

# Thermal Quench Dynamics and Heat Flux Distribution during Massive-Impurity-Injection Triggered Disruption in EAST



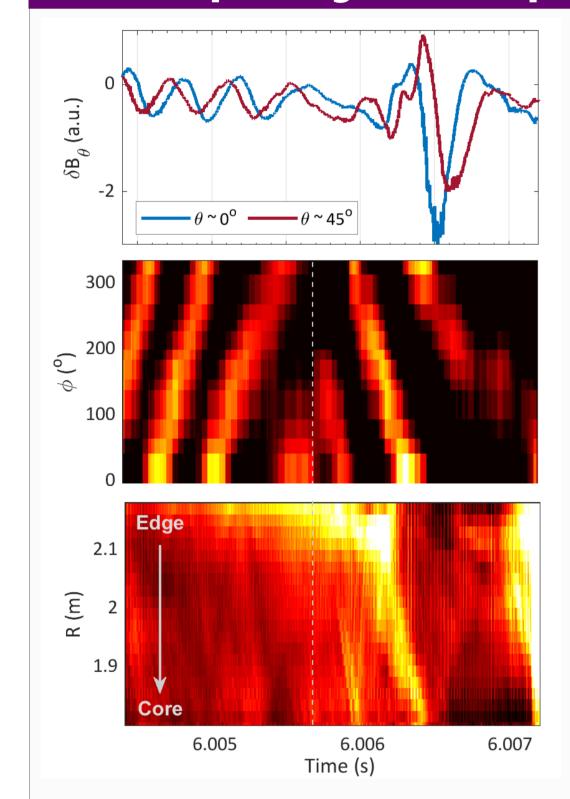
L. Zeng<sup>1\*</sup>, W. Xia<sup>2,3</sup>, D. Hu<sup>4</sup>, Y. Sun<sup>2</sup>, T. Tang<sup>2</sup>, X. Zhu<sup>5</sup>, S. Zhao<sup>2,3</sup>, T. Shi<sup>2</sup>, D. Chen<sup>2</sup>, Y. Duan<sup>2</sup>, L. Xu<sup>2</sup>, G. Li<sup>2</sup>, G. Zuo<sup>2</sup>, J. Hu<sup>2</sup>, X. Gao<sup>2</sup>, Z. Gao<sup>1</sup>, the JOREK Team<sup>6</sup>

<sup>1</sup>Department of Engineering Physics, Tsinghua University, Beijing 100084, China <sup>2</sup>Institute of Plasma Physics, Chinese Academy of Sciences, Hefei 230031, China <sup>3</sup>University of Science and Technology of China, Hefei 230031, China <sup>4</sup>School of Physics, Beihang University, Beijing 100191, China <sup>5</sup>Advanced Energy Research Center, Shenzhen University, Shenzhen 518060, China <sup>6</sup>Hoelzl et al 2024 (https://doi.org/10.1088/1741-4326/ad5a21) for the JOREK Team \*E-mail: zenglong@tsinghua.edu.cn

## Summary

- Thermal quench dynamics in massive-impurity-injection triggered EAST disruptions have been demonstrated in experiments and simulations.
- Sudden change in mode frequency: A sudden change of the radial electric field (Er), caused by collisionality transition from the banana to the P-S regime, can drive the poloidal rotation and enhance impurity influx.
- Double-stage TQ: The 3/1 mode couples with the 4/1 and 5/1 modes, contributing to edge stochasticity. The non-linear interaction between the 3/1 and 2/1 modes primarily leads to global stochastic.
- **3D** splitting of heat flux: The parallel convective particle flux at the upper-outer target, combined with the strike-point splitting, exhibits an n = 1 structure.

# Impurity Transport during the pre-TQ Phase

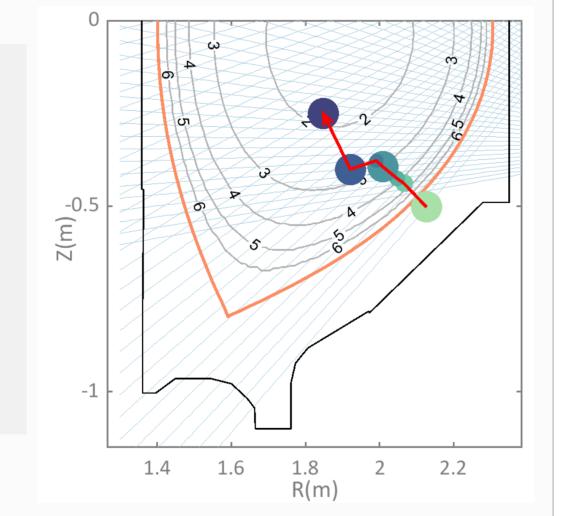


- The intrinsic 2/1 mode's rotation slows down and then reverses from the ion diamagnetic to the electron diamagnetic direction after neon injection into cocurrent NBI-heating plasmas.
  - The mode frequency changes from 2
    kHz to −3 kHz in < 0.5 ms.</li>
  - The radiation in the core region significantly increases following the rotation direction of the mode reversed.

# Observed in both SPI and MGI triggered disruptions

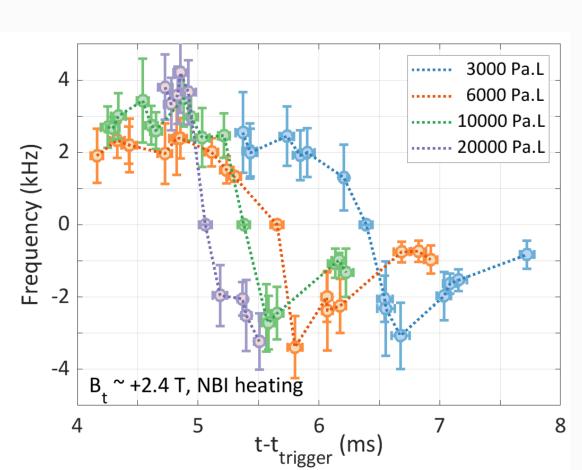
### Two stages for impurity transport

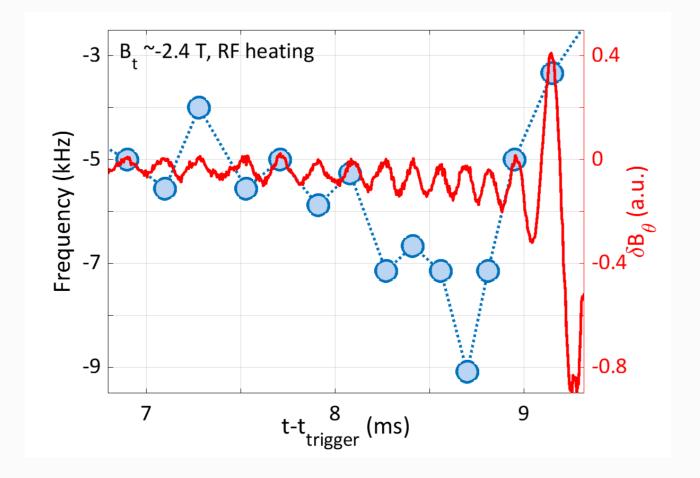
- STAGE 1: impurities drift along ion diamagnetic direction; Outside of q=2 surface. Inward pinch speed ~ 60 m/s.
- STAGE 2: impurities drift along electron diamagnetic direction; Across q=2 surface. Inward pinch speed ~ 300 m/s.



#### Characteristics of mode-frequency change

- The phenomenon has been observed by scanning species and numbers of impurities, and heating methods.
  - Amount scan: Earlier and faster reversion of rotation direction.
  - The mode frequency changes from -5 kHz to -9 kHz in ~ 0.5 ms after argon injection into RF-heating plasmas.





- o Change of mode frequency is mainly caused by poloidal angular frequency.
  - Time scale of frequency change (< 0.5 ms) is much less than  $\tau_{\rm E}$  (50-100 ms) on EAST.

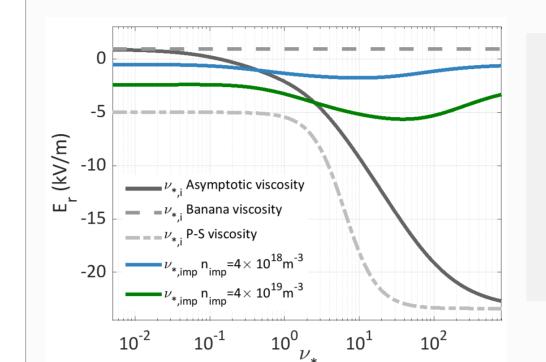
Poloidal rotation velocity

$$\delta U_{\theta} = -\frac{B_T}{B^2} \delta E_r$$

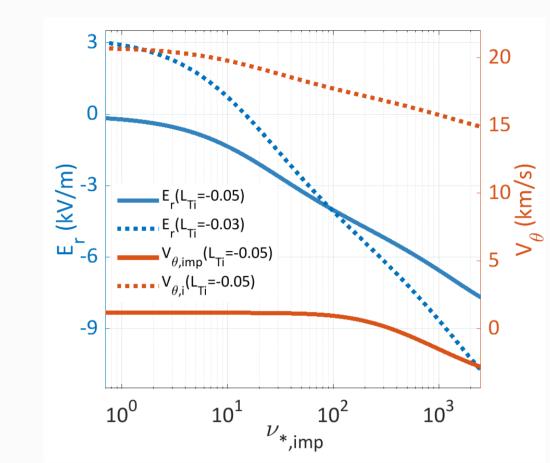
# Effect of Impurity on Poloidal Rotation

- The neoclassical theory is applied due to the short time scale.
  - The poloidal rotation can be driven by the viscosity associated with the collisionality and impurity.
  - The parallel momentum and heat flow balance equations for the impurity and ion species in the form of matrix equations are derived.

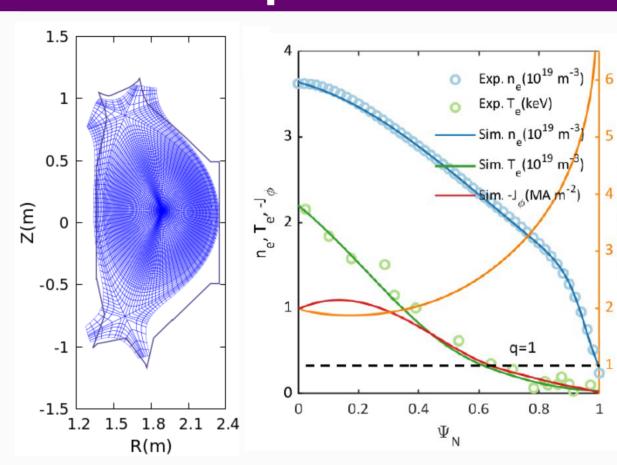
# A clear change in radial electric field is observed as collisionality transitions from the banana to the P-S regime.



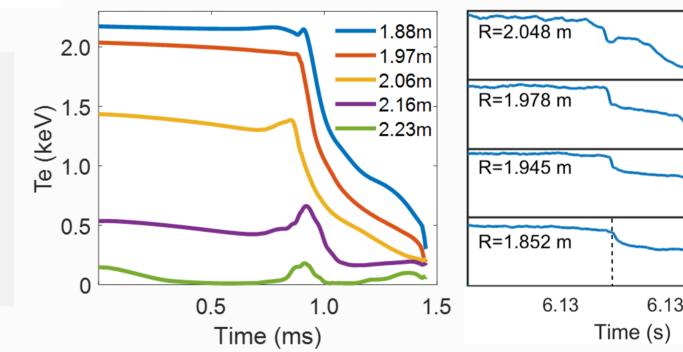
- As ion collisionality increases, the radial electric field decrease nonlinearly.
- The viscosity is sensitive to impurity strength: impurity orbits perturb ion orbits and thereby alter the effective viscosity.
- Global effects with massive impurity injection show variations in ion poloidal flow are governed by the E×B drift.
  - − E<sub>r</sub> is estimated in the range from −12 kV/m to −8 kV/m according to the typical parameters, consistent with the experimental value of ~ −10 kV/m.

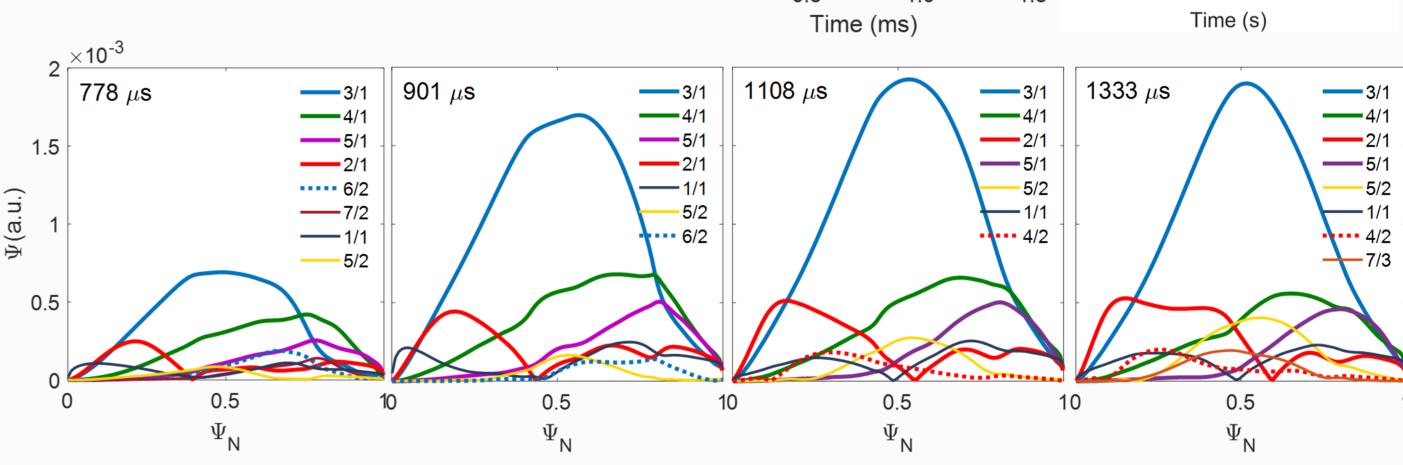


## Interpretive Simulation of TQ with JOREK



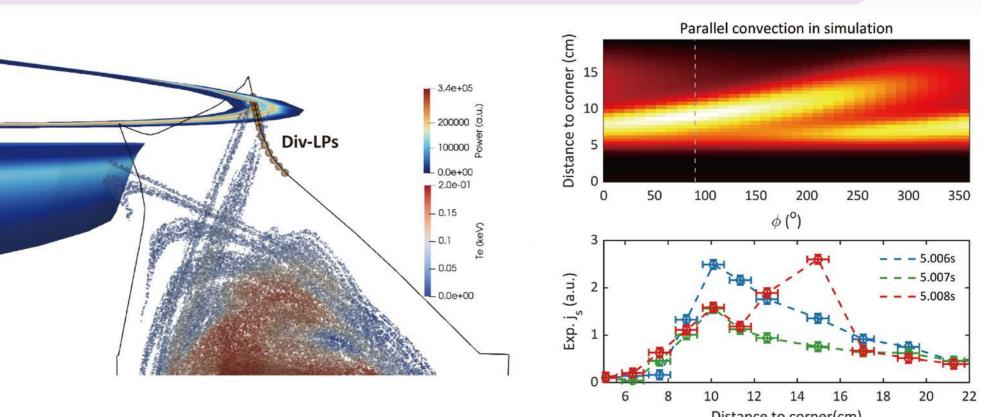
- Simulation setup: A typical L-mode plasma with  $I_P \sim 0.4$  MA and  $\mathbf{q}_{min} \sim 1.6$  in the core region.
  - The initial core  $T_e$  is ~2.2 keV, electron and ion temperatures are assumed to be equal, and the core  $n_e$  is ~3.6×10<sup>19</sup> m<sup>-3</sup>.
- o The process of a doublestage TQ has been represented with neon MGI, consistent with the experimental observation.





- > The 3/1 mode is dominant in the whole collapse process.
  - First collapse is from the outer region (q>2) and the coupling between 3/1 and 4/1 is the main reason.
  - **Second collapse**: The coupling among 2/1, 3/1 and higher harmonic mode 4/2 is dominant in the final collapse.

#### Strike point splitting on the divertor



Saturated ion flux measurement by Langmuir probes consistent with simulations

- The simulated parallel convective particle flux on the divertor, combined with the strike-point splitting, exhibits an n = 1 structure.
  - This result provides an explanation for the broadening of energy deposition width during the TQ in ASDEX Upgrade and JET.