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Experimental analyses and numerical modelling of trace neon shattered pellet injection discharges on JET

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Highlights

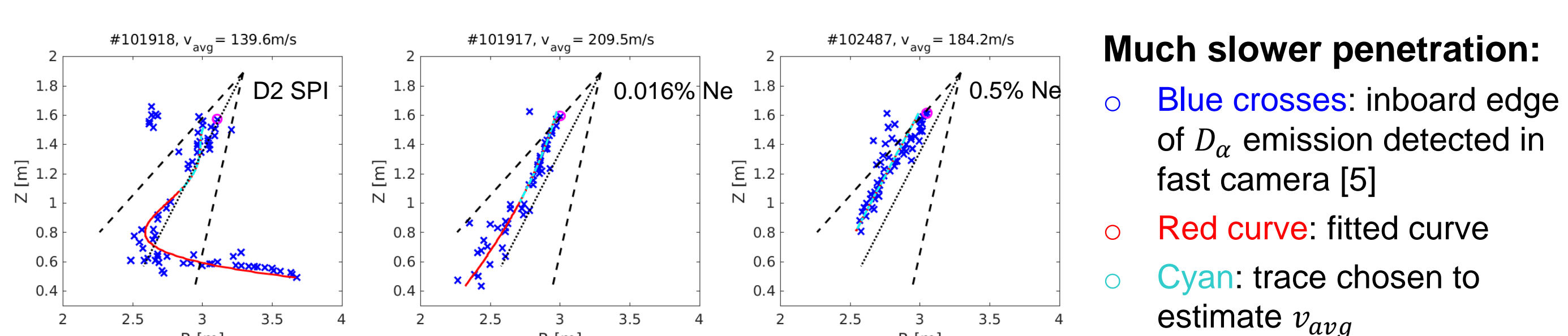
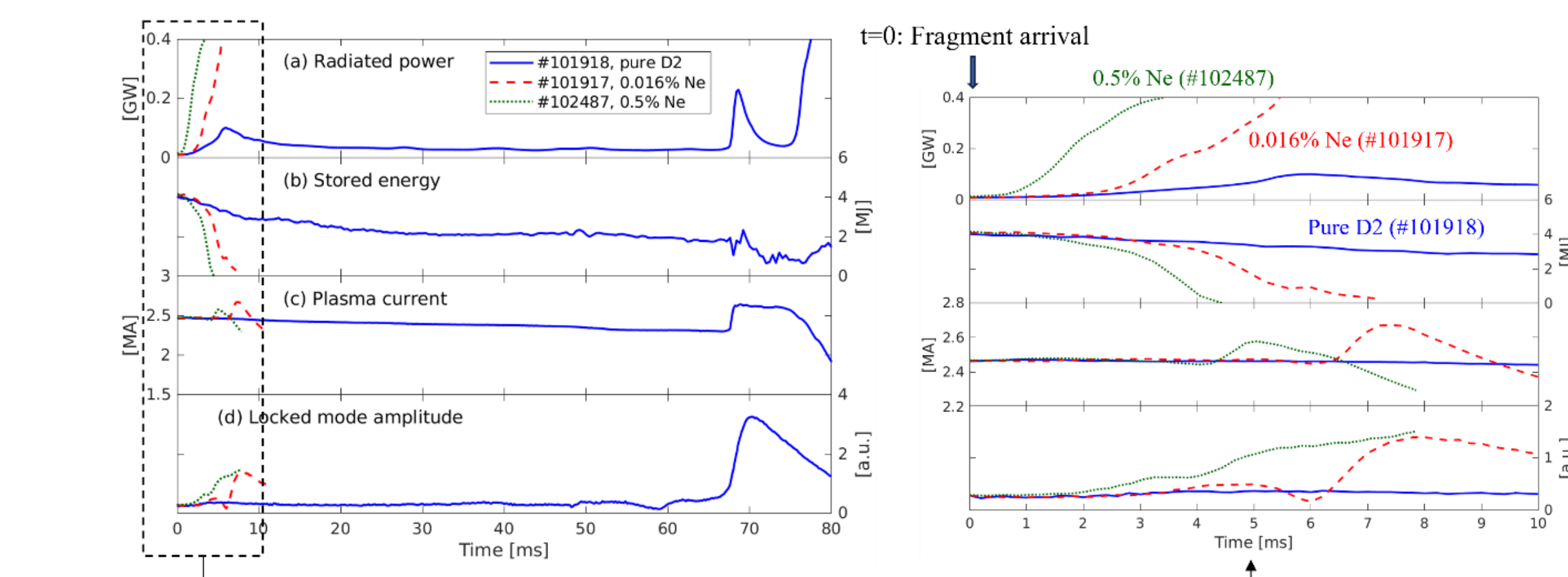
- Trace neon shattered pellet injection (SPI) discharges performed in final JET campaigns, demonstrating suppressed rocket motion of SPI fragments up to thermal quench (TQ) onset but much shorter pre-TQ compared to deuterium SPI
- Interpretative modelling on-going with the 3D non-linear MHD code JOEKE [1] and the 1.5D transport code INDEX [2]

Background & Motivation

- SPI**: technique used by the ITER disruption mitigation system (DMS) to mitigate thermal loads, mechanical forces & runaway electrons (REs) [3]
- Hydrogen isotope** (such as deuterium -- D2) SPI: envisaged for **RE avoidance** by strongly increasing core density before thermal quench (TQ) [4]
- However, JET SPI experiments & JOEKE modelling demonstrated that D2 SPI suffers from plasmoid drifts, hindering material assimilation & RE avoidance [5]
- Plasmoid drift** [6]: ∇B drift that leads to charge separation and $E \times B$ drift towards the tokamak LFS \rightarrow imbalanced ablation rate on the HFS and LFS of fragments \rightarrow outward motion of solid fragments (**rocket motion/effect**) [7]
- Trace neon SPI**: a small percentage of neon added to hydrogen-isotope SPI; envisaged to reduce pressure imbalance of the plasmoid via stronger local radiation, thus suppressing plasmoid drifts and facilitating RE avoidance [8]

JET trace neon SPI experiments

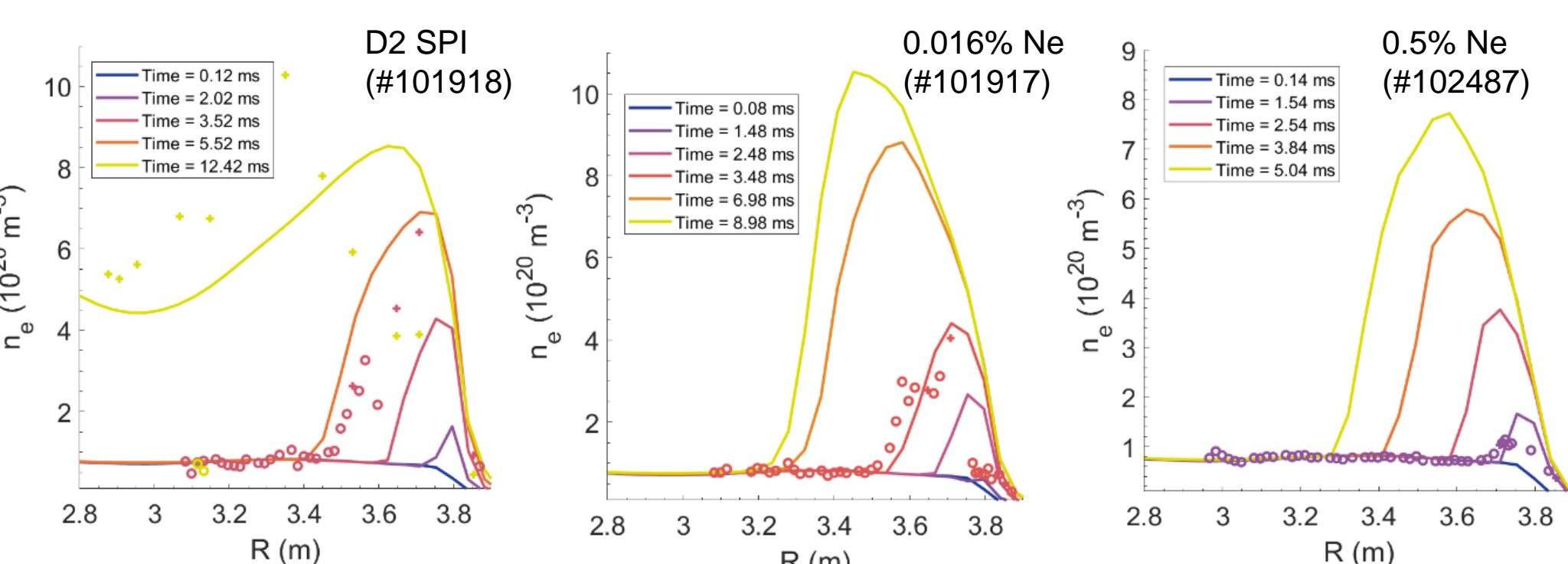
- Before SPI: H-mode plasma, 2.5 MA / 2.5 T, 4 MJ
- Pellet: 10 mm diameter, **350 m/s**, neon atomic mixture ratio 0.016% to 2% ($\sim 1.26 \times 10^{19}$ neon atoms in 0.016% neon SPI); D2 SPI performed as reference



- Two key observations**: 1) evident rocket motion in D2 SPI, but not in trace neon SPI with 0.016% neon up to TQ onset, indicating weaker plasmoid drifts; 2) cooling time, defined as the duration between SPI arrival at the plasma edge ($t = 0$) and the I_p spike, drops from ~ 70 ms in D2 SPI to ~ 7 ms in 0.016% neon

Modeling with 1.5D INDEX code

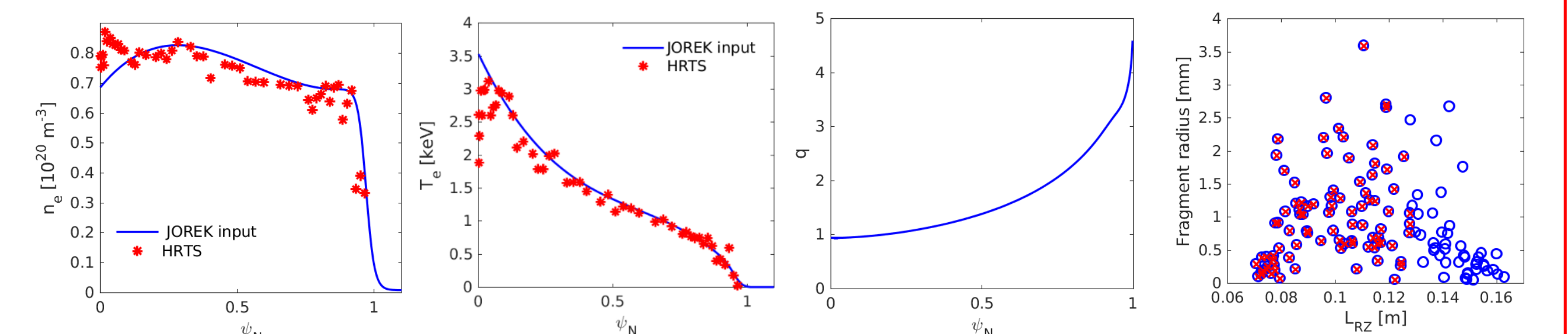
- Same input profiles and SPI setup as in JOEKE simulations
- A back-averaging model [9-10] used for plasmoid drifts: ablated material distributed uniformly between fragment and $\psi_N = 1$ here; fraction of ablated material shifted to fit TS data is 50%, 20% & 0 for D2, 0.016% & 0.5% neon SPI



- Comparing with Thomson scattering (TS) data shown in symbols:
 - Fragments reach HFS before the rocket motion in D2 SPI \rightarrow increased core fueling

JOEKE model & simulation setup

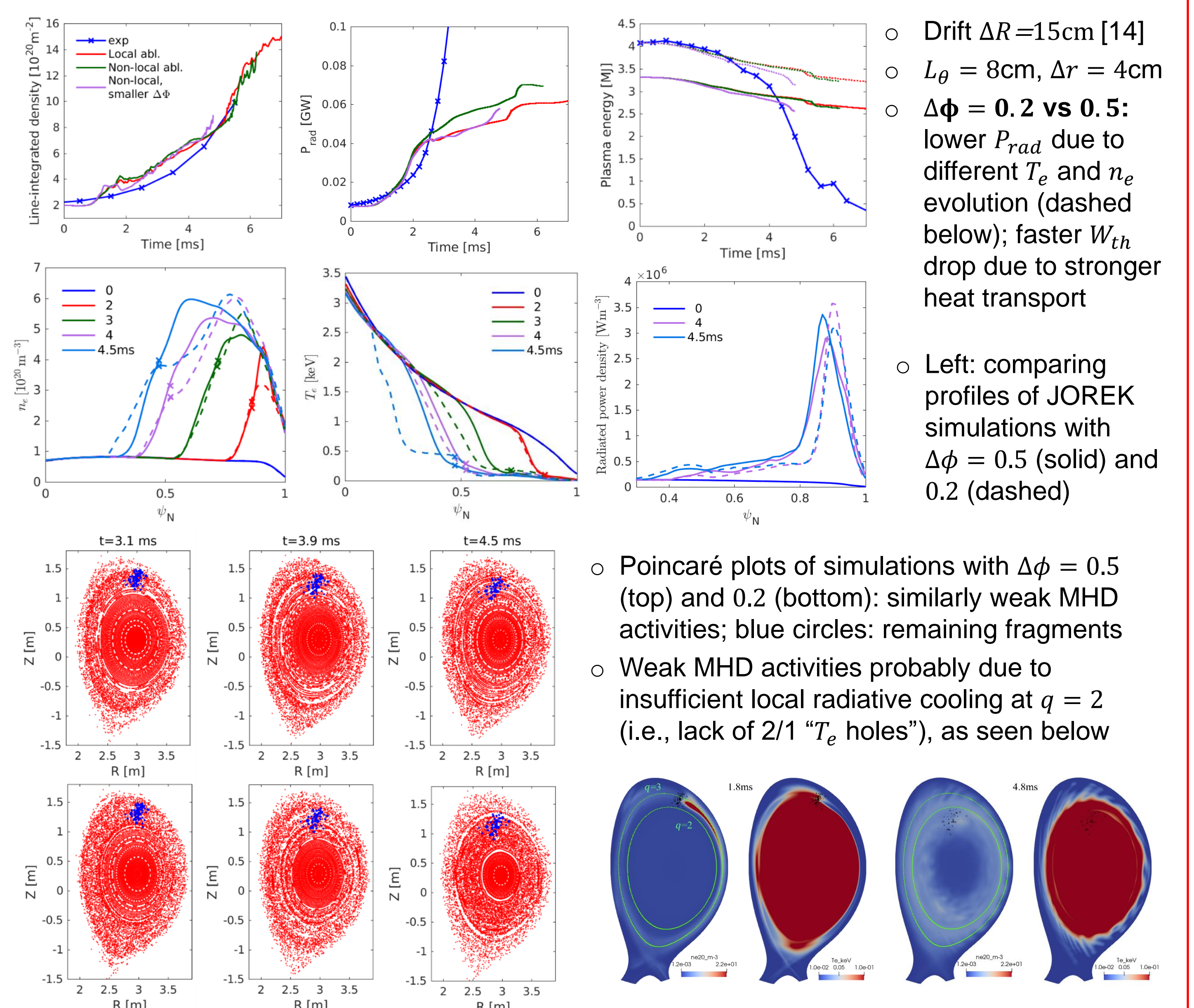
- Reduced MHD with parallel flows and extension to impurities [1]
- Coronal equilibrium (CE) for impurity radiation, including uniform **background tungsten & neon**; collisional-radiative (CR) model [11] available



- Ablation rate $\propto T_e^{5/3} \tau_p^{4/3} n_e^{1/3}$, where T_e and n_e taken from fragment location (**local-ablation model**) or flux-surface-averaged (**non-local ablation**)
- Ablated source deposited around each fragment (no drift) or radially shifted by ΔR (with drifts) via a "teleportation" model [5], both with Gaussian shape
- Fragment size distribution follows a **fragmentation model** [12]; front fragments removed to consider material loss in shattering process, where the gas generated can take up to 20% of initial pellet mass [13]
- SPI **penetration velocity** is only about half of the injection velocity, possibly due to braking by shock waves generated in the gas created in the shattering process [15]

3D JOEKE modeling results

- Modelling of #101917 (0.016% neon SPI) as an example**



Conclusions & Outlook

- JET trace neon SPI experiments** performed, exhibiting suppressed rocket motion at least up to TQ onset & $\sim 10\times$ shorter cooling time compared to D2 SPI
- INDEX and JOEKE simulations** suggest that plasmoid drift is weaker in 0.016% neon than in D2 SPI and there is no evident drift with 0.5% neon SPI
- Difficult to directly recover the observed early TQ onset in trace neon SPI in 3D JOEKE simulations, hypothesized to result from insufficient local cooling with toroidally oversized plasmoid typically used in 3D non-linear MHD modelling
- Outlook**: artificially condensed radiation sinks to investigate the susceptibility of plasma to T_e holes and the necessary conditions to trigger the observed TQ

References

- [1] M. Hoelzl et al. NF 61 (2021) 065001
 [2] A. Matsuyama et al. PPCF 64 (2022) 105018
 [3] S. Jachmich et al. NF 62 (2022) 026012
 [4] E. Nardon et al. NF 60 (2020) 126040
 [5] M. Kong et al. NF 64 (2024) 066004
 [6] P.B. Parks et al. POP 7 (2000) 1968
 [7] H.W. Mueller et al. NF 42 (2002) 301-9
 [8] A. Matsuyama et al. POP 29 (2022) 042501
 [9] A. Patel et al. NF 65 (2025) 086031
 [10] A. Lvovskiy et al. NF 64 (2024) 016002
 [11] D. Hu et al. PPCF 63 (2021) 125003
 [12] T. Gebhart et al. IEEE TPS. 48 (2020) 1598-605
 [13] T. Gebhart et al. FST 76 7 (2020) 831-835
 [14] M. Kong et al. EPS 2025; to be submitted to PPCF
 [15] E. Nardon et al. Nucl. Fusion 57 (2017) 016027

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