

FAST ION TRANSPORT INDUCED BY EDGE LOCALIZED MODES

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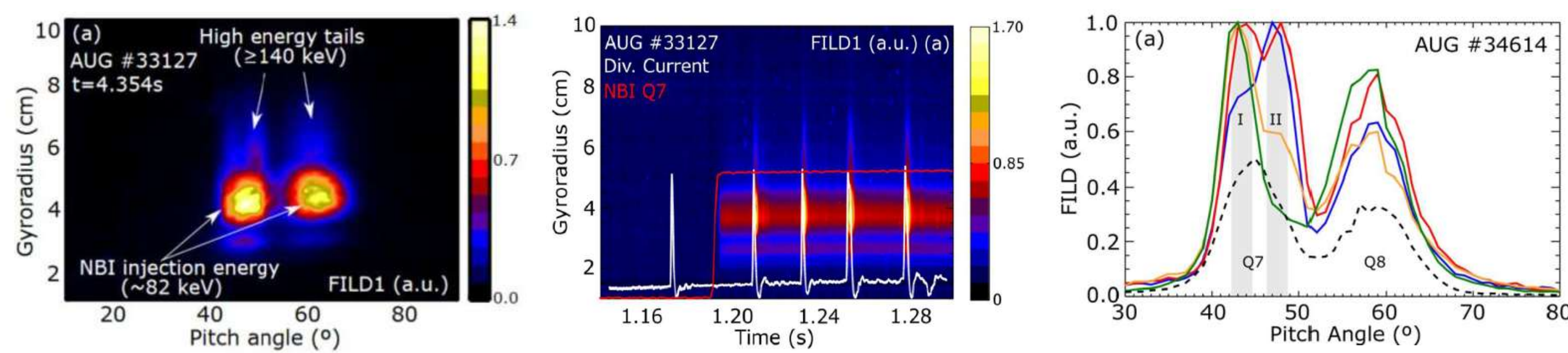
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Abstract

- Fast ions exhibit a notable acceleration during edge localized modes (ELMs) in tokamak devices.
- This paper presents an analytical investigation into the phase-space transport of fast ions driven by ELMs.
- Contrary to previous simulation results, it is shown that ELMs with low-frequency characteristics are inefficient at accelerating fast ions.
- Instead, the transport of fast ions is dominated by radial particle transport, resulting from the exchange of canonical toroidal angular momentum.
- The associated diffusivity increases sharply for high-energy particles, making fast-ion loss measurements in velocity space appear as an acceleration process

Background

- Velocity space measurements of fast-ion losses reveal a population at energies well above the main NBI injection energy during ELMs [Galdon-Quiroga PRL, 2018;NF, 2019].
- Two populations can be observed at two main different pitch angles of $\pi/4$ (Q7,passing) and $\pi/3$ (Q8, trapped).
- NBI injects neutrals at three different energies $E_0 = 82\text{keV}$, $E_0/2$ and $E_0/3$.
- The observation of high-energy tail is reproducible and well correlated with the NBI heating and the occurrence of ELMs.
- High-energy population exhibits a pitch angle structure (‘spikes’) that depends on the beam source and q.



Methods

Gyrokinetic Model

- The gyrocenter Hamiltonian in 5D phase space

$$H(\mathbf{X}, \mu, w, t, \tau) = \mu B + \frac{w^2}{2} + \frac{e}{m} \langle (\delta\phi) \rangle - \frac{w}{c} \langle \delta A_{\parallel} \rangle + \frac{v_{\perp} \rho}{c} \langle \delta B_{\parallel} \rangle_{*}.$$

Fluctuations

- During the ELM crash a broad spread in frequency is generally measured in AUG but low frequencies $\omega \sim 10\text{kHz}$ are dominant, with $n \sim 5$ and $\delta B_{\perp}/B \sim O(10^{-3})$.
- The inter-ELM modes in a high frequency range $\omega \sim 100\text{kHz}$ appears in the linear and early nonlinear phases, with $n \sim 10$ and $\delta B_{\perp}/B \sim O(10^{-5} - 10^{-4})$.
- Therefore, the perpendicular magnetic perturbation dominates the perturbed Hamiltonian of fast ions.

Fast Ion Acceleration

- An estimate for the time required for the particle acceleration

$$\Delta t \sim \frac{H_0}{\omega \delta H} \sim \frac{k_{\perp} \rho_0}{\omega} \frac{B_0}{\langle \delta B_{\perp} \rangle}$$

- Substituting into parameters of low-n modes and inter-ELM modes, respectively, one can estimate the required time as

$$\Delta t_{\text{low}} \sim O(10^3)/\omega_{\text{low}} \sim 1\text{s}, \quad \Delta t_{\text{int}} \sim O(10^4 - 10^5)/\omega_{\text{int}} \sim 1 - 10\text{s}.$$

- one concludes that both the inter-ELM modes and the dominant low-n modes cannot account for the observed fast ion ‘acceleration’ on $\sim 100\mu\text{s}$.
- Fast-ion is not accelerated by ELMs.

Lagrangian Perspective

Wave-Particle Resonance

- Near a single island, the maximum energy exchange is $\delta E = \frac{2\omega_n H_{m,l}}{\omega_B}$
- The maximum canonical momentum exchange, meanwhile, is $\delta P_{\zeta} = \frac{2n H_{m,l}}{\omega_B}$
- Qualitatively, for both the low-n modes and inter-ELM modes during ELM crash, the canonical momentum exchange dominates the transport.

Lagrangian Perspective

Fluctuations

An arbitrary perturbation, say δQ , can be decomposed into the Lagrangian description by introducing the fast (t_0) and slow (t_1) time scales

$$\delta Q_n = e^{-in[\omega_{\zeta} t_0 + \partial_{P_{\zeta}} \omega_{\zeta} \int_0^{t_1} \delta P_{\zeta} dt' + \partial_E \omega_{\zeta} \int_0^{t_1} \delta E dt'] - in \int_0^{t_1} \delta \zeta dt'} \sum_{m,l} e^{i(m\delta_{\zeta} + l)[\omega_b t_0 + \partial_{P_{\zeta}} \omega_b \int_0^{t_1} \delta P_{\zeta} dt' + \partial_E \omega_b \int_0^{t_1} \delta E dt'] + im \int_0^{t_1} \delta \theta dt'} c_{m,l,s}$$

where

$$c_{m,l}(E, \mu B_0, P_{\zeta}, t_1) = \oint \frac{\omega_b dt_0}{2\pi} e^{-il\omega_b t_0 - in\zeta + im\tilde{\theta}} A_n(\bar{r} + \bar{r} + \underbrace{\int_0^{t_1} \delta \tilde{r} dt'}_{\text{FDOW}}).$$

- Due to the finite drift orbit width (FDOW) effect, the orbit-averaged fields for trapped particles are typically smaller than those for circulating particles, leading to weaker cross-field transport.
- This is consistent with the experimental observation that trapped particles has a weaker FILD signal.

Eulerian Perspective

Quasilinear Theory

In the linear limit, the gyrokinetic equation can solved as

$$\delta G = i\delta\Phi L^{-1} [\omega \frac{\partial F_0}{\partial E} + n \frac{\partial F_0}{\partial P_{\zeta}}]$$

where the propagator L accounts for the wave-particle interaction.

- The effective potential

$$\delta\Phi = J_0(\delta\phi - \delta\psi) + J_0 \frac{\omega_d \delta\psi}{\omega} + \frac{v_{\perp} J_1}{k_{\perp} c} \delta B_{\parallel}$$

is respectively due to three forces: the parallel electric field which vanishes in the ideal MHD limit; $\mathbf{v}_d \times \delta \mathbf{B}_{\perp}$; and the mirror force.

- The explicit quasilinear transport equation

$$\partial_t F_0 = \frac{\partial}{\partial \mathbf{J}} \cdot [D \cdot \frac{\partial F_0}{\partial \mathbf{J}}]$$

- The interaction of fast ions with turbulence is constrained to a two-dimensional $\mathbf{J} = (P_{\zeta}, E)$ manifold embedded in the full 6D phase space.
- The associated transport coefficients are defined as

$$D_{P_{\zeta} P_{\zeta}} = -\text{Im}[\frac{n^2 |\delta\Phi|^2}{L}] \quad D_{EE} = -\text{Im}[\frac{\omega^2 |\delta\Phi|^2}{L}] \quad D_{P_{\zeta} E} = D_{EP_{\zeta}} = -\text{Im}[\frac{n\omega |\delta\Phi|^2}{L}]$$

- In the low-frequency ELM scale, the diffusion will be constrained to a 1D manifold in P_{ζ} .

‘Acceleration’

- For circulating particles, as an example, the radial diffusion coefficient $D_{rr} \propto v^3$ due to $\delta\Phi \propto \omega_d$ and the summation over p.
- In contrast to the microturbulence case [Zhang, PRL, 2009], high energy particles will be transported faster, leading to FILD signals seems like an acceleration process.
- For typical AUG parameters, the required time for cross-field diffusion at edge is on the order of $\Delta t \sim 100\mu\text{s}$, consistent with experimental observations.

‘Spikes’

- The resonance condition for circulating and trapped particles are, respectively,

$$\begin{aligned} \omega + (m - n\bar{q} + l)\omega_b - n\omega_d &= 0, \\ \omega + l\omega_b - n\omega_d &= 0. \end{aligned}$$

- Note that the circulating particle resonance condition depends on poloidal mode number, spikes correspond to multiple phase space islands for circulating particles.

CONCLUSION

- The ELM crash cannot effectively accelerate fast ions, it induces an efficient radial transport of fast ions with $D_{rr} \propto v^3$, yielding a strong FILD signal in high energy tail.
- Multiple spikes in pitch angle are due to multiple phase space islands for circulating particles.
- Theoretical predictions agree with the experimental observations:
 - Circulating particles are more easily transported, so the corresponding FILD signal is stronger.
 - The required time for cross-field diffusion process is estimated as $t \sim 100\mu\text{s}$, consistent with experimental observations.

ACKNOWLEDGEMENTS

This work was supported in part by the ITER-CN under Grant No. 2019YFE03020000, by the National Natural Science Foundation of China under Grant Nos. 12375213 and 12125502.