Isotopic Fuel Tailoring as Actuator for Burn Control in Tokamak Reactors: Tritium Concentration Effects on Operations & Control

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Kinetic (Burn) Control

- What Type of Model Do We Need for Burn-Control Design?
- Why a Nonlinear Control Solution is Needed?
- Synthesis of Nonlinear Controller for the Regulation the Burn Condition
- How Do We Deal With Model Uncertainties and Unmodeled Dynamics?
- How Do We Select the Correct Reference for the Controller?
- How Do We Close the Loop if State Is Not Fully Measurable?
- How Do We Handle Actuator Dynamics and Constraints?

Robust Burn Control Against Drift/Biases in Fueling-line Tritium Fractions

- Impact of Fueling-line Tritium Fractions on ITER's Operational Space
- Summary of the Presentation

Burn Control \equiv Density and Temperature Control of All Species



- Fusion power regulation in DT plasma \leftrightarrow species density/temperature control
- $\bullet\,$ Dimensionality and nonlinearity of the burning plasma \rightarrow model-based control

Actuators Used to Control the Plasma Kinetic Condition



- Plasma current contributes to heating through Ohmic heating (small in reactor)
- Magnetic configuration affects burn condition through confinement time
- Neutral beam injectors and radio frequency waves heat the plasma
- Refueling at the plasma boundary is achieved through gas puffing
- Refueling in the plasma core is achieved through pellet injection
- Impurity injection dilutes the fuel content and increases radiation losses
- Gas pumping removes exhausted fuel, alpha particles, and impurities

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Burning Plasma Model (Dynamics): 1D Conservation Equations

Alpha Particle Density

$$\frac{\partial}{\partial t}n_{\alpha} = -\nabla \cdot \vec{\Gamma}_{\alpha} + S_{\alpha}$$

Deuterium Density:

$$\frac{\partial}{\partial t}n_D = -\nabla \cdot \vec{\Gamma}_D - S_\alpha + S_D^R + S_D$$

Tritium Density:

$$rac{\partial}{\partial t}n_T = -
abla \cdot ec{\Gamma}_T - S_lpha + S_T^R + S_T$$

Impurity Density:

$$\frac{\partial}{\partial t}n_I = -\nabla \cdot \vec{\Gamma}_I - S_I$$

Energy Density:

$$\frac{\partial}{\partial t} \left[\frac{3}{2} nT \right] = -\nabla \cdot \vec{\Gamma}_E + Q_\alpha S_\alpha + P_{ohm} - P_{brem} + P_{aw}$$

- *T*: Plasma Temperature ($T \triangleq T_i = T_e$)
- n: Plasma Density
- Paux: Auxiliary Power
- Pohm: Ohmic Power
- P_{rad}: Radiation Power
- S_D, S_T: Deuterium, Tritium Refueling
- $-S_D^R, S_T^R$: Deuterium, Tritium Recycling
- Fusion Reaction Rate: $S_{\alpha} = n_D n_T \langle \sigma \nu \rangle$
- Reactivity: $\langle \sigma \nu \rangle$
- $\vec{\Gamma}_{\alpha}, \vec{\Gamma}_{D}, \vec{\Gamma}_{T}, \vec{\Gamma}_{I}, \vec{\Gamma}_{E}$: Transport Fluxes

Particle Densities and Temperature are functions of time (t) and radial position ($\hat{\rho}$)

Burning Plasma Model (Controls): 0D Balance Equations

- 1D burning plasma dynamic model may be intractable for control synthesis
 - Transport models are too complex and still under development
- Control goal is $0D \Rightarrow 0D$ response model is what is needed for control synthesis
- $\bullet\,$ Control synthesis: 0D model $\rightarrow\,$ Performance assessment (simulations): 1D model



- 0D response model is based on energy/particle balance equations
- Particle balance equations are needed for all species

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Isotopic Fuel Tailoring as Actuator for Burn Control

Energy:
$$\frac{dE}{dt} = -\frac{E}{\tau_E} + \underbrace{Q_{\alpha}S_{\alpha} - P_{rad} + P_{Ohm} + P_{aux} + P_{aux}^{burn}}_{P}$$
Alpha particles:
$$\frac{dn_{\alpha}}{dt} = -\frac{n_{\alpha}}{\tau_{\alpha}} + S_{\alpha}$$
Deuterium:
$$\frac{dn_D}{dt} = -\frac{n_D}{\tau_D} - S_{\alpha} + S_D^{rec} + S_D + S_D^{inj}$$
Tritium:
$$\frac{dn_T}{dt} = -\frac{n_T}{\tau_T} - S_{\alpha} + S_T^{rec} + S_T^{inj}$$
Impurities:
$$\frac{dn_I}{dt} = -\frac{n_I}{\tau_I} + S_I^{sp} + S_I^{inj}$$
 (actuators/disturbances in red/blue)
uasi-neutrality:
$$n_e = n_D + n_T + 2n_{\alpha} + Z_I n_I$$

Quasi-neutrality: $n_e = n_D + n_T + 2n_\alpha + Z_I n_I$ Density: $n = n_\alpha + n_D + n_T + n_I + n_e = 2n_D + 2n_T + 3n_\alpha + (Z_I + 1) n_I$ Temperature: $T = \frac{2}{3} \frac{E}{n} = \frac{2}{3} \frac{2}{2n_D + 2n_T + 3n_\alpha + (Z_I + 1)n_I}$ $(T \triangleq T_i = T_e \text{ (can be relaxed))}$

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• Reaction rate: $S_{\alpha} = n_D n_T \langle \sigma \nu \rangle \rightarrow S_{\alpha} = \gamma (1 - \gamma) n_{DT}^2 \langle \sigma \nu \rangle$. Tritium fraction: $\gamma \triangleq n_T / n_{DT}, n_{DT} = n_T + n_D$.

• The DT reactivity $\langle \sigma \nu \rangle$ is a highly nonlinear function of the plasma temperature calculated as

$$\langle \sigma \nu \rangle = exp\left(\frac{a}{T^r} + a_2 + a_3T + a_4T^2 + a_5T^3 + a_6T^4\right)$$

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• Confinement scaling (IPB98(y,2)): $\tau_E = 0.0562 H_H (I_{coils}^{non-axi}) I_p^{0.93} B_T^{0.15} P^{-0.69} n_{e19}^{0.41} M^{0.19} R^{1.97} \epsilon^{0.58} \kappa_{95}^{0.78}$.

• Particle confinement assumed proportional to τ_E , i.e. $\tau_{\alpha} = k_{\alpha}\tau_E$, $\tau_D = k_D\tau_E$, $\tau_T = k_T\tau_E$, $\tau_I = k_I\tau_E$.

• Impurity sputtering source: $S_I^{sp} = f_I^{sp} \left(\frac{n}{\tau_I} + \dot{n} \right) \quad (0 \le f_I^{sp} << 1) \Rightarrow n_I^{sp} = f_I^{sp} n \quad (n_I = n_I^{inj} + n_I^{sp}).$

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• Fuel recycling is included in the model through nonlinear functions S_D^{rec} and S_T^{rec} of the states.

- *P_{rad}*, *P_{Ohm}* are also highly nonlinear functions of the states (and external variables like *I*).
- Paux, S_D represent effect of actuators (NBI, RF H&CD) under other competing controllers.
- P_{aux}^{burn} , S_D^{inj} , S_T^{inj} , S_I^{inj} are the actuators available for burn control.

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$$S_{D}^{R} = \frac{f_{eff}}{1 - f_{ref} \left(1 - f_{eff}\right)} \left\{ f_{ref} \frac{n_{D}}{\tau_{D}} + \left(1 - \gamma^{PFC}\right) \times \left[\frac{\left(1 - f_{ref} \left(1 - f_{eff}\right)\right) R^{eff}}{1 - R^{eff} \left(1 - f_{eff}\right)} - f_{ref} \right] \left(\frac{n_{D}}{\tau_{D}} + \frac{n_{T}}{\tau_{T}} \right) \right\}$$

$$S_{T}^{R} = \frac{f_{eff}}{1 - f_{ref} \left(1 - f_{eff}\right)} \left\{ f_{ref} \frac{n_{T}}{\tau_{T}} + \gamma^{PFC} \times \left[\frac{\left(1 - f_{ref} \left(1 - f_{eff}\right)\right) R^{eff}}{1 - R^{eff} \left(1 - f_{eff}\right)} - f_{ref} \right] \left(\frac{n_{D}}{\tau_{D}} + \frac{n_{T}}{\tau_{T}} \right) \right\}$$



[†]Ehrenberg J. 1996 Physical Processes of the Interaction of Fusion Plasmas with Solids (New York: Academic) [1] M.D. Boyer and E. Schuster, Nuclear Fusion 55 (2015) 083021 (24pp).

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Isotopic Fuel Tailoring as Actuator for Burn Control

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Impact of Fueling-line Tritium Fractions on ITER's Operational Space

Summary of the Presentation

Burn Control Challenges



• Potential for thermal instability \rightarrow excursions and quenching:

$$P_f = n_D n_T < \sigma v > Q_{DT} \propto \beta^2 B^4, \quad \beta \propto rac{nkT}{rac{B^2}{2\mu_o}}$$

 Even when operating at stable equilibria, system performance during transients and plasma response to disturbances could be undesirable without active control.

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Isotopic Fuel Tailoring as Actuator for Burn Control

Burn Control Needs and Objectives

- Burn condition demands effective feedback control scheme to avoid:
 - Undesirable transient performance due to nonlinear/coupled dynamics
 - Perturbations due to plasma changes (confinement, impurity content)
 - Potentially disruptive plasma conditions due to thermal instabilities
- Nonlinear coupling with other control problems and objectives is severe
 - Confinement: PF coils (shape, current), Non-axisymmetric coils (RWM/ELM)
 - Heating/Density: Non-inductive heating & current drive (q-profile, NTM)
- Wall heat/particle load tolerance may impose constraints on core burn regulation
 - Requires controller that can effectively change operating conditions (Q, P_f, etc.)
- Capability of controller designed based on linearized model:
 - ✓ Regulation around a desired burning equilibrium point
 - \times Drive plasma from one operating point to another (Modify Q or P_f)
 - $\times\,$ Access to and exit from the burning plasma mode
 - $imes\,$ Handle nonlinear coupling with other competing controllers
- Reactor-specific additional challenges for effective burn control:
 - $P_{\alpha} >> P_{aux}$: control by heating may not be effective
 - Wall recycling effects may also make control by fueling not effective
 - Limited and noisy set of diagnostics

Burn Control Solution: Overview of Proposed Approach



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- Nonlinear dynamics
- Multiple inputs and outpus

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- The approach embeds whole nonlinear dynamics of burning plasma in controller by avoiding linearization of the model around operating point.
 - Preserving nonlinear dynamics is key to achieve controller's goals.
- The approach uses combination of actuators (SISO \rightarrow MIMO).

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Isotopic Fuel Tailoring as Actuator for Burn Control

Isotopic Fuel Tailoring

• Alpha power depends on tritium fraction

$$P_{\alpha} = n_D n_T \langle \sigma \nu \rangle Q_{\alpha} = \gamma \left(1 - \gamma\right) n_{DT}^2 \langle \sigma \nu \rangle Q_{\alpha}$$

where $n_{DT} \triangleq n_T + n_D$ and $\gamma \triangleq n_T/n_{DT}$.

- The function $\gamma (1 \gamma)$ achieves its maximum of 0.25 at $\gamma = 0.5$ and decreases steeply for smaller/larger γ 's.
- Fueling system in ITER (gass puffing (D) and pellet injection (D + DT)) will allow fuel mix regulation.
- Diagnostics for measuring the tritium ratio in both the edge and core plasma should be available in real time.
- Feedback control of the tritium ratio in ITER plasmas through isotopic fuel tailoring should be feasible.
- Effectiveness may be limited due to particle recycling.



No unique approach (depends on actuators and goals) but usually leads to tandem approach:

Step 1: Calculate value of P_{aux}^{burn} that stabilizes E^r subject to saturation.

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Step 2: Calculate stabilizing values of tritium fraction, γ^* , and/or energy confinement, τ_E^* .

• If the auxiliary power was not saturated, then $\gamma^* = \gamma^r$ and $\tau_E^* = \tau_E^{max}$.

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Step 3: Calculate value of non-axisymmetric magnetic fields to track τ_E^* .

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- Step 3: Calculate value of non-axisymmetric magnetic fields to track τ_E^* .
- **Step 4**: Calculate fuel injection rates that stabilize γ^* and n^r subject to saturation.
 - If fueling rates saturate, γ^* may not be tracked and an excursion could occur.

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- Step 5: If $P_{aux}^{burn} = (P_{aux}^{burn})^{max}$, $\tau_E^* = \tau_E^{max}$ and γ^* cannot be tracked \rightarrow thermal quench. • Change magnetic plasma parameters to improve energy confinement.

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- Step 6: If $P_{aux}^{burn} = (P_{aux}^{burn})^{min}$, $\tau_E^* = \tau_E^{min}$ and γ^* cannot be tracked \rightarrow thermal excursion. • Calculate a stabilizing value of impurity density, n_I^* , and the fueling/impurity injection rates needed to track n_I^* , γ^* and n^r .

[1] E. Schuster, M. Krstic and G. Tynan, Fusion Engineering and Design, 63-64, pp. 569-575, 2002.

[2] E. Schuster, M. Krstic and G. Tynan, Fusion Science and Technology, vol. 43, no. 1, 2003.

[3] M.D. Boyer and E. Schuster, Nuclear Fusion 55 (2015) 083021 (24pp).

[4] A. Pajares and E. Schuster, Fusion Engineering and Design 123 (2017) 607–611.

Potential of Nonlinear Control: Burn Performance



- Comparative study is carried out generating initial perturbations around the equilibrium for *T* and n_e keeping $f_{\alpha} = n_{\alpha}/n_e$ constant.
- While the boundaries shown for the linear controllers are absolute, for the nonlinear controller they only indicate the test limits.
- Embedding nonlinear dynamics in control synthesis ightarrow higher levels of performance!

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Potential of Nonlinear Control: Burn Robustness



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- Embedding nonlinear dynamics in control synthesis \rightarrow higher levels of robustness!

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Adding Robustness to Controller by Specific Design Techniques

How Do We Handle Uncertainties and Time-variations In Control-oriented Models?

- One of the main characteristic of feedback is its ability to deal with model uncertainties:
 - Poorly understood and/or unmodeled dynamics in response models
- Moreover, there are specific tools within the body of mathematical theory of control to specifically deal with model uncertainties:
 - Adaptive Control
 - Robust Control

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Stage 2: Adaptation by Real-time Estimation of Model Parameters



- Many of the burning plasma model parameters may be uncertain (highlighted in orange).
- The control algorithm must make use of estimated model parameters.
- Adaptive control is proposed to ensure tracking despite uncertainty.

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Stage 2: Adaptation by Real-time Estimation of Model Parameters

We define a system observer as

$$\begin{split} \dot{E}^{ob} &= -\hat{\theta}_{1} \frac{E}{\tau_{E}} + P_{\alpha} - P_{rad} + P_{aux} + P_{Ohm} - K_{E}^{ob} \left(E^{ob} - E \right) \\ \dot{n}_{\alpha}^{ob} &= -\hat{\theta}_{2} \frac{n_{\alpha}}{\tau_{E}} + S_{\alpha} - K_{\alpha}^{ob} \left(n_{\alpha}^{ob} - n_{\alpha} \right) \\ \dot{n}_{D}^{ob} &= -\hat{\theta}_{3} \frac{n_{D}}{\tau_{E}} + \hat{\theta}_{4} \frac{n_{T}}{\tau_{E}} - S_{\alpha} + S_{D}^{inj} - K_{D}^{ob} \left(n_{D}^{ob} - n_{D} \right) \\ \dot{n}_{T}^{ob} &= \hat{\theta}_{5} \frac{n_{D}}{\tau_{E}} - \hat{\theta}_{6} \frac{n_{T}}{\tau_{E}} - S_{\alpha} + S_{T}^{inj} - K_{T}^{ob} \left(n_{T}^{ob} - n_{T} \right) \\ \dot{n}_{I}^{ob} &= -\hat{\theta}_{7} \frac{n_{I}}{\tau_{E}} + S_{I}^{inj} + S_{I}^{sp} - K_{I}^{ob} \left(n_{I}^{ob} - n_{I} \right) \end{split}$$



The dynamics of the error $\tilde{\theta} = \theta - \hat{\theta}$ can be asymptotically stabilized by taking

$$\dot{\hat{\theta}} = -\frac{1}{\tau_E} \Gamma \begin{bmatrix} \tilde{n}_{\alpha}^{ob} n_{\alpha} & \tilde{E}^{ob} E & \tilde{n}_D^{ob} n_D & -\tilde{n}_D^{ob} n_T & -\tilde{n}_T^{ob} n_D & \tilde{n}_T^{ob} n_T & \tilde{n}_I^{ob} n_I \end{bmatrix}^T, \Gamma > 0$$

where

$$\tilde{n}_{\alpha}^{ob} = n_{\alpha}^{ob} - n_{\alpha}, \tilde{E}^{ob} = E^{ob} - E, \tilde{n}_{I}^{ob} = n_{I}^{ob} - n_{I}, \tilde{n}_{D}^{ob} = n_{D}^{ob} - n_{D}, \tilde{n}_{T}^{ob} = n_{T}^{ob} - n_{T}.$$

[1] M.D. Boyer and E. Schuster, Plasma Physics and Controlled Fusion 56 104004 (2014).

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Stage 3: Reference Governor by Real-time Optimization



- Part of burn control problem is selection of controller references.
- References must be chosen to optimize figure of merit for performance.
- Convex optimization is proposed to ensure optimal reference selection.

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Stage 3: Reference Governor by Real-time Optimization



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Stage 3: Reference Governor by Real-time Optimization



• The cost function is user-defined! More sophisticated optimization problems are possible!

[1] M.D. Boyer and E. Schuster, Plasma Physics and Controlled Fusion 56 104004 (2014).

- Conditions: $\gamma^r = 0.5$ is kept constant in this simulation and $w_{\gamma} \equiv 0$
- Optimization weights: $w_T = 0.1$ and $w_{P_{\alpha}} = 1$.
- References for fusion heating and temperature modified twice (at t = 60s and t = 120s)
- Constraints: $53MW < P_{aux} < 73MW$
- Recycling: $\gamma^{PFC} = 0.5$, $f_{eff} = 0.3$, $f_{ref} = 0.5$, $R_{eff} = 0.95 \Rightarrow \text{poor } \gamma \text{ control}$
- Simulation conditions chosen to ensure impurity injection is needed







Nonlinear controller and online optimization scheme were able to switch between desired operating points.





Tritium fraction was reduced in response to saturation of auxiliary power.





Due to saturation of other actuators, impurity injection is briefly used to track energy reference.

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Summary of the Presentation

Stage 4: State Estimation by Observer From Noisy/Limited Data



- Lower quality (e.g., noise, drifts, biases, etc.) of diagnostics.
- Critical issue in fusion reactors!

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Isotopic Fuel Tailoring as Actuator for Burn Control

Stage 4: State Estimation by Observer From Noisy/Limited Data



Stage 4: State Estimation by Observer From Noisy/Limited Data

We define an observer as



- We consider a general nonlinear output map $y = h(n_{\alpha}, E, n_I, n_D, n_T)$.
- The system is augmented with an additional state, \check{z} , governed by $\dot{\check{z}} = \mathring{y} y = \check{y}$.
- Lyapunov analysis \rightarrow injection terms L_E , L_α , L_D , L_T , L_I adopt a proportional-integral form.

[1] M. D. Boyer, E. Schuster, International Federation of Automatic Control World Congress (2014).

E. Schuster (LU Plama Control Group)

Kinetic (Burn) Control

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Summary of the Presentation

Stage 5: Optimal Allocation of Actuators with Dynamics

- Two-temperature model ($T_i \neq T_e$). Heating and fueling as actuation.
- Virtual control inputs \leftrightarrow Effector System \leftrightarrow Physical control inputs.



[1] V. Graber and E. Schuster, Nuclear Fusion 62 (2022) 026016 (18pp).

Stage 5: Optimal Allocation of Actuators with Dynamics

The Effector System maps the control efforts v to the actuator efforts u: $v = [P_{aux,i} P_{aux,e} S_D S_T]^T \iff u = [P_{ic} P_{ec} P_{nbi_1} P_{nbi_2} S_{D_{pel}} S_{DT_{eel}} S_{DT_{ees}}]^T$

$$P_{aux,i} = \eta_{ic} P_{ic} + \eta_{nbi_{1}} \phi_{nbi} P_{nbi_{1}} + \eta_{nbi_{2}} \phi_{nbi} P_{nbi_{2}}$$

$$P_{aux,e} = \eta_{ec} P_{ec} + \eta_{nbi_{1}} \phi_{nbi} P_{nbi_{1}} + \eta_{nbi_{2}} \phi_{nbi} P_{nbi_{2}} \qquad (\text{where } \phi_{nbi} = 1 - \phi_{nbi})$$

$$S_{D} = \eta_{nbi_{1}} \frac{P_{nbi_{1}}}{\varepsilon_{nbi_{0}}} + \eta_{nbi_{2}} \frac{P_{nbi_{2}}}{\varepsilon_{nbi_{0}}} + \eta_{pel_{1}} S_{D_{pel}} + \eta_{pel_{2}} (1 - \gamma_{pel}) S_{DT_{pel}} + \eta_{gas} (1 - \gamma_{gas}) S_{DT_{gas}}$$

$$S_{T} = \eta_{pel_{2}} \gamma_{pel} S_{DT_{pel}} + \eta_{gas} \gamma_{gas} S_{DT_{gas}}$$
Uncertain Parameters

- Ion cyclotron, electron cyclotron & NBI heating: Pic, Pec, Pnbi1, Pnbi2
- DT pellet & gas injection with Tritium fractions γ_{pel} & γ_{gas}: S<sub>D_{pel}, S<sub>DT_{pel}, S<sub>DT_{gas}
 </sub></sub></sub>
- Efficiency factors: η_{ic} , η_{ec} , η_{nbi_1} , η_{nbi_2} , η_{pel_1} , η_{pel_2} , η_{gas}
- Pellet fueling efficiency decreases with increasing plasma energy: $\eta_{pel_i} = \rho_{pel_i}(1-E/E_0), i \in \{1,2\}$
- The NBI ion-heating fraction $\phi_{nbi} = \rho_{nbi} \phi^{\star}_{nbi}$ contains uncertainty (ρ_{nbi}).
- NBI thermalization delay contains uncertainty: ρ_{th}

$$\tau_{nbi}^{lag} = \rho_{th} \tau_{nbi}^{\star} = -\rho_{th} \frac{2}{3B} \ln \left[\frac{\left(\frac{\varepsilon_{nbith}}{\varepsilon_{nbi0}}\right)^{3/2} + \left(\frac{\varepsilon_c}{\varepsilon_{nbi0}}\right)^{3/2}}{1 + \left(\frac{\varepsilon_c}{\varepsilon_{nbi0}}\right)^{3/2}} \right] \qquad (\varepsilon_{nbith} = T_i)$$

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Summary of the Presentation

Unmeasurable DT Content in Fueling Lines May Vary Over Time

- Two fueling lines will be available in ITER:
 - D injector (S^{line}): pellets of approximately 100% D
 - DT injector (S^{line}): pellets of approximately 10% D - 90% T
- These nominal concentrations in the fueling lines may not be sustained over long discharges
- Moreover, drifts and biases in the tritium fractions would not be measurable in real time
- Can we make the burn controller robust against these unmeasurable drifts and biases?



Adding Robustness to Controller by Specific Design Techniques

How Do We Handle Uncertainties and Time-variations In Control-oriented Models?

- One of the main characteristic of feedback is its ability to deal with model uncertainties:
 - Poorly understood and/or unmodeled dynamics in response models
- Moreover, there are specific tools within the body of mathematical theory of control to specifically deal with model uncertainties:
 - Adaptive Control
 - Robust Control

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Burn Control Scheme: Overview of Proposed Approach



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Burn Control Scheme: Overview of Robust Proposed Approach



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Injection rates for D and T can be written as

$$S_D^{inj} = (1 - \gamma_{DT}) S_{DT}^{line} + (1 - \gamma_D) S_D^{line}$$
$$S_T^{inj} = \gamma_{DT} S_{DT}^{line} + \gamma_D S_D^{line}$$

 $-\gamma_{DT} \in [0, 1], \gamma_D \in [0, 1]$ characterize the tritium concentration in the DT and D fueling lines.

- In the nominal case, $\gamma_{DT} = \gamma_{DT}^{nom} = 0.9$ and $\gamma_{D} = \gamma_{D}^{nom} = 0$.
- Unknown variations over time in the tritium fractions are modeled as

$$\gamma_{DT} = \gamma_{DT}^{nom} + \delta_{DT}, \ \ \gamma_D = \gamma_D^{nom} + \delta_D,$$

 $-\delta_{DT}$ and δ_D are "model uncertainties" in the tritium fractions.

• From definition, $\delta_{DT} \in [-0.9, 0.1]$, $\delta_D \in [0, 1] \Rightarrow$ bounded uncertainties.

[1] A. Pajares, E. Schuster, Nuclear Fusion 59 (2019) 096023 (18pp).

(1)

- Controller tries to regulate the system around a nominal equilibrium point defined by $\bar{T} = 12 keV$, $\bar{\gamma} = 0.4$ and $\bar{\beta}_N = 1.5$
- The system starts from a perturbed initial condition of +5% in *E*.
- Initial biases and time drifts in γ_{DT} and γ_D are simulated (not known by controller).





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4 Summary of the Presentation

Coupling Two-Chamber Model and Two-Point Model \rightarrow CSD Model



[1] V. Graber and E. Schuster, Fusion Engineering and Design 171 (2021) 112516.

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Core Chamber: Radial Temperature & Density Profiles

Peaked Electron and Ion Temperature Profiles

$$T_e(t,\psi) = (T_{e,0} - T_u)(1 - \psi/\psi_0)^2 + T_u$$

$$T_i(t,\psi) = (T_{i,0} - T_u)(1 - \psi/\psi_0)^2 + T_u$$

Central Ion Temperature: $T_{i,0}$ Central Electron Temperature: $T_{e,0}$ Upstream Separatrix Temperature: T_u

Flat Density Profiles

$$n_e(t,\psi)=n_{e,0}=n_u$$

Central Electron Density: $n_{e,0}$ Upstream Separatrix Density: n_u

Radial profiles couple core conditions (T_0 / n_0) with conditions at the separatrix (T_u / n_u)



toroidal magnetic flux coordinate

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Two-Point Model: Straightening Out the SOL Plasma

- Relates upstream (separatrix) and downstream (target) conditions[†]
- **Upstream**: density n_u , temperature T_u , parallel power flux density q_{\parallel}
- Downstream: density n_t, temperature T_t

Two Point Model Equations:

$$2n_{t}T_{t} = f_{mom}n_{u}T_{u}, \qquad T_{u}^{\frac{7}{2}} = T_{t}^{\frac{7}{2}} + \frac{7}{2}\frac{f_{cond}q_{\parallel}L}{\kappa_{0}}, \qquad (1 - f_{pow})q_{\parallel} = \gamma_{s}n_{t}T_{t}c_{st}$$

The Two-Point Model can be solved in terms of the **electron density** and **power entering the SOL** which are controllable with core-plasma actuators (pellet injection and auxiliary power).



† P.C. Stangeby, "The Plasma Boundary of Magnetic Fusion Devices," IoP Publishing, 2000.

POPCON (Plasma OPeration CONtour) Analysis

Generating POPCONs:

First, the Core-SOL-Divertor Model is solved in steady-state.

Then, operational constraints and plasma conditions can be plotted in temperature-density space.

Correction Factors for SOL:

- Conduction factor $f_{cond} = 1$ models inclusion of convection.
- Momentum loss factor $f_{mom} = 0.4$ models frictional collisions with neutrals.
- Power loss factor $f_{pow} = 0.8$ models radiation losses below the X-point.

ITER Operational Constraints:

- External Fueling Saturation
 - 100% D Pellet Injector
 - 90%T-10%D Pellet Injector
 - Tritium concentration in fueling lines can fall during long pulses
- Auxiliary Heating Saturation
 - Neutral beam injection
 - Radio frequency heating
- Power Threshold for H-mode
 - Total power must exceed threshold power
- Maximum Divertor Heat Load
 Maximum: 10 MW/m²
- Divertor Detachment ($T_t < 7eV$)

POPCON: ITER Plasma With Significant DT Recycling



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POPCON: ITER Plasma Without Any DT Recycling



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Isotopic Fueling Can Play Key Role in Burn-Control Schemes

- Burn-control solution: feedback controller, state observer, actuator allocator, reference governor.
 - Control techniques are <u>available</u> to incorporate nonlinear dynamics into the design o Performance \uparrow
 - Adaptive & robust control techniques are available to deal with model uncertainties ightarrow Robustness \uparrow
- Burn control by fueling becomes critical when control by heating becomes ineffective.
 - Density control by fueling is constrained by efficiency of actuators (gas puffing, pellet injection)
- Burn control designs can effectively incorporate isotopic fueling as an "actuator."
- Control by (isotopic) fueling may lose effectiveness as particle recycling increases.
- Burn controllers can be robustified against tritium-concentration drifts in fueling lines.
- Tritium fraction in fueling lines could impose limits on operational space in low recycling scenarios but drops of tritium fraction below nominal values are expected to be small.
- - Further work on actuator/diagnostic/transport modeling is needed
 - Need for multi-zone response model for control synthesis could be determined from simulation results
- Testing of proposed density-burn-control algorithms in present devices is needed.
 - Emulation of α heating, and even particle recycling, is possible through different mechanisms.
 - Emulation of ITER's actuators and diagnostics is also possible.