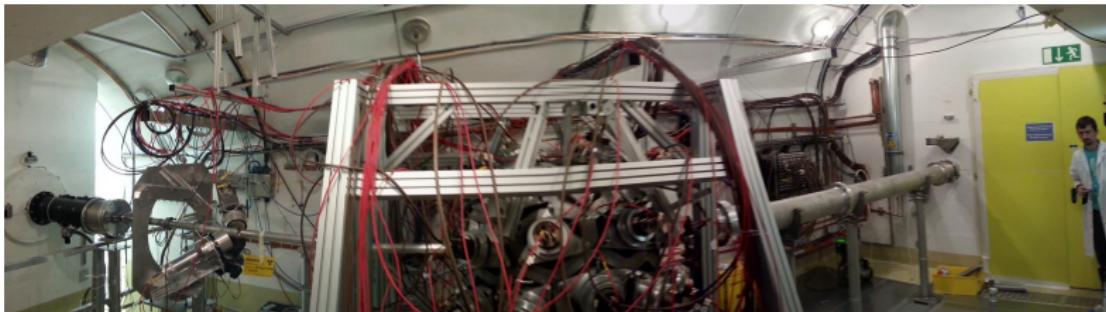


The dipole photon strength of uranium isotopes from TAC measurements at n_TOF

Standa Valenta
for n_TOF collaboration

July 10, 2024



The dipole photon strength of uranium isotopes from TAC measurements at n_TOF

PHYSICAL REVIEW C 105, 024618 (2022)

Constraints on the dipole photon strength for the odd uranium isotopes

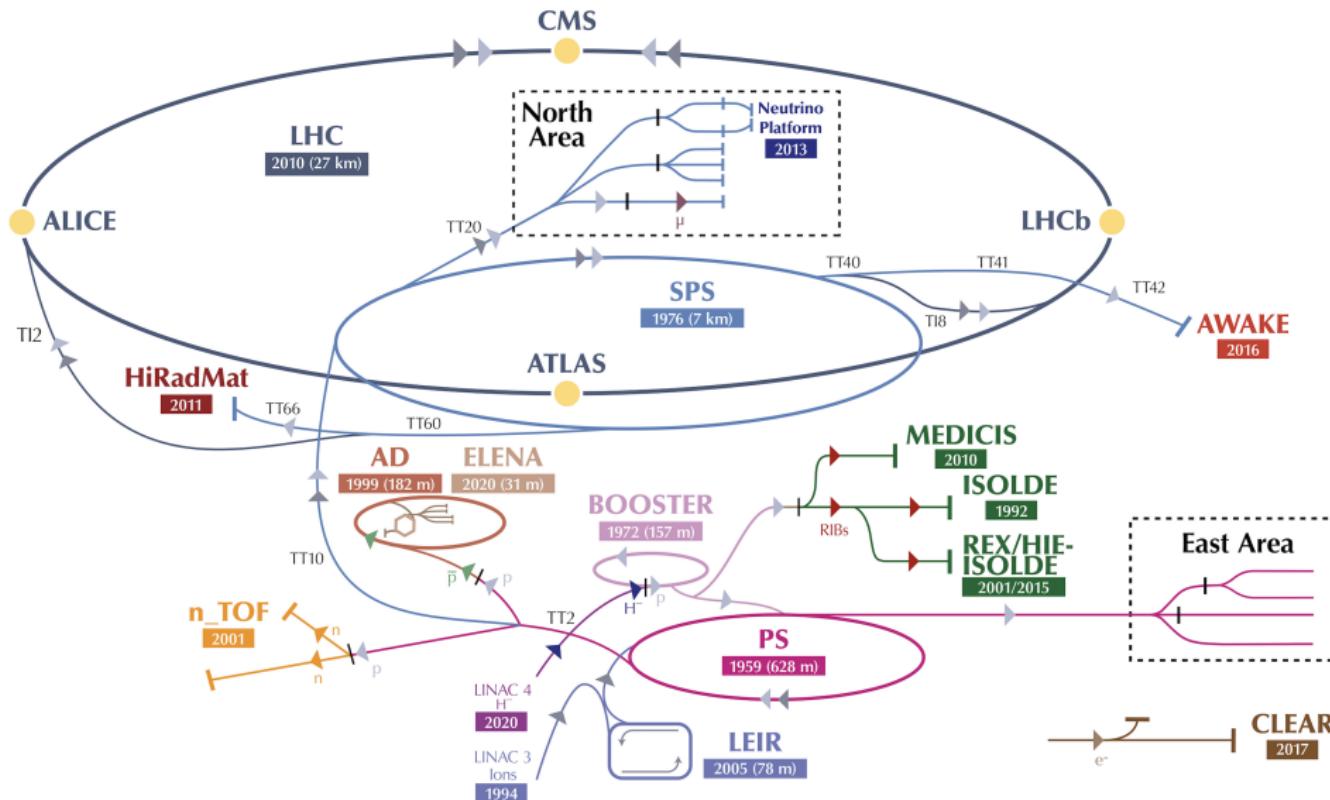
- J. Moreno-Soto,^{1,*} S. Valenta,^{2,†} E. Berthoumieux,¹ A. Chebboubi,³ M. Diakaki,³ W. Dridi,⁴ E. Dupont,¹ F. Gunsing,^{1,‡} M. Krtička,² O. Litaize,³ O. Serot,³ O. Aberle,⁵ V. Alcayne,⁶ S. Amaducci,⁷ J. Andrzejewski,⁸ L. Audouin,⁹ V. Bécares,⁶ V. Babiano-Suarez,¹⁰ M. Bacak,^{5,11,1} M. Barbagallo,^{5,12} Th. Benedikt,¹³ S. Bennett,¹⁴ J. Billowes,¹⁴ D. Bosnar,¹⁵ A. Brown,¹⁶ M. Busso,^{17,18} M. Caamaño,¹⁹ L. Caballero-Ontanaya,¹⁰ F. Calviño,²⁰ M. Calviani,⁵ D. Cano-Ott,⁶ A. Casanovas,²⁰ F. Cerutti,⁵ E. Chiaveri,^{5,14} N. Colonna,¹² G. Cortés,²⁰ M. A. Cortés-Giraldo,²¹ L. Cosentino,⁷ S. Cristallo,^{17,22} L. A. Damone,^{12,23} P. J. Davies,¹⁴ M. Dietz,²⁴ C. Domingo-Pardo,¹⁰ R. Dressler,²⁵ Q. Ducasse,²⁶ I. Durán,¹⁹ Z. Eleme,²⁷ B. Fernández-Domínguez,¹⁹ A. Ferrari,⁵ P. Finocchiaro,⁷ V. Furman,²⁸ K. Göbel,¹³ A. Gawlik-Ramiega,⁸ S. Gilardoni,⁵ I. F. Gonçalves,²⁹ E. González-Romero,⁶ C. Guerrero,²¹ S. Heinitz,²⁵ J. Heyse,³⁰ D. G. Jenkins,¹⁶ A. Junghans,³¹ F. Käppeler,^{32,§} Y. Kadi,⁵ A. Kimura,³³ I. Knapová,² M. Kokkoris,³⁴ Y. Kopatch,²⁸ D. Kurtulgil,¹³ I. Ladarescu,¹⁰ C. Lampoudis,³⁵ C. Lederer-Woods,²⁴ S. J. Lonsdale,²⁴ D. Macina,⁵ A. Manna,^{36,37} T. Martínez,⁶ A. Masi,⁵ C. Massimi,^{36,37} P. Mastinu,³⁸ M. Mastromarco,⁵ E. A. Maugeri,²⁵ A. Mazzone,^{12,39} E. Mendoza,⁶ A. Mengoni,⁴⁰ V. Michalopoulou,^{5,34} P. M. Milazzo,⁴¹ F. Mingrone,⁵ A. Musumarra,^{7,42} A. Negret,⁴³ R. Nolte,²⁶ F. Ogállar,⁴⁴ A. Oprea,⁴³ N. Patronis,²⁷ A. Pavlik,⁴⁵ J. Perkowski,⁸ L. Piersanti,^{17,22} C. Petrone,⁴³ E. Pirovano,²⁶ I. Porras,⁴⁴ J. Praena,⁴⁴ J. M. Quesada,²¹ D. Ramos-Doval,⁹ T. Rauscher,^{46,47} R. Reifarth,¹³ D. Rochman,²⁵ M. Sabaté-Gilarte,^{21,5} A. Saxena,⁴⁸ P. Schillebeeckx,³⁰ D. Schumann,²⁵ A. Sekhar,¹⁴ A. G. Smith,¹⁴ N. V. Sosnin,¹⁴ P. Sprung,²⁵ A. Stamatopoulos,³⁴ G. Tagliente,¹² J. L. Tain,¹⁰ A. Tarifeño-Saldivia,²⁰ L. Tassan-Got,^{5,34,9} P. Torres-Sánchez,⁴⁴ A. Tsinganis,⁵ J. Ulrich,²⁵ S. Urllass,^{31,5} G. Vannini,^{36,37}



neutron_Time Of Flight facility



neutron_Time Of Flight facility



What is neutron time of flight?



neutron time of flight



Prihlásiť sa

Všetko

Obrázky

Videá

Nákupy

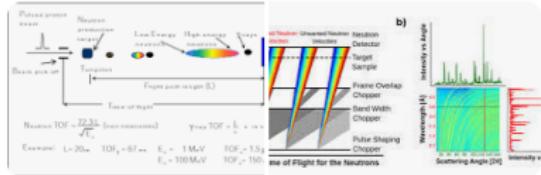
Správy

Viac

Nástroje

Približne 6 720 000 výsledkov (0,46 sekundy)

The Neutron Time Of Flight (n-TOF) facility is a neutron spectrometer at CERN. It consists of a pulsed source, a flight path of 200 m length, and a detector systems. Neutron energies are deduced from the time of flight between source and detector; hence the name of the facility.



wikipedia.org

https://en.wikipedia.org/wiki/Neutron_Time_of_Flight

[Neutron Time Of Flight - Wikipedia](#)

Vybrané úryvky • Spätná väzba

Podobné dopady

How fast is a neutron?



What is the relation between the neutron wavelength and the time of flight over a particular distance?



What is the temperature of thermal neutron?



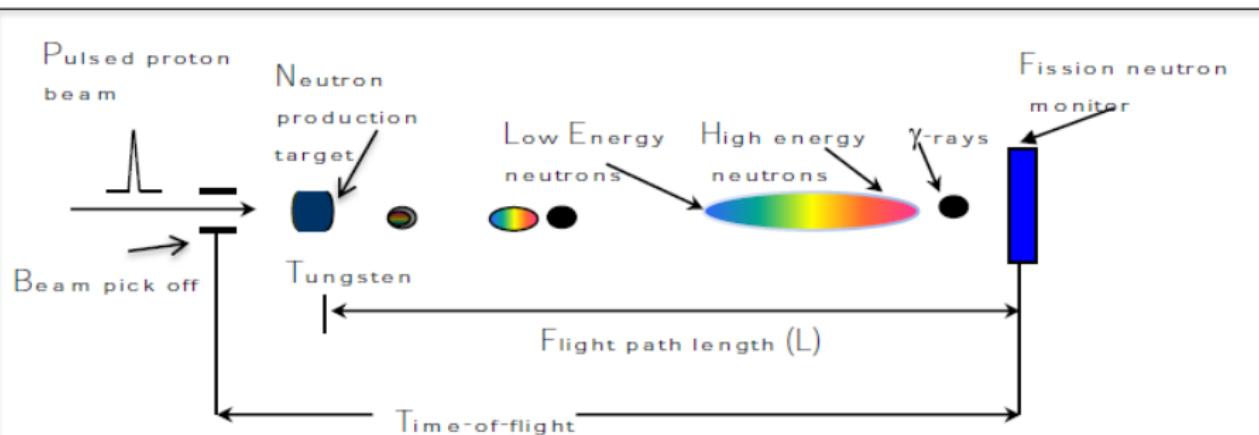
Spätná väzba

Neutron Time Of Flight

Preložené z angličtiny - Zariadenie Neutron Time Of Flight je neutrónový spektrometer v CERN-e. Pozostáva z impulzného zdroja, dráhy letu s dĺžkou 200 m a systému detektorov. Energie neutrónov sú odvodené z času letu medzi zdrojom a detektorm; odtiaľ názov zariadenia. [Wikipedia \(angličtina\)](#)



What is neutron time of flight?



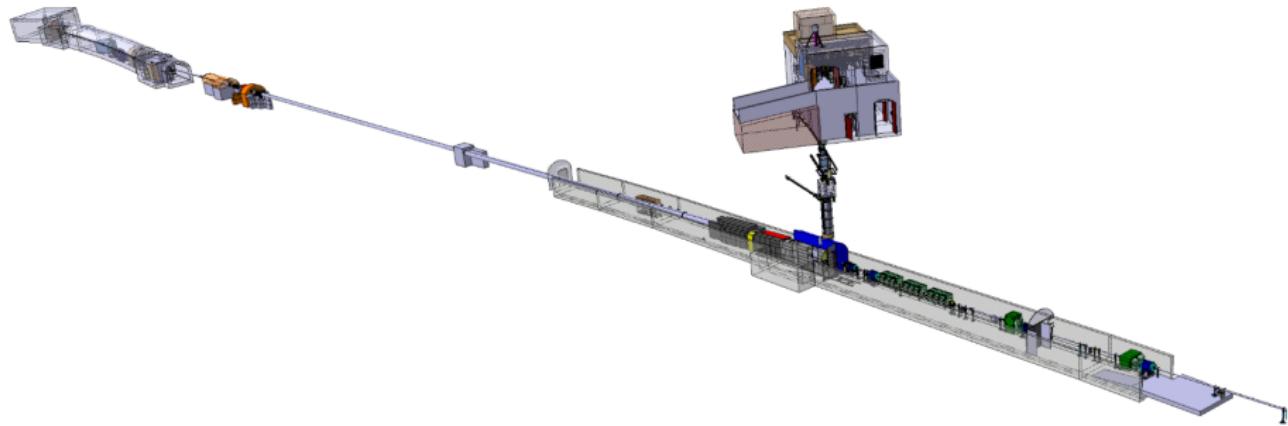
$$\text{Neutron TOF} = \frac{72.3L}{\sqrt{E_n}} \quad (\text{non-relativistic})$$

$$\gamma\text{-ray TOF} = \frac{L}{c} \quad c \text{ is velocity of light}$$

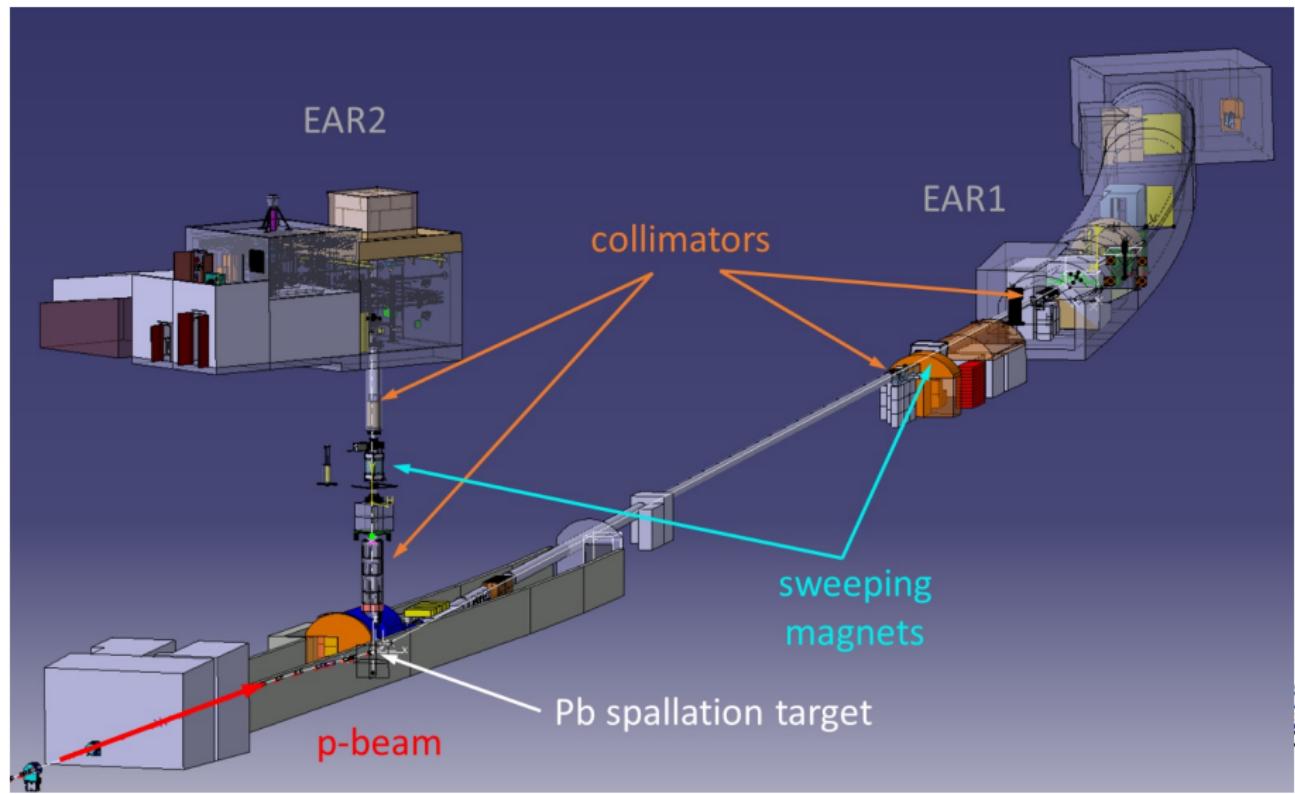
Example:	$L = 20 \text{ m}$	$\text{TOF}_\gamma = 67 \text{ ns}$	$E_n = 1 \text{ MeV}$	$\text{TOF}_n = 1.5 \mu\text{s}$
			$E_n = 100 \text{ MeV}$	$\text{TOF}_n = 150 \text{ ns}$

from <https://lansce.lanl.gov/facilities/time-of-flight.php>

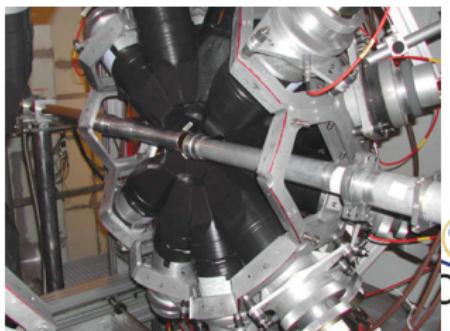
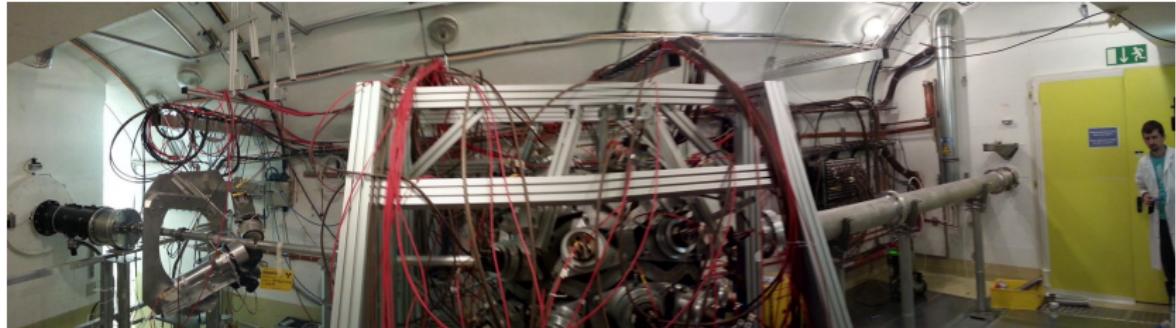
neutron_Time Of Flight facility



neutron_Time Of Flight facility

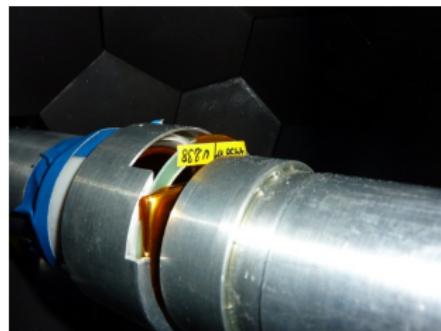
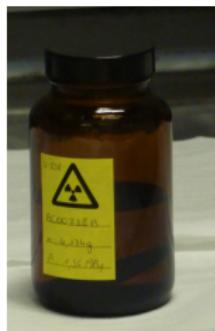


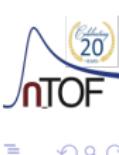
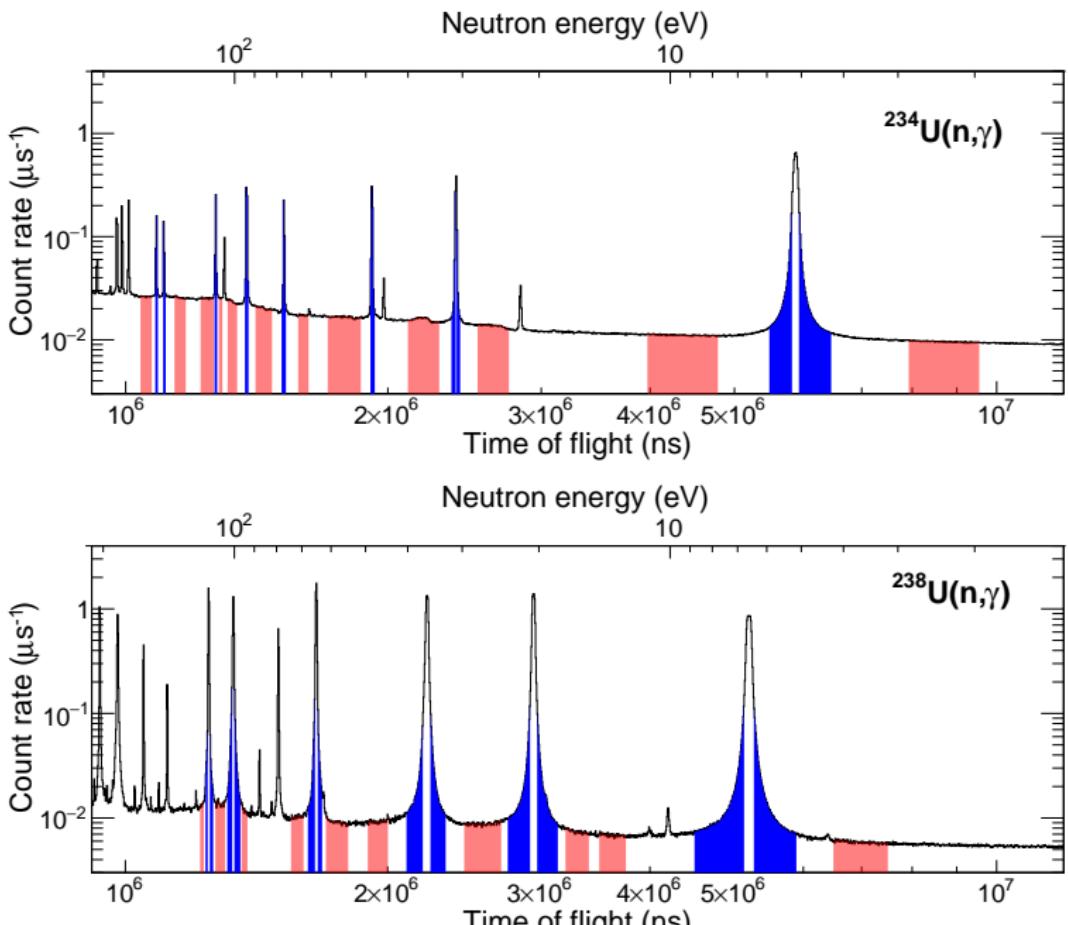
Experimental ARea 1



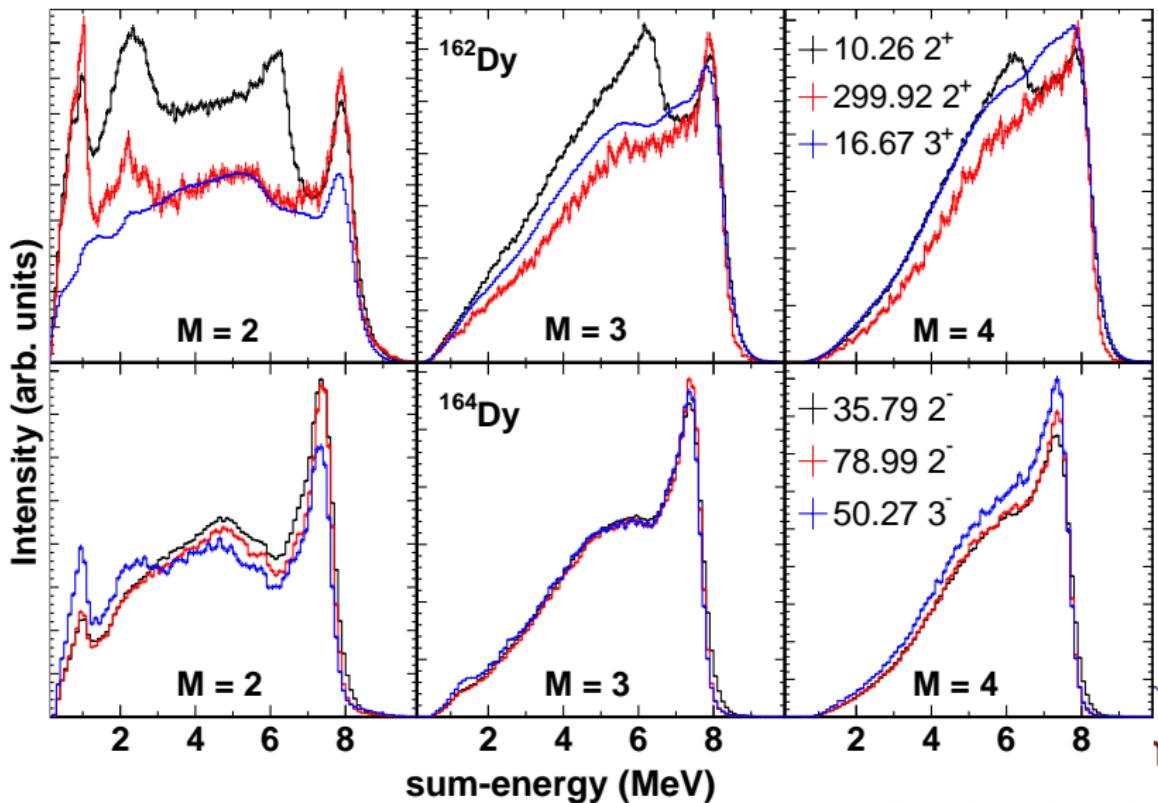
Experimental conditions & samples

	^{234}U	^{236}U	^{238}U
Mass (mg)	32.7	338	6125
Areal density (10^{-4} atoms/b)	1.07	10.9	9.56
Canning	Ti	Al	Al
$S_n^{\text{compound}} \approx Q$ (MeV)	5.297	5.126	4.806
Resolution @ 0.9 MeV (%)	14.5	16.5	16.5
Absorber	^6Li -salt	B-PE	B-PE



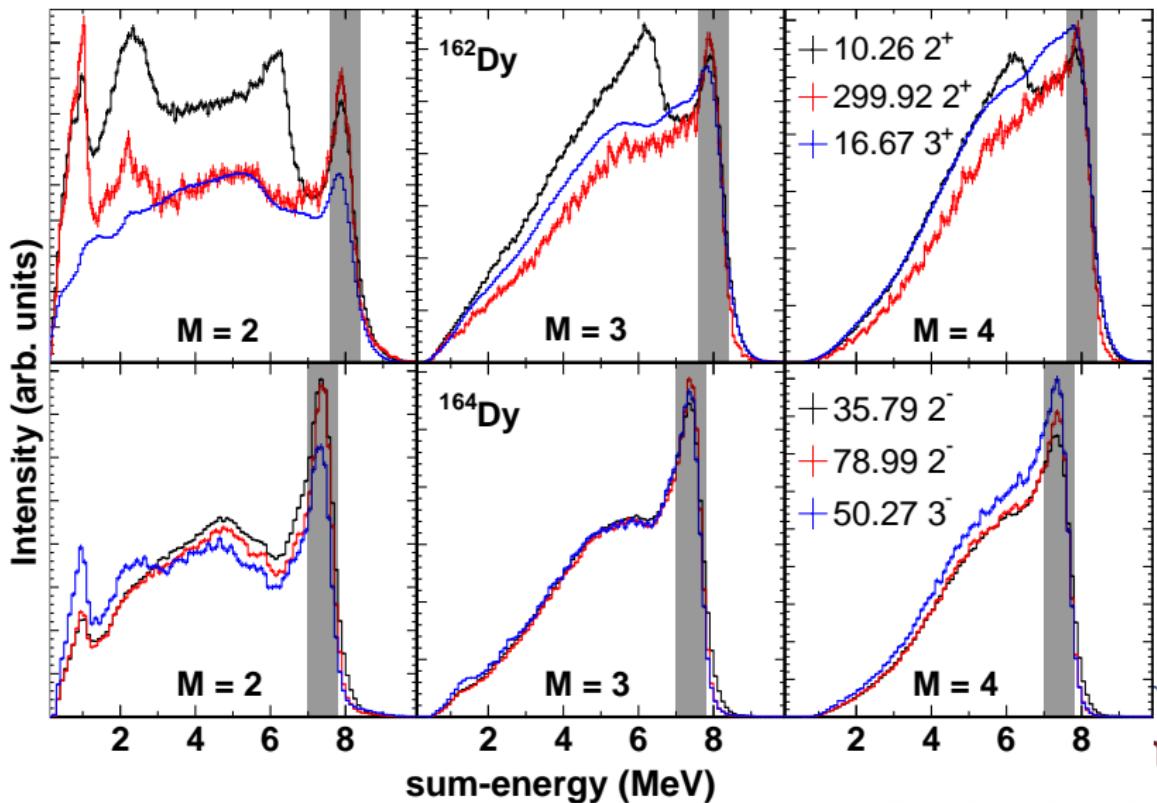


Spectra of energy sums for strong isolated resonances



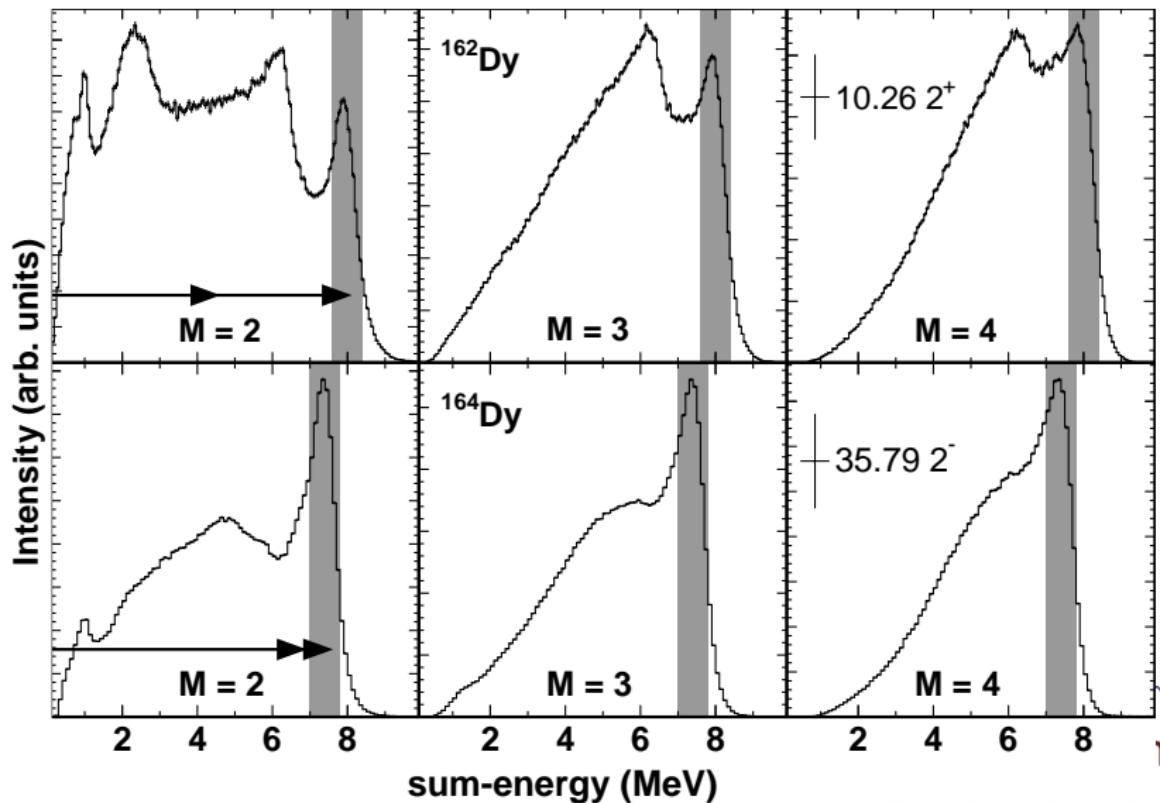
TOF

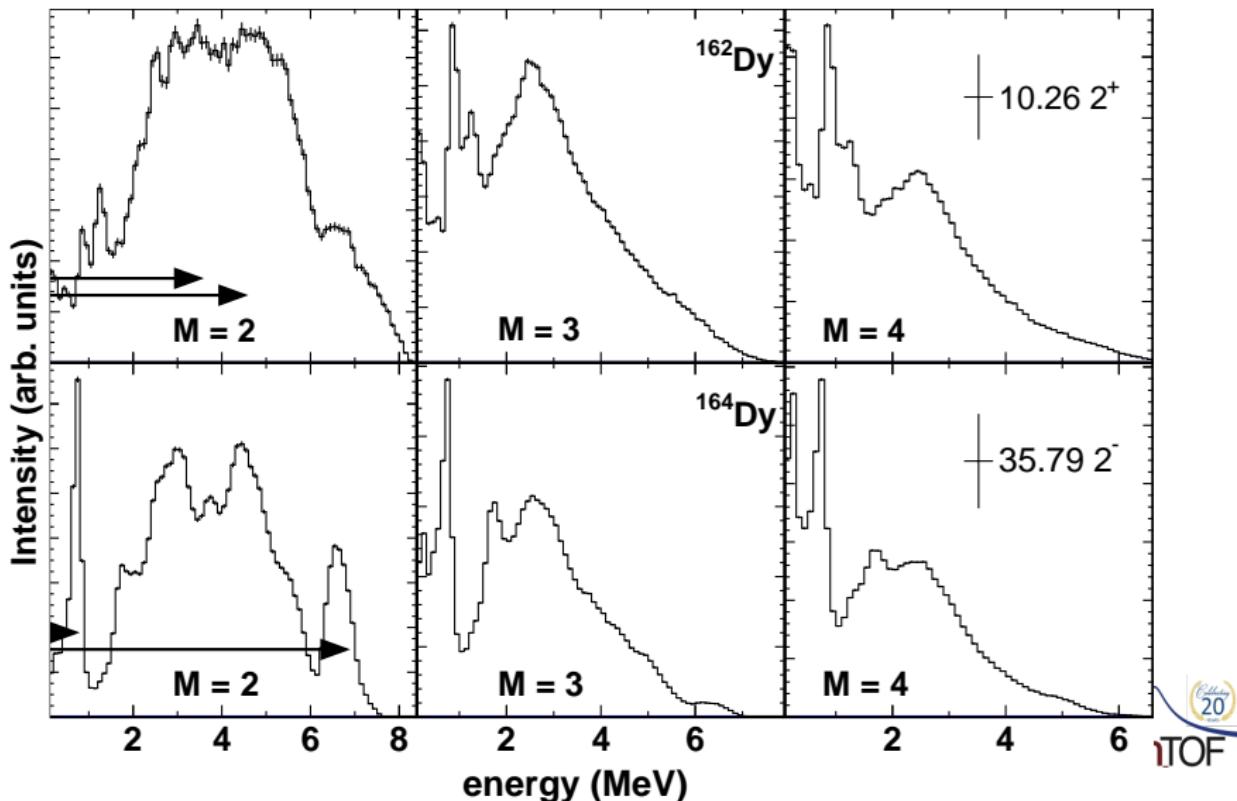
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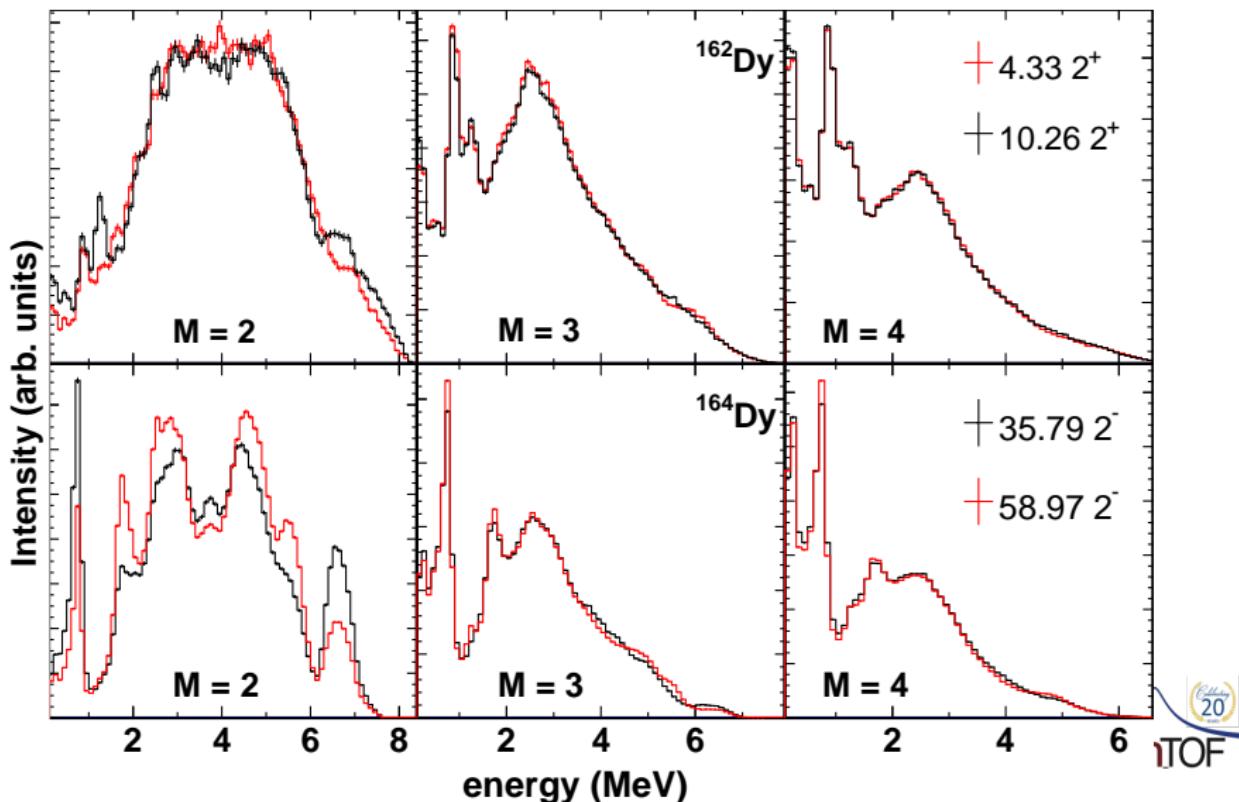
TOF

The energies forming the cascades ev-by-ev

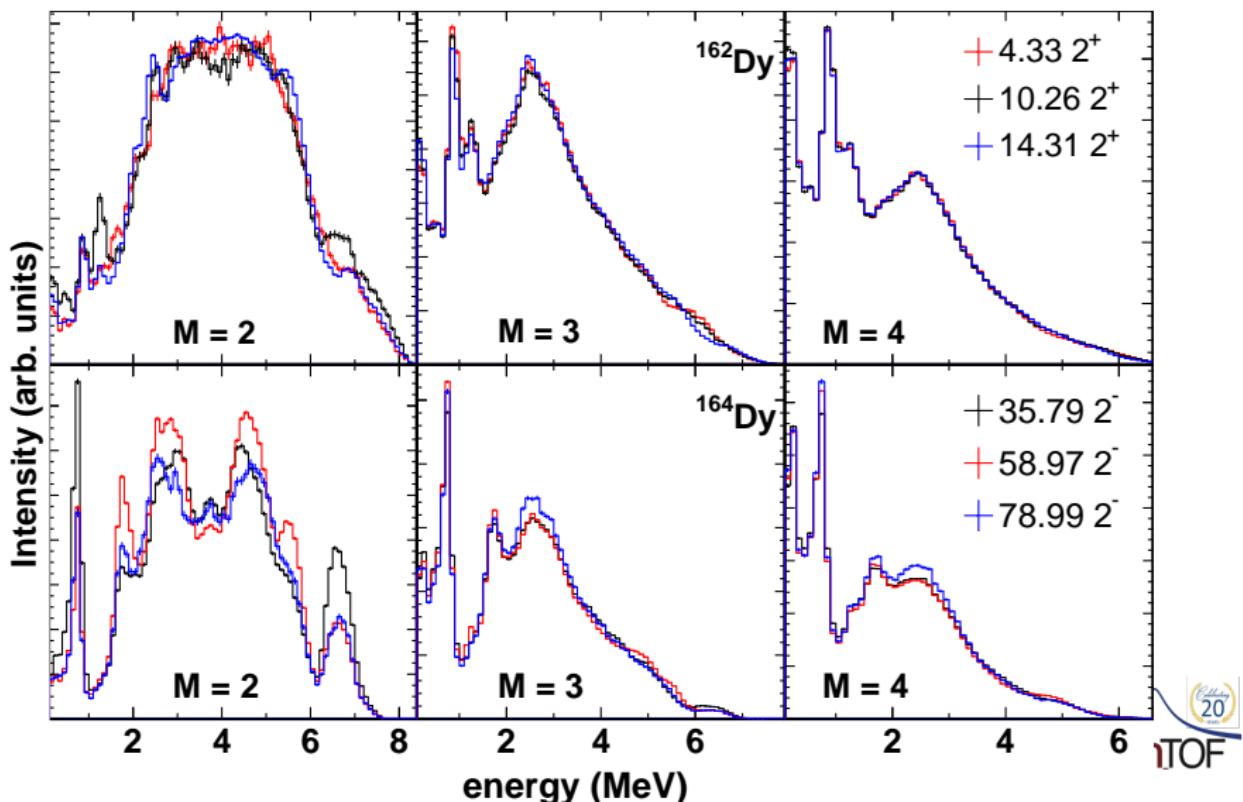


Experimental multi-step γ cascades (MSC) spectra

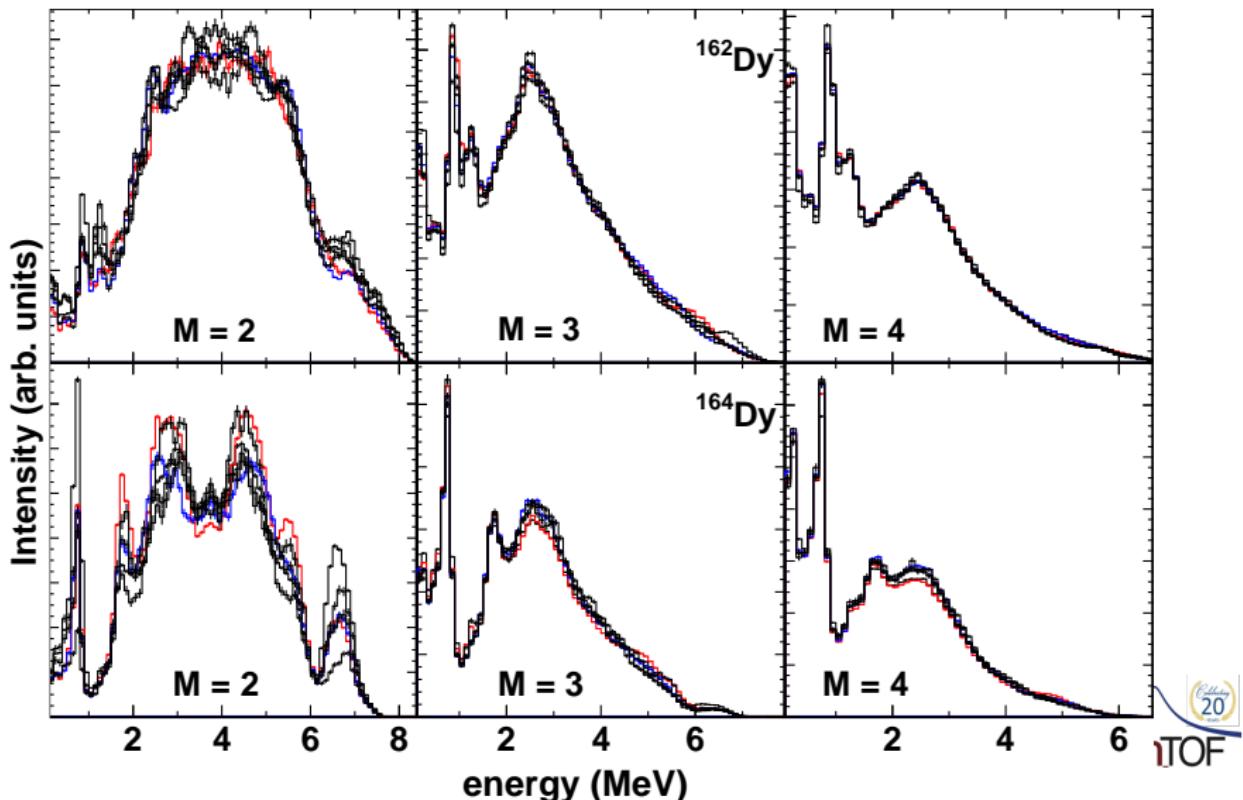
Experimental MSC spectra



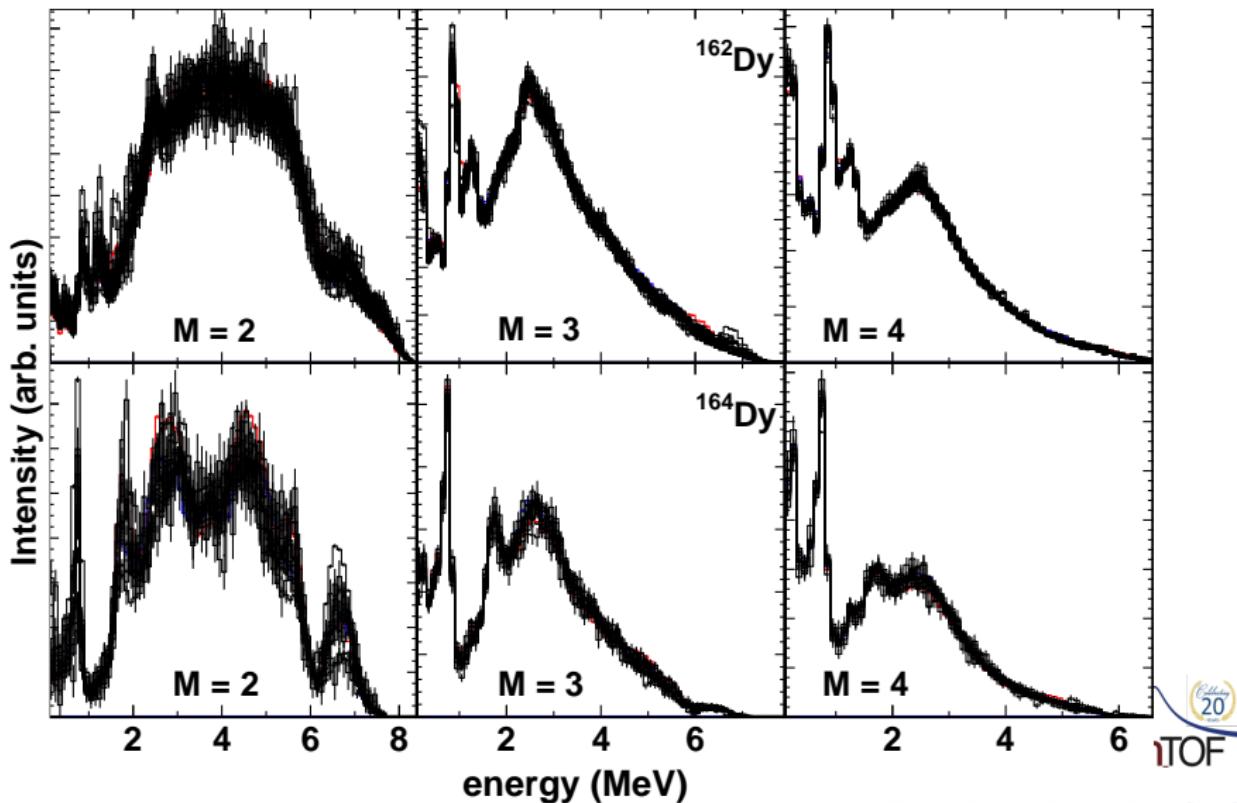
Experimental MSC spectra



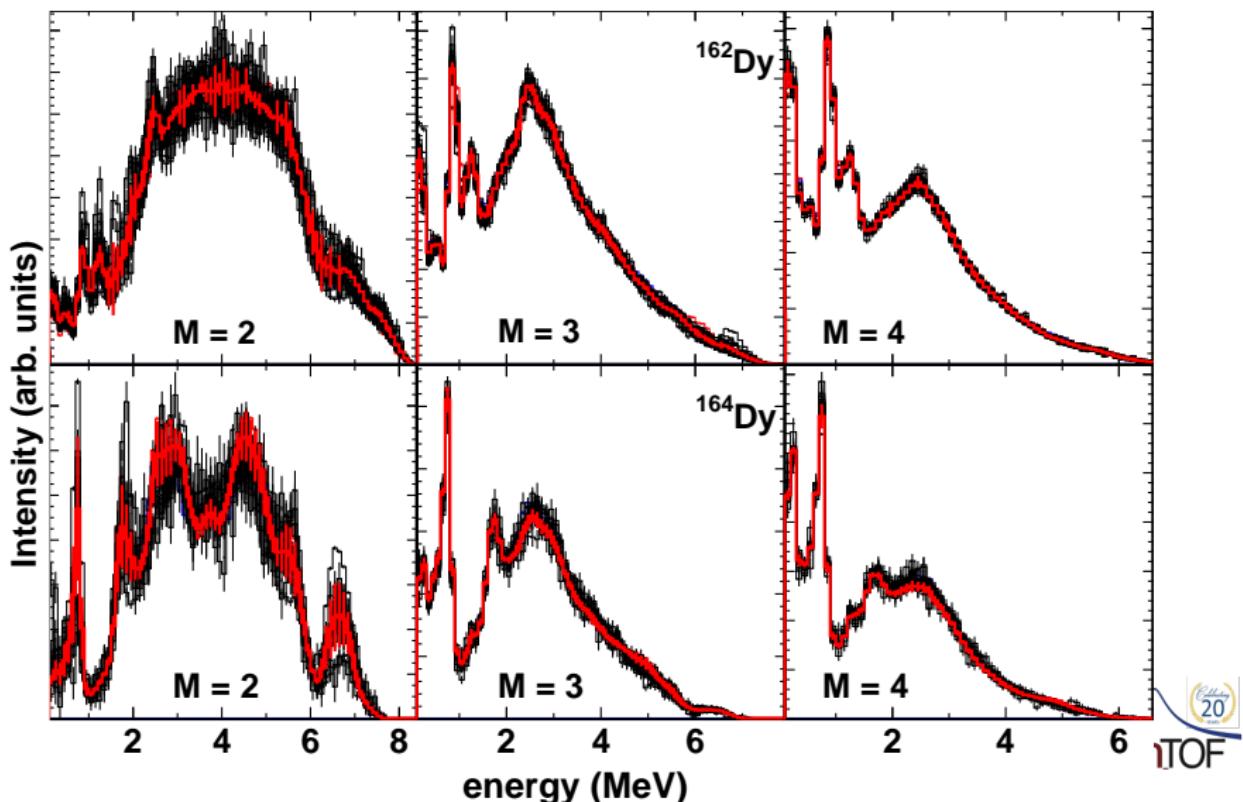
Experimental MSC spectra



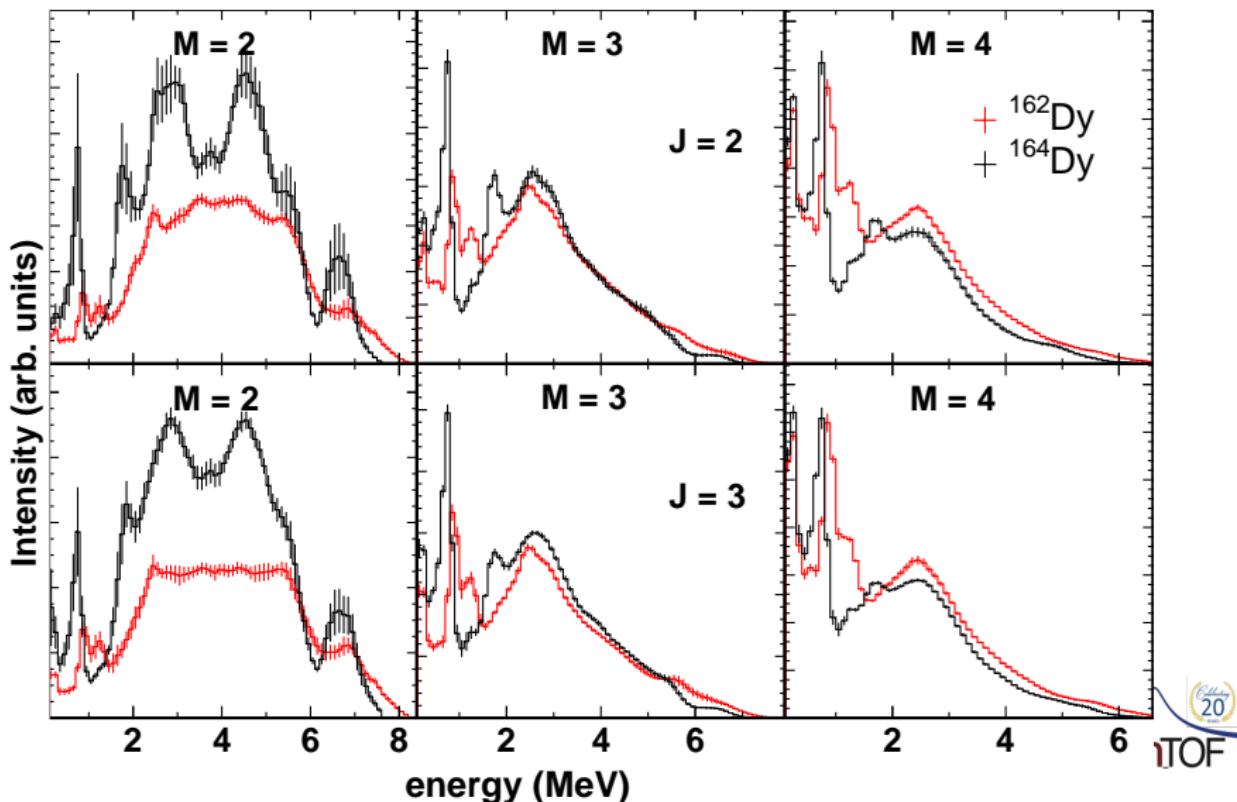
Experimental MSC spectra

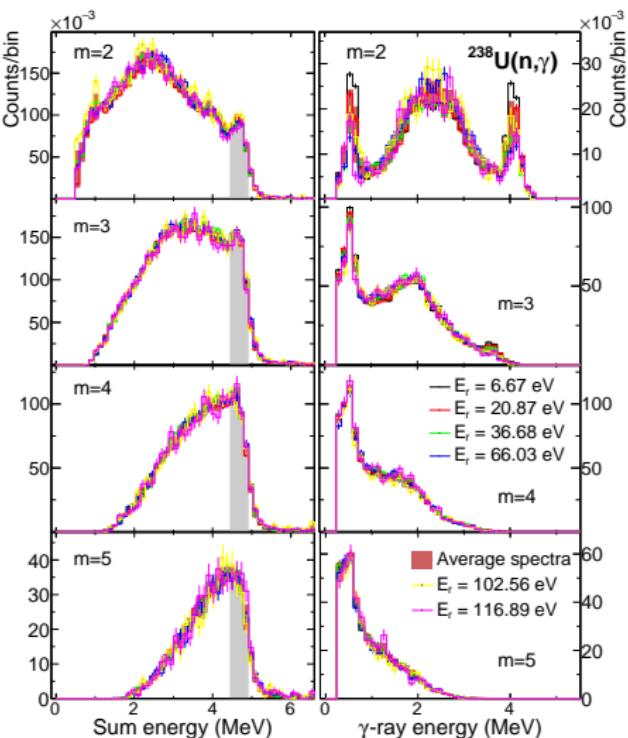
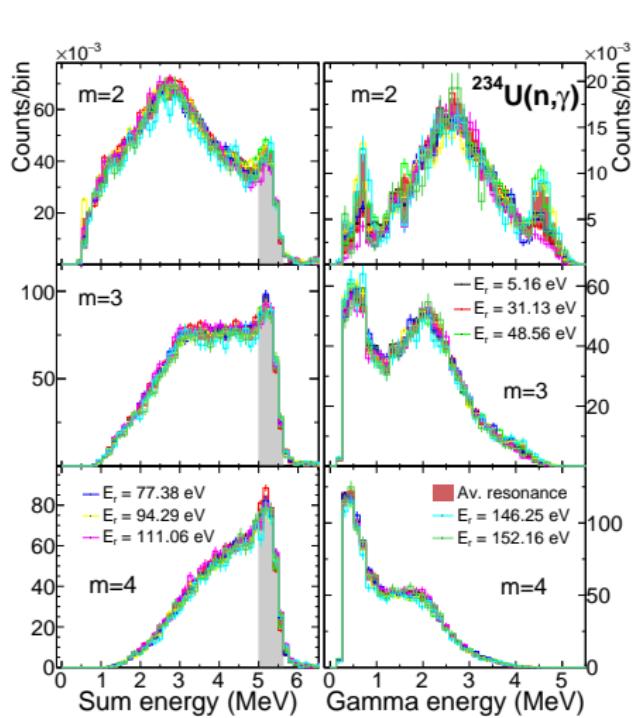


Average experimental MSC spectra



Average experimental MSC spectra

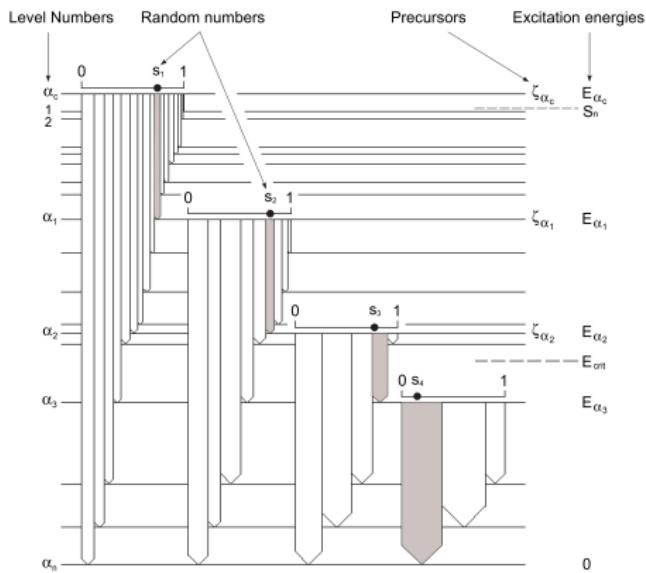




Simulations of MSC and sum-energy spectra

Cascades from DICEBOX simulation

$$\Gamma_{i\gamma f} = \sum_{XJ} y_{ifXJ}^2 (E_i - E_f)^{2J+1} \frac{f^{(XJ)}(E_i - E_f)}{\rho(E_i, J_i, \pi_i)}$$

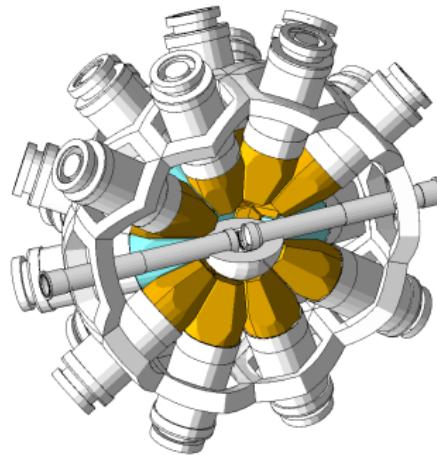


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are fed to GEANT4 detector response of TAC

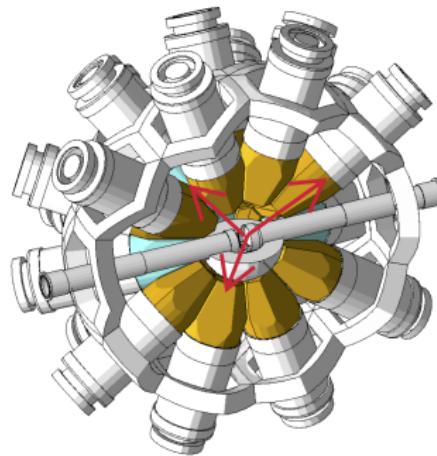


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Nuclear levels and their decay/excitation

Lane and Lynn defined the photon strength function $S^{(XL)}$ as rescaled smoothed cross section:

$$S^{(XL)}(E_\gamma) = \frac{1}{(\pi\hbar c)^2} \frac{\bar{\sigma}_\gamma^{(XL)}(E_\gamma)}{(2L+1)E_\gamma^{2L-1}}, \quad (1)$$

or related to the partial radiation width according to Bartholomew:

$$S^{(XL)}(E_\gamma) = \frac{\bar{\Gamma}_{i\gamma f}^{(XL)}(E_\gamma)\rho(E_i, J_i, \pi_i)}{E_\gamma^{2L+1}}, \quad (2)$$

These two definitions are connected by detailed-balance principle:

$$\bar{\Gamma}_{i\gamma f}^{(XL)}(E_\gamma)\rho(E_i, J_i, \pi_i) = \frac{E_\gamma^2}{(\pi\hbar c)^2} \frac{1}{2L+1} \bar{\sigma}_{\gamma, f \rightarrow E_i}^{(XL)}(E_\gamma),$$

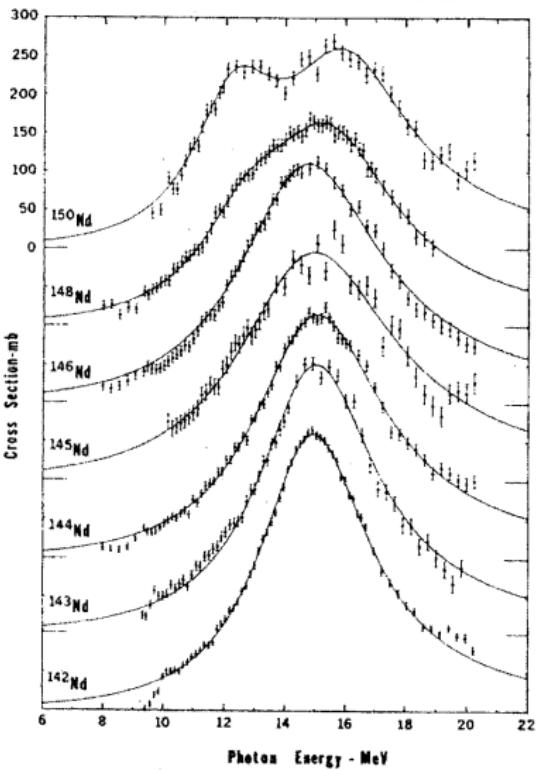


Giant Electric Dipole Resonance

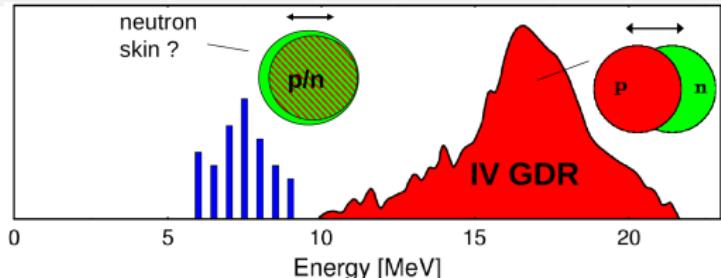
- measured in (γ, x_n) experiments
 - from classical electrodynamics – this collective mode in $E1$ PSF should be described by Lorentzian shape

$$S_{SLO}^{E1} = \frac{1}{3(\pi hc)^2} \frac{\sigma_G E_G \Gamma_G^2}{(E_\gamma^2 - E_G^2)^2 + E_\gamma^2 \Gamma_G^2}$$

- positions E_G , widths Γ_G and cross sections σ_G obtained from fits of data
 - extrapolation below S_n uncertain \Rightarrow phenomenological models: KMF, GLO, EGLO, MGLO, ... and calculations



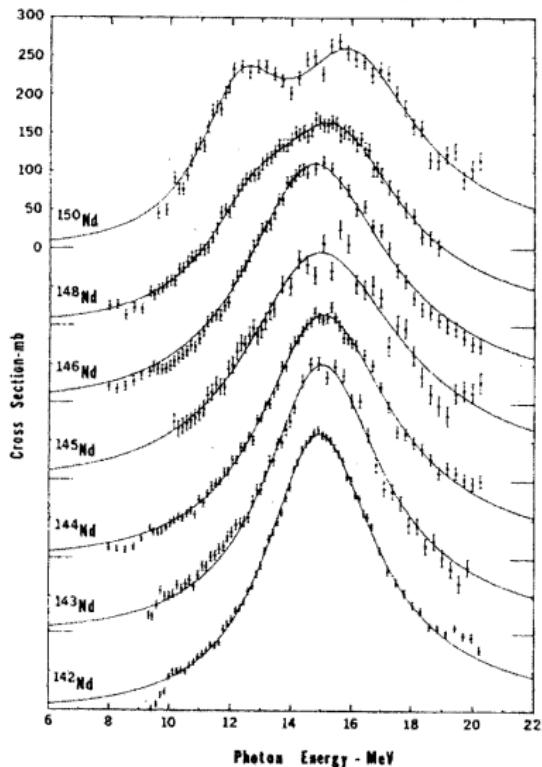
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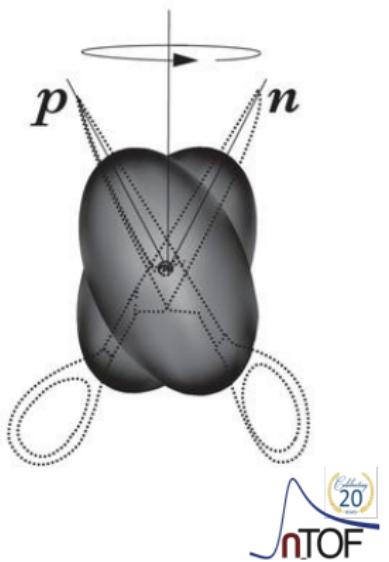
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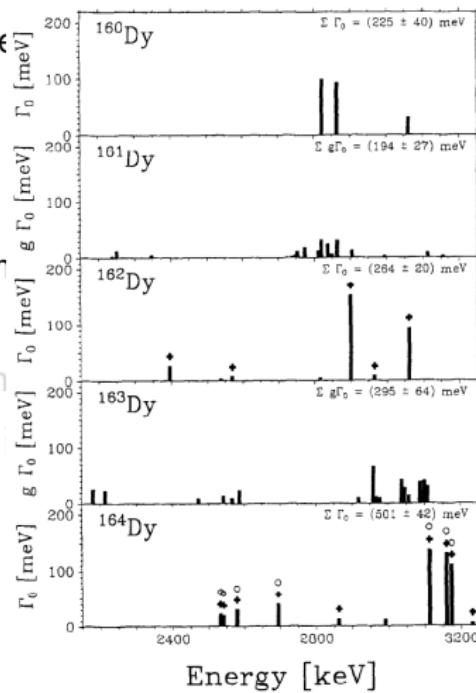
Scissors Mode

- $M1$ collective states in deformed nuclei predicted by theory: Interacting Boson Model, Two-Rotor Model
- experimentally discovered in (e, e') on ^{156}Gd
- measured in nuclear resonance fluorescence scattering (NRF; (γ, γ')) for many rare-earths and later for actinides
- \Rightarrow experimental SM strength dependence on ...
- in decay confirmed from neutron capture & in light ion induced reactions – Oslo method
- the results from different reactions usually agree on the position of SM, but sometimes differ in strength



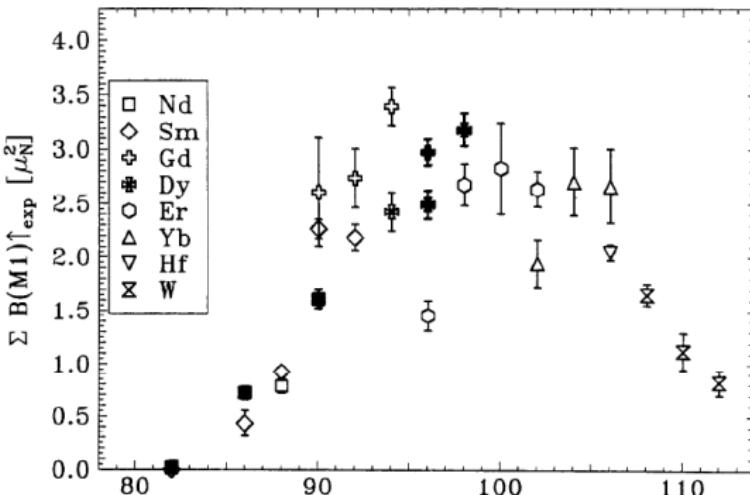
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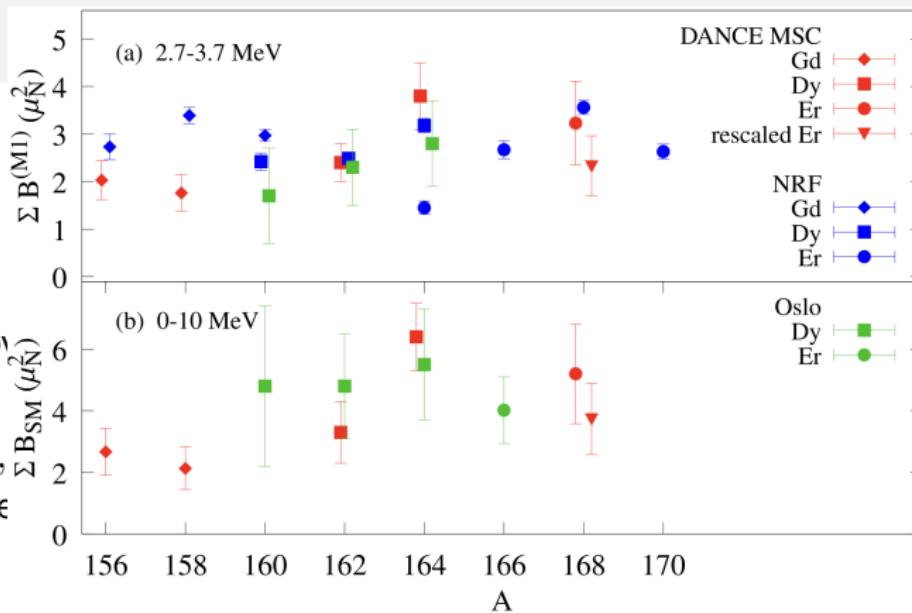
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Model combinations from literature

- from **RIPL-3** database: the analytical **GLO** model for $E1$ PSF and the spin-flip **SF** Lorentzian for $M1$ PSF coupled with the constant-temperature **CT** NLD model,
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- the original model interpretation of the **Oslo** PSF and NLD data: **EGLO** $E1$ model, the **SM** and **SF** $M1$ Lorentzians, and the **CT** NLD,
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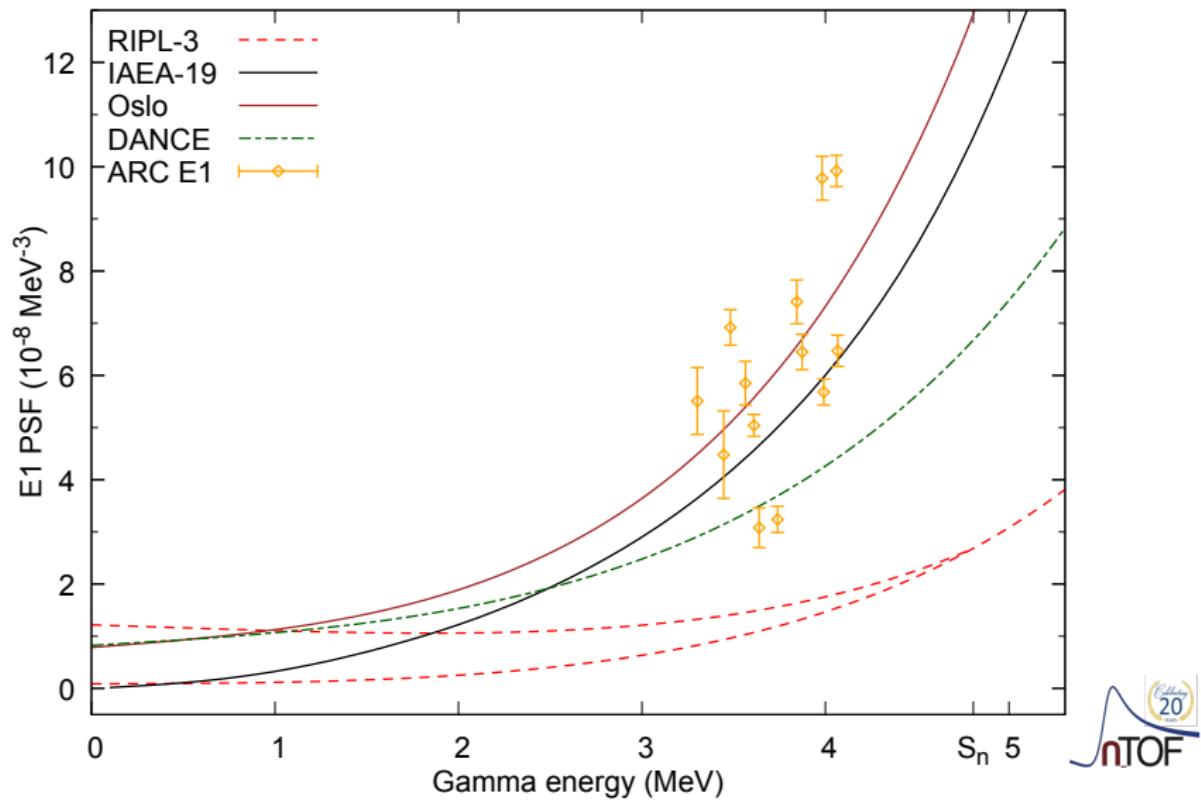
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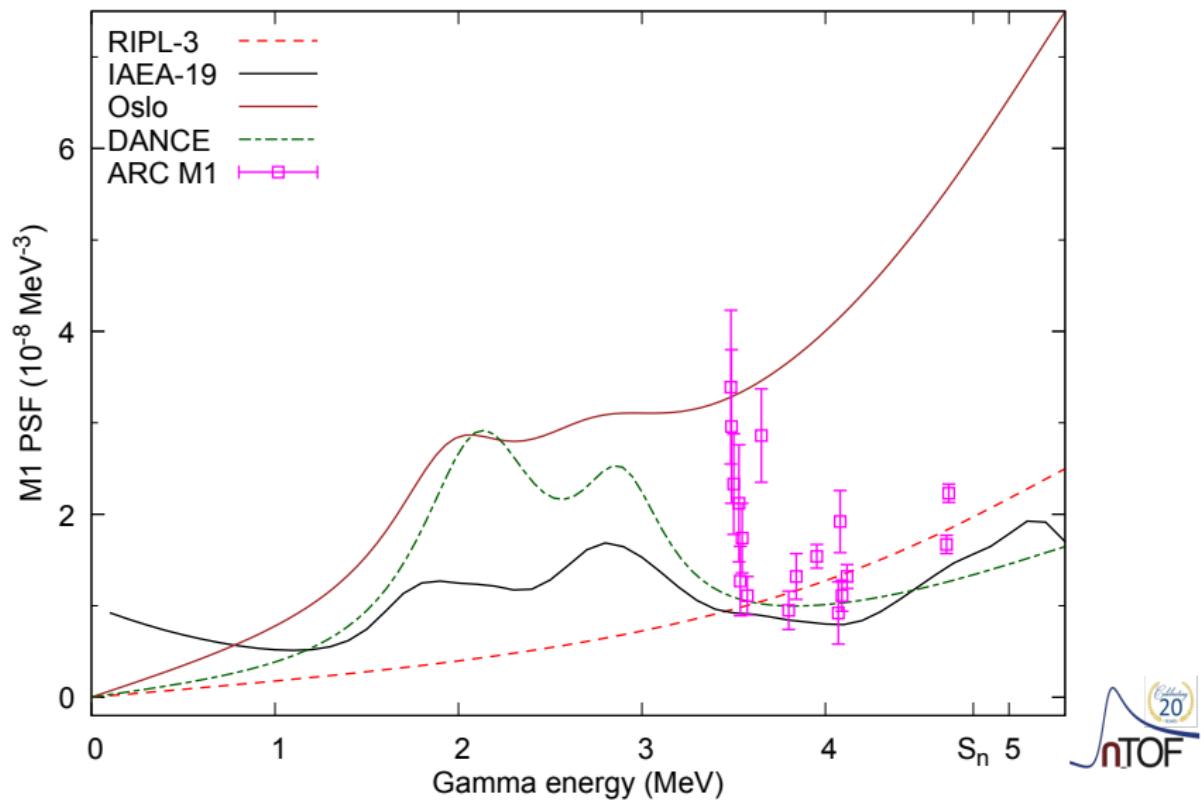


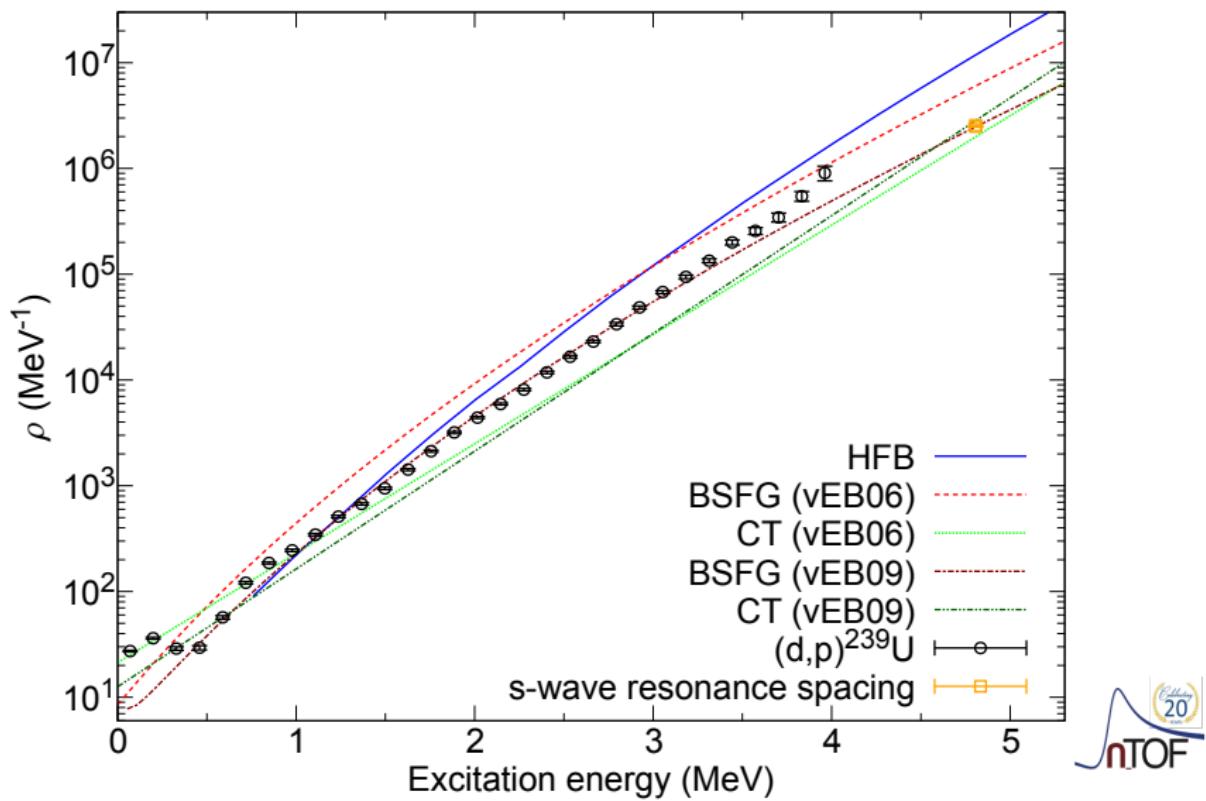
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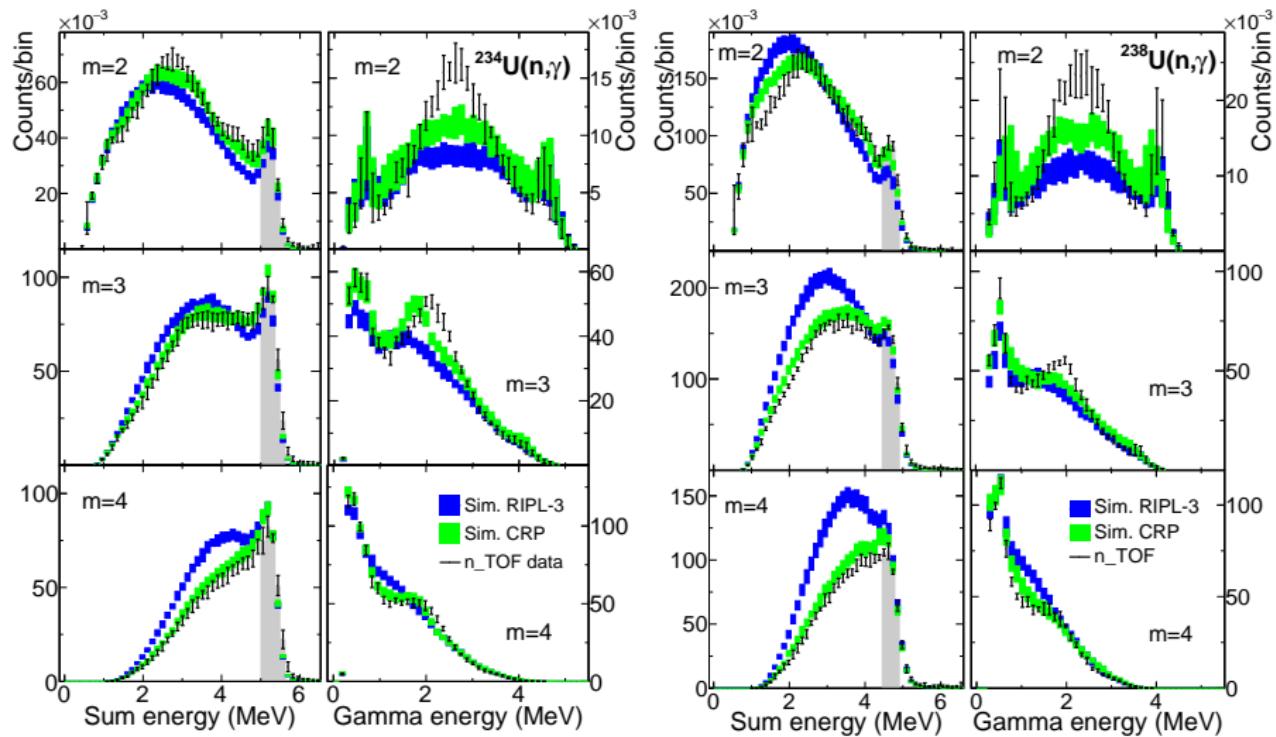


Literature $E1$ PSFs of ^{239}U 

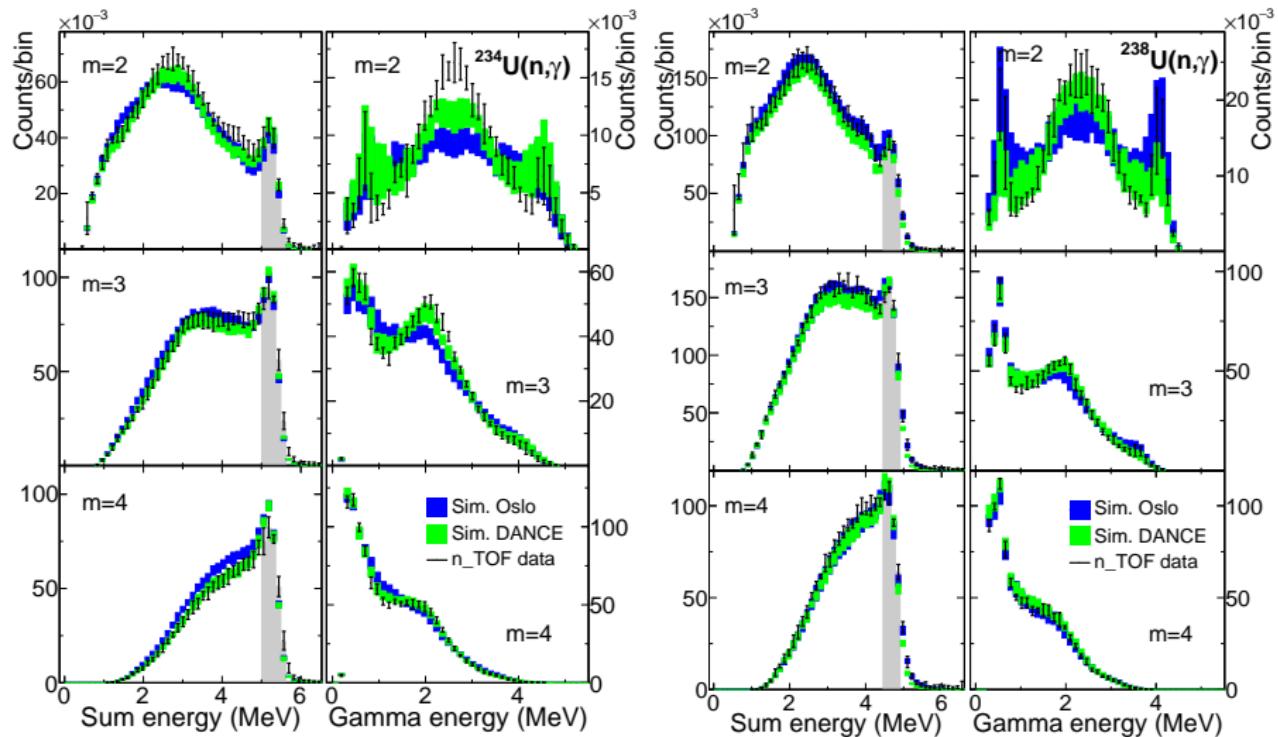
Literature $M1$ PSFs of ^{239}U 

Literature NLD of ^{239}U 

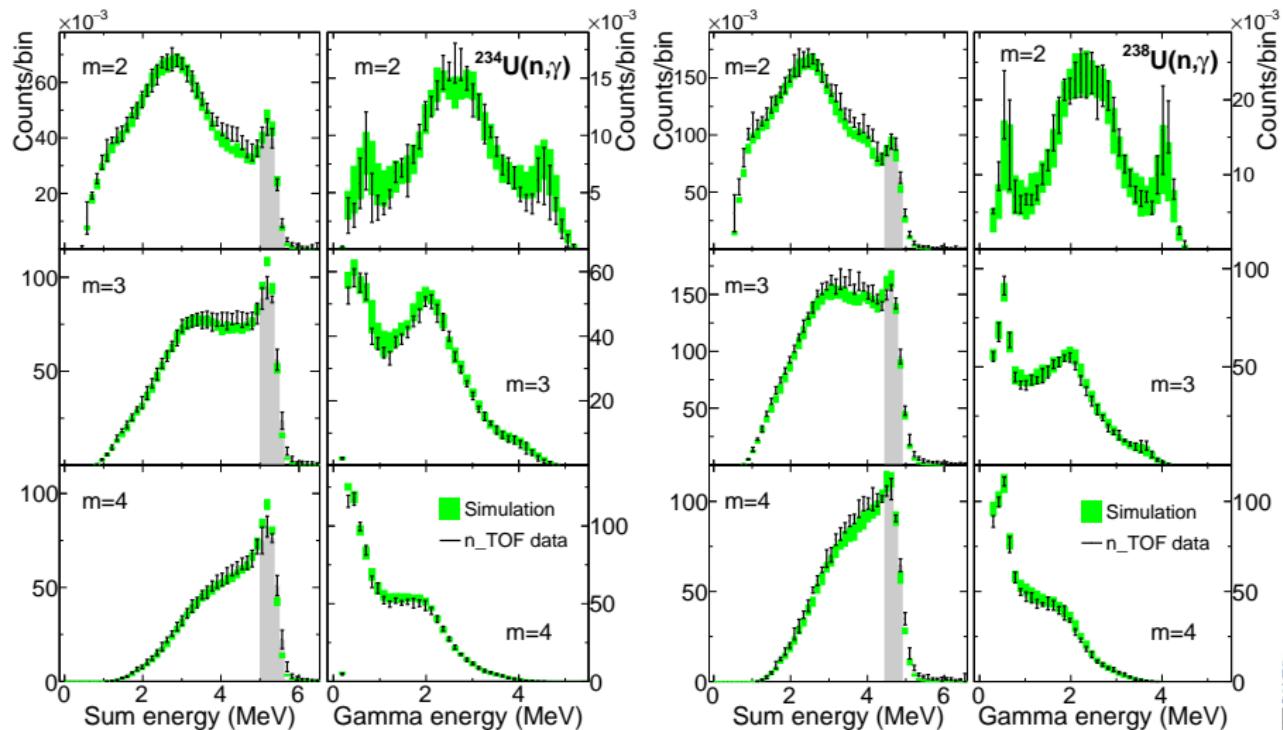
^{235,239}U – RIPL-3 & IAEA-19



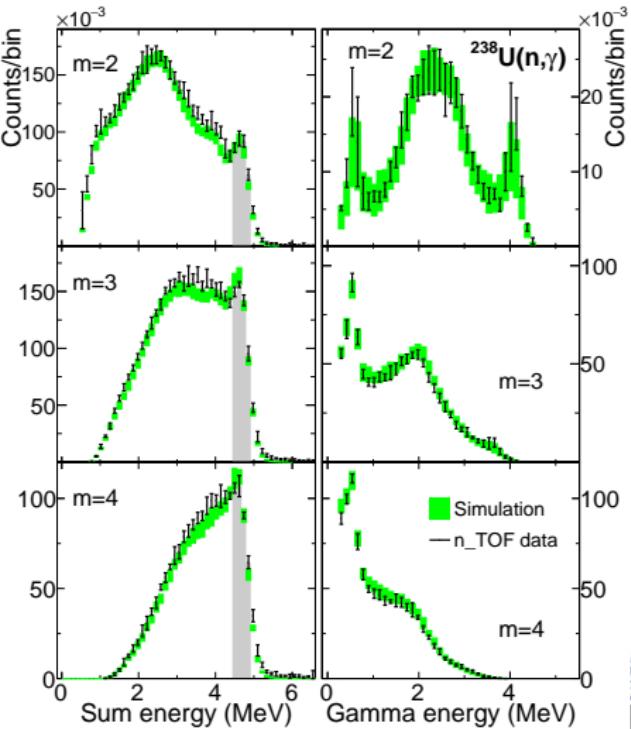
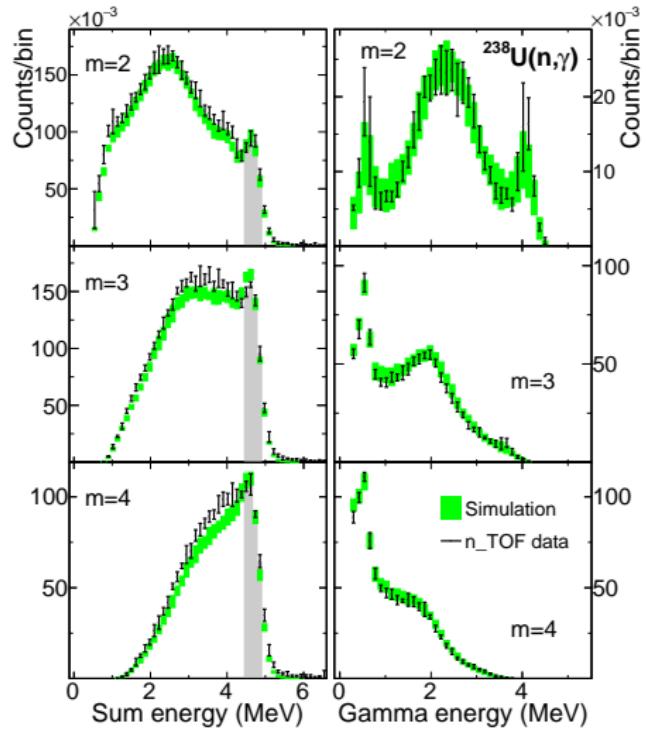
$^{235,239}\text{U} - \text{Oslo} \& \text{DANCE}$



$^{235,239}\text{U}$ – adjusted PSFs from n_TOF analysis



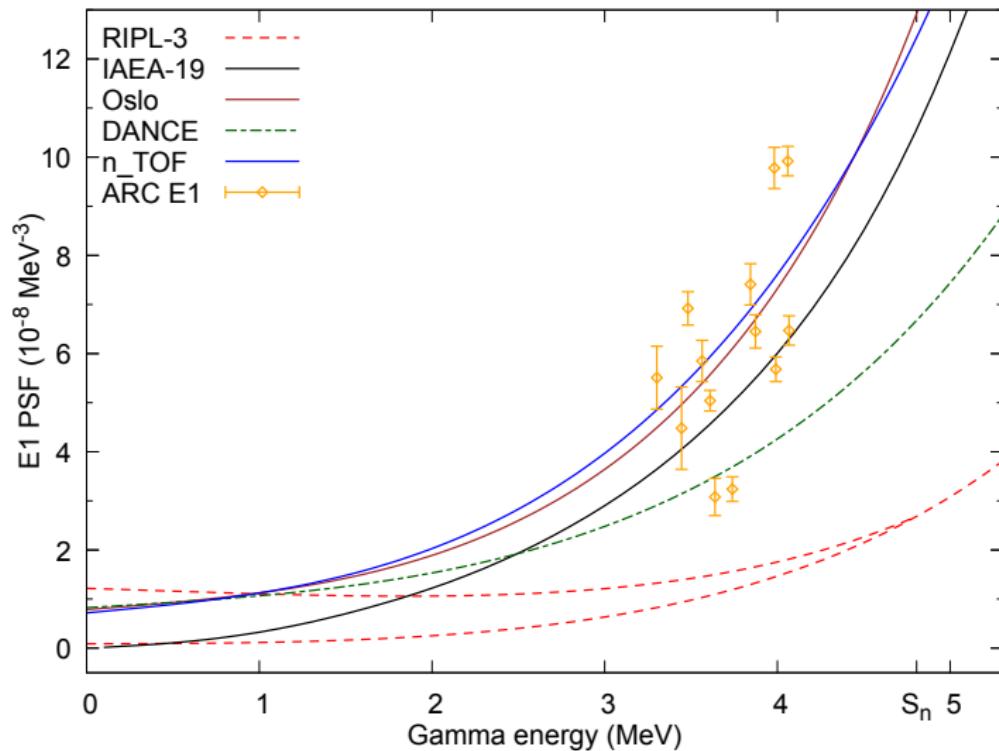
$^{239}\text{U} - \text{MGLO}$ with $k = 3$ vs $k = 1.8$

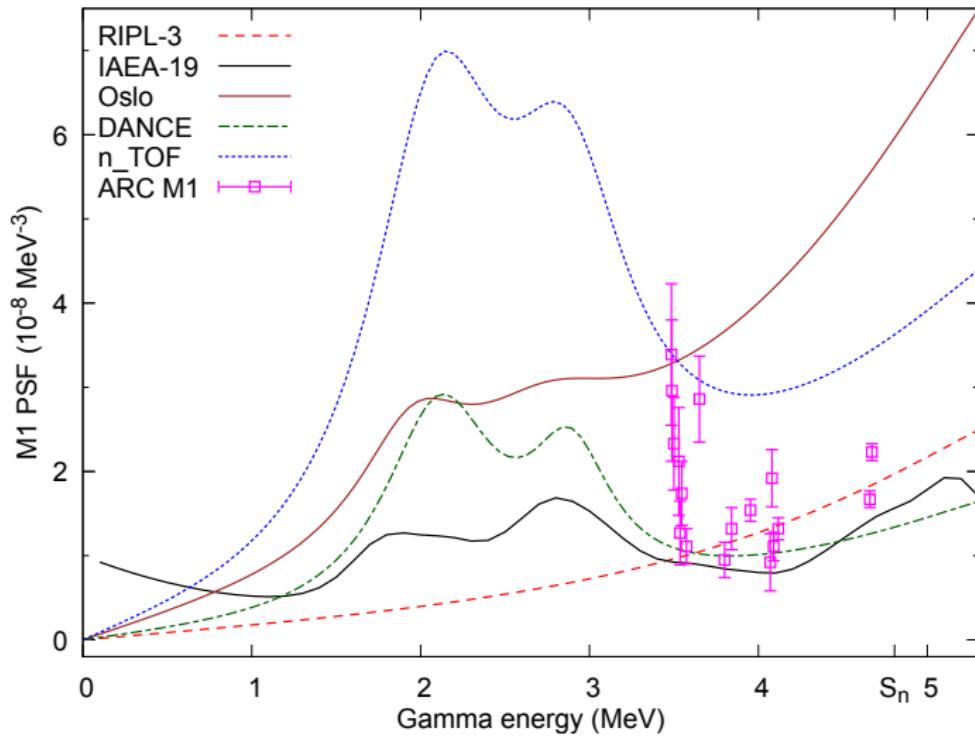


Model combination PSF-LD	Γ_γ (meV)	$^{236}U(n, \gamma)$	$^{238}U(n, \gamma)$
	$^{234}U(n, \gamma)$		
RIPL-3	16.1(2)	12.9(2)	9.5(2)
IAEA-19	29.4(6)	19.3(5)	13.9(5)
Oslo	19.9(4)	20.4(6)	18.6(8)
DANCE	22.0(5)	17.2(4)	15.9(6)
MGLO(1.8)	25.4(7)	20.1(5)	15.9(6)
MGLO(2.5)	30.5(10)	23.9(7)	18.8(7)
MGLO(3.0)	39.0(12)	30.9(9)	24.3(9)
MGLO($k, T(E)$)	26.7(7)*	24.5(6)†	19.2(7)‡
Mughabghab's atlas 2006	25.3(10)	23.4(8)	23.36(31)
Mughabghab's atlas 2018	36.7(7)	23.4(8)	22.9(4)
JEFF-3.3	26.0	23.0	22.5
ENDF/B-VIII.0	26.0	19.5	22.5

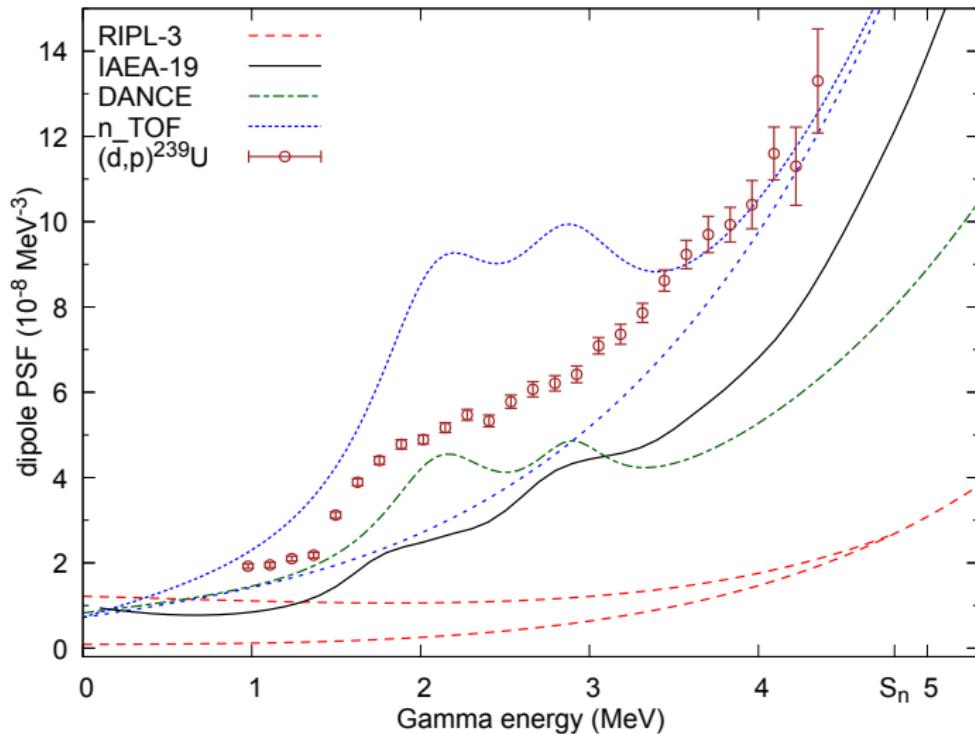
* with $k = 1.8$; † with $k = 2.5$; ‡ with $k = 3.0$



Resulting $E1$ PSF for ^{239}U 

Resulting $M1$ PSF for ^{239}U 

Resulting dipole PSF for ^{239}U



LD and PSFs in odd U's

- sum-energy and MSC spectra of ^{odd}U measured with TAC at n_TOF
- no available model combination is able to describe our data
- we have performed extended search to find a satisfactory reproduction of our experimental spectra
- and we succeeded, see 10.1103/PhysRevC.105.024618 ;)
- CT LD is strongly favored (with unique parameters for each isotope)
- MGLO model is able to describe the $E1$ PSF for energies from 0 to GEDR region
- scissors mode strongly influences decay below S_n , two-resonance form is strongly favored
- taking Γ_γ 's into account implies there are no universal PSFs for odd U's, i.e. SM strength increases with mass



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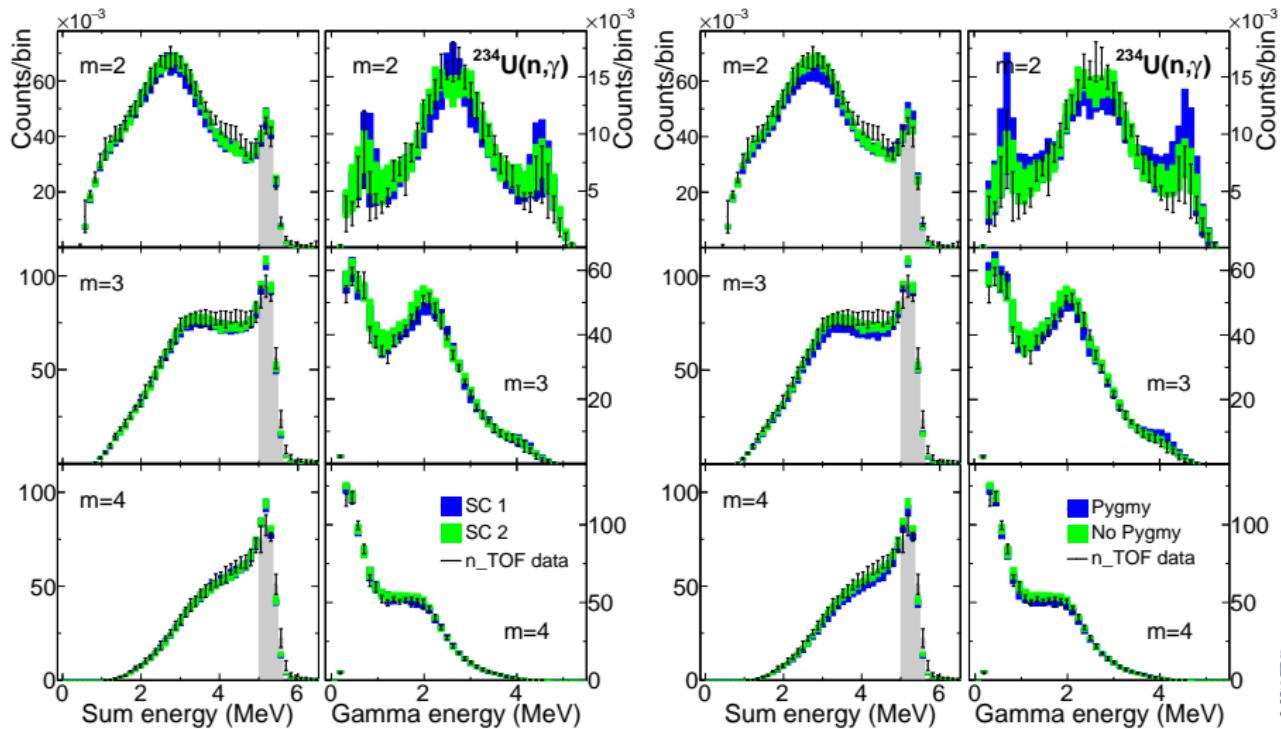
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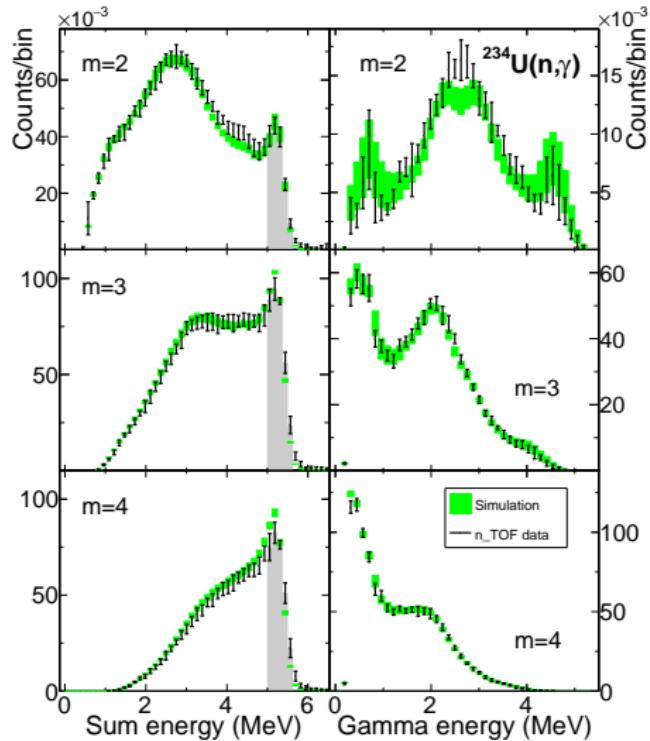
Thank You for listening!



SM two- or single- resonance? $E1$ Pygmy?



LEE? For moderate one we can not say.



The influence of LEE depends on its shape and magnitude.

In general LEE shifts the multiplicity to higher values, but the footprint in spectra is not so straightforward due to gamma energy threshold.

Note that if LEE is sizable and “squished” to very low energies, the internal conversion will hide its presence in MSC spectra. The question would then be if there are so many electrons flying out?



Footprint of individual SM resonance terms

