

# Interplay between particle transport, zonal flows and zonal density in Dissipative Trapped-Electron Mode turbulence

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# Outline

## 1 Introduction

## 2 Motivation

- Turbulent flux and transport crossphase
- Basic Mechanism of crossphase effects on the particle flux

## 3 Our Results: Analytical theory

- Wave kinetic approach including crossphase
- Schematic derivation of the model
- DTEM wave-kinetic model

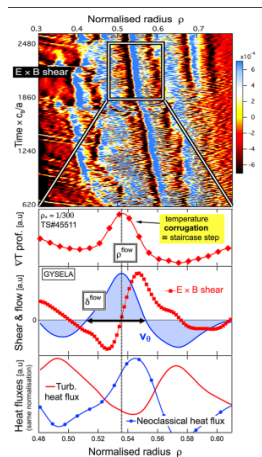
## 4 Our Results: Simulations of the Hasegawa-Wakatani model with BOUT++ (or DTEM for $f_t \rightarrow 1$ )

- zonal flows & zonal density corrugations
- Transport crossphase radial modulation

## 5 Summary and conclusions

# General Intro

- Transitions to enhanced confinement regimes → key for future devices (ITER)
- Recently,  $E \times B$  staircase observed in gyrokinetic simulations of ITG (GYSELA).
- Associated with **profile corrugations** in  $T_i$  and other quantities.
- Are profile corrugations universal, independent of the type of turbulence?
- What is the role of the **transport crossphase**?



ExB staircase [**Dif-Pradalier '15**]

## General Intro (ct'd)

- H-mode regime well studied both experimentally and theoretically  
→ L-H transition explained via 'flow-shear paradigm'  
**[Biglari Diamond Terry '90], [Moyer '95], [Diamond Itoh<sup>2</sup> Hahm '05]**: flow shear reduces transport by **suppressing** turbulence intensity.
- Crossphase may play a role in the L-H transition [**Kobayashi '17**].
- Zonal density effect on turbulent transport was observed in gyrokinetic simulations of Trapped-electron-Mode (TEM) turbulence [**Lang '08**] and in fluid simulations of resistive drift-wave turbulence [**An '17**].
- Here we identify a new mechanism for **self-organization of crossphase and zonal density corrugations**.
- In this work, we focus on the **particle transport** channel

# Turbulent flux and crossphase

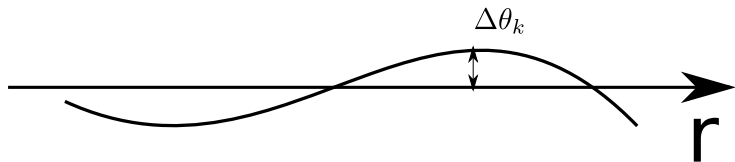
- the turbulent particle flux can be written:

$$\Gamma_{turb} = \sqrt{\langle n^2 \rangle} \sqrt{\langle \phi^2 \rangle} \sin \theta_{n\phi}$$

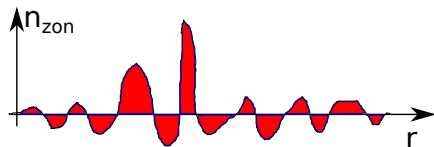
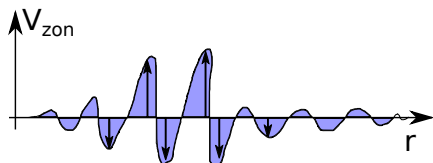
- with  $\theta_{n\phi}$  : transport crossphase
- the particle flux can be suppressed by:
  - suppressing the turbulence amplitudes  
 $\sqrt{\langle n^2 \rangle}, \sqrt{\langle \phi^2 \rangle} \searrow$  (eddy-shearing paradigm)
  - and/or
  - suppressing the crossphase  $\theta_{n\phi} \searrow$

## Basic Mechanism:

$E \times B$  nonlinearity  $\rightarrow$  transport crossphase modulation  
 $\rightarrow$  drives zonal density corrugations



modulation of **transport crossphase**  $\theta_k$  between density and potential fluctuations



Cartoon of zonal flows and zonal density

# Model

- fluid model for dissipative trapped-electron mode turbulence (with  $\frac{L_n}{L_{Te}} \rightarrow 0$ ) including **zonal modes**:<sup>1</sup>

$$\begin{aligned} \frac{\partial n}{\partial t} + v_E \cdot \nabla n + v_* \frac{\partial \phi}{\partial y} &= -\nu(\tilde{n} - \tilde{\phi}) \\ \frac{\partial}{\partial t} \left[ (1 - f_t) \tilde{\phi} - \nabla_{\perp}^2 \phi \right] - v_E \cdot \nabla \nabla_{\perp}^2 \phi + (1 - f_t) v_* \frac{\partial \phi}{\partial y} \\ &= f_t \nu (\tilde{n} - \tilde{\phi}) \end{aligned}$$

$n$  : effective density  
 $\phi$  : electric potential  
 $\nu = \nu_{ei}/f_t$  : de-trapping rate ( $\sim$  e-i collision frequency)

normalizations : space ( $\rho_S$ ), time: ( $L_n/c_S$ )

For  $f_t \rightarrow 1$  : reduces to the Hasegawa-Wakatani model with  $\frac{\nu_{ei}}{f_t}$

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<sup>1</sup>based on [Baver '02]

## Model (ct'd)

- This model has two nonlinearities:

polarization nonlinearity

$$v_E \cdot \nabla \nabla_{\perp}^2 \phi = \sum_{k=k'+k''} (k_{\perp}'^2 - k_{\perp}''^2) (\hat{\mathbf{z}} \times \mathbf{k}') \cdot \mathbf{k}'' \phi_{k'} \phi_{k''}$$

ExB convective nonlinearity

$$v_E \cdot \nabla n = \frac{1}{2} \sum_{k=k'+k''} (\hat{\mathbf{z}} \times \mathbf{k}') \cdot \mathbf{k}'' (n_{k'} \phi_{k''} - \phi_{k'} n_{k''})$$



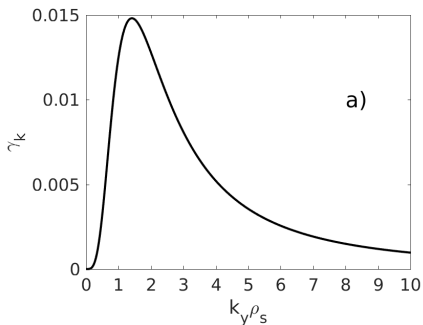


Figure: growth-rate v.s.  $k_y$

dispersion relation

$$\omega \left[ 1 + \frac{k_{\perp}^2}{1 + i\omega/\nu} \right] - \omega_* = 0$$

DTEM frequency  $\omega \ll \nu$

$$\omega_k = \omega_* / (1 + k_{\perp}^2)$$

growth-rate

$$\gamma_k \simeq \frac{1}{\nu} \frac{k_{\perp}^2 \omega_*^2}{(1 + k_{\perp}^2)^3}$$

with  $\omega_* = k_y v_*$

## Zonal flows & zonal density

- zonal flows  $V_{zon}(x, t) = \langle v_E \rangle$   
& zonal density  $n_{zon}(x, t) = \langle n \rangle$   
are nonlinearly driven by DTEM turbulence:

$$\frac{\partial V_{zon}}{\partial t} = -\frac{\partial}{\partial x} \sum_k k_r k_\theta |\phi_k|^2$$
$$\frac{\partial n_{zon}}{\partial t} = -\frac{\partial}{\partial x} \sum_k k_\theta \frac{|n_k|}{|\phi_k|} |\phi_k|^2 \sin \theta_k$$

$V_{zon}$	:	zonal flows
$n_{zon}$	:	zonal density
$ \phi_k ,  n_k $	:	amplitudes
$\theta_k$	:	crossphase

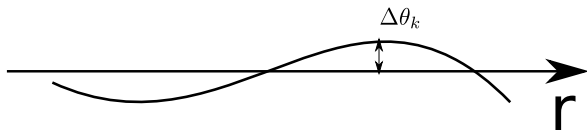
- - zonal flows driven by the polarization nonlinearity
- zonal density driven by  $E \times B$  nonlinearity

# Wave kinetic approach

- The turbulent particle flux at wavenumber  $k$  can be written:

$$\Gamma_k \propto W_k \sin \theta_k$$

- with  $W_k = (1 + k^2)|\phi_k|^2/\omega_k$  the wave action density (Wigner function) for DW.
- Define  $W_k = \langle W_k \rangle + \Delta W_k(r, t)$  and  $\theta_k = \theta_k^0 + \Delta\theta_k(r, t)$
- $\Delta\theta_k(r, t)$  : **crossphase modulation**



## Wave kinetic approach (ct'd)

### Wave Kinetic equation (WKE) [Malkov, Diamond, Rosenbluth '01]

$$\frac{\partial W_k}{\partial t} + \frac{\partial \omega}{\partial k_r} \frac{\partial W_k}{\partial r} - \frac{\partial \omega}{\partial r} \frac{\partial W_k}{\partial k_r} = \hat{\gamma}_{NL} W_k$$

- with the non-linear growth-rate operator satisfying:

$$\hat{\gamma}_{NL} W_k = 2\gamma_k W_k - \Delta\omega W_k^2$$

- We extend the wave kinetic equation to include the dependence of growth-rate on **crossphase**

$$\gamma_k \simeq (\theta_k^0 + \Delta\theta_k)\omega_k, \quad (\text{valid for } |\theta_k^0|, \Delta\theta_k \ll 1),$$

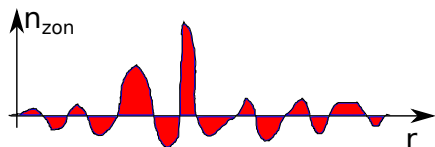
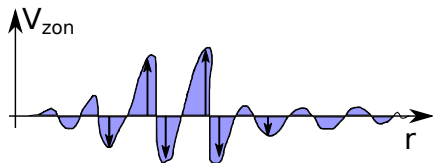
- crossphase modulation  $\rightarrow$  growth-rate modulation

# Wave Kinetic Equation

- and we obtain the Extended Wave-Kinetic Equation (EWKE):

$$\frac{\partial W_k}{\partial t} + v_{gr} \frac{\partial N_k}{\partial k_r} - k_\theta \nabla_r U \frac{\partial W_k}{\partial k_r} = 2\gamma_k^0 W_k - 2c_k W_k \nabla_r N$$

- 1st term on RHS: linear growth  $\gamma_k^0 = \theta_k^0 \omega_k$
- 2nd term: **nonlinear stabilization** due to **radial modulation of the crossphase**  $\Delta\theta_k \propto -\nabla_r N$  [Leconte & Kobayashi '21]
- What causes this radial modulation?  
↪ zonal density corrugations (as shown in the following)



# Transport crossphase modulation

- consider the weakly-nonlinear density response [Zhou, Zhu & Dodin '19]:

## Weakly-nonlinear Density Response

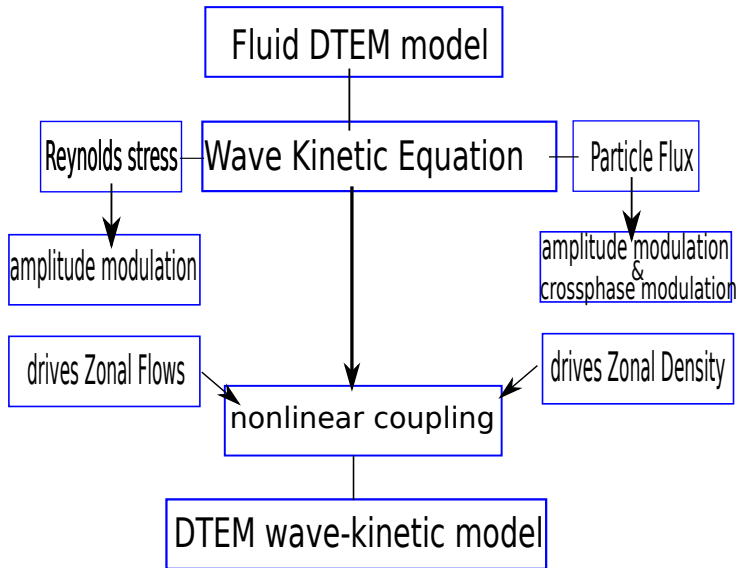
$$\frac{n_k}{\phi_k} \simeq 1 - i\theta_k$$

- $\theta_k \rightarrow \theta_k^0$  in the linear regime
- with  $\theta_k^0 = \nu^{-1}[k_{\perp}^2 \omega_*/(1 + k_{\perp}^2)]$  the linear crossphase
- Writing the nonlinear crossphase:  $\theta_k = \theta_k^0 + \Delta\theta_k$ , we have:

crossphase modulation similar to [Holland & Diamond '01] (ETG)

$$\Delta\theta_k = -\frac{\partial\theta_k^0}{\partial\omega_*} k_{\theta} \nabla_r N$$

# Schematic derivation of the model



# DTEM wave-kinetic model

- Wave-kinetic analysis leads to the DTEM wave-kinetic model :

$$\begin{aligned}\frac{\partial W_k}{\partial t} + \frac{\partial \omega_k}{\partial k_r} \nabla_r W_k - k_\theta \nabla_r U \frac{\partial W_k}{\partial k_r} &= 2\gamma W_k - 2c_1 W_k \nabla_r N \\ \frac{\partial U}{\partial t} &= -c_2 \nabla_r \int W_k dk_r - \mu U \\ \frac{\partial N}{\partial t} &= -c_3 \nabla_r \int W_k dk_r + c_4 \nabla_r \left[ \int W_k dk_r \nabla_r N \right] + D_0 \nabla_r^2 N\end{aligned}$$

$W_k$  : turbulent wave action density

$U$  : zonal flows

$N$  : zonal density

- 2 'predators': zonal flows and zonal density corrugations
- numerical implementation of wave-kinetic model left for future work



# Results from direct fluid turbulence simulations

Hereafter, we present results that support the theory, using **direct numerical simulations** of the fluid **modified Hasegawa-Wakatani model** using BOUT++ framework [**Dudson '09**]

- fluid model for collisional drift-waves (modified Hasegawa-Wakatani) including **zonal modes**
- : limit of DTEM for  $f_t \rightarrow 1$  (strongly unstable DTEM)

$$\frac{\partial n}{\partial t} + v_E \cdot \nabla n + \kappa \frac{\partial \phi}{\partial y} = \nu(\tilde{\phi} - \tilde{n})$$
$$\frac{\partial}{\partial t} \nabla_{\perp}^2 \phi + v_E \cdot \nabla \nabla_{\perp}^2 \phi = \nu(\tilde{\phi} - \tilde{n})$$

# Results: zonal flows & zonal density corrugations

- 2D BOUT++ Simulations on a grid of size  $[L_x \cdot L_y] = 256 \times 256$ ,
  - constant equilibrium density gradient  $\kappa = 0.5$  ( $L_n = 2\rho_s$ ),
  - de-trapping rate  $\nu = 1$ .

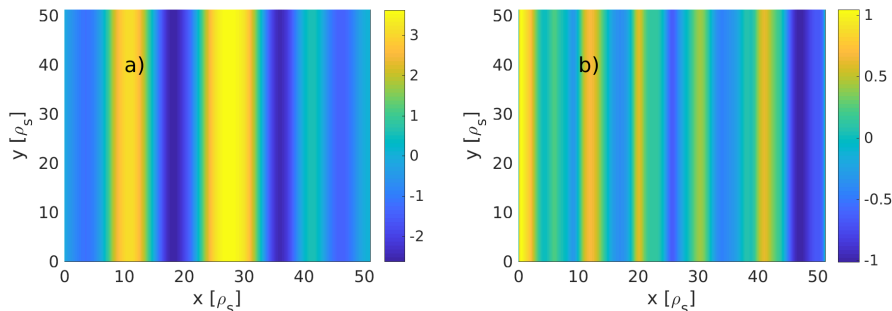
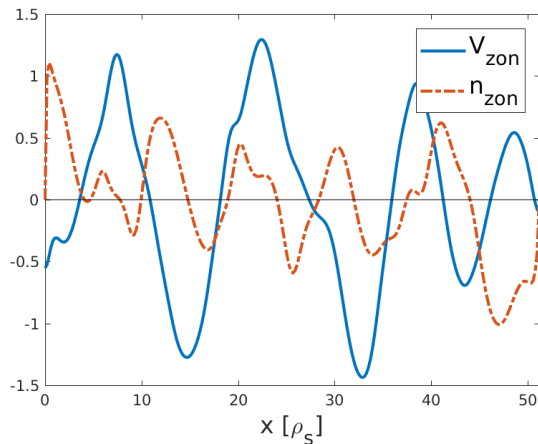


Figure: Snapshot of the simulations in saturated state:  $\phi$ ,  $n$

- Contours of electric potential  $\phi \rightarrow$  **zonal flows**
- Contours of density  $n \rightarrow$  **zonal density corrugations**

# Results: Evidence of staircase



- zonal density corrugations  $n_{zon}$  have different radial scale compared to zonal flows  $V_{zon}$  !  $\rightarrow$  suppress different scales?

# Results: Transport crossphase radial modulation

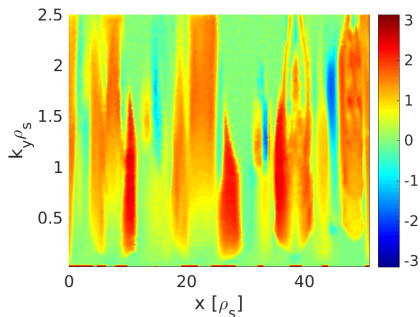


Figure: Time-averaged **nonlinear crossphase**  $\theta_k(x)$  v.s. radius and  $k_y \rho_s$

- turbulence spectrum divided into the  $|\phi_k|^2$ ,  $|n_k|^2$  power spectrum, and cross-spectrum
- cross-spectrum  $\rightarrow n_k^* \phi_k$  is complex-valued

crossphase spectrum (crossphase)

$$\theta_k = \arg n_k^* \phi_k$$

- Time-averaged **nonlinear crossphase**  $\theta_k(x)$  shows **radial modulations**
- This result supports the theory proposed in this talk

# How efficient is zonal density at suppressing transport?

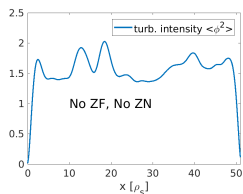


Figure: Turbulence intensity profile

No ZF, No ZN

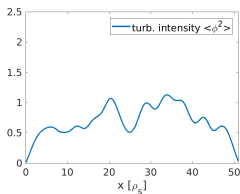


Figure: Turbulence intensity profile

ZN only

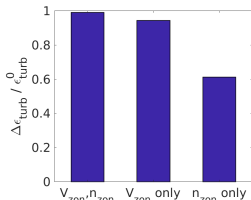


Figure: Zonal efficacy for 3 cases

zonal efficacy

$$\chi = \frac{|I - I_0|}{I_0}$$

↪ Zonal density has a significant effect on turbulence suppression

## Summary and conclusions

- In the framework of a fluid DTEM model, we showed that zonal density corrugations can be nonlinearly driven by the DTEM turbulence
- The wave kinetic analysis predicts that DTEM turbulence can nonlinearly drive zonal density corrugations that can suppress the transport crossphase by radially modulating it, thus reinforcing the zonal density corrugations.
- An extended wave-kinetic equation (EWKE) is derived. Nonlinear coupling to zonal flows and zonal density yields a novel DTEM wave kinetic model.
- Direct numerical simulations of the related Hasegawa-Wakatani model confirm that crossphase modulation can drive zonal density, and zonal density stabilizes turbulence.
- Open Questions
  - Can zonal density corrugations be observed experimentally?
  - Are radial modulations of crossphase observed in nonlinear gyrokinetic simulations? (ongoing work)
  - work in progress: extension to collisionless modes (CTEM)

# THANK YOU

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