Impcat of Stark broadening on ion temperature measurement in the ITER divertor plasma

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1. Introduction

- DIM measures impurity ion spectra in the ITER divertor plasma.
- Precise ion temperature $(T_{\rm i})$ measurement is essential to evaluate divertor sputtering.
- Impurity flow direction and magnitude are also key for detachment control.
- Low-Z materials (Be, B) provide useful emission lines for diagnostics.
- Line broadening by Zeeman effect and Stark broadening must be clarified.

2. Method

- The Stark broadening is modeled by ion (static) and electron (impact) contributions.
- Electron broadening follows Griem's model with a Lorentzian line shape.
- The width is proportional to electron density:
- Ion microfields are described by the Holtsmark distribution assuming

$$n_{\rm i} = n_{\rm e}$$
.

- The plasma is isotropic, but the magnetic field introduces anisotropy.
- The total Hamiltonian under arbitrary external fields is

$$H = H_0 - \mu_B B(g_L L_z + g_S S_z) - E d_{\xi}.$$

- The quantization axis is taken along the magnetic field direction.
- The electric-field axis is defined by Euler rotation $R(\theta_1,\theta_2)$:

$$|LJM\rangle_{\xi} = \sum_{m} |LJm\rangle D_{Mm}^{(J)*}(R).$$

$$\langle LJM| - Ed_{\xi} |L'J'M'\rangle =$$

$$-E\sum_{j} D^{(J)}D^{(J')*}\langle LJm |d|L'J'm'\rangle$$

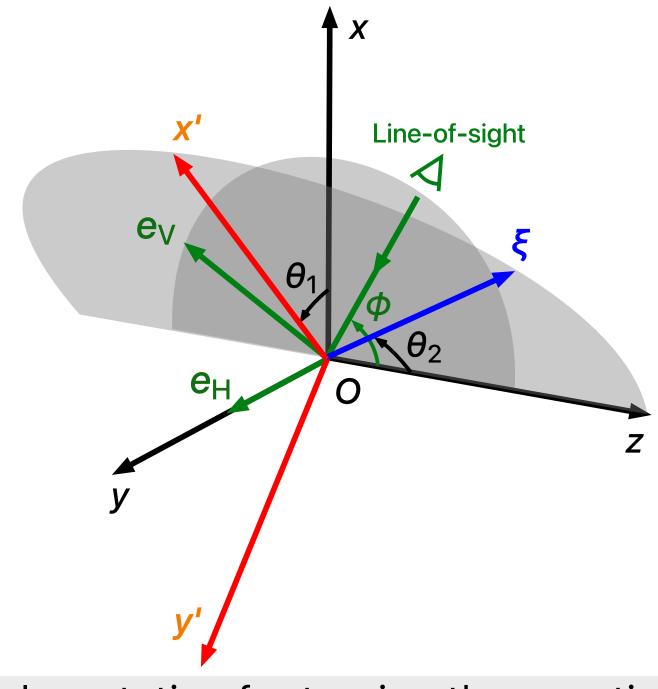
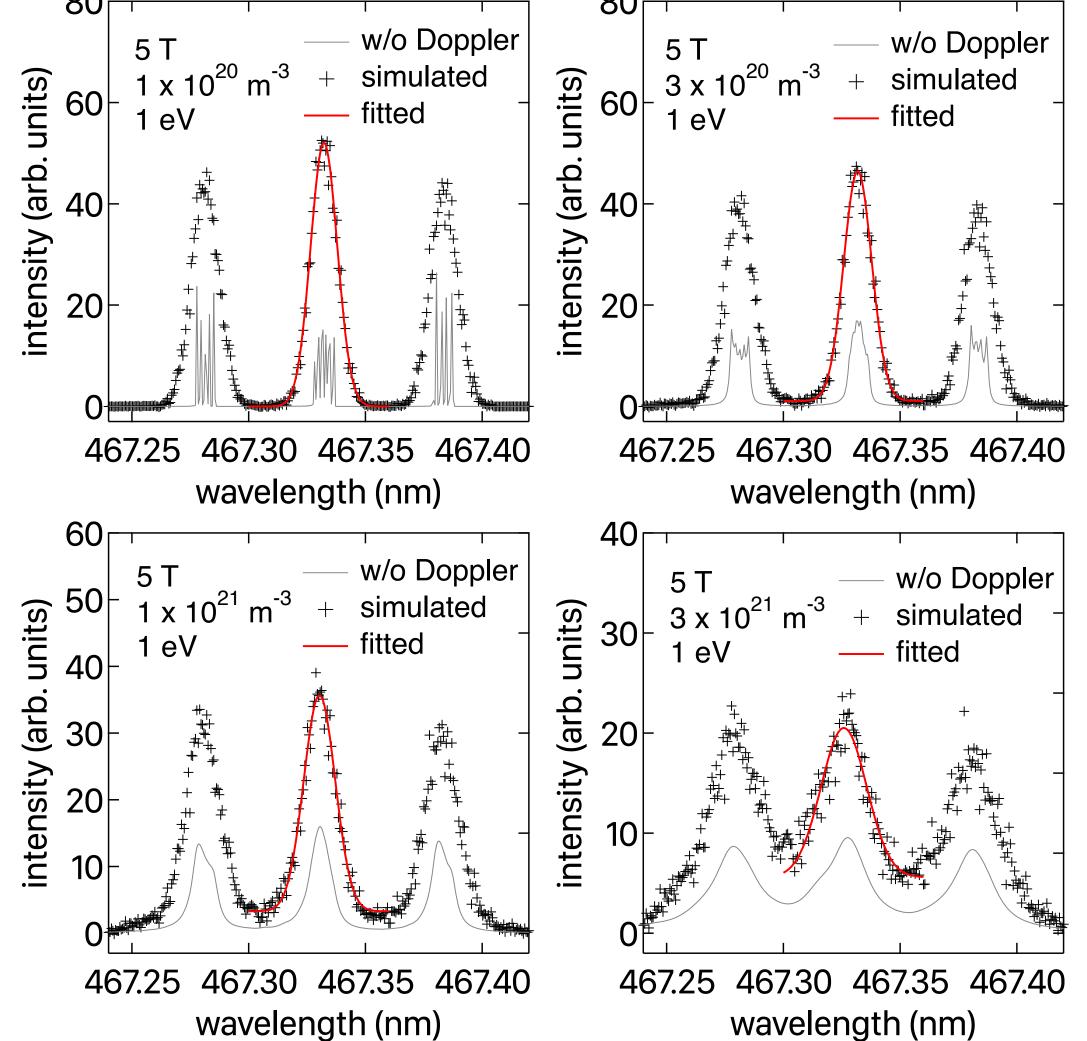
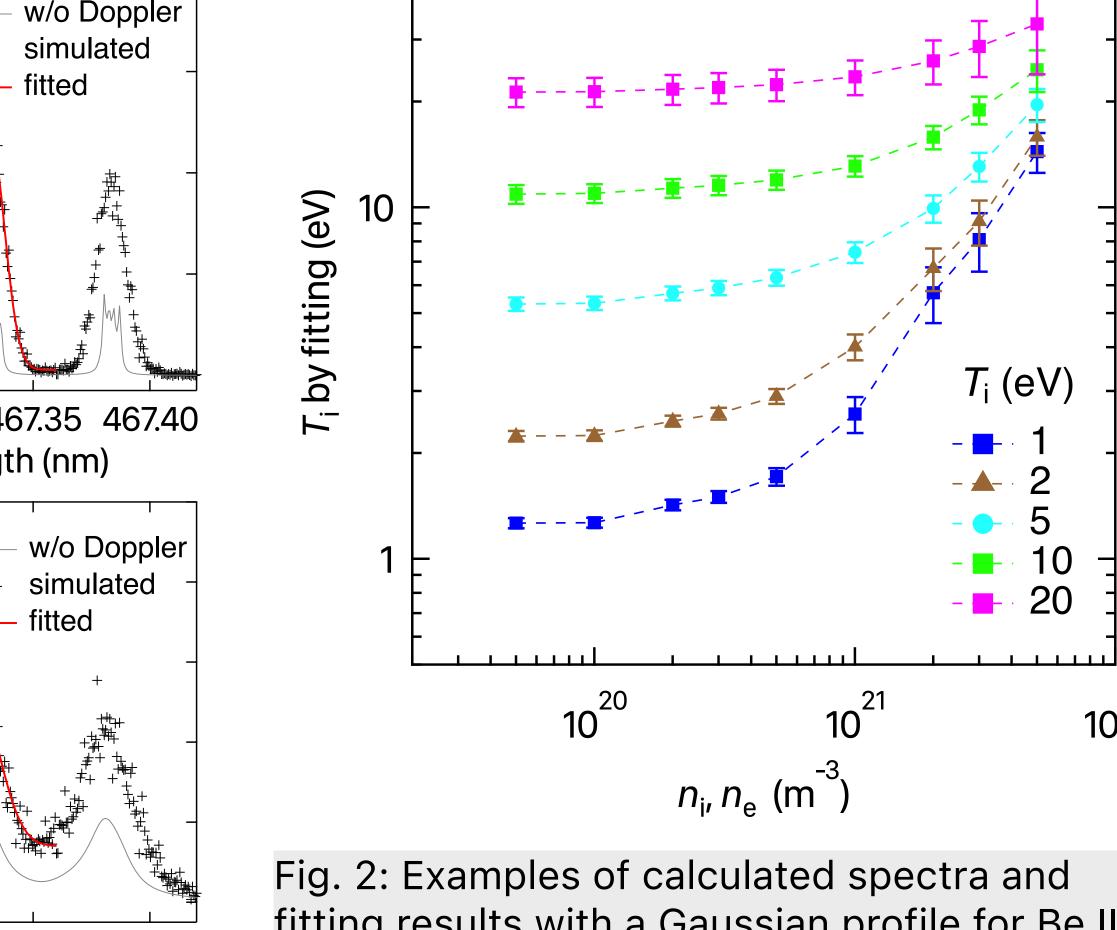


Fig. 1: Euler rotation for turning the quantization axis (z-axis) into the ξ -axis direction by first rotating θ_1 around the z-axis and then θ_2 around the y'-axis which is the new y-axis after the first rotation. The definition of observation vectors \mathbf{e}_V and \mathbf{e}_H are also shown: The former is in the x-z plane and is perpendicular to the line-of-sight, and the latter is on the y-axis.





fitting results with a Gaussian profile for Be II 467.3 nm line.

- Matrix elements include Zeeman and Stark interactions between sublevels.
- Perturbed energies are obtained by diagonalizing the total Hamiltonian.
- Transition intensities are evaluated from $I = |\langle l | \textbf{\textit{e}} \cdot \textbf{\textit{d}} | u \rangle|^2.$
- Observed polarization is decomposed into horizontal ($e_{
 m H}$) and vertical ($e_{
 m V}$) components.
- The final line shape is obtained by integrating over all microfields and convolving with the impact profile.

3. Results: Be II (467 nm)

- Line splits into π and σ components under 5 T magnetic field.
- At $n_{\rm e} > 10^{21}\,{\rm m}^{-3}$, Stark broadening becomes significant.
- Gaussian fitting to central peak overestimates $T_{\rm i}$.

wavelength (nm)

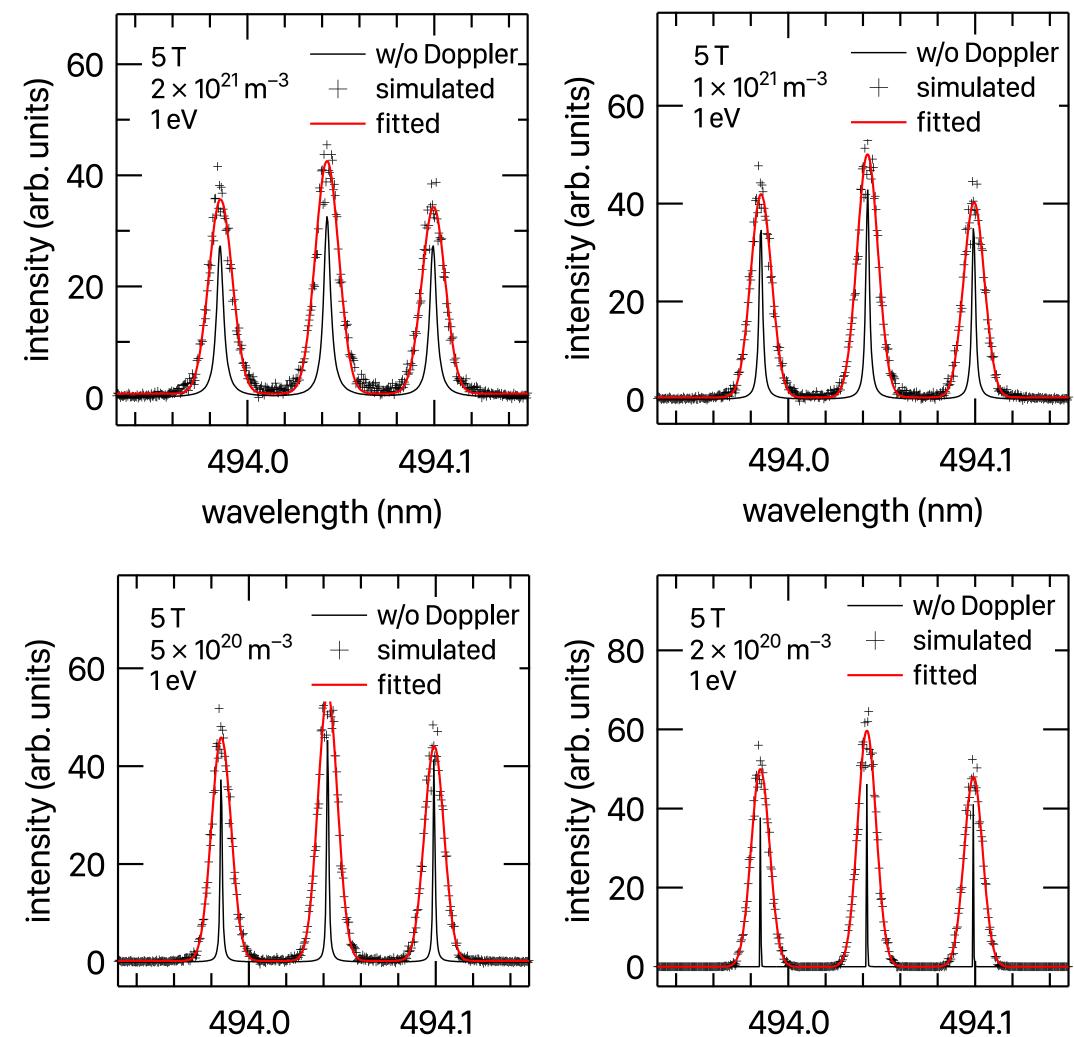
• Even at low density, Zeeman splitting inflates apparent $T_{\rm i}$.

4. Results: B II lines

- 412 nm line: complex triplet structure, difficult for $T_{\rm i}$ diagnostics.
- 494 nm line: Stark broadening visible above $5 \times 10^{20} \, \mathrm{m}^{-3}$.
- 345 nm line: negligible Stark effect, dominated by Doppler width.
- Combining 494 nm and 345 nm lines enables $T_{\rm i} + n_{\rm e}$ diagnostics.

5. Summary

- Stark broadening affects Be II (467 nm) at high density.
- B II (494 nm, 345 nm) offer complementary diagnostic capability.
- Simultaneous $T_{\rm i}$ and $n_{\rm e}$ measurements possible for ITER divertor plasma.



wavelength (nm)

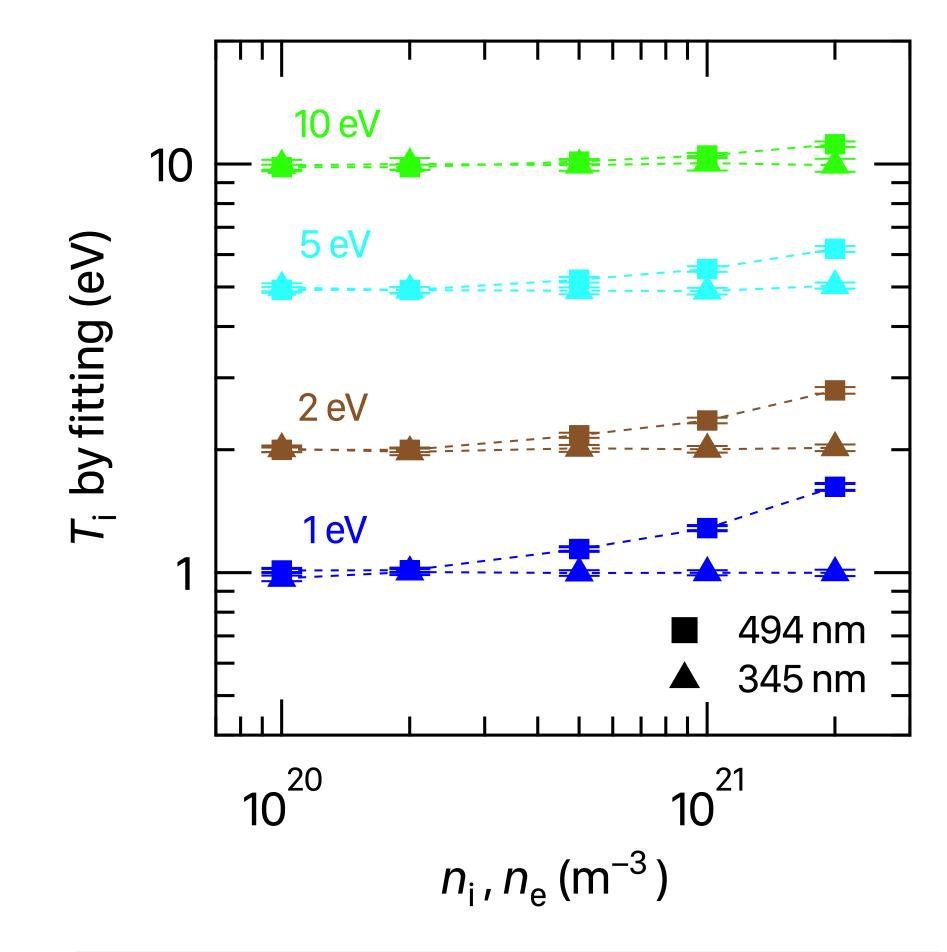


Fig. 3: Examples of calculated spectra and fitting results with Gaussian profiles for B II 494.0 nm.