

# Global calculation of neoclassical impurity transport

## including the variation of electrostatic potential

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### ABSTRACT

- We investigate neoclassical impurity transport in an impurity hole plasma using a global neoclassical simulation code FORTEC-3D.
- By the global simulation, we find the ambipolar radial electric field ( $E_r$ ) which changes the sign from negative to positive along the minor radius.
- With such an  $E_r$ , the radial carbon flux ( $\Gamma_C$ ) can be outwardly directed even where  $E_r < 0$  and the carbon density profile is hollow.
- The sign-changing property of  $E_r$  in an impurity hole plasma has been experimentally confirmed and our result is qualitatively consistent with it.

### BACKGROUND

- Since the presence of impurity ions in the core plasma can degrade the performance of fusion reactor, understanding the impurity transport process to prevent the impurity accumulation is a crucial task.
- According to neoclassical theory and experiments, impurity ions are usually expected to accumulate in the core of stellarators when  $E_r < 0$ .
- However, a notable exception, the formation of "impurity hole", is observed in LHD (Large Helical Device).
- An impurity hole is a hollow structure of an impurity density profile in the core where  $E_r < 0$ .
- An experiment shows that, in an impurity hole plasma,  $E_r < 0$  at  $r/a < 0.55$  but  $E_r > 0$  at the outer region [1],

$r$ : radial coordinate  
 $a$ : minor radius

[1] T. Ido et al., *Plasma Phys. and Control. Fusion*, 52(12), 2010.

### METHODS AND SETUP

#### THE PLASMA PROFILE

- We investigate three different cases, **A**, **B** and **C**, each corresponding to a different carbon density ( $n_C$ ) profile (Fig.1).
- For case **A**, the  $n$ - $T$  profiles, including the  $n_C$  profile, are the same as those used in previous studies[2,3]: the hollow structure of the  $n_C$  profile is formed at an off-axis region ( $r/a > 0$ ).
- However, some measurement data indicates impurity holes are formed around the magnetic axis.
- For case **B**, the carbon density gradient  $\nabla n_C$  near the axis is thus flattened.
- For case **C**, on the other hand,  $\nabla n_C$  near the axis is steepen.

#### SELF-CONSISTENT CALCULATION OF $\Phi_1$ , $E_r$ and $\Gamma_I$

- We solve the quasi-neutrality condition to evaluate the variation of electrostatic potential on the flux surface  $\Phi_1$ , the ambipolar condition to determine the ambipolar  $E_r$  and a drift-kinetic equation for multiple ion species including  $\Phi_1$  simultaneously.
- Local  $\Gamma_e$  is used to determine  $E_r$  and  $\delta f_e = 0$  is assumed to evaluate  $\Phi_1$ .

[2] A. Mollén, et al., *Plasma Phys. and Control. Fusion*, 60(8), 2018.

[3] M. Nunami, et al., *Phys. of Plasmas*, 27(5), 2020.

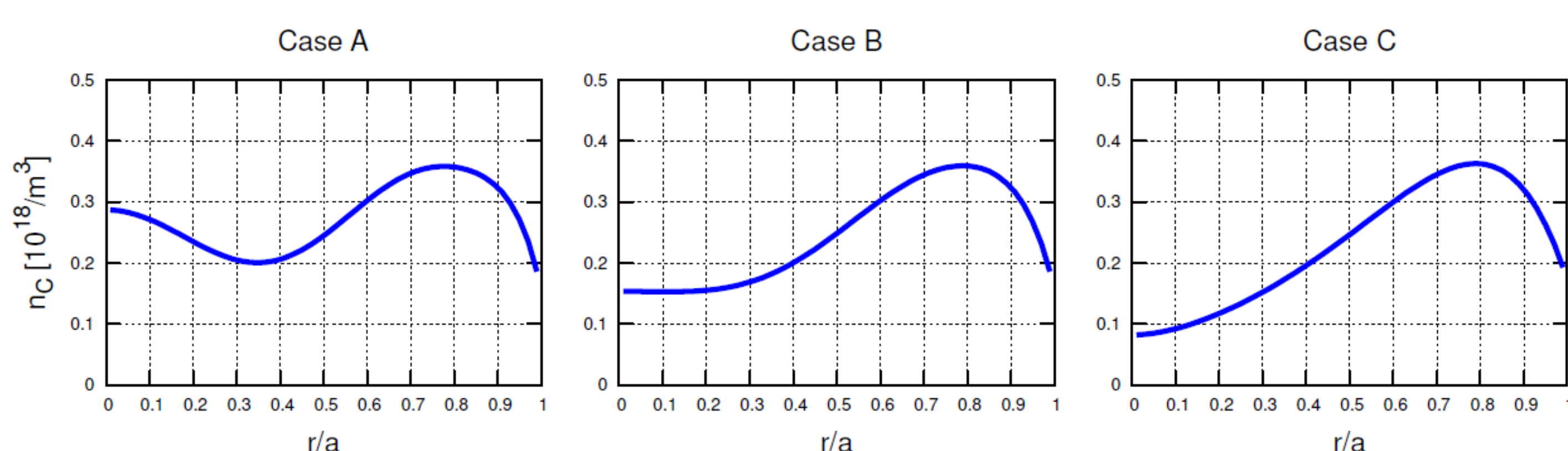


Fig.1 The radial profiles of carbon density for each case.

### OUTCOME

#### THE AMBIPOLAR $E_r$

- For case **A**, the sign of the local solution (cyan points ●) of the ambipolar condition is negative for almost the entire radius.
- The sign of the global solutions (red line — for the case w/ $\Phi_1$  and green lines — for the case w/o  $\Phi_1$ ) is negative near the axis but transits to positive around  $r/a = 0.25$  (Fig.2).
- The local calculation also finds an electron root (blue points ●) but only partially at  $r/a > 0.8$ .
- Qualitatively similar solutions are obtained for case **B** and **C**

#### THE CARBON FLUX $\Gamma_C$

##### ► Case **A**

- Regardless of the effect of  $\Phi_1$ ,  $\Gamma_C$  is outwardly directed where the sign of  $E_r$  is negative ( $r/a < 0.25$ ) and where the  $n_C$  profile is hollow (from  $r/a \sim 0.2$  to  $r/a \sim 0.6$ ) (Fig.2).
- The effect of  $\Phi_1$  tends to drive the flux more.

##### ► Case **B**

- The  $\Gamma_C$  without  $\Phi_1$  is close to zero, but when  $\Phi_1$  is considered, the flux is driven outwardly for the entire radius (Fig.3).

##### ► Case **C**

- $\Gamma_C$  near the axis is negative regardless of the  $\Phi_1$ -effect (Fig.3).
- This indicates the steep  $\nabla n_C$  becomes a dominant driving force for  $\Gamma_C$ .

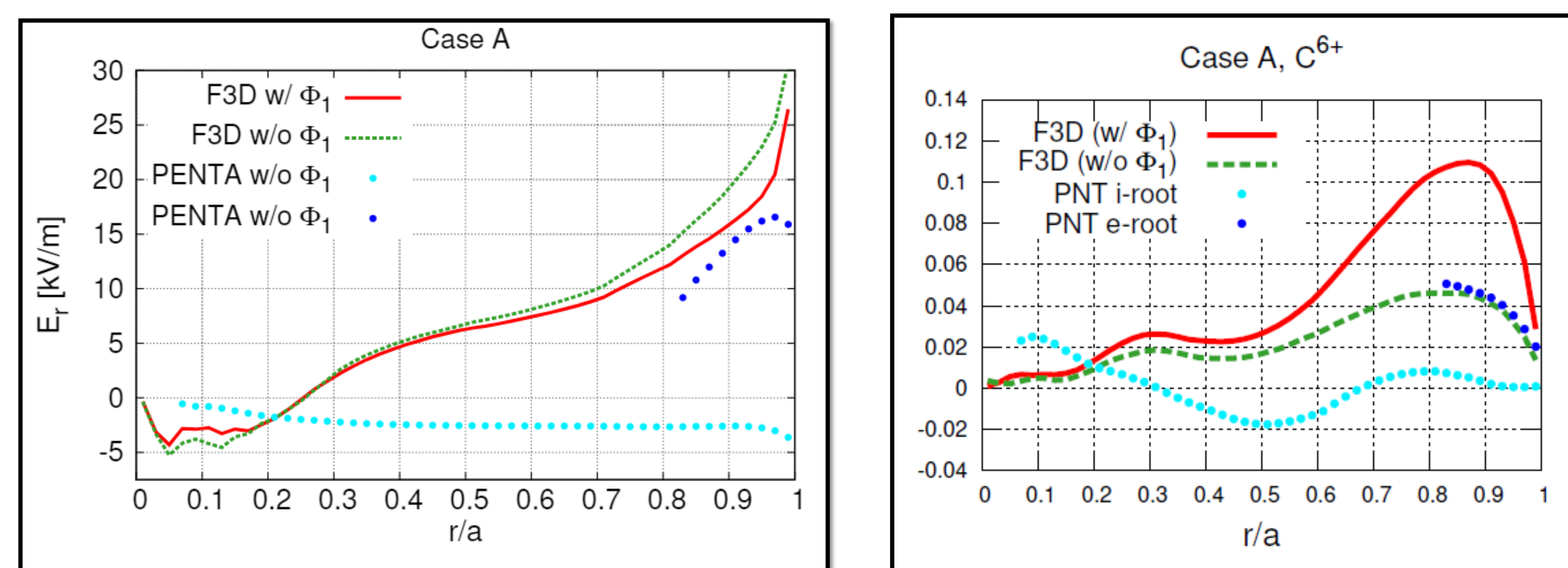


Fig.2 The radial profiles of  $E_r$  (left) and the carbon flux (right) for case **A**.

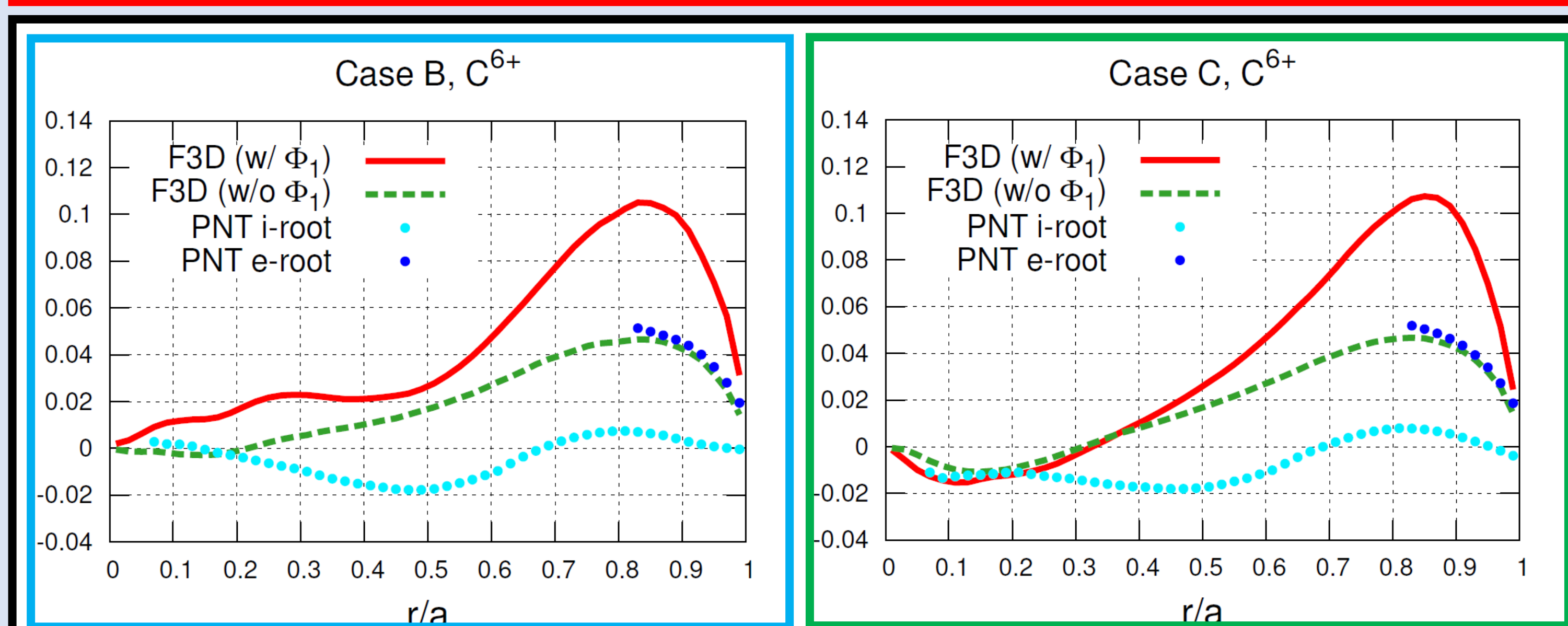


Fig.3 The radial profiles of carbon flux for case **B** (left) and **C** (right).

### CONCLUSION

- By global simulation, we find ambipolar  $E_r$  which changes the sign from negative to positive along the radius and the carbon flux can be outward.
  - Further, it is found that the outward carbon flux balances with the inward turbulent flux within factor 2 ~ 3 (omitted in this presentation).
  - Yet, several points remain to be improved or further investigated, e.g.,
1. Validity of adiabatic electron is yet to be checked by global simulation.
  2. Discrepancy in the transitioning point needs further investigation:  $r/a < 0.25$  in the simulation while  $r/a < 0.55$  in the experiment.