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On the possible injection schemes with the ITER SPI system

E. Nardon, A. Matsuyama, M. Lehnen,
D. Hu, M. Hoelzl, D. Bonfiglio,
F. Wieschollek, JOREK team

IAEA Technical Meeting on Plasma Disruptions and
their Mitigation, 20-23 July 2020 (Virtual Meeting)



- We will propose 2 « original » SPI schemes focused on the objective of avoiding RE beams in ITER
 - The compatibility of these schemes with other DMS objectives will not be addressed

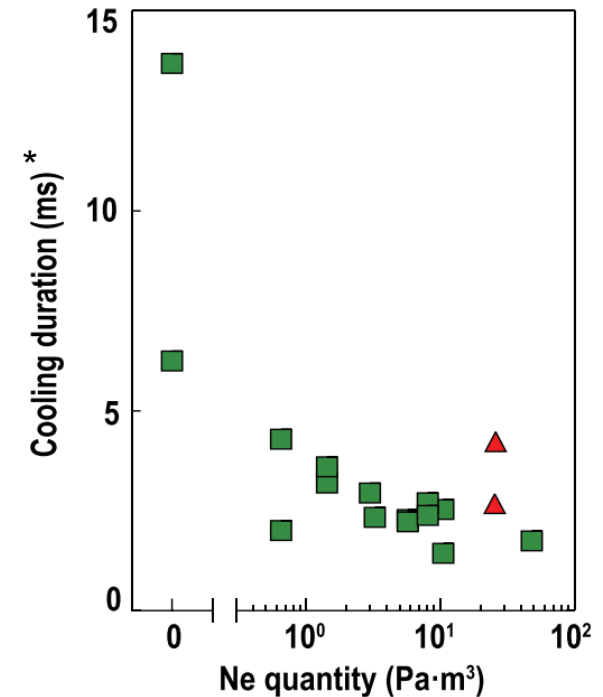
- 2 main reasons why RE generation could be large in ITER:
 - 1) Strong hot tail due to high T_e [Paz-Soldan 2020]
 - 2) Enormous avalanche gain + small continuous seeds (Tritium β decay, Compton scattering of γ 's from the activated wall)

- The current RE avoidance strategy is based on raising n_e by a large factor (~20-40) [Martín-Solís, NF 2017]
 - Hot tail generation assumed to be small
- Scheme 1: Pure H₂ SPI before TQ could be the best hope to raise n_e ...
 - ...and to avoid hot tail generation thanks to pre-TQ dilution cooling
 - Impurity SPI would follow to trigger the TQ and radiate the energy
- !! But the validity of the current RE avoidance strategy for the Baseline 15 MA DT scenario is **questioned by recent work** !! [T. Fülöp, this conf.]
- An alternative idea to deal with the small seed + large avalanche issue:
Scheme 2: Post-TQ SPI to continuously deplete RE seeds

**Pure H₂ (or D₂) SPI before the TQ
for fast plasma dilution**

- Pre-TQ (or 'cooling') phase is observed to be longer for pure D_2 SPI than for impurity SPI on several machines (DIII-D, JET, ...)

D_2+Ne SPI in DIII-D

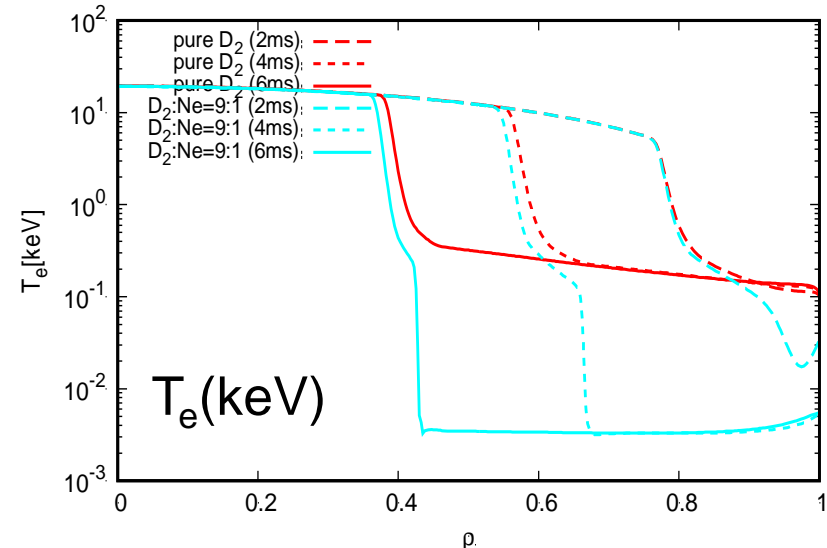
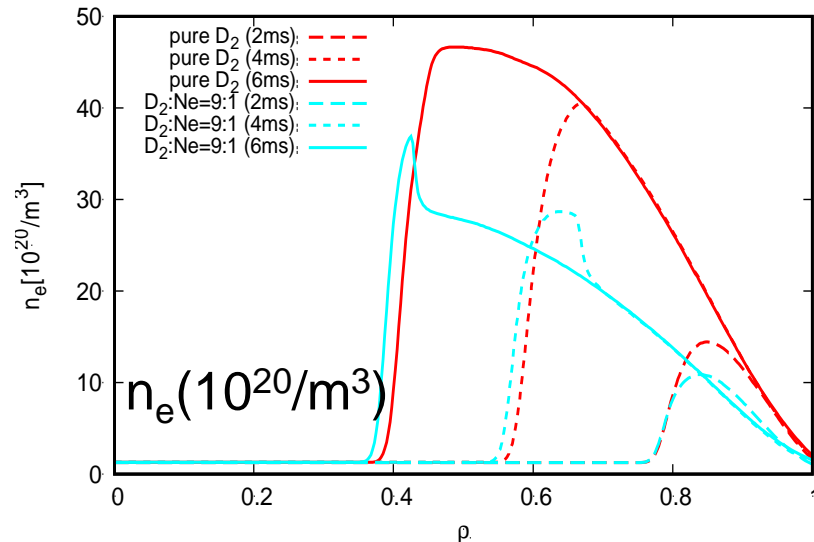


From D. Shiraki et al., Phys. Plasmas **23** 062516 (2016)

*Time between the arrival of initial pellet fragments and the I_p spike

Pure D₂ vs. D₂:Ne = 9:1

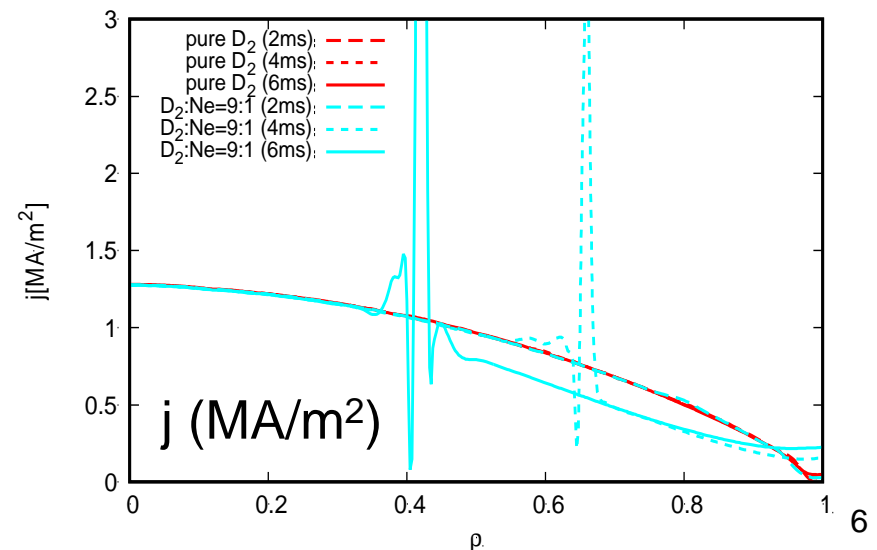
INDEX 1.5D simulations, ITER baseline 15 MA H-mode plasma
Single 28 mm pellet, $V_p=200$ m/s, $N_{shards} = 300$



With pure D₂ SPI:

Only dilution cooling (no radiative collapse)

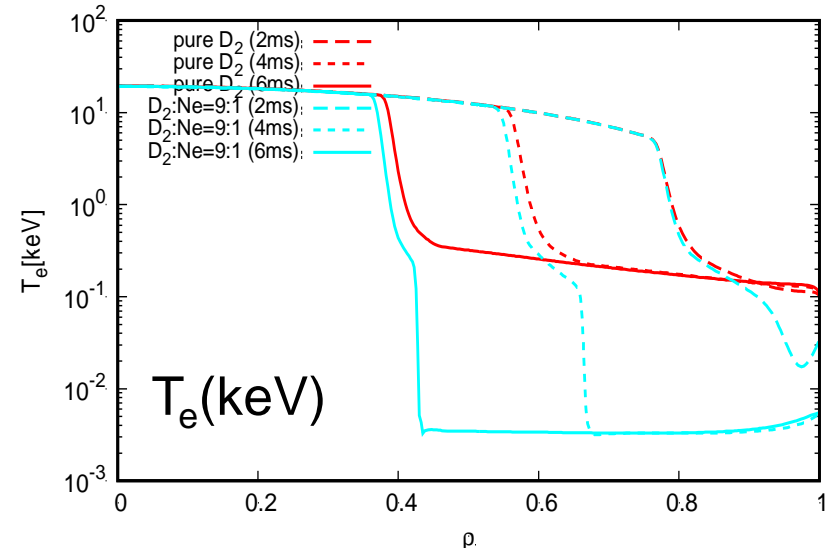
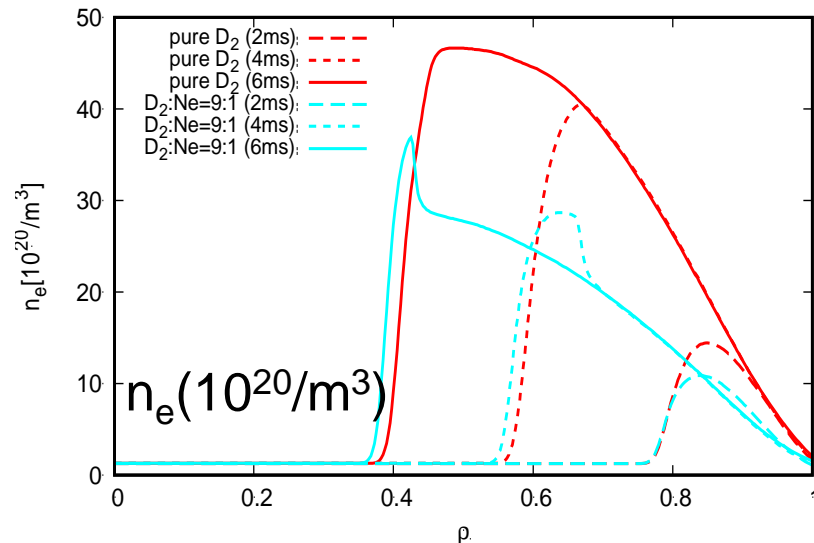
- ⇒ T_e remains > 100 eV
- ⇒ Resistive j decay time $\gg a/V_p$
- ⇒ No modification of j profile



Interpretation: cold front vs. 'lukewarm' front

Pure D₂ vs. D₂:Ne = 9:1

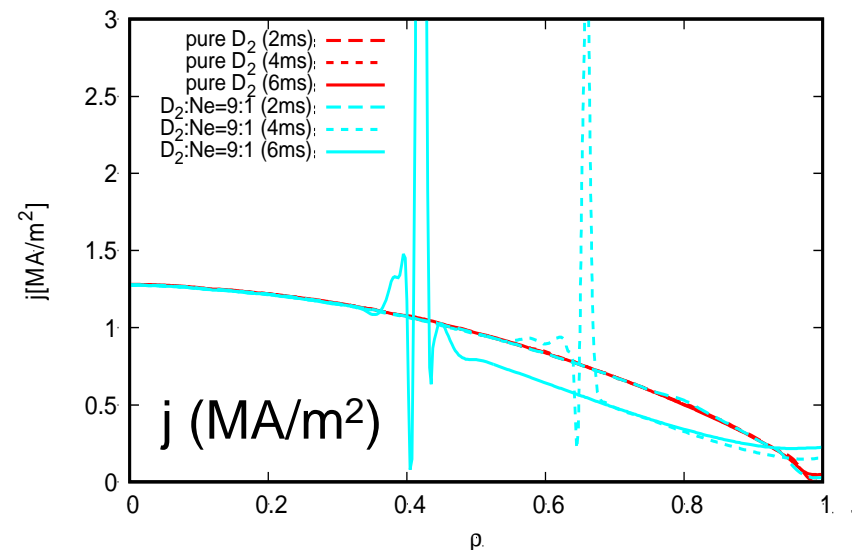
INDEX 1.5D simulations, ITER baseline 15 MA H-mode plasma
Single 28 mm pellet, V_p=200 m/s, N_{shards} = 300



With Ne+D₂ SPI:

Radiative collapse in cold front

- ⇒ T_e goes down to a few eV
- ⇒ Resistive j decay time < a/V_p
- ⇒ Modification of j profile
- ⇒ Likely to trigger an early TQ



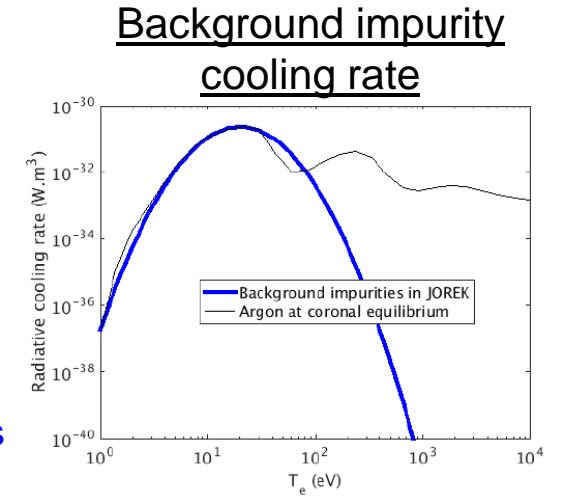
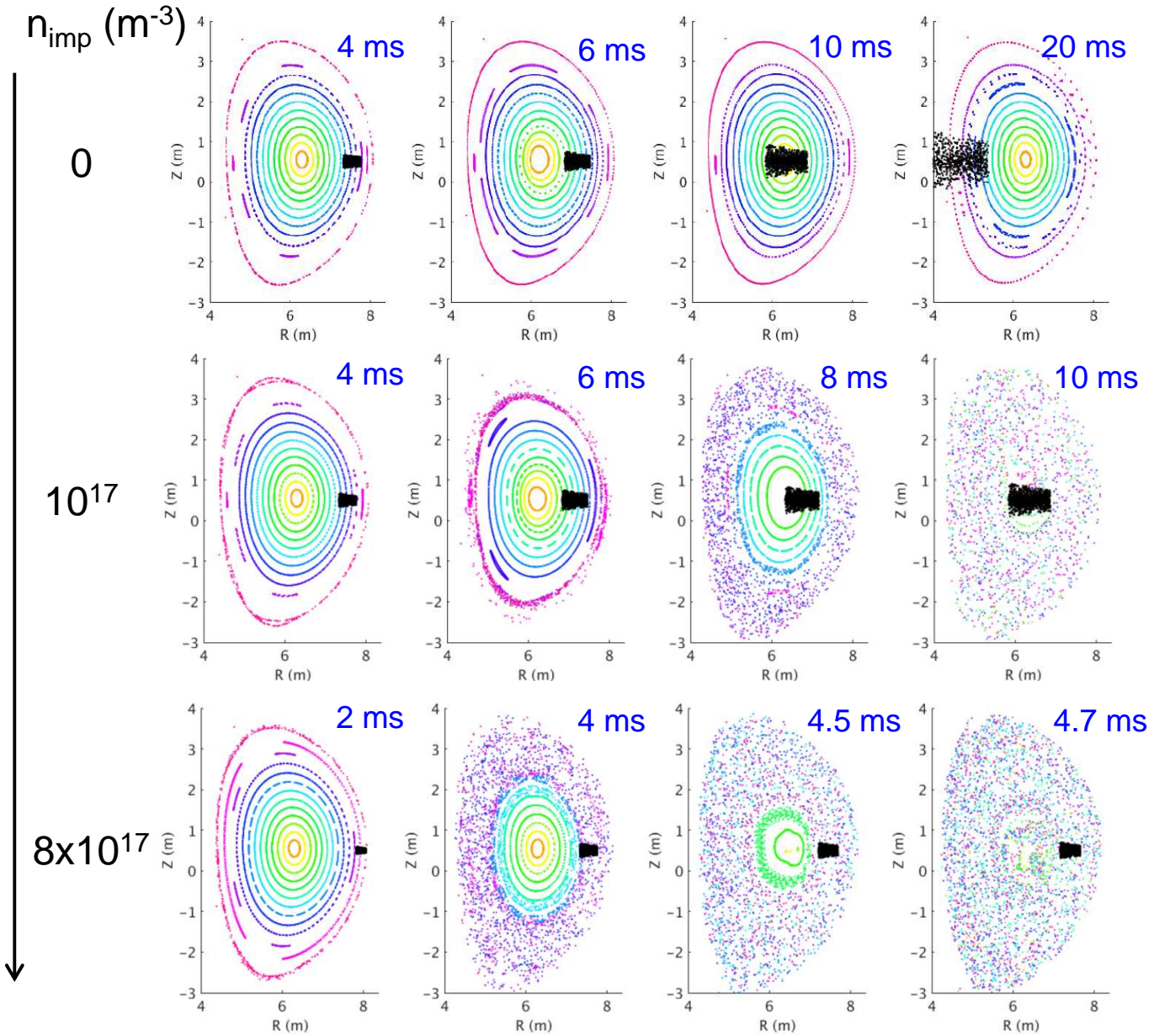
⇒ Experimental and early modelling results suggest that with pure H₂ (or D₂) SPI, shards may have the time to reach the plasma core and dilute the whole plasma before triggering a TQ

■ What about the effect of:

- Background impurities in the target plasma?
 - Lukewarm front → cold front?
- Toroidal localization of ablation clouds?
 - Enhances MHD modes?

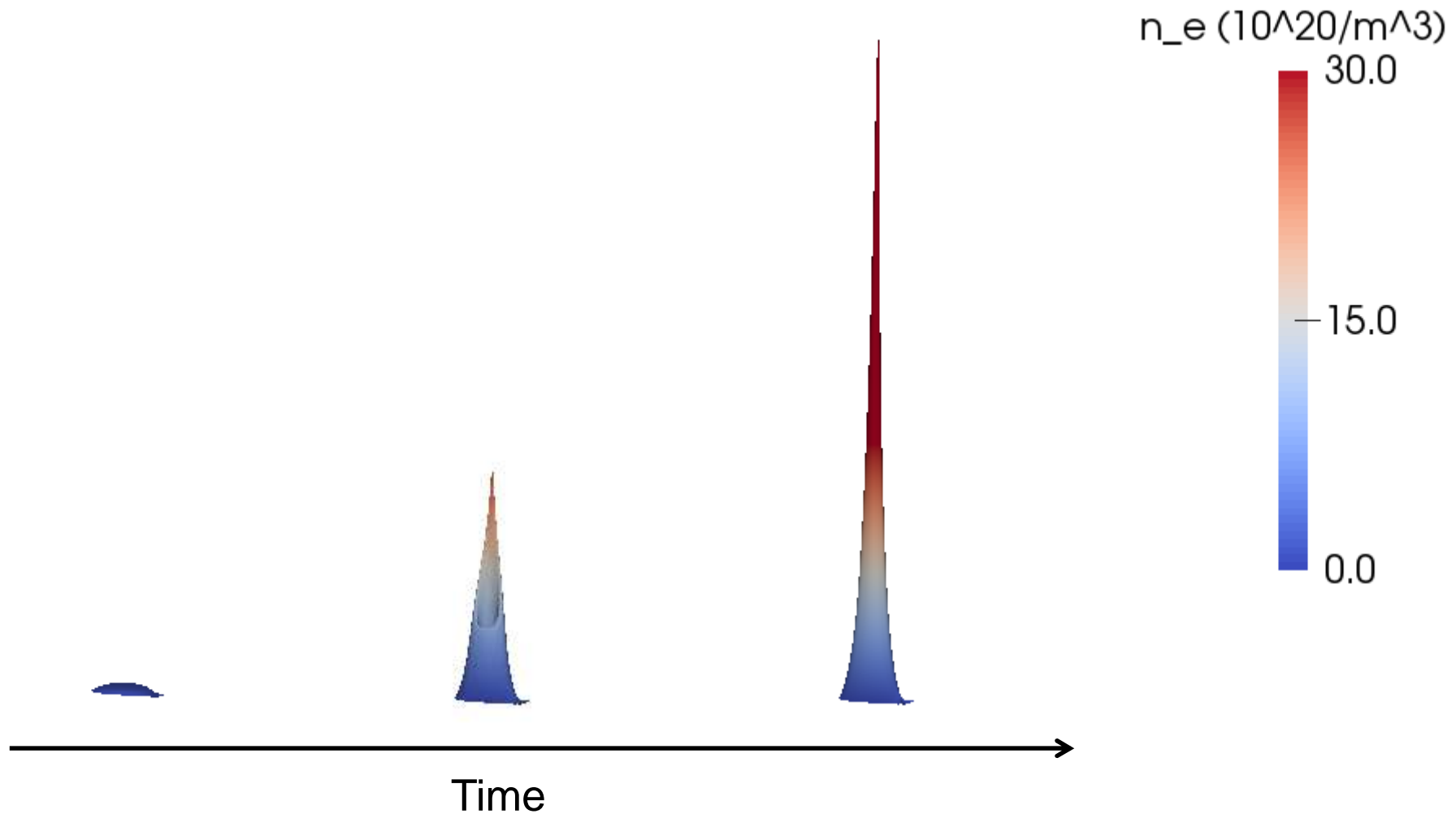
⇒ JOREK simulations

- Target plasma = 15 MA L-mode
- Injection of a 28 mm D₂ pellet with $V_p=200$ m/s, $N_{\text{shards}}=1000$

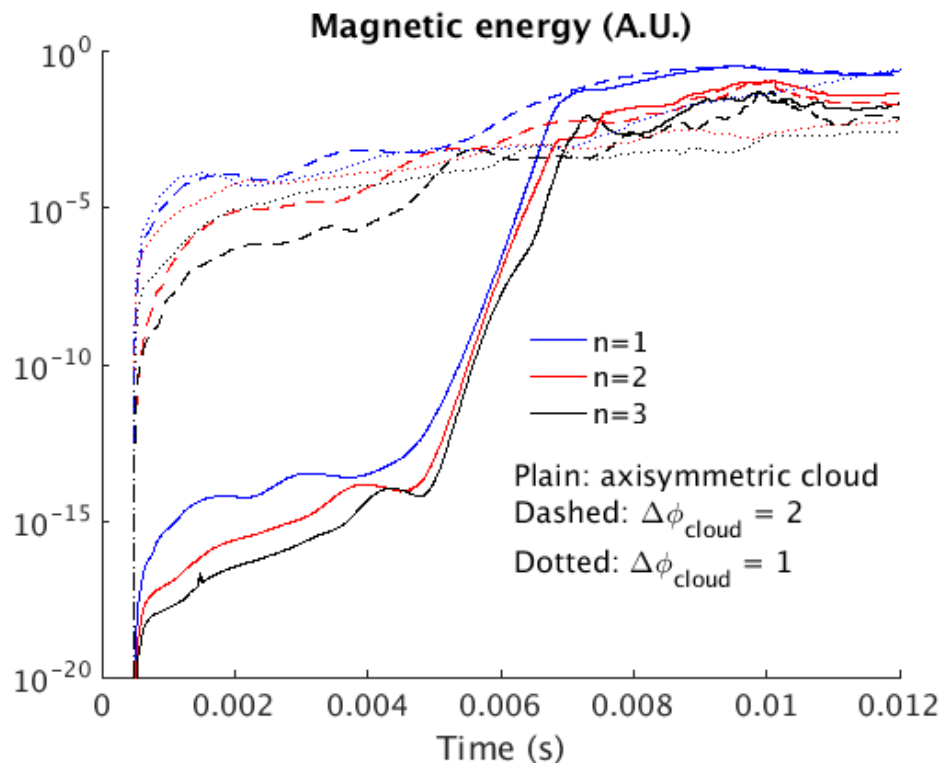


⇒ Motivates further studies with a more realistic description of background impurities

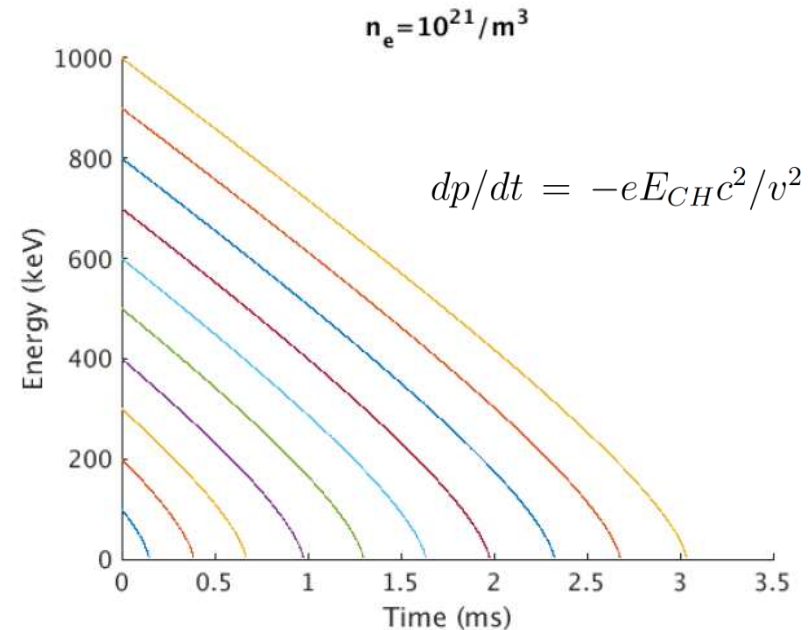
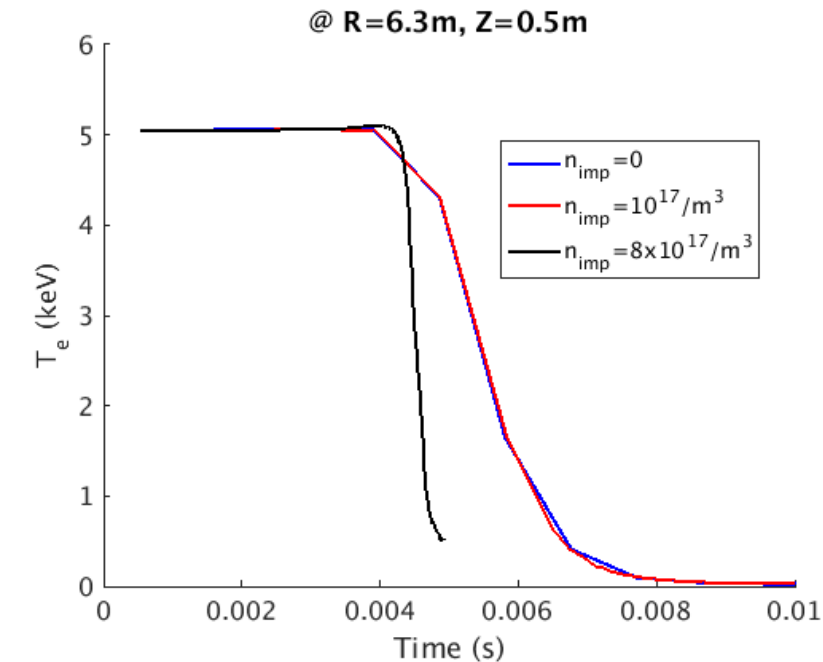
A dilution by a factor > 10 is obtained throughout the plasma



- Due to limited toroidal resolution in JOREK, ablation clouds toroidal extension $\Delta\phi_{\text{cloud}}$ is much larger than in reality
- A scan in $\Delta\phi_{\text{cloud}}$ shows that low-n mode excitation decreases with $\Delta\phi_{\text{cloud}}$
- Interpretation: low-n modes are primarily destabilized via a current profile modification, which is stronger for larger $\Delta\phi_{\text{cloud}}$
 - Direct mode drive by helical cooling does not compensate



⇒ JOREK predictions with artificially large $\Delta\phi_{\text{cloud}}$ probably overpredict low-n mode excitation, i.e. are « pessimistic »



- At low n_{imp} , dilution cooling in the core takes place on a **multi-ms timescale**, allowing the electron distribution to remain Maxwellian
 - Hot tail generation during the TQ is very sensitive to τ_e/τ_{TQ}
 - Electron collision time $\tau_e \sim T_e^{3/2}/n_e \sim D^{-5/2}$, where D = dilution factor
 - TQ duration (radiative collapse) $\tau_{TQ} \sim 1/n_e \sim 1/D$
- $\Rightarrow \tau_e/\tau_{TQ} \sim D^{-3/2}$, i.e. **dilution may dramatically reduce hot tail generation**

- JOREK simulations suggest that plasma can be diluted by a large factor (>10) all the way to the core without triggering a global MHD crash, provided n_{imp} is low enough
 - Best way to raise n_e (baseline RE avoidance strategy)
 - Dilution cooling may suppress the risk of hot tail RE generation

- For more detail: <http://arxiv.org/abs/2006.16020>

- Motivates more detailed studies, and in particular:
 - More accurate modelling of background impurities
 - Assessment of the effect of pre-existing islands
 - Validation against present experiments
 - Assessment of compatibility with other DMS objectives

**Post-TQ SPI
to deplete RE seeds**

Principle: REs travel across shards. When they do so, they lose an energy $\Delta E_{RE} = p\delta$, where p = stopping power, δ = distance travelled across shard.

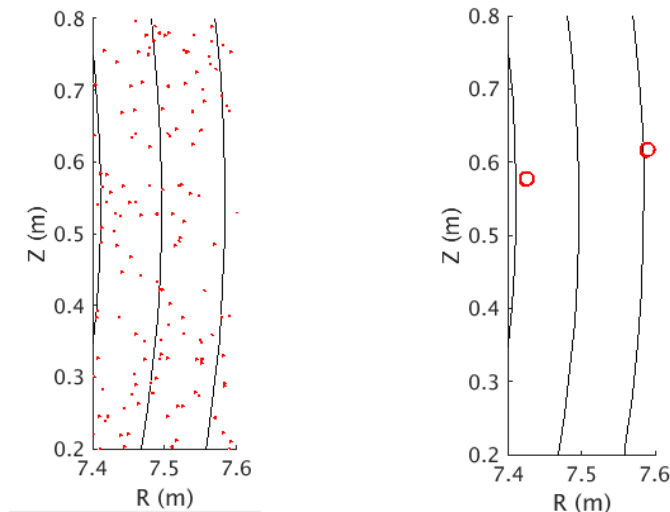
⇒ Try to use this effect to deplete small RE seeds **before they avalanche**

Since seeds are continuously produced, need **repeated or continuous injection**.

Required conditions for success:

- 1) Shards should be able to **stop RE seeds**
- 2) Shards should be able to **penetrate**
- 3) Additional assimilated material should **not make the CQ too short**

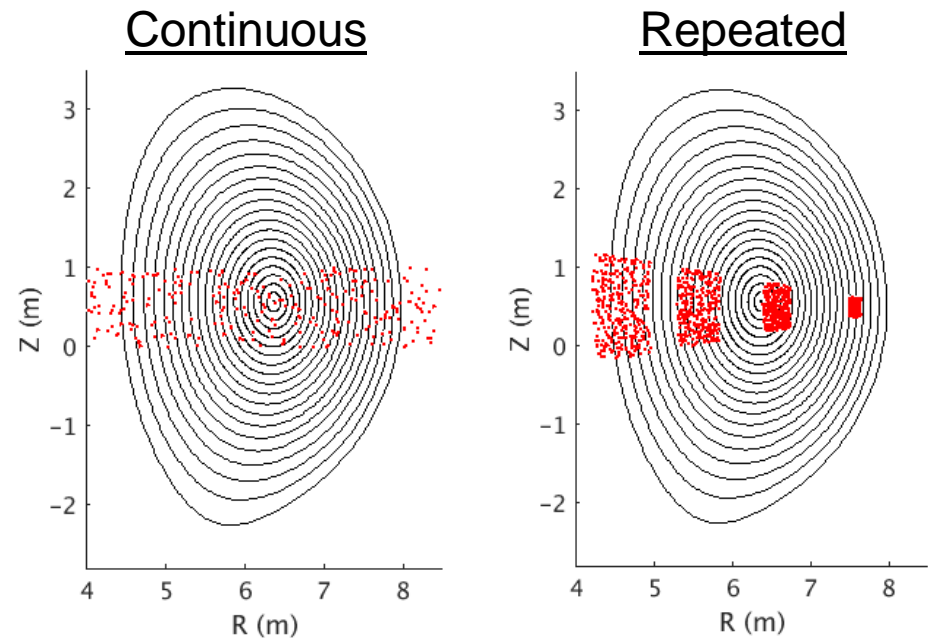
■ Many small shards or a few large shards?



■ Continuous passage of small shards vs. discrete passages of large shards

■ Continuous or repeated injection?

■ Note: « Few large shards » case assimilated to shard by shard injection at given freq.



■ Pellet material?

- Hydrogen → lower stopping power p but less risk to make CQ too fast
- Neon or Argon → larger p but more risk to make CQ too fast

- **Braking condition:** $E_{\text{gain/turn}}/E_{\text{loss/turn}} < 1$, where:

- $E_{\text{gain/turn}} = eV_{\text{loop}}$

- $E_{\text{loss/turn}} =$ energy lost per toroidal turn due to interaction with shards

- For many small shards: $E_{\text{loss/turn}} \simeq \frac{V_s p}{2\pi r_\psi} \frac{dN_s}{dr_\psi}$

- Hydrogen: need more than 5.3 large (28mm) pellet/meter

- Neon or Argon: need more than 0.9 pellet/meter

- For a few large shards: $E_{\text{loss/turn}} \simeq \frac{V_s p}{2\pi r_\psi} \frac{2}{\pi r_s}$

- Hydrogen: need $r_s > 8.3$ mm

- Neon or Argon: need $r_s > 3.4$ mm

($p_{\text{H}_2}=38\text{MeV/m}$, $p_{\text{Ne/Ar}}=230\text{MeV/m}$; Assumptions: $V_{\text{loop}}=1.1\text{kV}$, $r_\psi=1\text{m}$)

- For case of repeated injections with many small shards, each injection needs to fully stop REs, otherwise they would re-accelerate between 2 injections

- **Stopping condition:**
$$\frac{t_{stop}}{t_{pass}} \simeq \frac{4\pi^2 R r_\psi E_{RE}^0 \langle v_s \rangle}{c N_s V_s p} < 1$$

(assuming $E_{gain/turn}/E_{loss/turn} \ll 1$)

→ Hydrogen: need more than 5 large (28mm) pellet / injection

→ Neon or Argon: need more than 0.9 pellet / injection

(Assumptions: $R=6\text{m}$, $r_\psi=1\text{m}$, $E_{RE}^0=10\text{MeV}$, $\langle v_s \rangle=100\text{m/s}$)

- For case of a few large shards, each shard needs to fully stop REs, otherwise they would re-accelerate between 2 shards

- **Stopping condition:**
$$\frac{t_{stop}}{t_{pass}} \simeq \frac{\pi^3 R E_{RE}^0 r_\psi v_s}{c V_s p} < 1$$

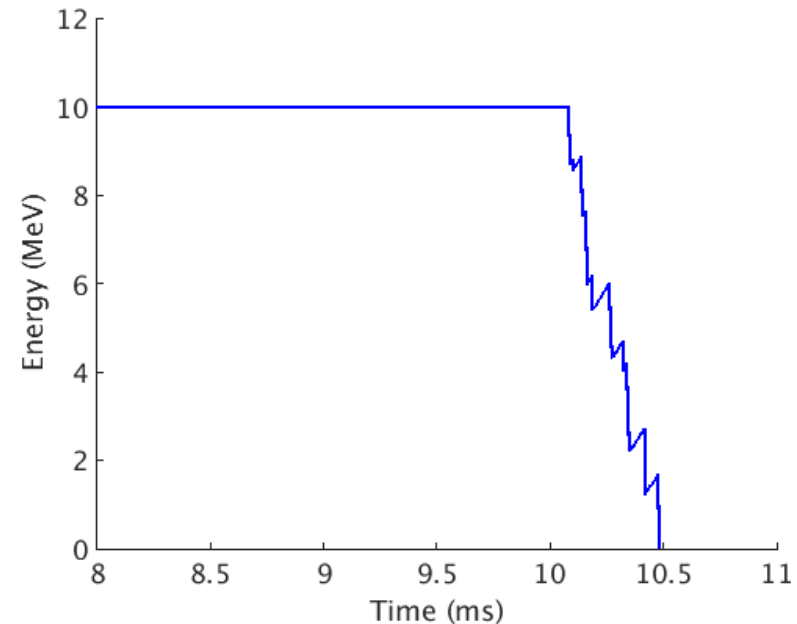
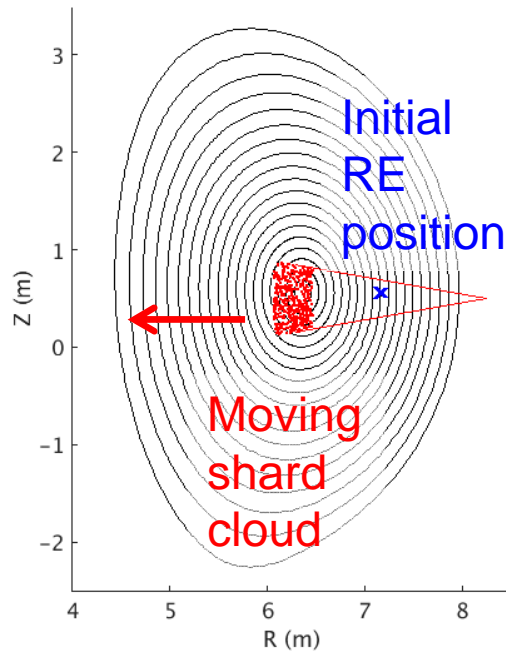
(assuming $E_{gain/turn}/E_{loss/turn} \ll 1$)

→ Hydrogen: need $r_s > 16$ mm

→ Neon or Argon: need $r_s > 8.6$ mm

⇒ More stringent than braking condition

(Assumptions: $R=6$ m, $r_\psi=1$ m, $E_{RE}^0=10$ MeV, $v_s=100$ m/s)

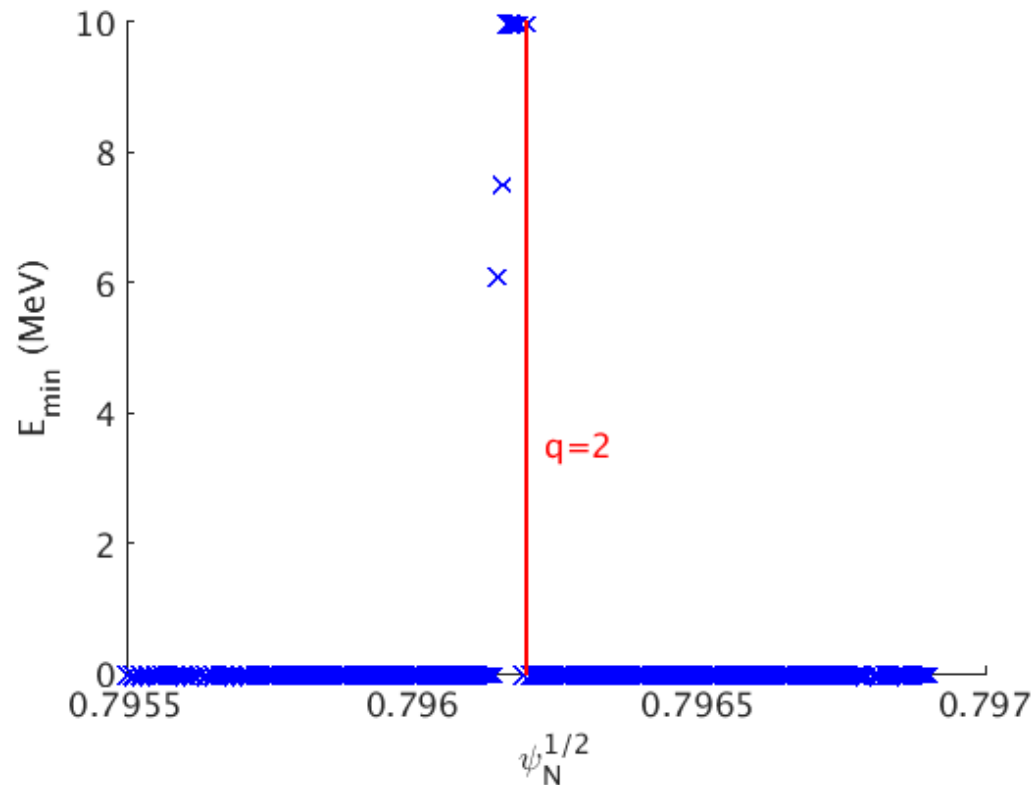


■ Model:

- REs assumed to follow field lines and tracked by steps of 1 toroidal turn
- RE energy evolved at each turn: $\Delta E_{RE} = eV_{loop} - p\bar{\delta}$, where $\bar{\delta}$ = distance travelled across shards (calculated from position of shards at current time)
 - E_{RE} saturated at 10 MeV

⇒ Trends from analytical estimates confirmed, but less optimistic results due to realistic equilibrium geometry

Effect of rational surfaces is very limited in the absence of islands



- REs perform a **large number of toroidal turns** over the passage of shards
- ⇒ **Magnetic shear** makes the effect of rational q surfaces very localized
- The presence of **islands** could be a concern
- Another possible concern is the **vertical plasma motion**

- When there are gaps between (clouds of) shards, the **avalanche during gaps** may lead to a large enough RE population to **vaporize shards**
- Energy into shard from REs (per unit volume):

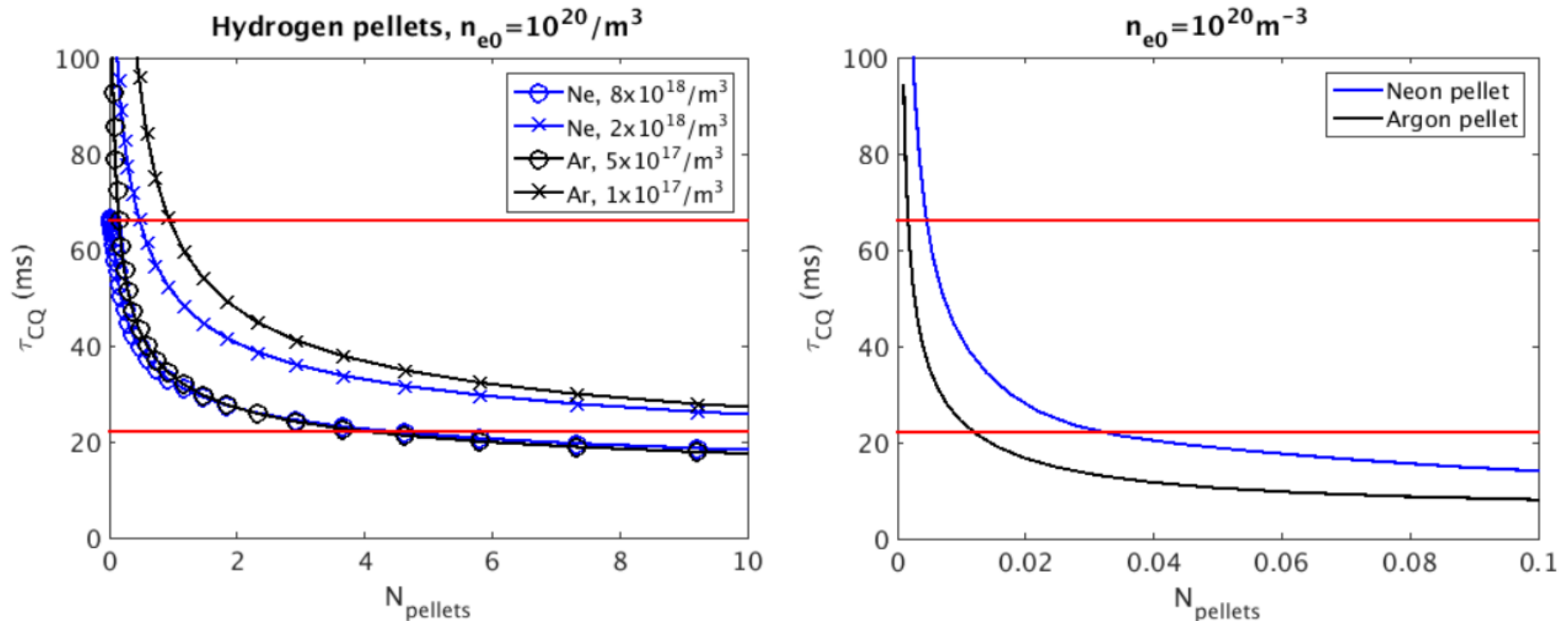
$$\exp(\Delta t_{\text{gap}}/\tau_{\text{e-fold}}) \times (j_{\text{RE}}^{\text{seed}}/e) \times p \times 2a/v_s$$

- Comparison to shard heat of vaporization → **Upper limit on Δt_{gap}** → Lower limit on injection rate

(Input numbers:

- $j_{\text{RE}}^{\text{seed}} \sim 0.1 \text{ A/m}^2$ [Boozer PPCF 2019; Martín-Solís NF 2017]
- $\tau_{\text{e-fold}} \sim 2 \text{ ms}$ ($\Delta\psi_{\text{e-fold}} = 2.3 \text{ V.s}$ [Boozer PPCF 2019] and $V_{\text{loop}} = 1.1 \text{ kV}$))

Exponential CQ timescale vs. amount of assimilated material



- Estimate based on a simple 0D model, assuming balance between Ohmic heating and radiative losses
- Note: H (or D) radiation not considered but could shorten CQ drastically when injecting several H pellets (recombination)...

Injection rate required to penetrate and stop RE seeds

(Unit = equivalent # of 28 mm pellets over 100 ms)

	H	Ne	Ar
Many small shards, continuous	60	9	9
Many small shards, repeated	100	18	18
A few large shards	30	6	6

Tolerable amount of assimilated material

(Unit = equivalent # of 28 mm pellets)

H	Ne	Ar
5-10	0.03	0.01

(Main assumptions: $E_{RE}^0=10\text{MeV}$, $v_s=100\text{m/s}$)

- Pellet number appears **prohibitive for H**, but **might be within DMS capabilities for Ne or Ar** (up to 24 pellets of 28 mm)
- **Ablation from thermal plasma** would have to be extremely small
 - Need a model for ablation in the CQ plasma ($T_e \sim 5\text{ eV}$)
 - Experiments on SPI during CQ in present machines would be useful
- Look for **better pellet material** (i.e. with large stopping power, poor radiation at low T_e , small ablation rate, ...)
- For more detail: [paper on arXiv](#) (lookup 'Nardon')