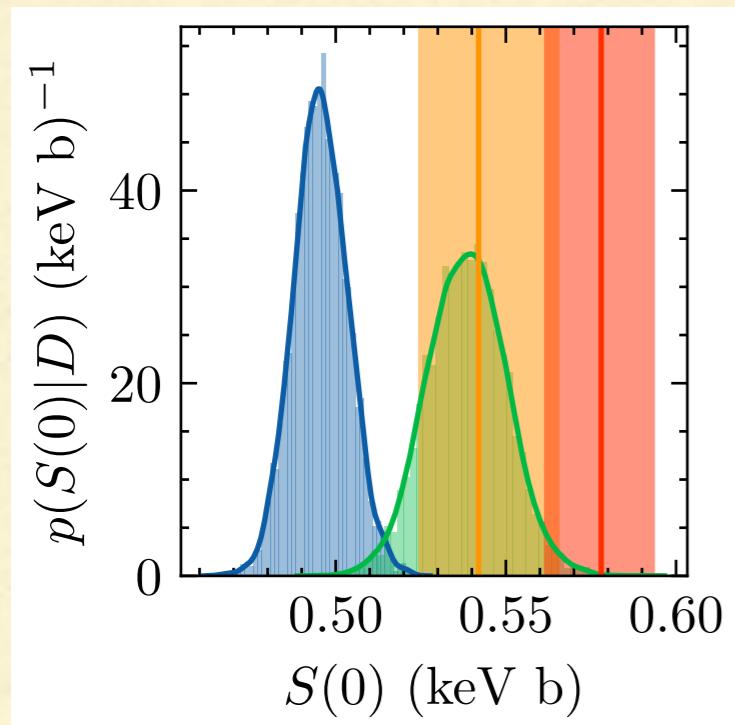
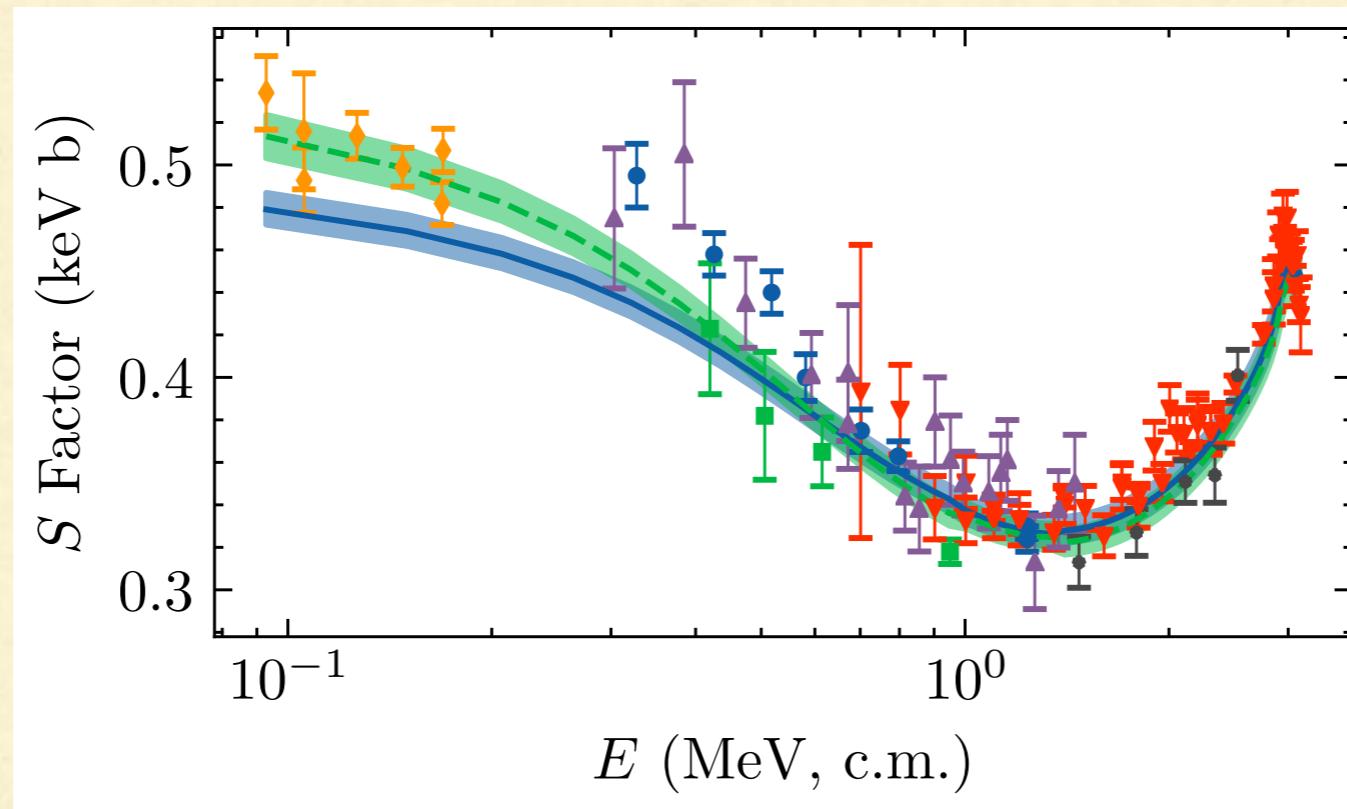
Speaking of which: SONIK data looks good



What about S-factor at solar energies?

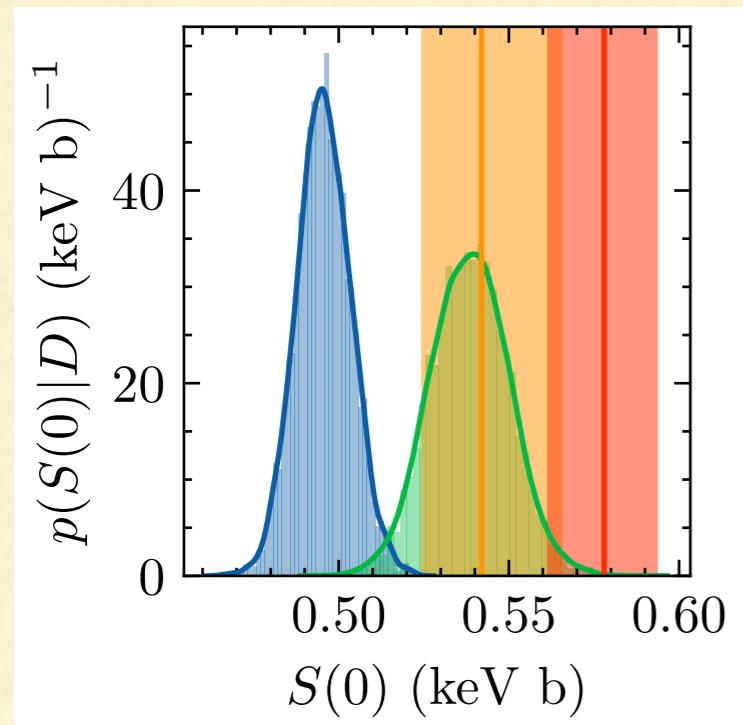


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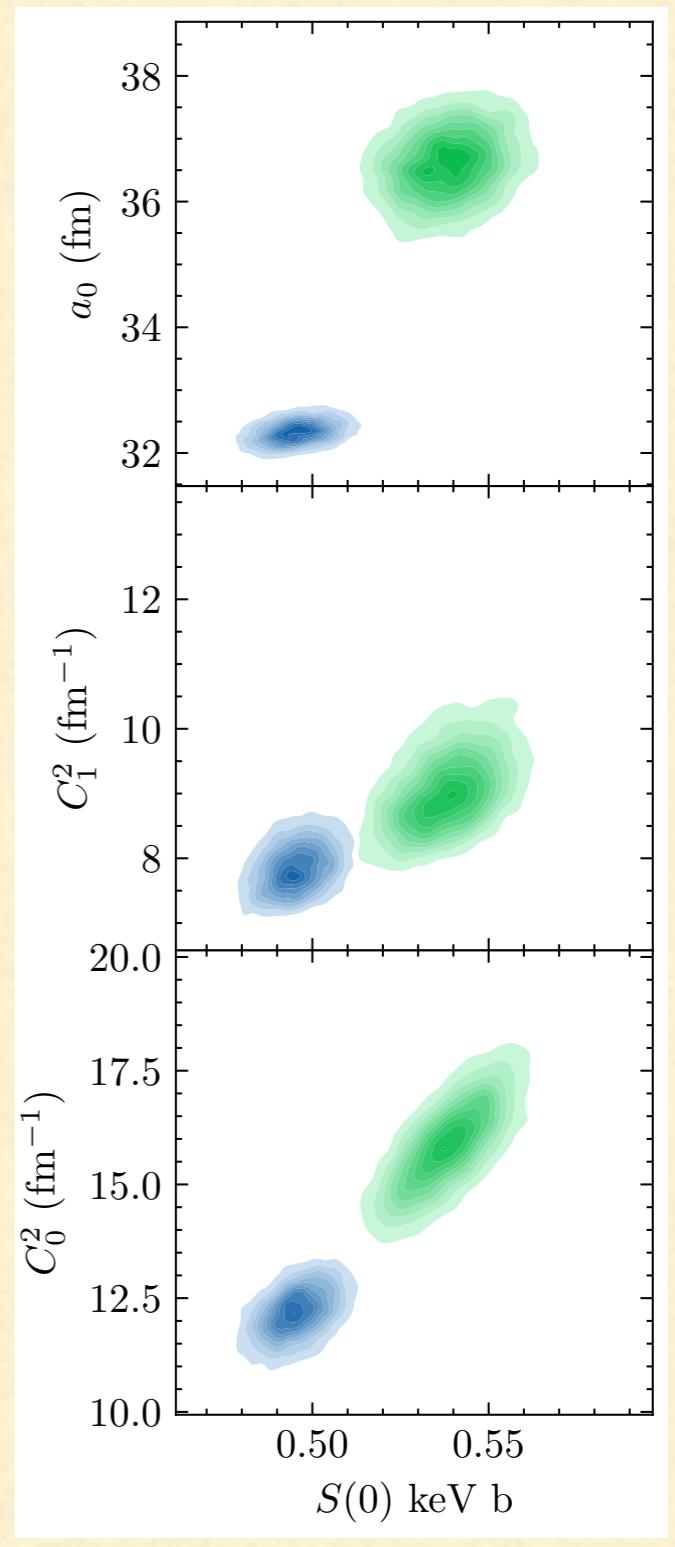
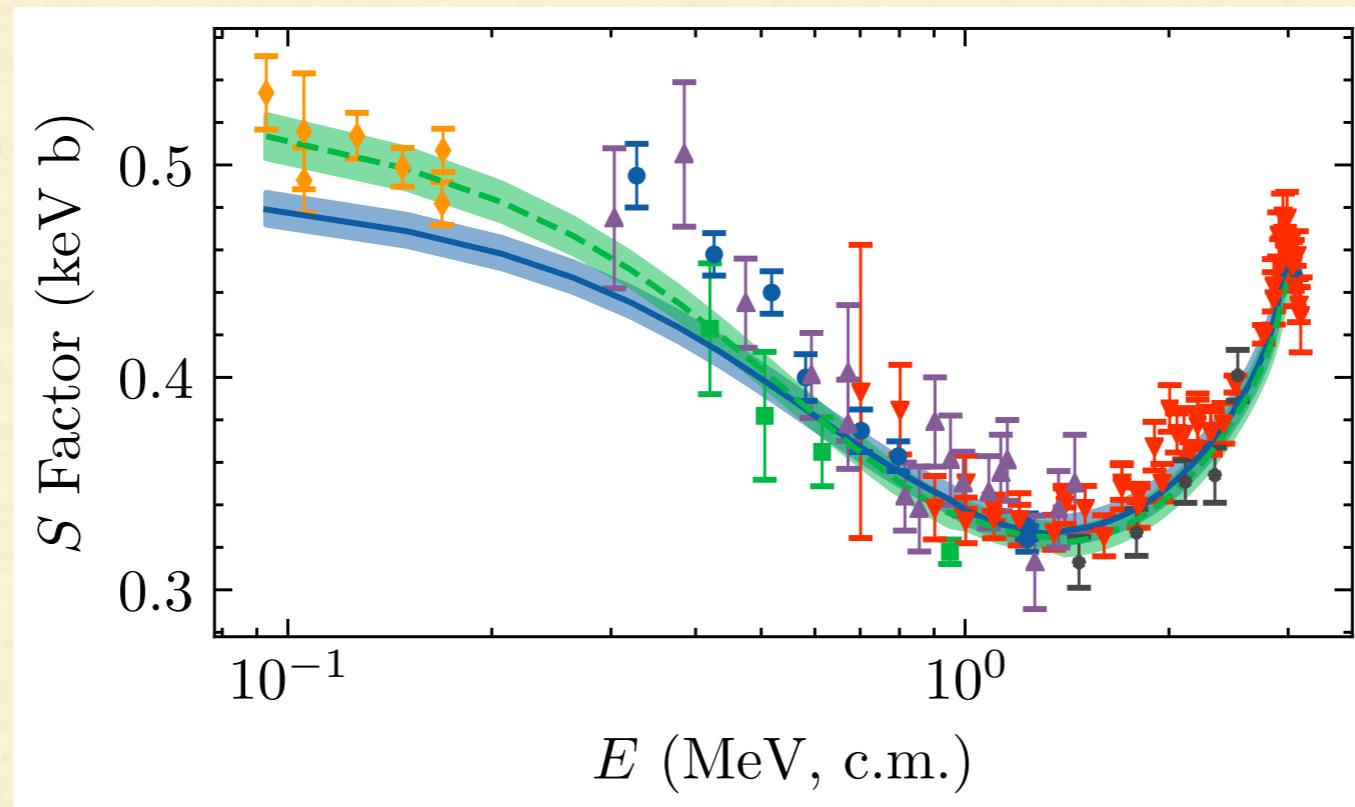


- Blue: CSB
- Green: CS
- Orange: de Boer et al.
- Red: Zhang, Nollett, DP

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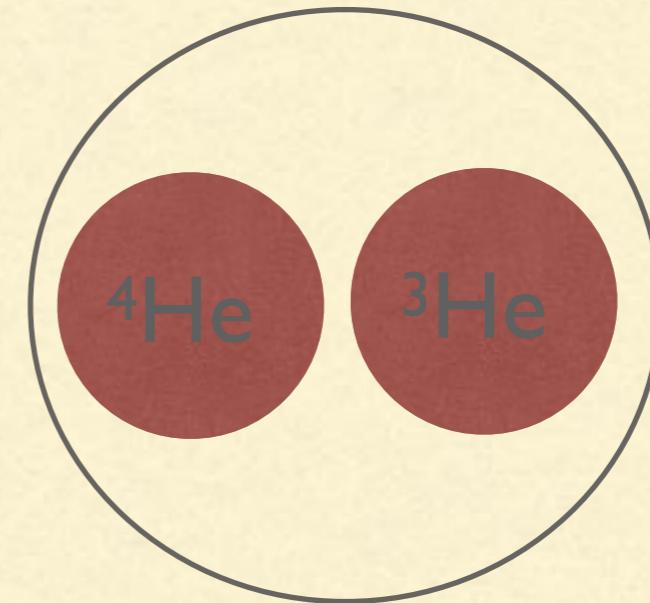
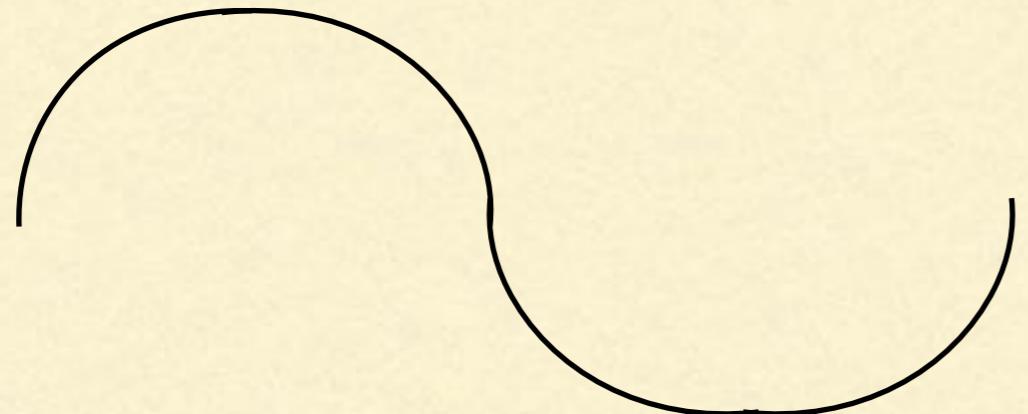
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Halo EFT

Bertulani, Hammer, van Kolck, NPA (2003);
Bedaque, Hammer, van Kolck, PLB (2003);
Reviews: Hammer, Ji, DP, J. Phys. G 44, 103002 (2017);

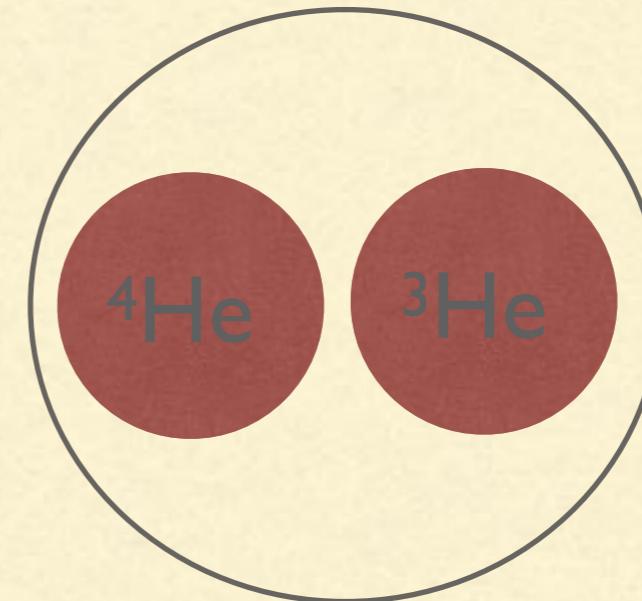
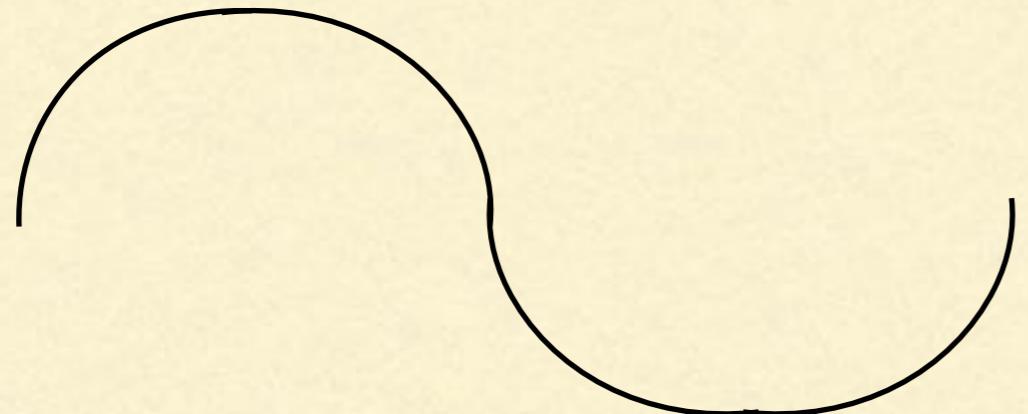
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$$\lambda \gg R_{\text{core}}; \lambda \lesssim R_{\text{halo}}$$



- Consider photo disintegration of ^7Be at long wavelength
- Define $R_{\text{halo}} = \langle r^2 \rangle^{1/2}$. Seek EFT expansion in $R_{\text{core}}/R_{\text{halo}}$. Valid for $\lambda \lesssim R_{\text{halo}}$
- Here $R \equiv R_{\text{core}} \sim 1.5$ fm. So this approach should be valid up to momenta of order 100 MeV
- Updates and systematizes cluster models

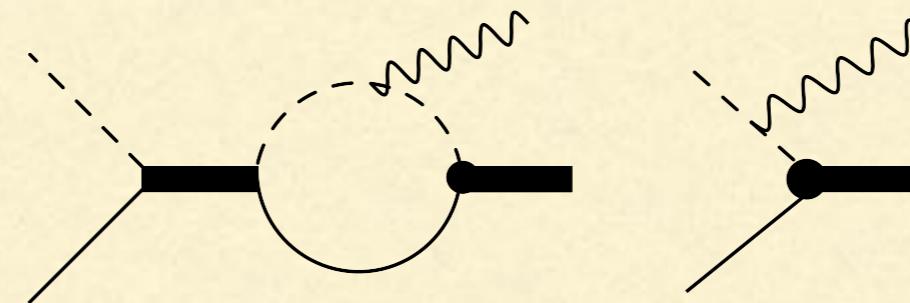
p-wave bound states and capture thereto

Hammer & DP, NPA (2011)

- At LO p-wave In halo described solely by its ANC and binding energy

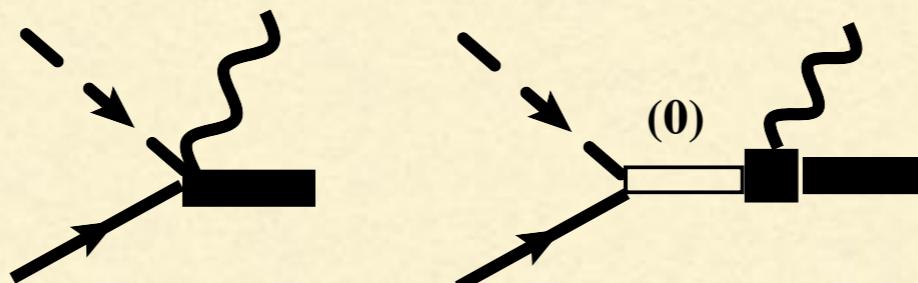
$$u_1(r) = A_1 \exp(-\gamma_1 r) \left(1 + \frac{1}{\gamma_1 r} \right) \quad \gamma_1 = \sqrt{2m_R B}$$

- Capture to the p-wave state proceeds via the one-body E1 operator: “external direct capture”



$$E1 \propto \int dr u_1(r) r (\cos(kr) + \sin(kr) \cot \delta); k \cot \delta \text{ from ERE}$$

- NLO: piece of the amplitude representing capture at short distances, represented by a contact operator \Rightarrow there is an LEC that must be fit



$^3\text{He} + ^4\text{He} \rightarrow ^7\text{Be} + \gamma_{\text{EI}}$ at LO in Halo EFT

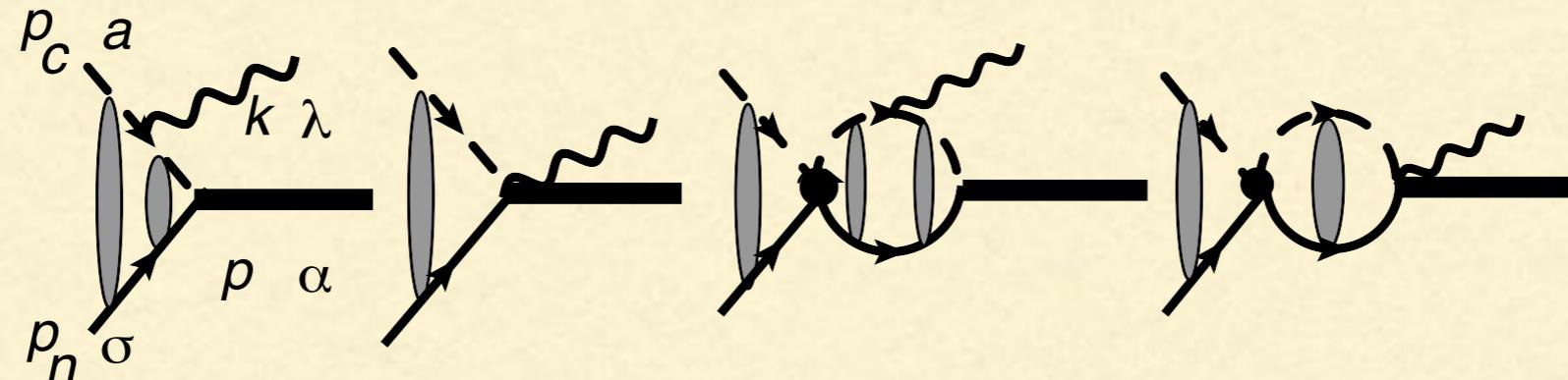
Zhang, Nollett, DP, JPG (2019), cf. Rupak, Higa, Vaghani, EPJA (2018)

- In this system $R_{\text{core}} \sim 1.5 \text{ fm}$, $R_{\text{halo}} \sim 3 \text{ fm}$

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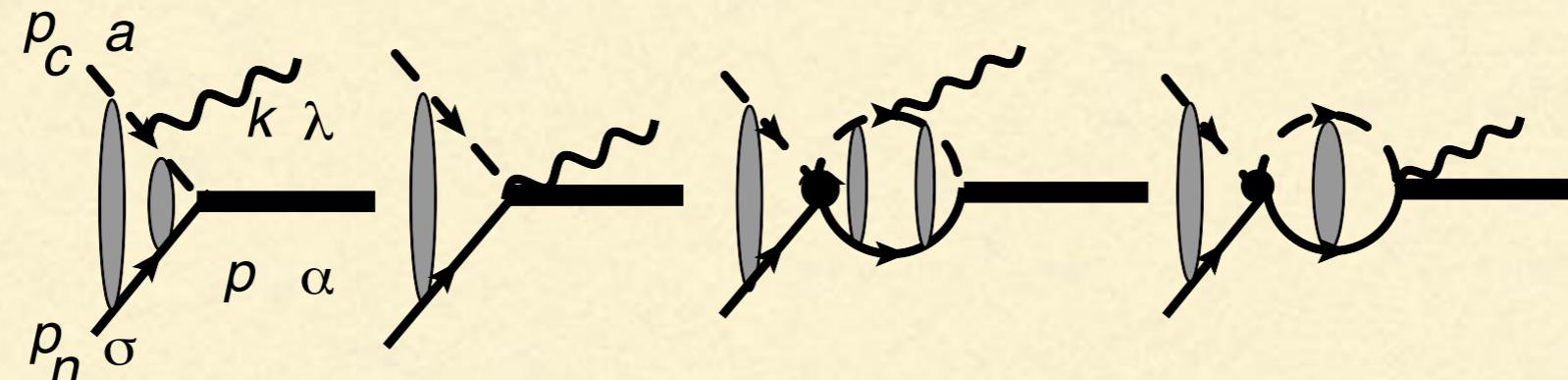
- In this system $R_{\text{core}} \sim 1.5 \text{ fm}$, $R_{\text{halo}} \sim 3 \text{ fm}$
- Also need to include Coulomb interactions non-perturbatively:
 $k_C = Q_c Q_n a_{\text{EM}} M_R = 17 \text{ MeV}$; $a \sim 10s$ of fm, both $\sim R_{\text{halo}}$



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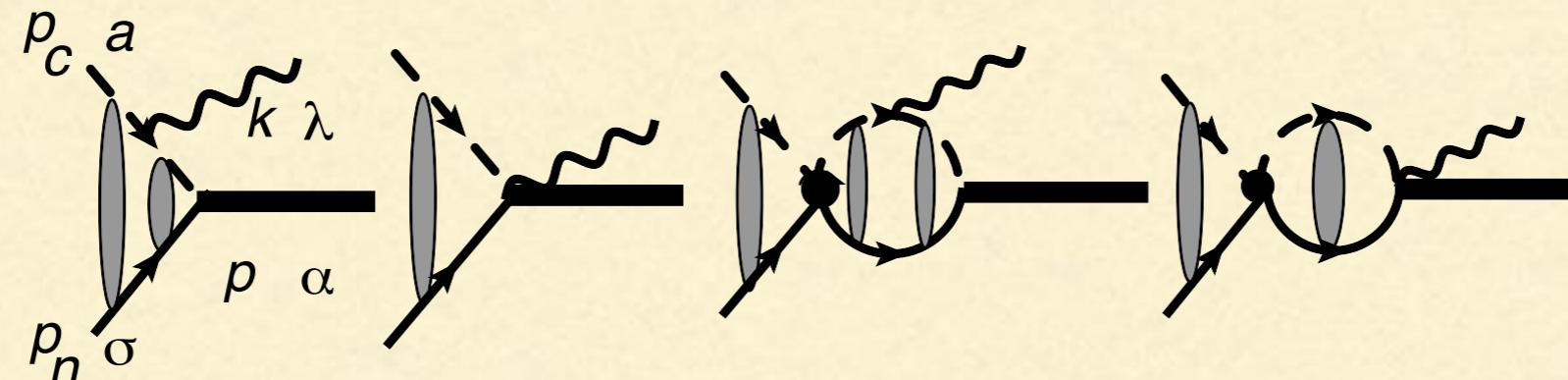


- Scattering wave functions are linear combinations of Coulomb wave functions F_0 and G_0 . Bound state wave function = the appropriate Whittaker function.

$^3\text{He} + ^4\text{He} \rightarrow ^7\text{Be} + \gamma_{\text{EI}}$ at LO in Halo EFT

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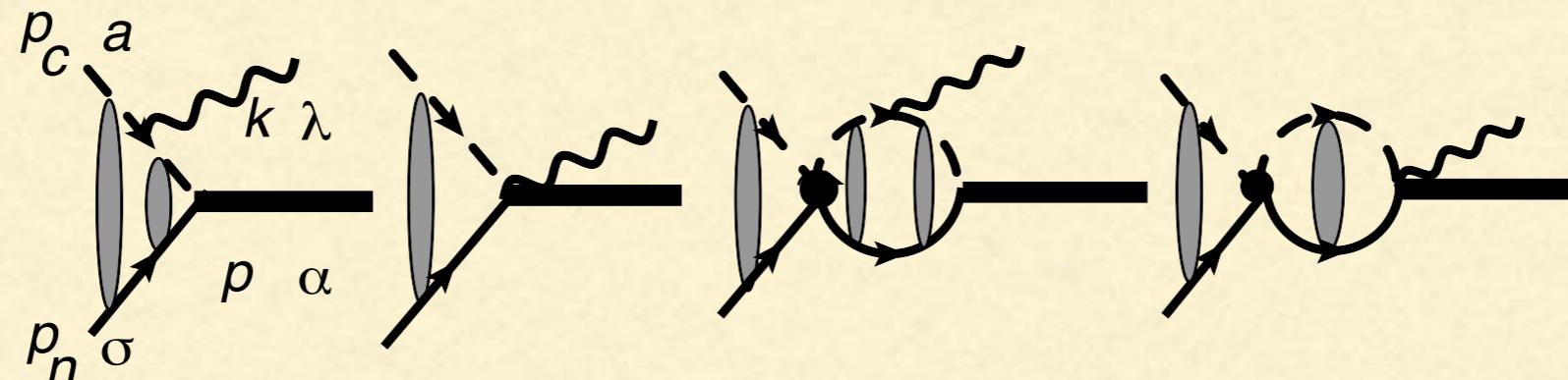
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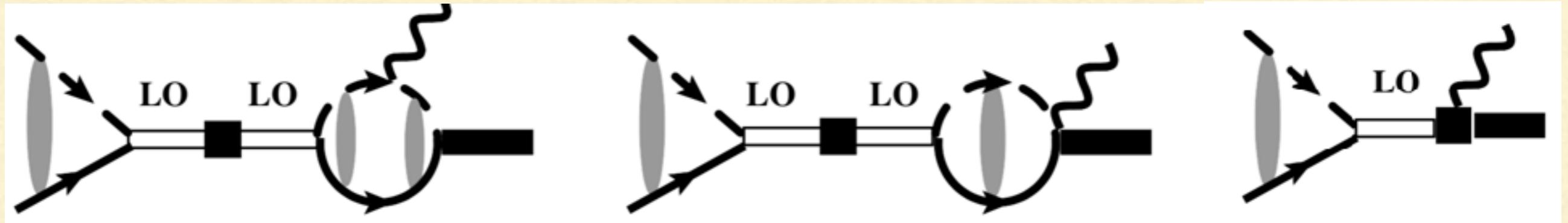
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- Can also predict capture to the excited $1/2^-$ in ^7Be

Three parameters at leading order

Additional ingredients at NLO

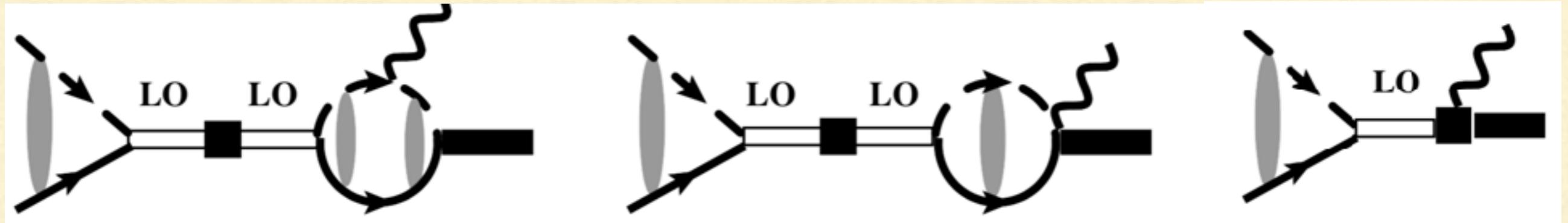


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Three more parameters at NLO

- Effective range (can add shape parameter which enters at N^3LO)
- LECs associated with contact interaction, \bar{L} and \bar{L}_*
- Can also consider contact interaction for D-wave capture, \bar{L}_D (enters at N^4LO)

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Pick data sets

- 88 S-factor data
 - Seattle (S)
 - Weizmann
 - Luna (L)
 - Erna
 - Notre Dame
 - Atomki
 - Plus 34 branching-ratio data
-

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 - Seattle (S)
 - Weizmann
 - Luna (L)
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 - Atomki
- Specify CMEs
 - S-factor: by set
 - Branching ratio: none
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$^3\text{He}(^4\text{He},\gamma)$ results

Zhang, Nollett, DP, JPG (2019)
cf. Higa, Rupak, Vaghani, EPJA (2018)

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- EI external direct capture to a shallow p-wave bound state
- Only one spin channel

$^3\text{He}(^4\text{He},\gamma)$ results

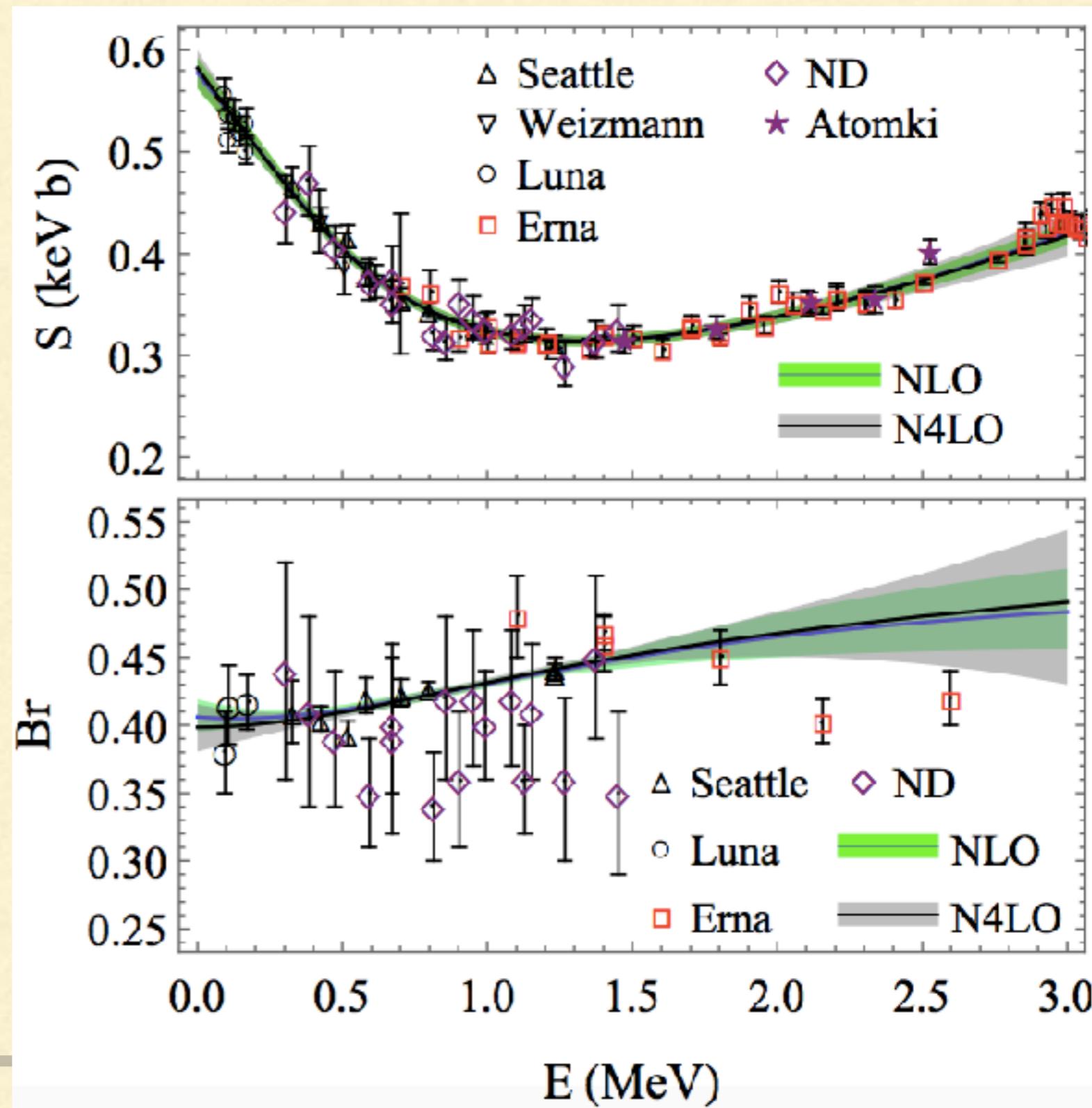
Zhang, Nollett, DP, JPG (2019)
cf. Higa, Rupak, Vaghani, EPJA (2018)

- E1 external direct capture to a shallow p-wave bound state
- Only one spin channel
- Integral is not dominated by as large r as in $^7\text{Be}(\text{p},\gamma)$
- More sensitivity to ^3He - ^4He scattering parameterization

$^3\text{He}(^4\text{He},\gamma)$ results

Zhang, Nollett, DP, JPG (2019)
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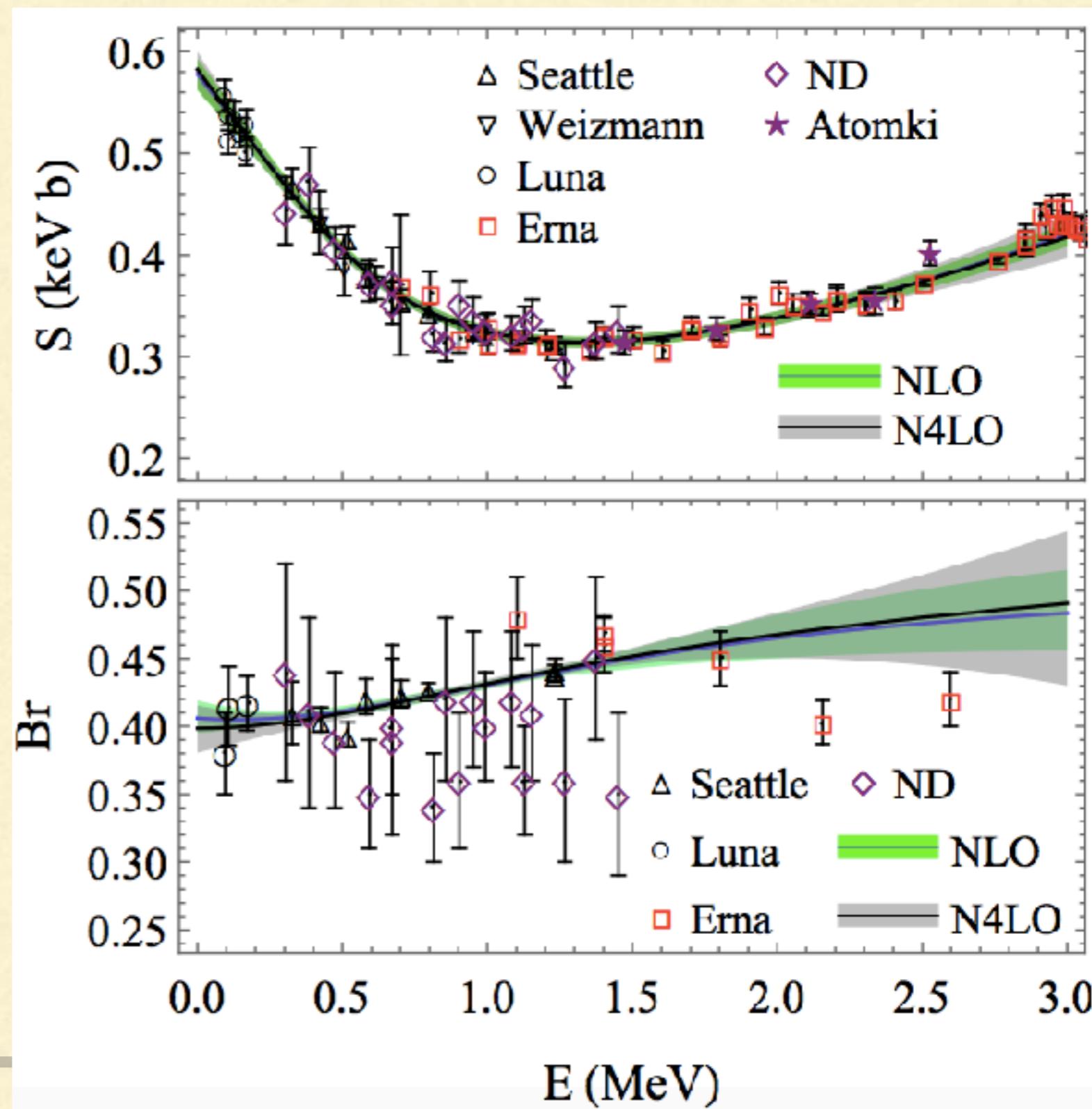
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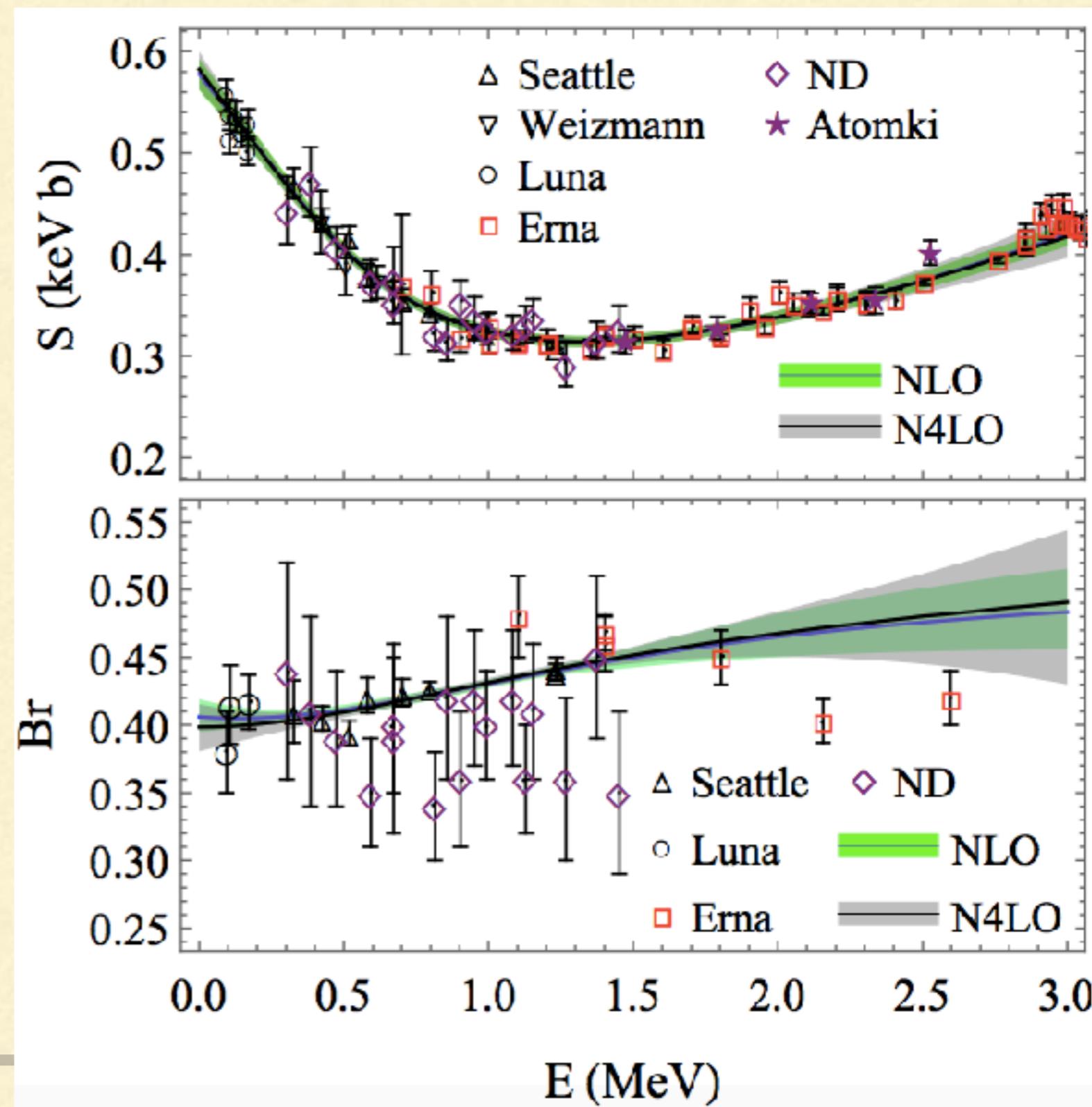
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- Distribution peaks at $\chi^2 = 82$



$^3\text{He}(^4\text{He},\gamma)$ results

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- Distribution peaks at $\chi^2 = 82$
- Bayesian evidence ratio ≈ 6 for NLO cf. N⁴LO



EFT treatment of $^3\text{He} + ^4\text{He}$ scattering

Poudel, Phillips, JPG (2022)

- Analyze SONIK data, Barnard data, and Boykin et al. A_y data
- Using Halo EFT to N2LO, $\mathcal{O}(Q^2)$
- $1/2^+$: a_0, r_0
- $1/2^-, 3/2^-$: a_1, r_1, P_1 ($\Leftrightarrow E_{^7\text{Be}}, \text{ANC}, P_1$) $\Lambda = 200 \text{ MeV}$
- $7/2^-$: Resonance at $E_{\text{cm}} = 2.98 \text{ MeV}$ with fitted Γ (R-matrix form)
- Likelihood: includes theory uncertainty based on convergence pattern of EFT expansion.

$$Q = \frac{(p, q)}{\Lambda}$$

$$\Lambda = 200 \text{ MeV}$$

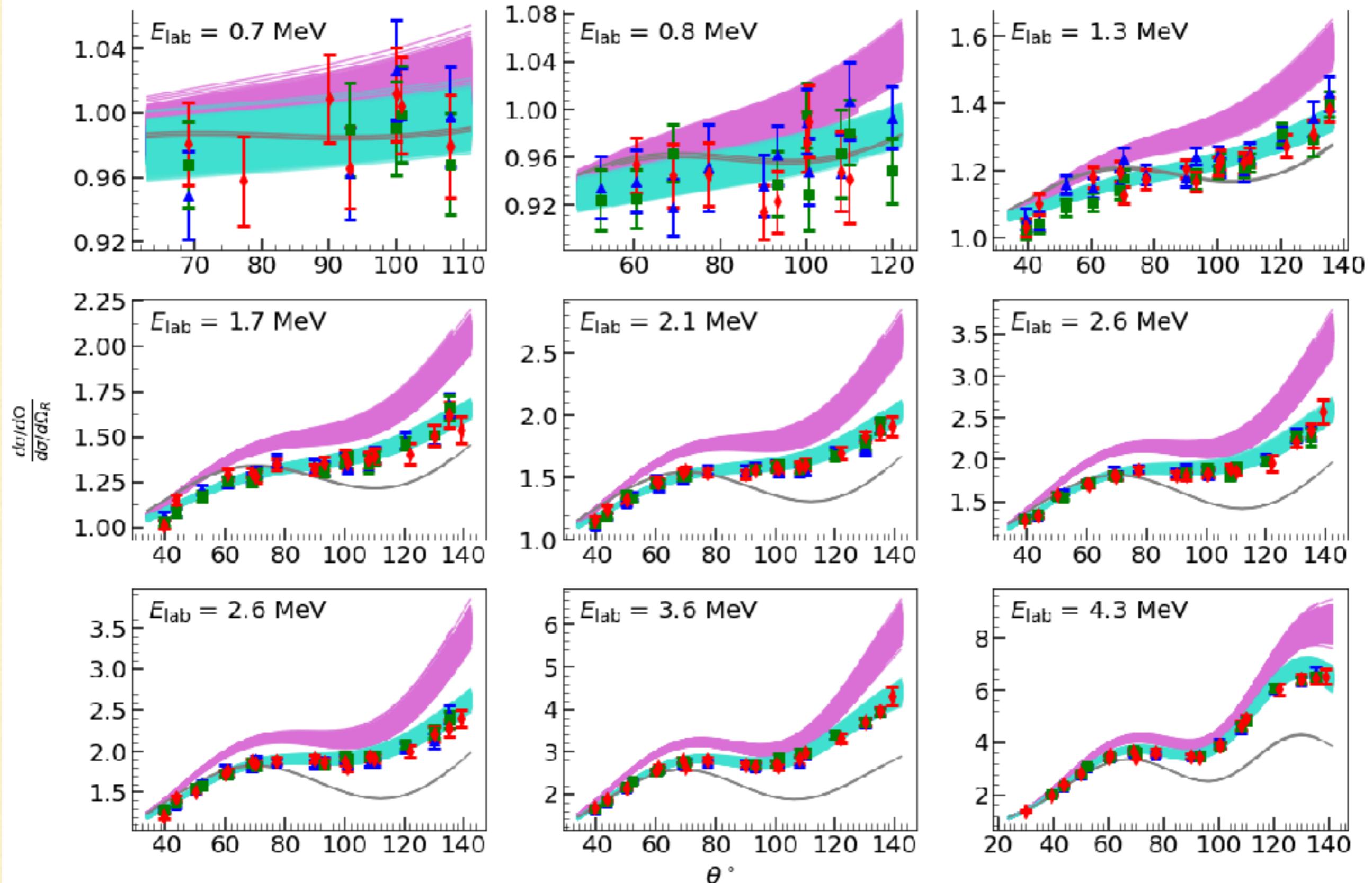
$$\mathcal{L} \propto \exp \left(-\frac{1}{2} \vec{r}^T (\Sigma^{\text{th}} + \Sigma^{\text{exp}}) \vec{r} \right)$$
$$\Sigma_{\alpha\beta}^{\text{th}} = (\Delta y)_\alpha (\Delta y)_\beta$$
$$r_{j\alpha} = y_{j\alpha} - f_\alpha \mu(x_{j\alpha}; \theta_{\text{EFT}})$$
$$\Delta y_\alpha = y_{\text{ref}} \bar{c} Q_\alpha^{\nu+1}$$

Pick data sets

- Scattering data
 - SONIK*: 451 from 0.385 to 2.7 MeV
 - Barnard: 646 from 1.49 to 2.7 MeV
 - Boykin et al.: 9 A_y data from 2.1 to 2.7 MeV
- Specify CMEs
 - SONIK: by energy
 - Barnard: 5%

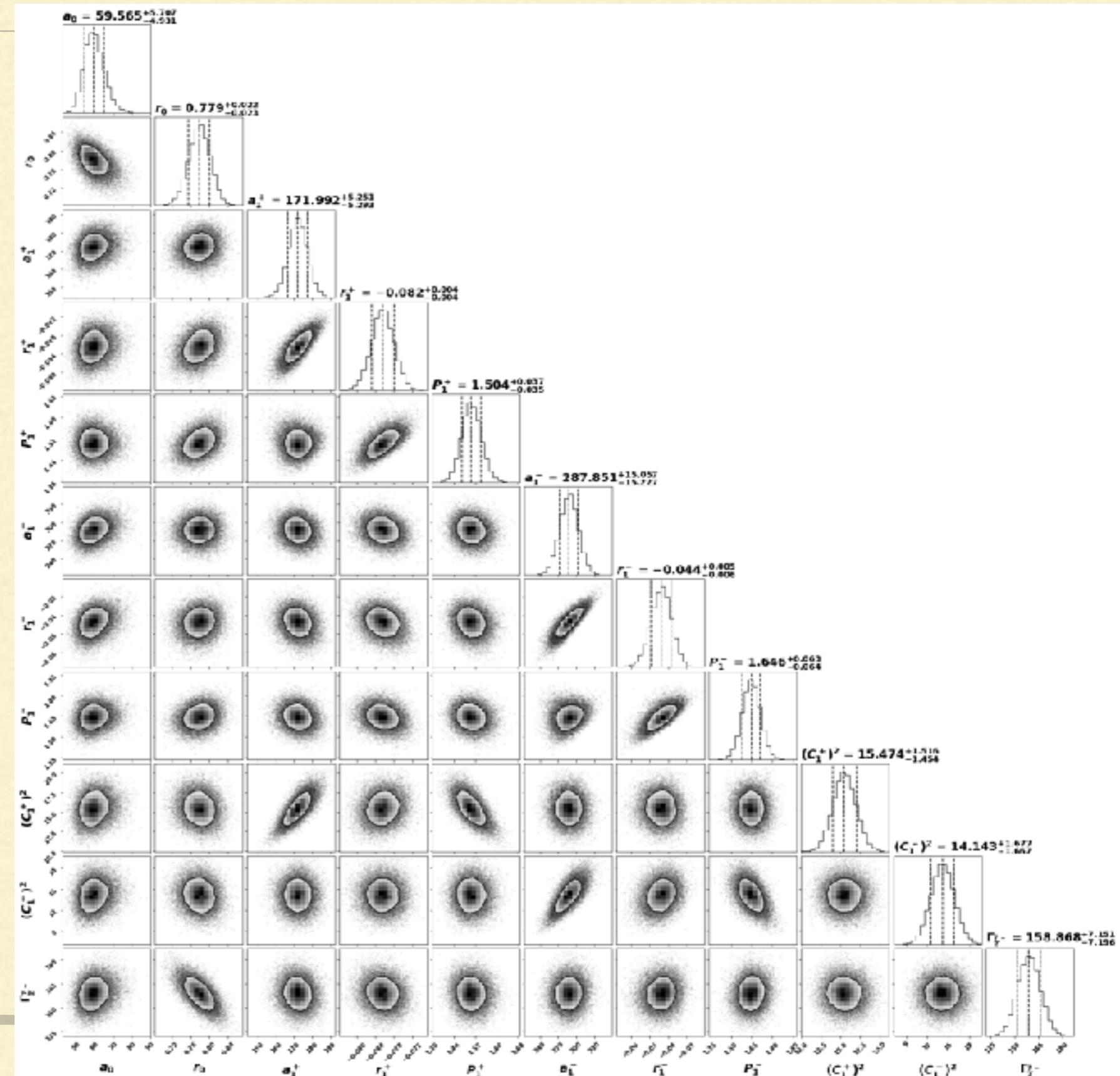
* Paneru et al., submitted to Phys. Rev. C; Paneru, Ohio University Ph.D. thesis, 2020

SONIK data, with truncation uncertainty



ERT parameters from scattering data

- Imposed prior on ANC_s from capture data, so not solely from scattering data
- Consistent values:
 $C_1^{+2} = 15.5 \pm 1.5$ fm;
 $C_1^{-2} = 14.1 \pm 1.7$ fm
- $a_0 = 60_{-5}^{+6}$ fm cf.
 $a_0 = 50_{-6}^{+7}$ fm from
capture and lower
number from R-
matrix

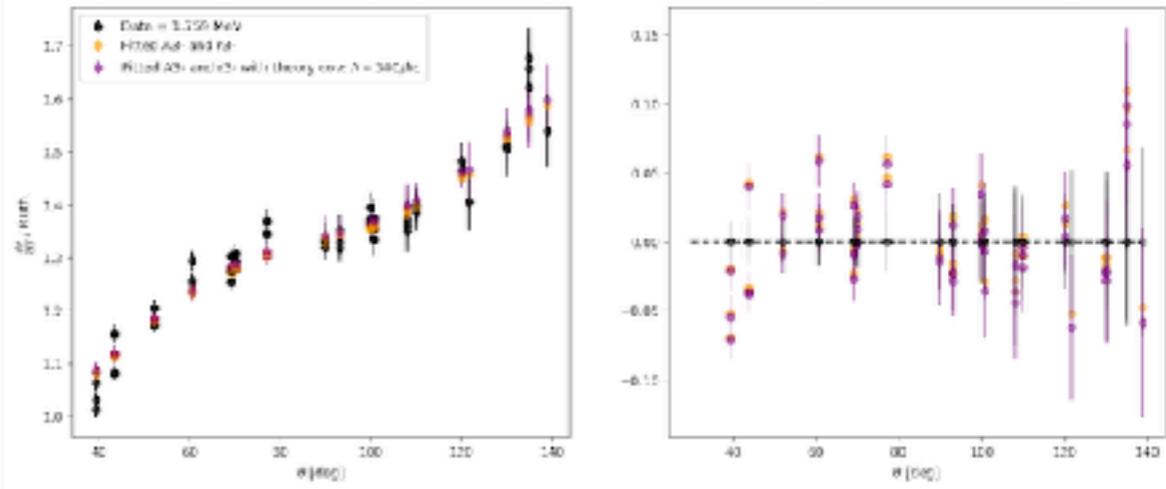
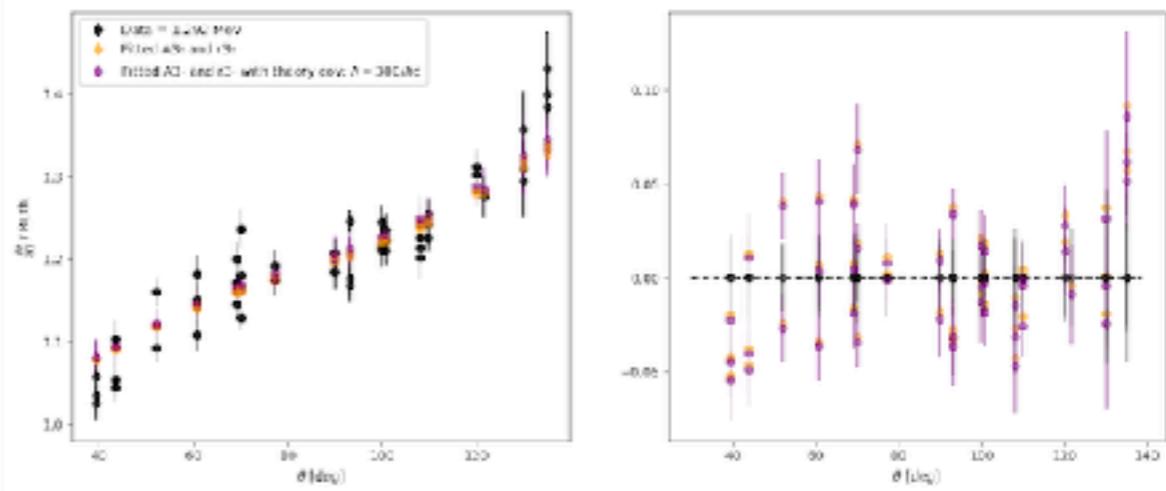
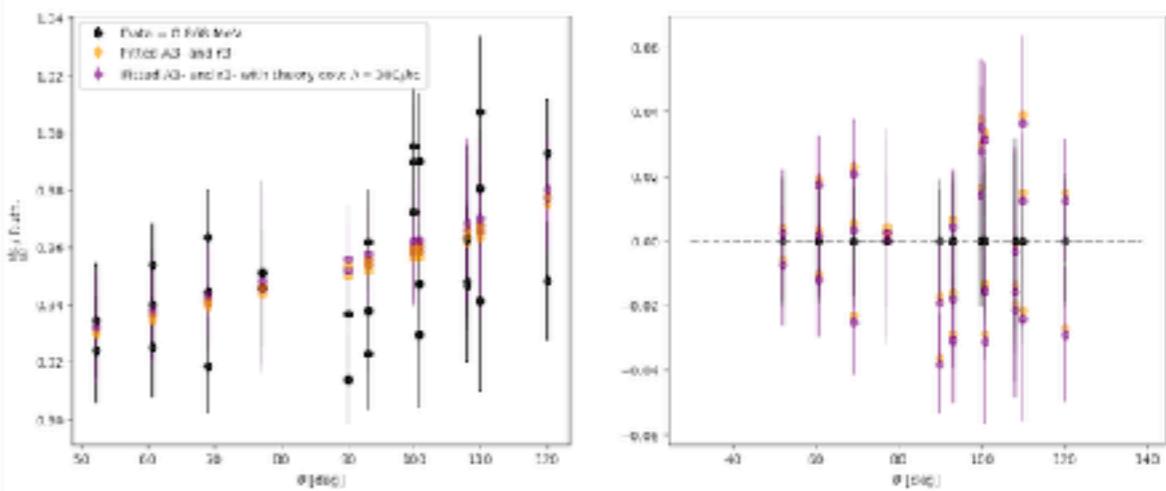
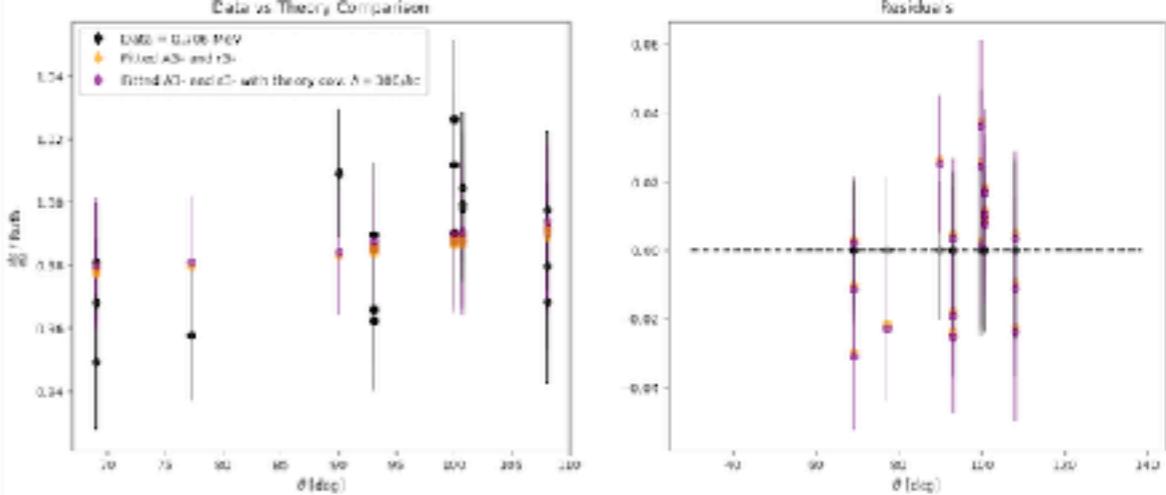


Inclusion of 7/2- and 5/2- resonances in Halo EFT

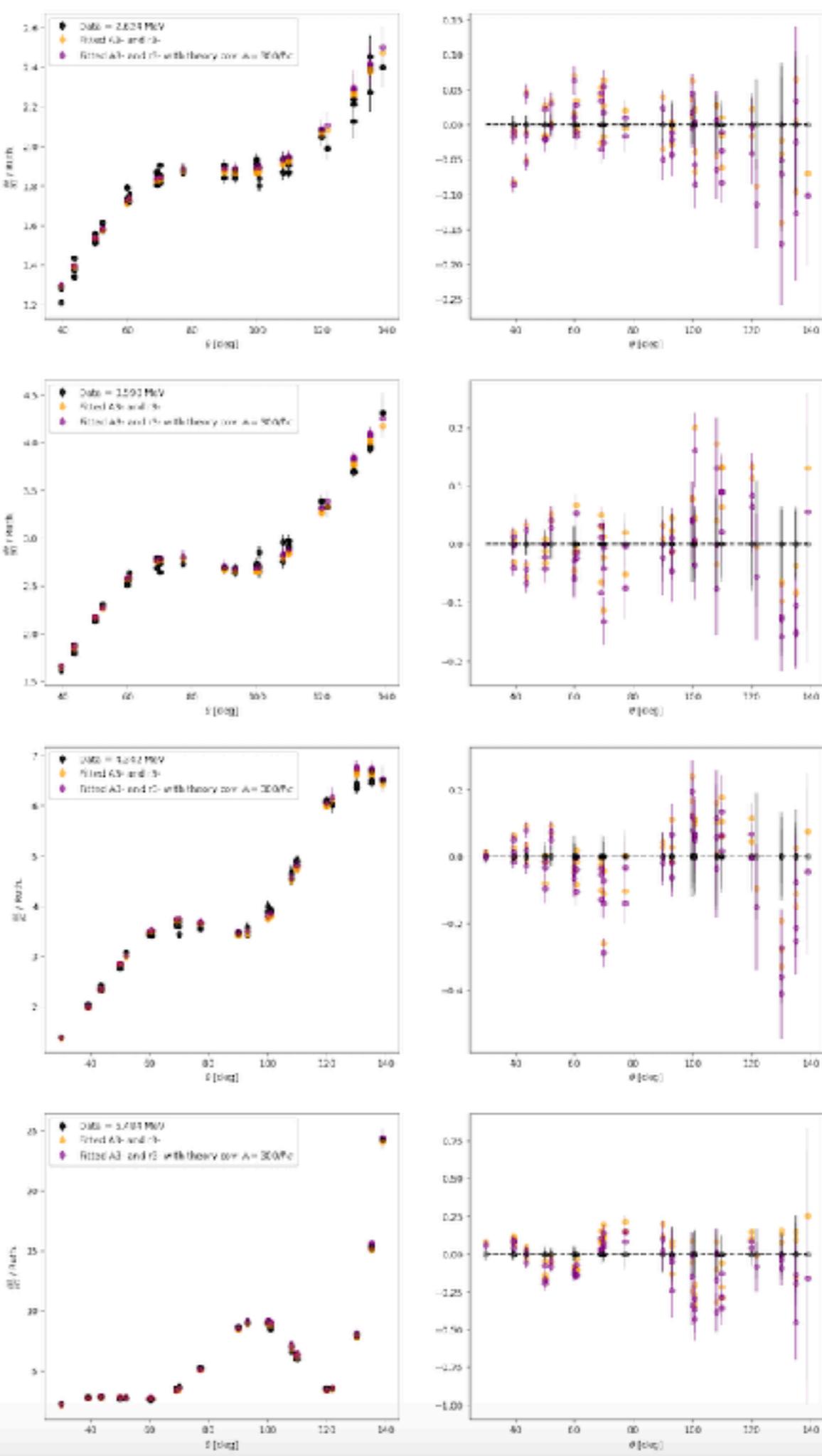
Burnelis, Phillips

- Analyze SONIK data
- Using Halo EFT to N2LO, $\mathcal{O}(Q^3)$
- $1/2^+$: a_0, r_0, P_0
- $1/2^-, 3/2^-$: a_1, r_1, P_1 ($\Leftrightarrow E_{7\text{Be}}, \text{ANC}, P_1$)
- 7/2-: Resonance at $E_{\text{cm}}=2.98$ MeV with form given by effective-range theory up to fourth order \Rightarrow width fitted to data
- 5/2-: fit effective-range theory up to second order to Boykin phase shifts and take as fixed
- Likelihood: includes theory uncertainty based on convergence pattern of EFT expansion.

Cross sections

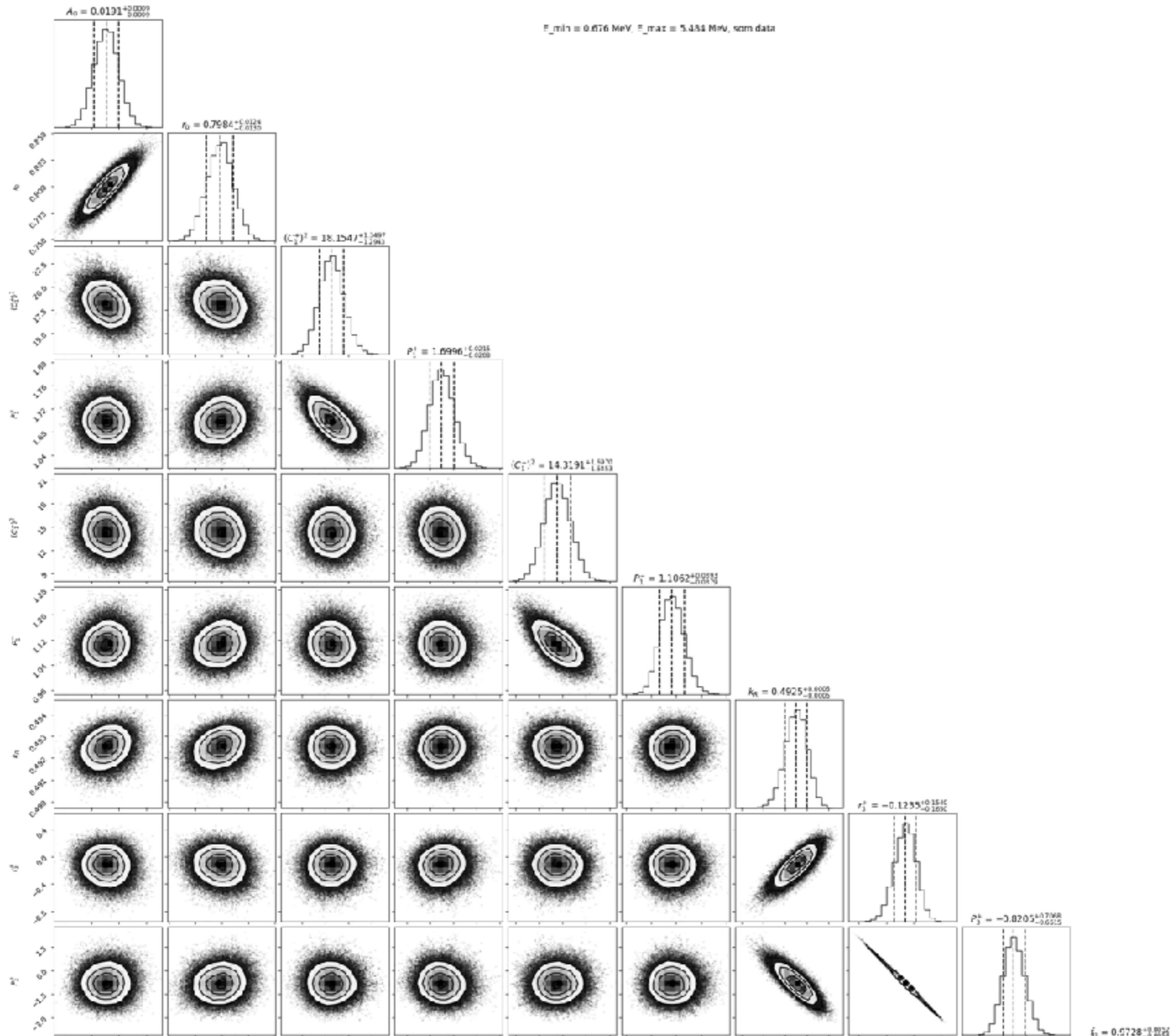


Cross sections



$$A_0 = 0.0191^{+0.0009}_{-0.0009}$$

E_min = 0.676 MeV, E_max = 5.484 MeV, sum data



Further applications of BRICK



- $^{19}\text{F}(p, \gamma)^{20}\text{Ne}$ – Zhang et al. (incl. deBoer, Odell) **Nature** **610**, 656-660 (2022)
 - Low-energy resonance opens up possibility of “warm” CNO breakout
- $^{10}\text{B}(p, \alpha)^7\text{Be}$ – Van de Kolk et al. (incl. deBoer, Odell) **PRC** **105**, 055802 (2022)
 - possible temperature probe for $^{11}\text{B}(p, 2\alpha)^4\text{He}$ – aneutronic plasma fusion source
- $^{23}\text{Na}(p, \gamma)^{24}\text{Mg}$ – Boeltzig et al. (incl. deBoer, Odell) **PRC** **106**, 045801 (2022)
 - breakout reaction linking NeNa and MgAl cycles
- $^{13}\text{C}(\alpha, n_1)^{16}\text{O}$ – deBoer et al. (incl. deBoer, Odell) **PRC** **106**, 055808 (2022)
 - partial cross section measurement, improves BG modeling

Summary

- Parametric uncertainties in R-matrix analyses can be quantified by MCMC sampling of the Bayesian posterior and evaluating derived quantities
-  <https://github.com/odell/brick>
- Multiple examples of successful application to different reactions
- Enables more sophisticated modeling of experimental imperfections
- Knowledge of full posterior provides access to parameter correlations, allows diagnosis of which parameters are not needed, shows where there is multimodality, non-Gaussianity, and more
- Error propagation to derived quantities is straightforward with samples in hand
- Model checking (residuals, coverage, etc.) needs to be done at end
- Model uncertainties of R-matrix analysis? Comparison to EFT, ab initio, etc.

