

# Linking optical potentials to resonances at lower energies

Checking the accuracy of common approximations

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# Summary

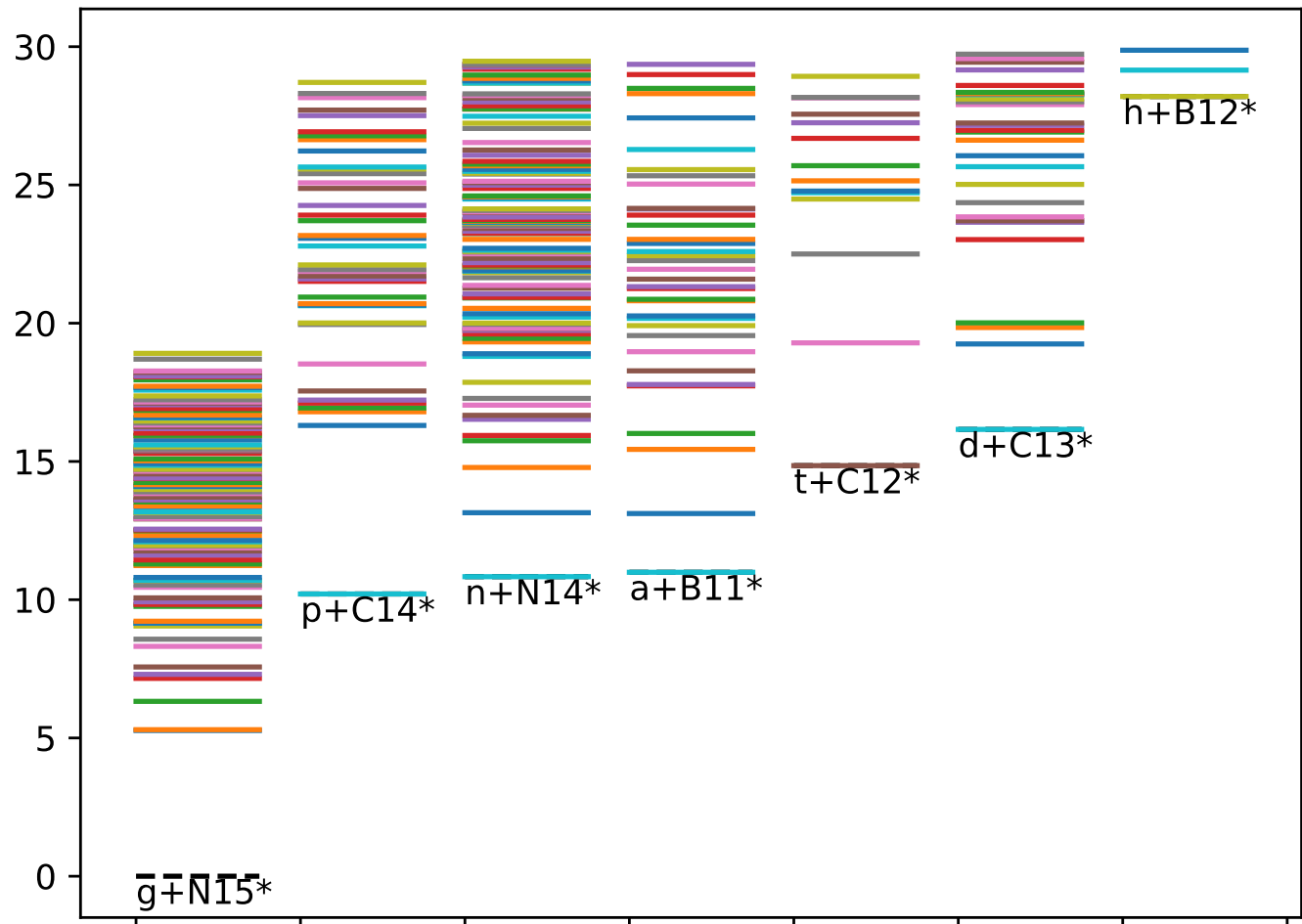
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- Neutrons +  $^{14}\text{N}$  reactions as a test case
- Known levels and level densities
- Going from R-matrix theory to Hauser-Feshbach models
- How to check the approximation used
- Try making R-matrix parameters to reproduce optical results, at least in middle range when resonances overlap more.

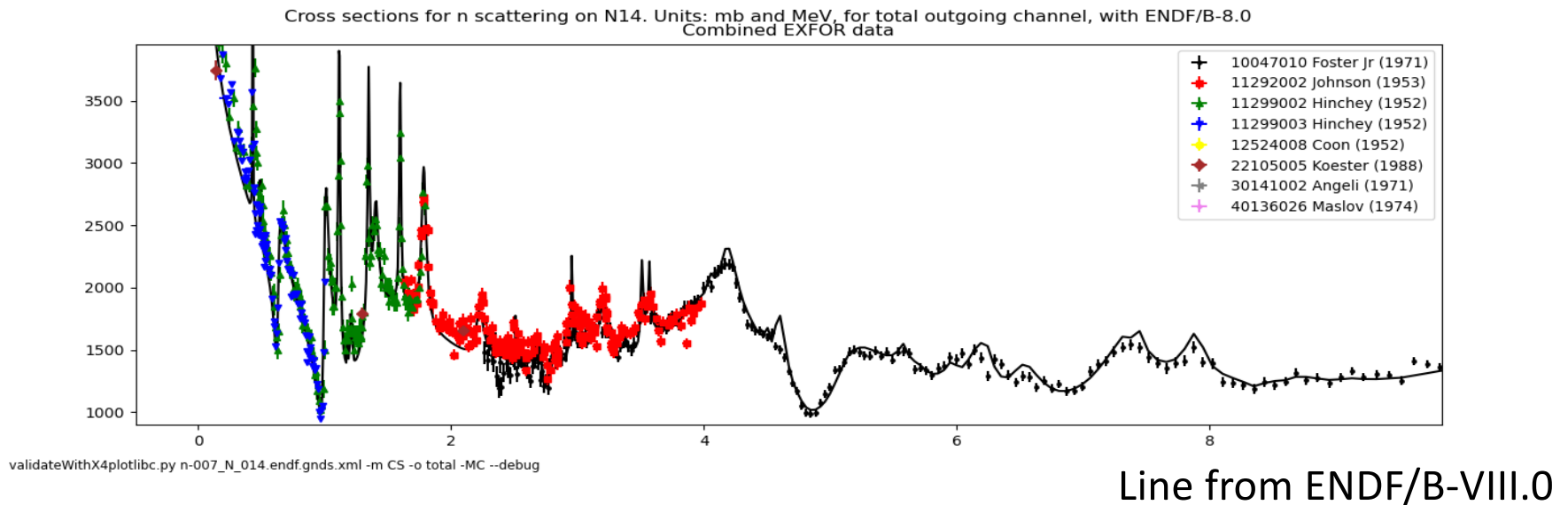


# Known levels in $n+N14 = N15^*$ reactions

- Known levels from RIPL3:
- $^{15}\text{N}^*$  levels for incident neutrons
- n, p and  $\alpha$  channels near.
- Many gamma decays from excited states.
- Needs a new evaluation!

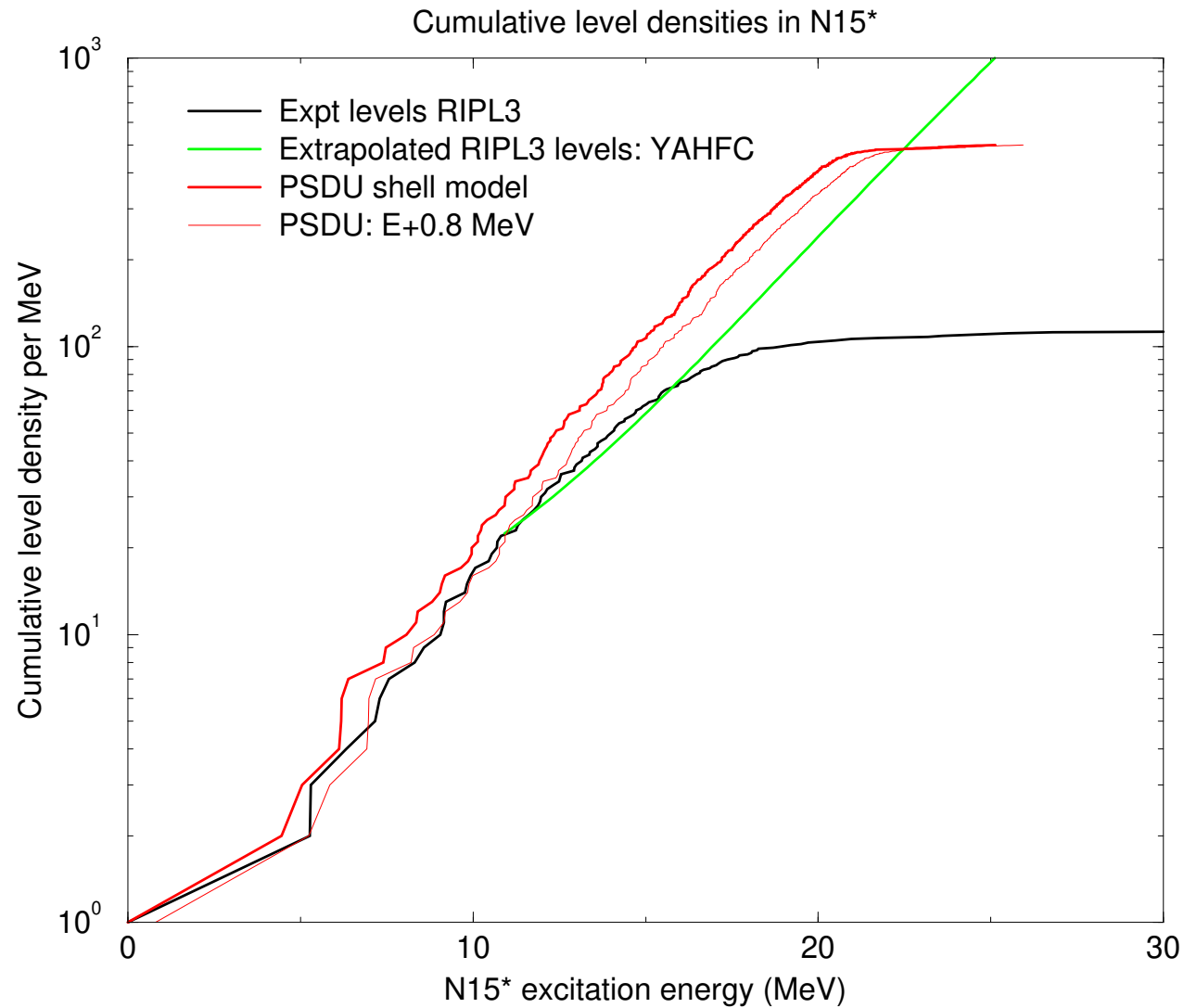


# Neutron 'total cross-section' data from EXFOR



- Pronounced resonances < 4 MeV.
- Smoother > 4 MeV: unresolved resonances with higher densities
- Cannot search to fit individual resonances to this data > 4 MeV.

# Level Densities: levels per MeV



RIPL3 levels  
for N15\* as on  
level diagram.

# R-matrix parameters

1. Diagonalize Hamiltonian inside R-matrix radius  $a$ 
  - (with fixed  $\psi/\psi'$  to make orthonormal basis).
2. Energy eigenvalues  $\varepsilon_p$  for level  $p$ .
3. eigenstate wave functions  $\psi_{p\alpha}(r)$
4. Wf values at  $r = a$ :  $\gamma_{pa} = \sqrt{\frac{\hbar^2}{2\mu_\alpha a}} \psi_{pa}(a)$ 
  - Called 'reduced width amplitudes'
5. Formal width  $\Gamma_{pa} = 2 \gamma_{pa}^2 P_\alpha$  . Penetrability  $P_\alpha = \text{Im} \left( \frac{aH^{+'}}{H^+} \right)$
6. Total width  $\Gamma_p = \sum_\alpha \Gamma_{pa}$

# Exact R-matrix theory

For each specific spin and parity  $J_{\text{tot}}^\pi$ , the multichannel R matrix with  $N$  poles and channels  $\alpha$  is

$$\mathbf{R}_{\alpha\alpha'}(E) = \sum_{p=1}^N \frac{\gamma_{p\alpha}\gamma_{p\alpha'}}{e_p - E}. \quad (1)$$

Using shift functions  $S$  and penetrability  $P$ , define shift and width matrices

$$\hat{\Delta}_{pq} + \frac{i}{2}\hat{\Gamma}_{pq} = \sum_{\alpha} \gamma_{p\alpha}[S_{\alpha} + iP_{\alpha}]\gamma_{q\alpha} \quad (2)$$

Construct the symmetric level matrix  $\mathbf{A}$  is

$$(\mathbf{A}^{-1})_{pq} = \delta_{pq}(e_p - E) - \hat{\Delta}_{pq} - \frac{i}{2}\hat{\Gamma}_{pq}, \quad (3)$$

that is

$$\mathbf{A} = - \begin{pmatrix} E - e_1 + \hat{\Delta}_{11} + \frac{i}{2}\hat{\Gamma}_{11} & \hat{\Delta}_{12} + \frac{i}{2}\hat{\Gamma}_{12} & \dots \\ \hat{\Delta}_{21} + \frac{i}{2}\hat{\Gamma}_{21} & E - e_2 + \hat{\Delta}_{22} + \frac{i}{2}\hat{\Gamma}_{22} & \dots \\ \dots & \dots & \dots \end{pmatrix}^{-1}, \quad (4)$$

The exact multi-channel multi-level  $\tilde{\mathbf{S}}$  matrix is

$$\tilde{\mathbf{S}}_{\alpha'\alpha} = \Omega_{\alpha} \left[ \delta_{\alpha'\alpha} + i \sum_{\lambda\lambda'} \Gamma_{\alpha\lambda}^{1/2} \mathbf{A}_{\lambda\lambda'} \Gamma_{\alpha'\lambda'}^{1/2} \right] \Omega_{\alpha'}. \quad (5)$$



# Single-level Breit-Wigner approximation (SLBW)

Neglect the off-diagonal terms of  $\hat{\Gamma}_{pq}$  for *all* level interference terms:

$$\mathbf{A}_{pq} = \delta_{pq} \frac{1}{e_p - E - \hat{\Delta}_{pp} - \frac{i}{2}\hat{\Gamma}_{pp}}, \quad (6)$$

from which the angle-integrated cross section to channel  $\alpha'$  from  $\alpha$  is

$$\sigma_{\alpha'\alpha}(J_{\text{tot}}^\pi; E) = \frac{\pi}{k^2} g_{J_{\text{tot}}} \sum_p \frac{\Gamma_{\alpha p} \Gamma_{\alpha' p}}{(E - E_p)^2 + \Gamma_p^2/4} \quad (7)$$

for spin coefficient  $g_{J_{\text{tot}}}$ . Each level  $p$  has a total width of  $\Gamma_p = \sum_{\alpha} \Gamma_{\alpha p}$ . The  $p$ -averaged cross section for average level spacing  $D$  is

$$\langle \sigma_{\alpha'\alpha}(E) \rangle = \frac{\pi}{k^2} g_{J_{\text{tot}}} \left\langle \frac{\Gamma_{\alpha p} \Gamma_{\alpha' p}}{\Gamma_p} \right\rangle \frac{2\pi}{D} \sim \frac{\pi}{k^2} g_{J_{\text{tot}}} \frac{\langle \Gamma_{\alpha} \rangle \langle \Gamma_{\alpha'} \rangle}{\langle \Gamma \rangle} \frac{2\pi}{D} \quad (8)$$

where the average sum  $\langle \Gamma \rangle = \sum_{\alpha} \langle \Gamma_{\alpha} \rangle$ .

The reaction cross-section: sum of all outgoing channels:  $\sigma_R = \sum_{\alpha'} \sigma_{\alpha'\alpha}(E)$ .



# Connecting widths with optical potentials

The reaction cross-section from Eq. (8), the sum of all outgoing channels, is

$$\sigma_R = \frac{\pi}{k^2} g_{J_{\text{tot}}} \frac{2\pi \langle \Gamma_\alpha \rangle}{D} \quad (9)$$

This reaction cross section should equal the reaction cross section  $\sigma_R$  given by a one-channel optical potential and its  $S$ -matrix elements,

$$\sigma_R = \frac{\pi}{k_i^2} g_{J_{\text{tot}}} (1 - |\mathbf{S}_{\alpha_i \alpha_i}^{J_{\text{tot}} \pi}|^2) \quad (10)$$

arising from the flux leaving an entrance channel  $\alpha_i$ .

Comparison of two expressions (9) and (10) gives

$$1 - |\mathbf{S}_{\alpha\alpha}^{\text{opt}}|^2 = \frac{2\pi \langle \Gamma_\alpha \rangle}{D} \quad \text{so} \quad \langle \Gamma_\alpha \rangle = \frac{D}{2\pi} (1 - |\mathbf{S}_{\alpha\alpha}^{\text{opt}}|^2). \quad (11)$$

Alternative expression (Simonius, 1974) giving larger widths when  $|\mathbf{S}| \ll 1$ :

$$\langle \Gamma_\alpha \rangle = -\frac{D}{2\pi} \ln(|\mathbf{S}_{\alpha\alpha}^{\text{opt}}|^2), \quad (12)$$

# Checking the approximation used

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1. The Single-level Breit-Wigner approximation
  - a) Check effects of off-diagonal terms
  - b) Check neglect of interferences for diagonal terms
2. Check that the width-fluctuation correction  $W_{\alpha\alpha'}$  is near 1 in

$$\left\langle \frac{\Gamma_{\alpha p} \Gamma_{\alpha' p}}{\Gamma_p} \right\rangle = W_{\alpha\alpha'} \frac{\langle \Gamma_{\alpha} \rangle \langle \Gamma_{\alpha'} \rangle}{\langle \Gamma \rangle}$$

3. Check conversion from optical  $|S_{\alpha}|^2$  to transmission coef  $T_{\alpha}$
4. Check overall Hauser-Feshbach models give cross-sections close to those from R-matrix models,  
at least in some transition region when both should work ok.

# First: Compare R-matrix with HF cross-sections

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- Use a range of known levels from RIPL3
  - Total cross-section has many resonances.
- Preliminary R-matrix fit to make plausible distributions.
- Use Koning-DelaRoche global optical potential for  $n+^{14}\text{N}$
- Use YAHFC Hauser-Feshbach code to also predict cross-sections
- Compare:
  - Total cross-sections for neutrons: reaction + elastic
  - Transfer cross-sections to  $\alpha + ^{11}\text{B}$
  - Transfers to excited state (measured by their gamma decays)

# Hauser-Feshbach models

The reaction cross-section from Eq. (8), the sum of all outgoing channels, is

$$\sigma_R = \frac{\pi}{k^2} g_{J_{\text{tot}}} \frac{2\pi \langle \Gamma_\alpha \rangle}{D} \quad (9)$$

This reaction cross section should equal the reaction cross section  $\sigma_R$  given by a one-channel optical potential and its  $S$ -matrix elements,

$$\sigma_R = \frac{\pi}{k_i^2} g_{J_{\text{tot}}} (1 - |\mathbf{S}_{\alpha_i \alpha_i}^{J_{\text{tot}} \pi}|^2) \quad (10)$$

arising from the flux leaving an entrance channel  $\alpha_i$ .

Comparison of two expressions (9) and (10) gives

$$1 - |\mathbf{S}_{\alpha\alpha}^{\text{opt}}|^2 = \frac{2\pi \langle \Gamma_\alpha \rangle}{D}. \quad (11)$$

So we calculate *transmission coefficients* as

$$\mathcal{T}_\alpha = 1 - |\mathbf{S}_{\alpha\alpha}^{\text{opt}}|^2, \quad (12)$$

in terms of which the channel cross sections of Eq. (8) become

$$\langle \sigma_{\alpha' \alpha}(J_{\text{tot}}^\pi; E) \rangle = \frac{\pi}{k^2} g_{J_{\text{tot}}} \frac{\mathcal{T}_\alpha \mathcal{T}_{\alpha'}}{\sum_{\alpha''} \mathcal{T}_{\alpha''}}. \quad (13)$$

The *Hauser-Feshbach branching ratios* here are

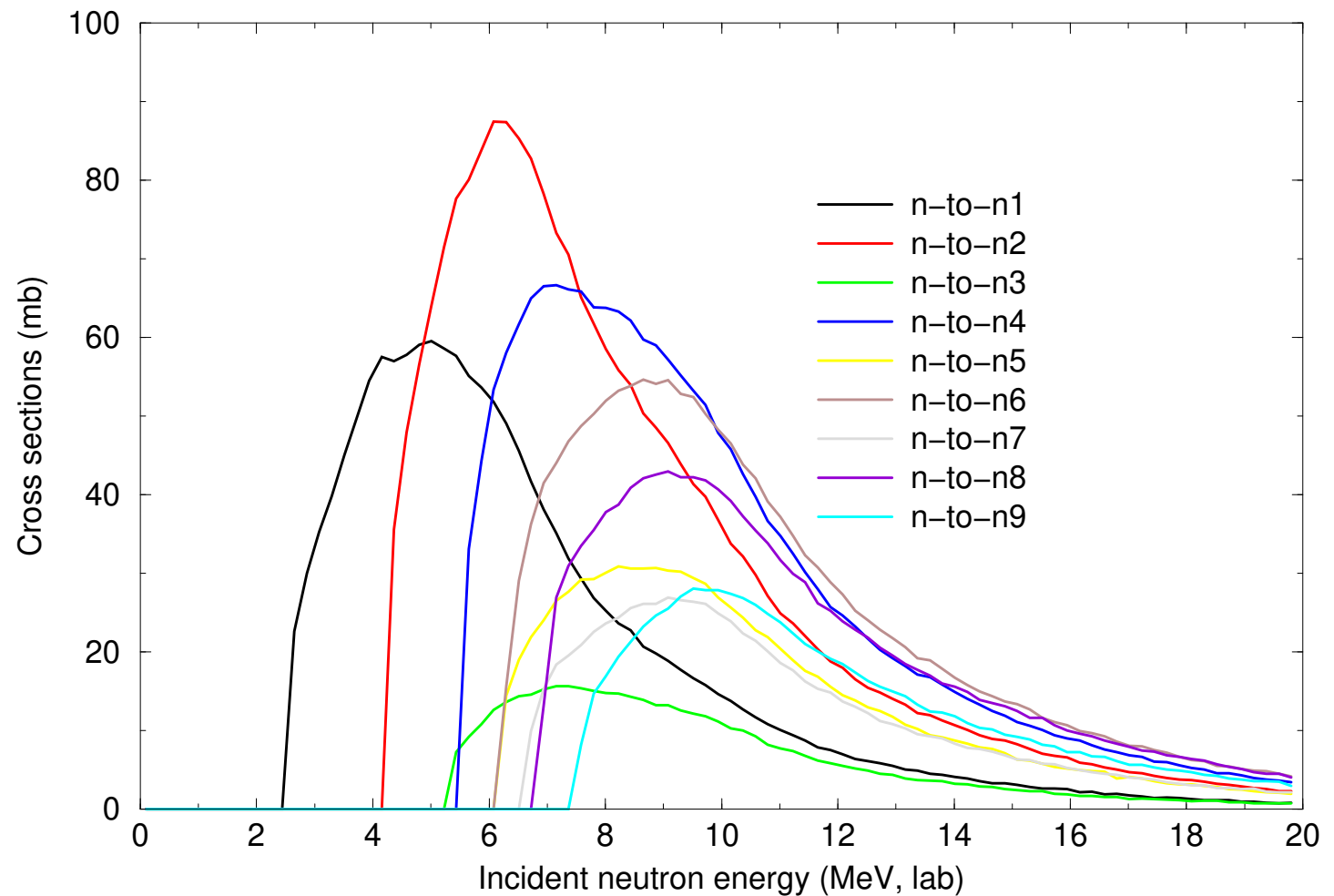
$$\mathcal{B}_{\alpha'} = \frac{\mathcal{T}_{\alpha'}}{\sum_{\alpha''} \mathcal{T}_{\alpha''}}. \quad (14)$$

# Hauser-Feshbach cross-sections (smooth!)

## Inelastic ( $n, n'$ )

At higher energies even more inelastic channels.

Fusion neutrons at 14 MeV will require all these inelastic channels and more !

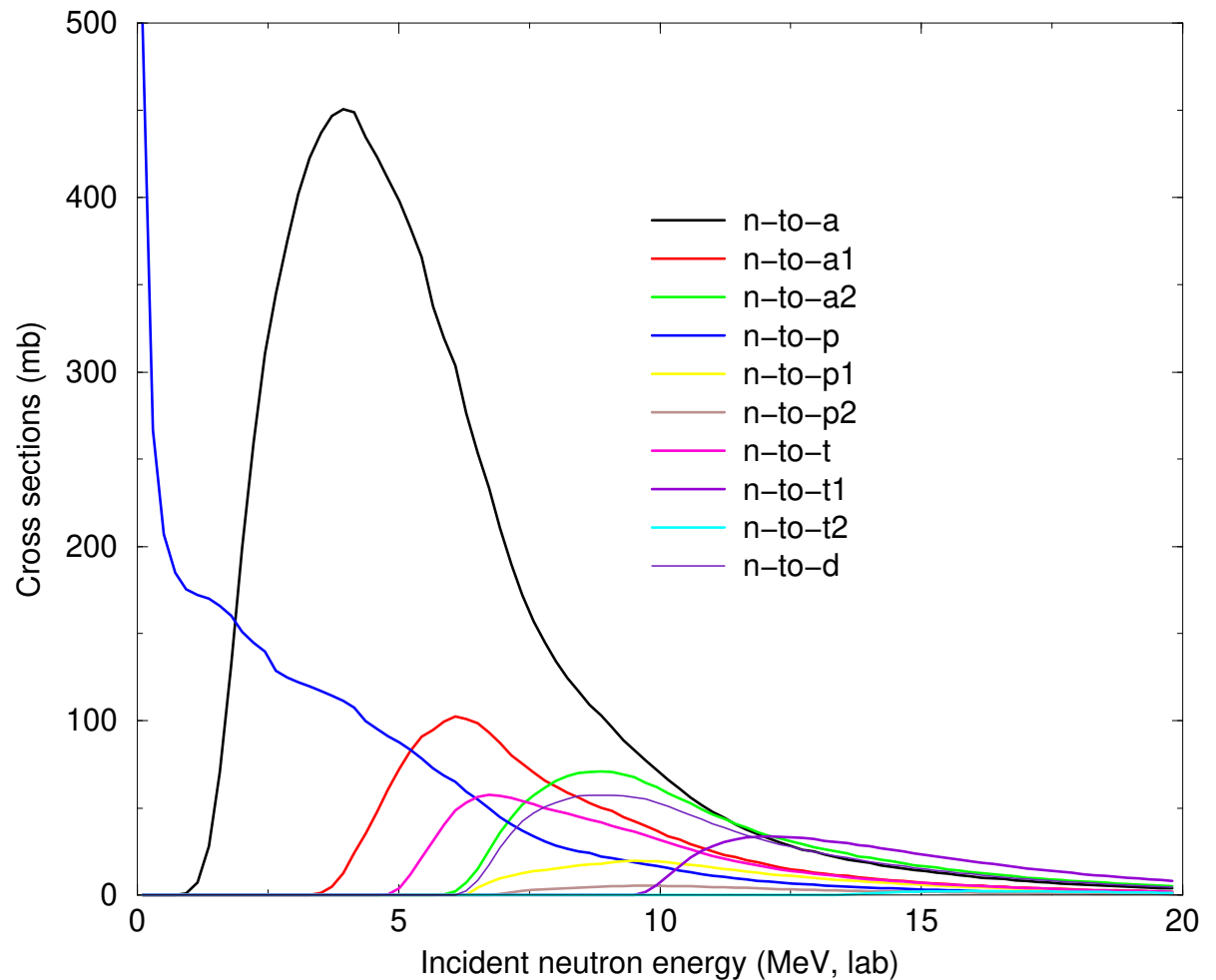


# Hauser-Feshbach transfer cross-sections

Transfers  
(n,a), (n,p), (n,t),  
(n,d)

Many larger than  
(n,n')

Fusion neutrons  
at 14 MeV will  
require all these  
transfer channels  
and more !

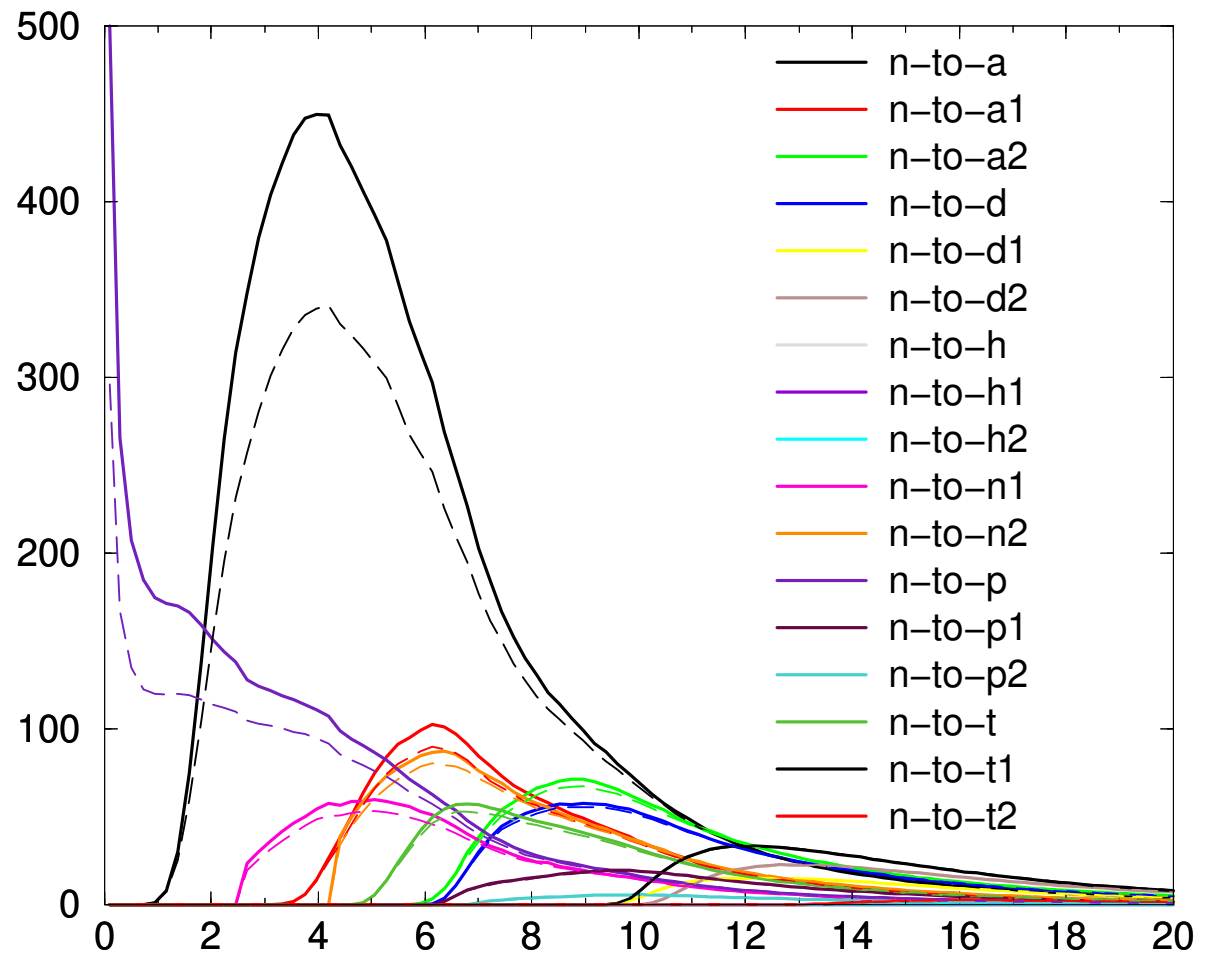


# Hauser-Feshbach width-fluctuation corrections

Dashed lines use WFC from Moldauer PRC (1976), NPA (1980).

For neutrons, usually supposed to be small above 1 MeV, but here we see effects up to 9 MeV.

Use WFC calculated factor in YAHFC

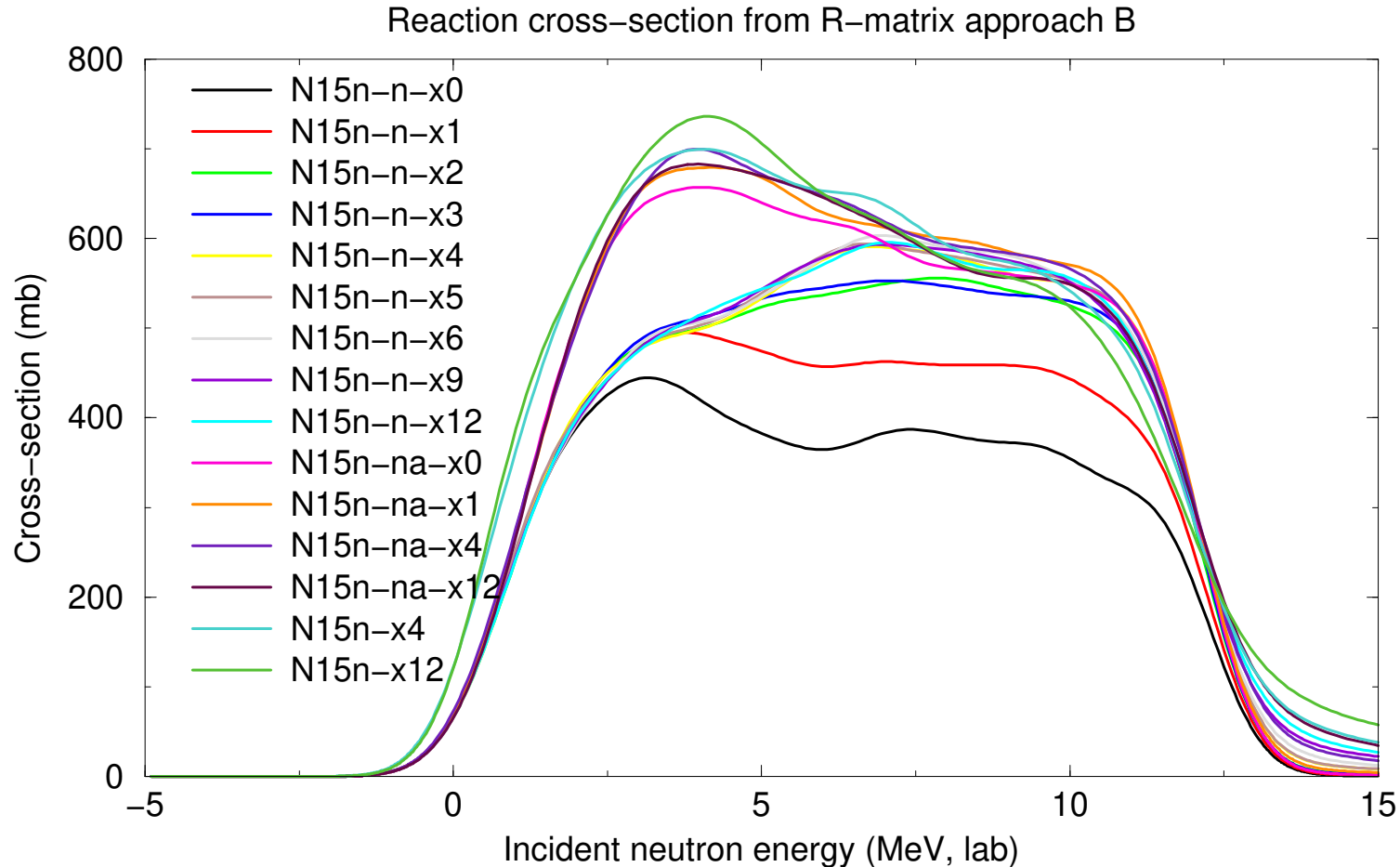




# Generate R-matrix poles from optical potential and level densities

- Start from optical potential for projectiles n, p,  $\alpha$  (etc)
- Choose which nuclear excited states to include (x1, x4, ... , x12)
- Use a level-density to generate spacings  $D$  up to 12 MeV.
- Find partial widths by A:  $\langle \Gamma_{\alpha} \rangle = -\frac{D}{2\pi} \ln(|\mathbf{S}_{\alpha\alpha}^{\text{opt}}|^2)$
- Or by method B::  $\langle \Gamma_{\alpha} \rangle = \frac{D}{2\pi} (1 - |\mathbf{S}_{\alpha\alpha}^{\text{opt}}|^2)$
- Or  $A_p, B_p$ : reduced width *amplitudes* have gaussian fluctuations
- Generate discrete levels with above statistics (like Dicebox)
- Find exact cross-sections from R-matrix theory
- Compare with HF results after smoothing (e.g. 1 MeV Gaussian)

# Reaction cross-sections from R-matrix method B



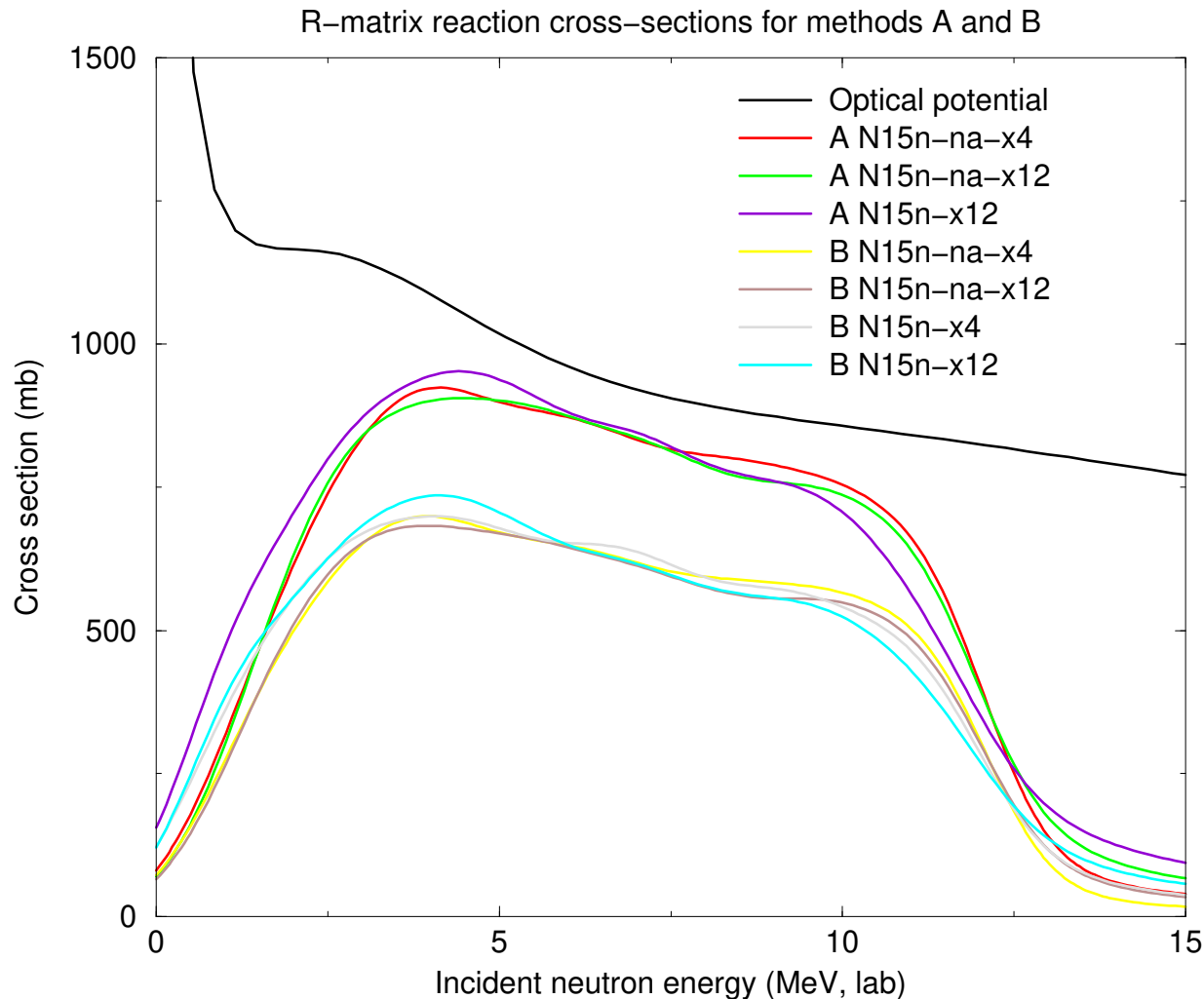
N15n-x is  
 $n + \alpha + p + d + t$

Curves larger  
as more excited  
levels included.

Seems to converge

R-matrix poles  
only to 12 MeV

# Comparing width methods A and B



**R-matrix A method:**  $\gamma$   
magnitude normal  $N(0, \text{mean})$ :  
Porter-T.

**R-matrix B method:**  $\gamma$   
magnitude fixed to mean,  
sign random

Method A gives reaction cross section nearest to that of the optical potential.

But strange short-coming at low energies.  
May be compound-elastic.

R-matrix poles only to 12 MeV

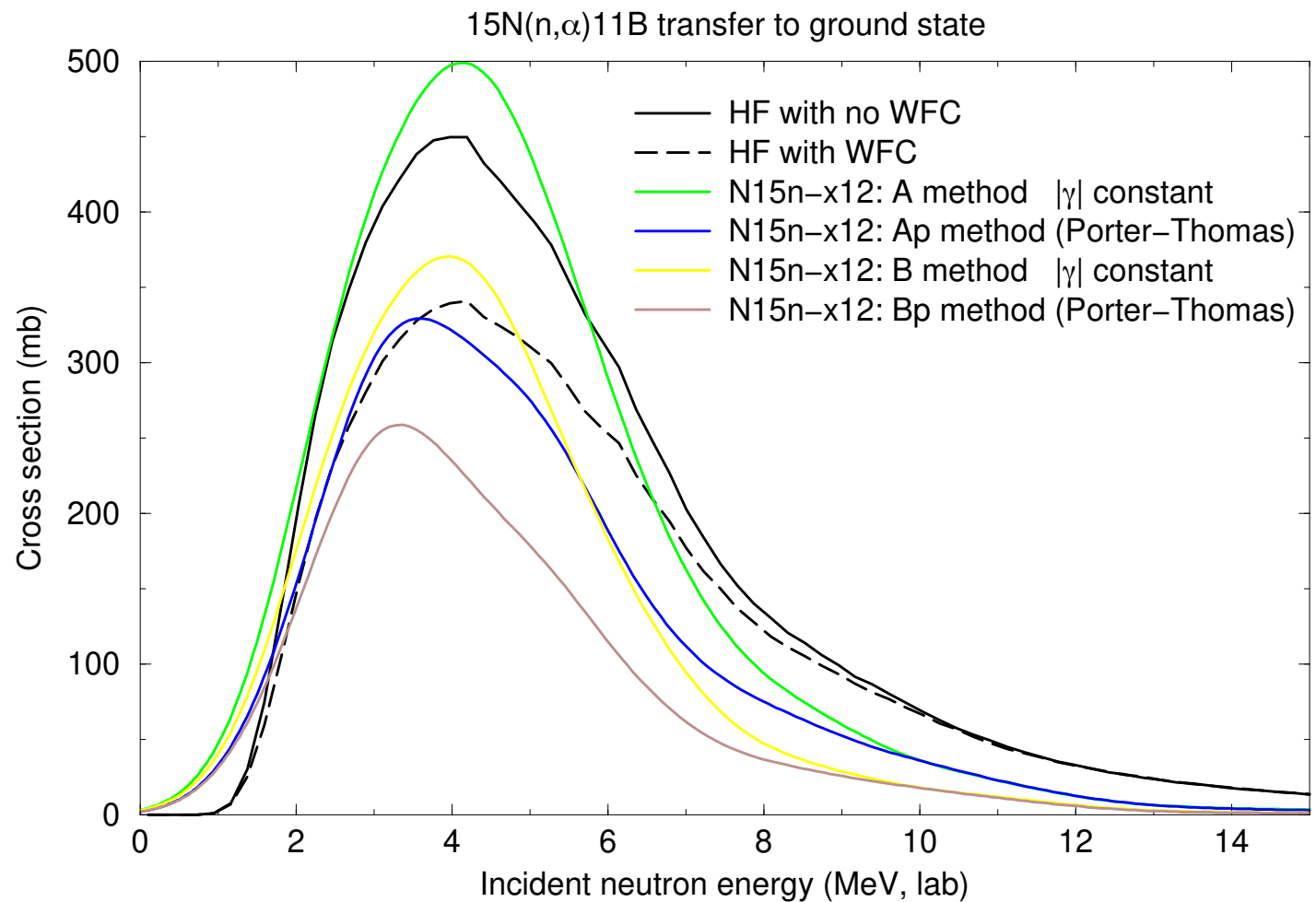
# Channel Comparisons: HF and full R-matrix (A,B) - (n, $\alpha$ ) channel

Most accurate  
are expected to be  
HF+WFC and  
Ap methods.

These two (blue and  
black-dashed) agree  
the best, at least  
up to  $\sim 5$  MeV.

R-matrix A method:  
 $\gamma$  magnitude normal  
N(0,mean): Porter-T.

R-matrix B method:  
 $\gamma$  magnitude fixed  
to mean, sign random

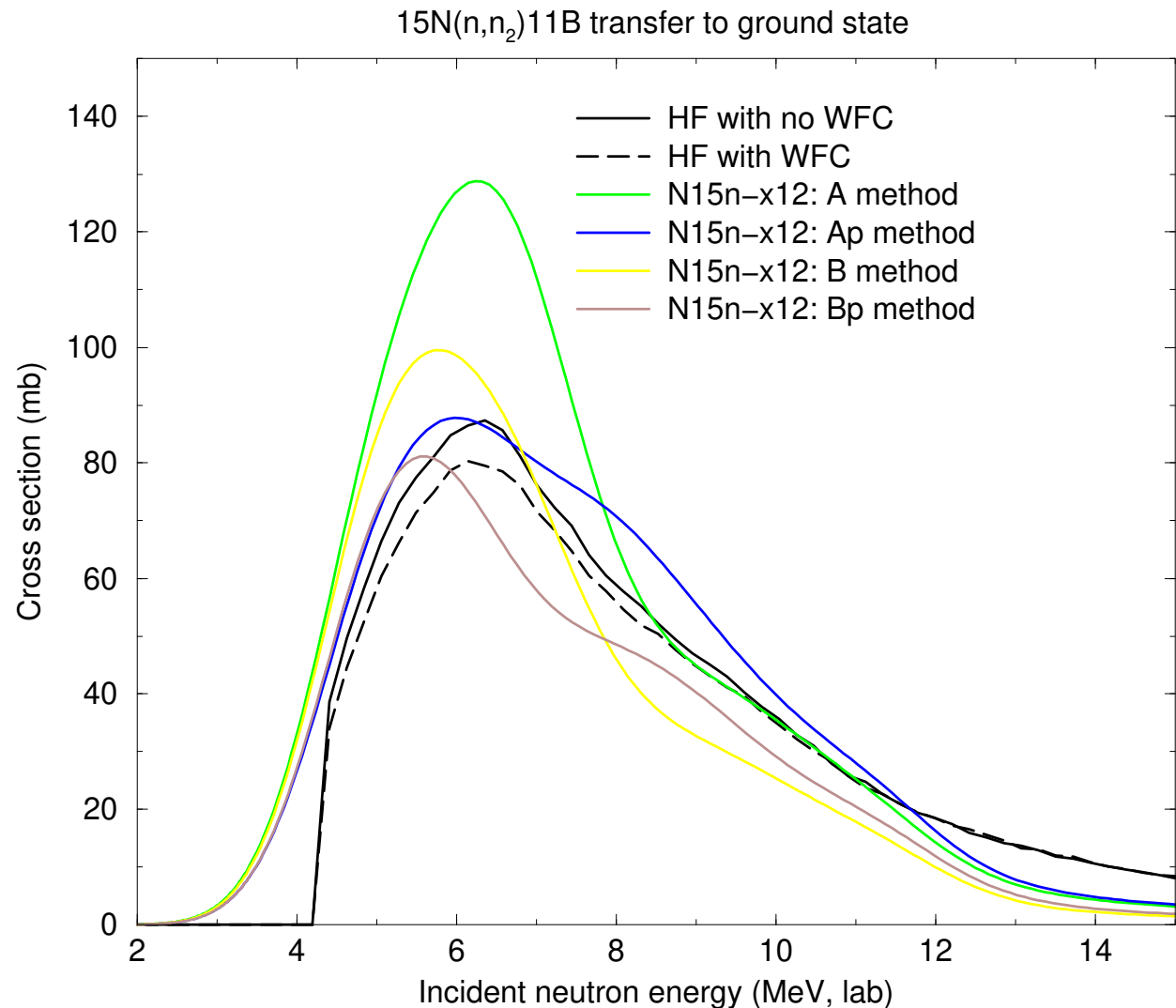


# Channel Comparisons: HF and full R-matrix (A,B) - $(n,n_2)$ channel

Most accurate  
are expected to be  
HF+WFC and  
Ap methods.

HF and Ap methods  
(blue and black)  
agree the best, at  
least up to  $\sim 7$  MeV.

WFC does not have  
large effects here.



# Summary

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- Possible to compare full R-matrix and Hauser-Feshbach models in transition region of unresolved resonances
- Necessary to include all excited residual states up to incident energy. This is well-known for HF, but not so well for R-matrix. (Needed anyway to predict gamma production cross-sections)
- Makes large R-matrix model: I use tensorflow of GPUs.
- Demonstrate best comparison agreements when including
  - Width-fluctuation corrections in HF (up to higher energies)
  - Full Porter-Thomas statistics of reduced width amplitudes  $\gamma_{p\alpha}$ .
- These are well known for high-A targets, but here for N15 too.



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