



# Mechanism of toroidal flow generation by electron cyclotron heating in HSX and LHD plasmas

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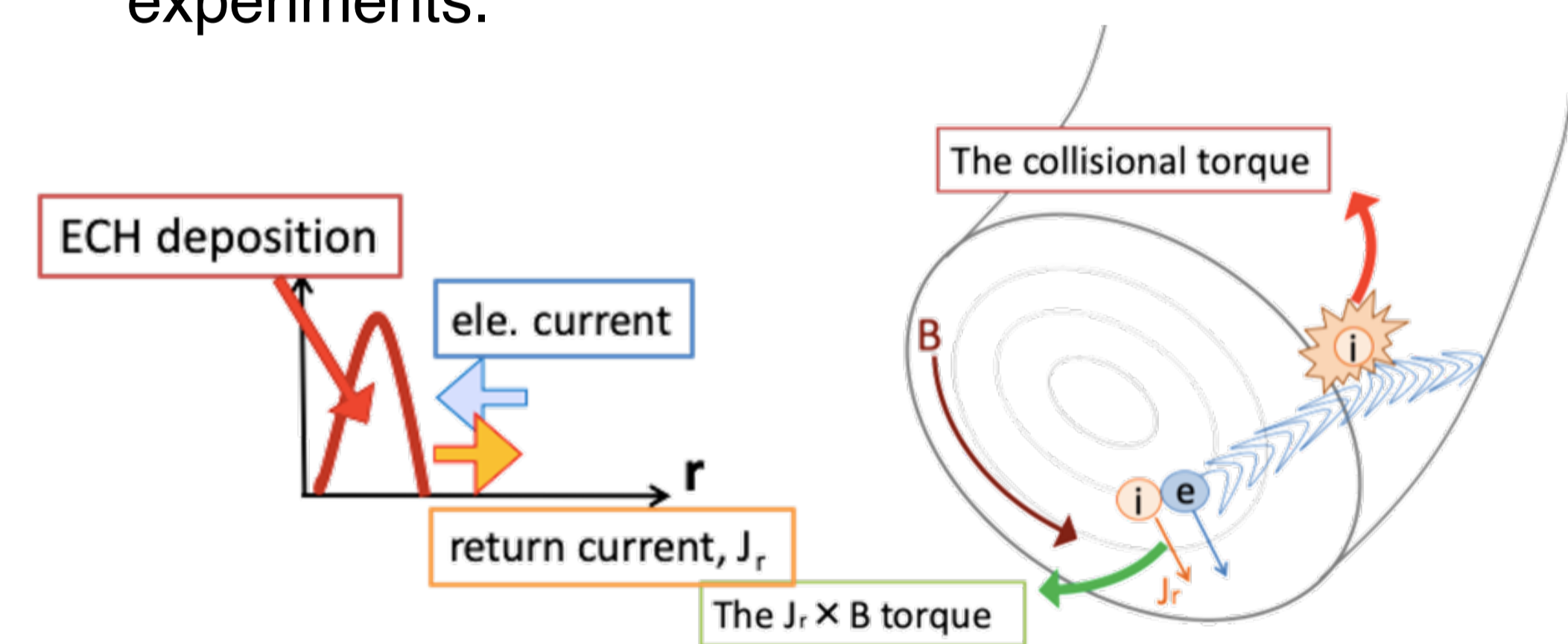


## Abstract

- Spontaneous plasma flows have been observed in electron cyclotron heating (ECH) plasmas in HSX and LHD.
- We uncover that supra-thermal electrons generated by ECH can apply forces on the plasma through  $J \times B$  and collisions, and the  $J \times B$  force overcomes the collisional force in HSX and LHD.
- We evaluate the parallel flow, solving the momentum balance equations and the diffusion equation.
- As a result, the obtained flows have reasonable agreement with the experiments.

## Introduction

- It has been observed that the toroidal flow and its shear are important for plasma confinement and MHD stability in many experiments.
- In a future reactor, NBI heating is not enough to drive a toroidal flow because of its high density and large size.
- Another method to drive the toroidal flow is required.
- Recently, the spontaneous toroidal flows driven by ECH has been observed in many tokamak and helical devices, e.g. JT-60U, LHD, HSX.
- In the previous studies, we have found that ECH can make the  $J \times B$  & collisional forces, and the  $J \times B$  force overcomes the collisional force in a non-symmetric configuration.
- In this study, we evaluate the toroidal flow driven by ECH force, and compare simulation results with experiments.



- Trapped supra-thermal electrons generate electron current due to the radial diffusion.
- Return current,  $J_r$ , cancels the electron current by ambipolar condition.
- The coupling with  $B$  makes the  $J_r \times B$  force.
- The supra-thermal electrons drift toroidally due to the precession motion.
- During the slowing down, they transfer their momenta to the bulk plasma due to collisions.
- This is the collisional force.

## Simulation model

GNET code [S.Murakami, NF (2000)]

- GNET code solves a linearized drift kinetic equation for energetic electrons in 5-D phase space based on the Monte Carlo technique.

$$\frac{\partial \delta f}{\partial t} + (\mathbf{v}_d + \mathbf{v}_{\parallel}) \cdot \frac{\partial \delta f}{\partial \mathbf{r}} + \dot{\mathbf{v}} \cdot \frac{\partial \delta f}{\partial \mathbf{v}} - C^{coll} = S^{ql}$$

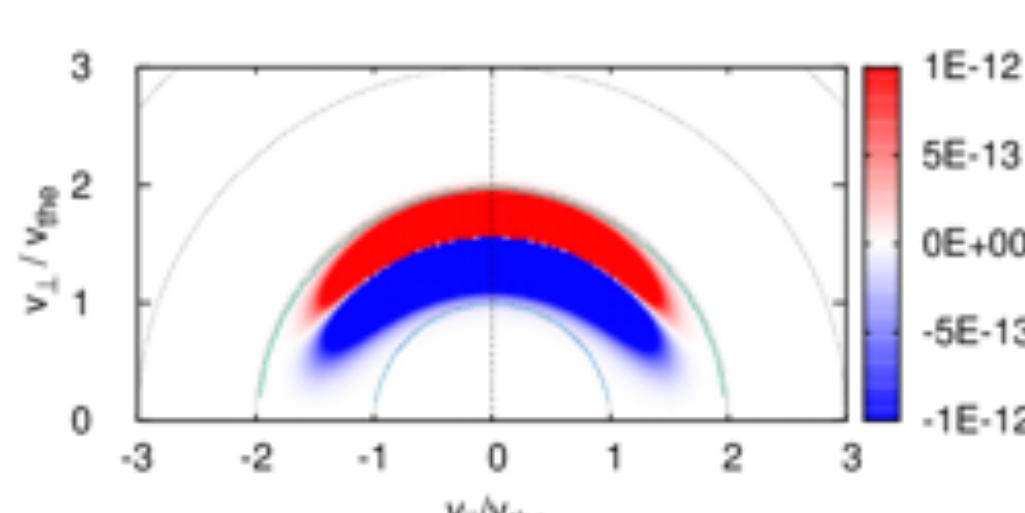
$\delta f$  : oscillation part of velocity distribution  
 $\mathbf{v}_d$  : drift velocity,  
 $C^{coll}$  : Collision operator,  
 $\mathbf{v}_{\parallel}$  : parallel velocity,  
 $S^{ql}$  : ECH quasi-linear diffusion operator

## ECH quasi-linear source term

$$S^{ql}(f_{Max}) = \frac{\partial}{\partial v_i} D_{ij}^{ql} \frac{\partial f_{Max}}{\partial v_j}$$

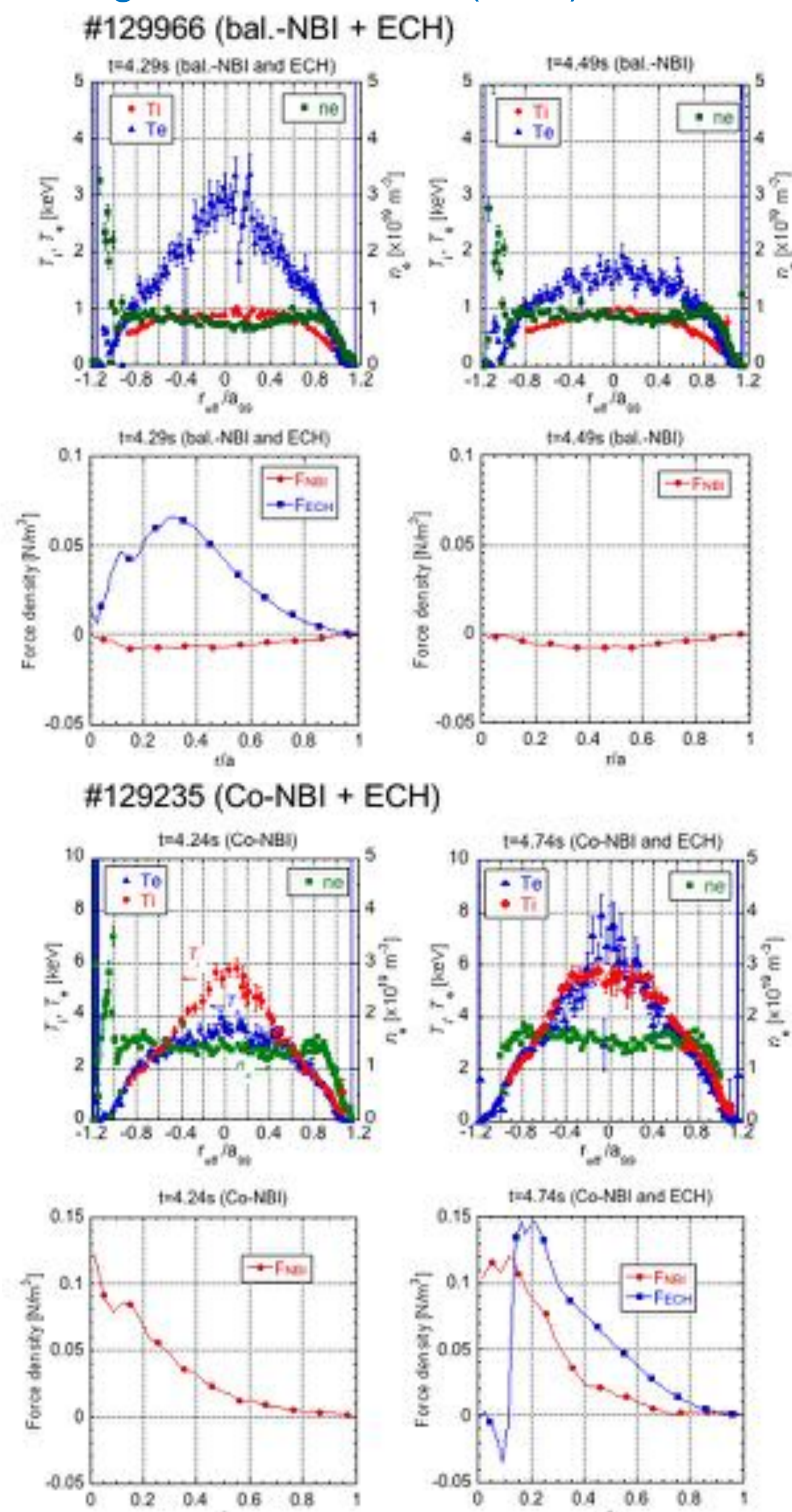
$r_0$  : heating point  
 $D$  : Quasi-linear diffusion tensor  
 $f_{Max}$  : Maxwell distribution

Example of the source term.  
 X-mode,  $n\Omega_{ce}/\omega=1.02$ ,  $N_{\parallel}=0.0$



## Simulation results

The Large Helical Device (LHD)



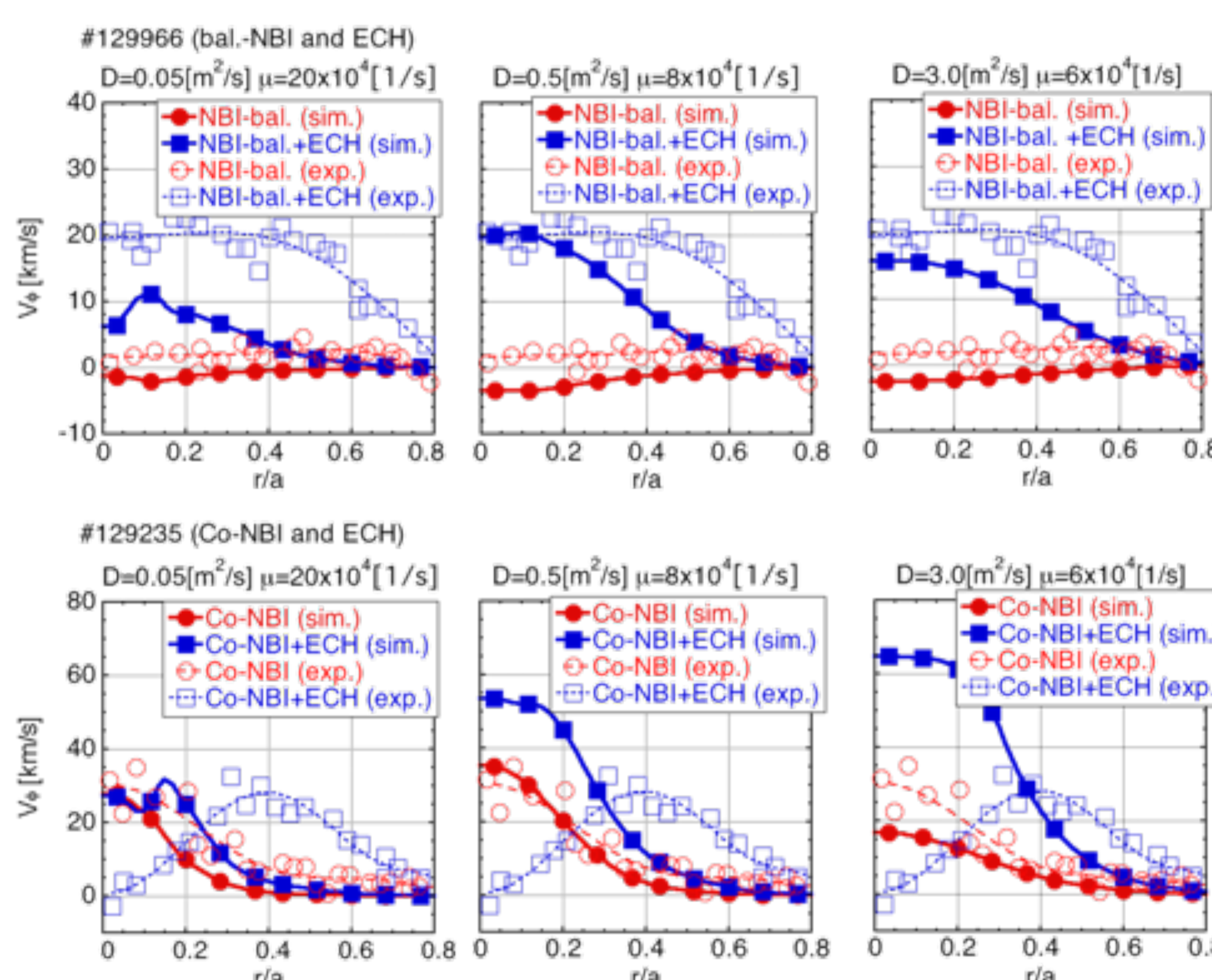
- The direction of the force is counter (co) direction in the inner (outer) region of the ECH heating location, and it qualitatively agrees with the experimental toroidal flow change.
- We can see the ECH force can be comparable with the NBI force.

## Comparison with experiments

We evaluate the toroidal flow velocity in the steady state by solving momentum diffusion equation.

$$\frac{\partial V}{\partial t} = \frac{1}{r} \frac{\partial}{\partial r} \left( r D \frac{\partial V}{\partial r} \right) + \frac{1}{m_i n_i} (F_{ECH} + F_{NBI} + F_{NTV})$$

$$F_{NTV} = -m_i n_i \mu (\delta B / B)^2 V$$

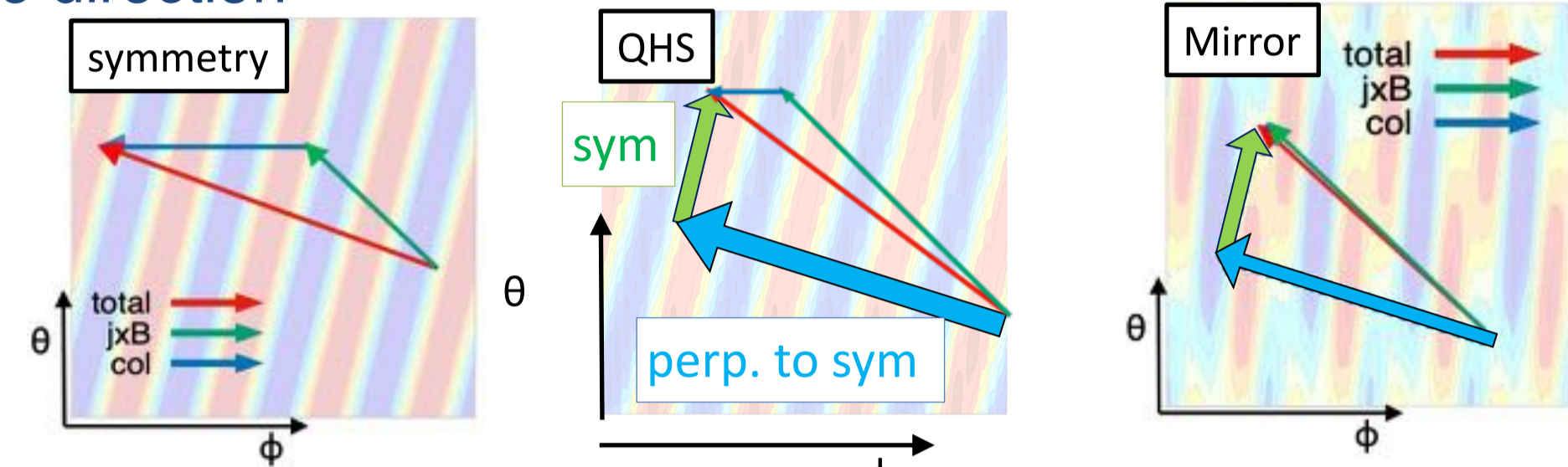


- We reproduce the co-rotating toroidal flow quantitatively in the balanced-NBI+ECH heated case.
- With the relatively large coefficient ( $D = 0.5$  &  $3.0$ ), the toroidal flow velocity increases over the entire minor radius.
- We see a difference in the toroidal flow profiles in the co-NBI+ECH heated case.

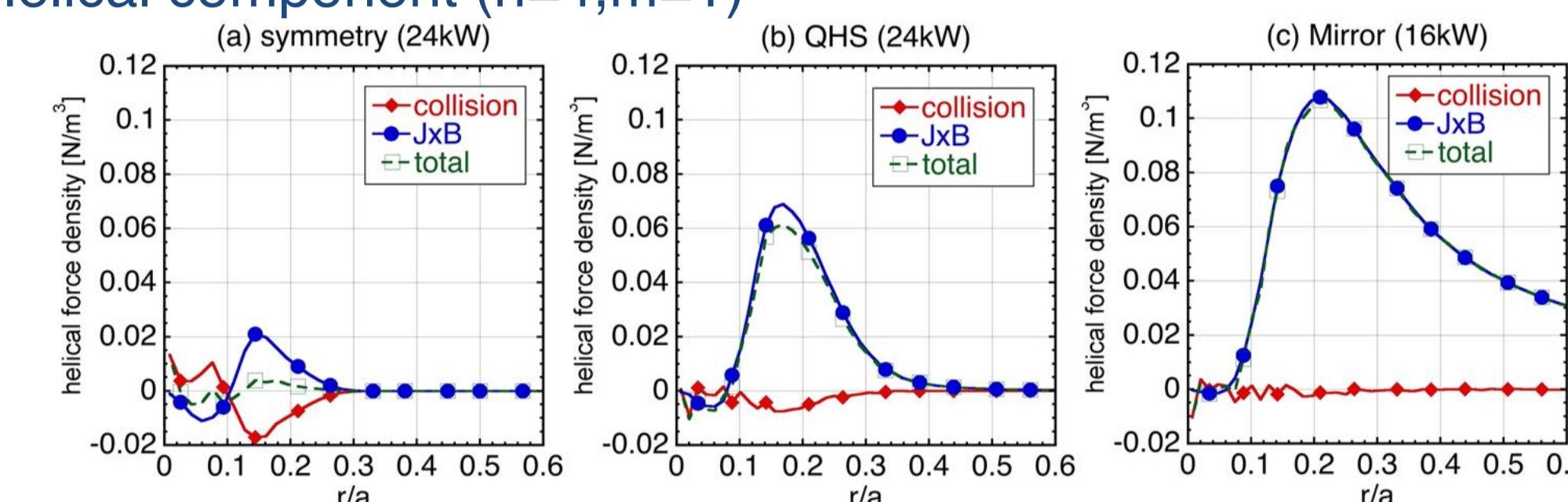
## The helically symmetric experiment (HSX)

- There are two typical configurations for HSX.
- Quasi-Helical Symmetry (QHS) configuration  
 It has a single dominant helical component,  $B(4,1)$ .
- Mirror (FL14) configuration  
 Two toroidal mirror terms,  $B(4,0)$  &  $B(8,0)$ , break the helical symmetry. Neoclassical viscosity of Mirror config. is larger than that of QHS config.

## Force direction



## Helical component (n=4, m=1)



## Comparison with experiments

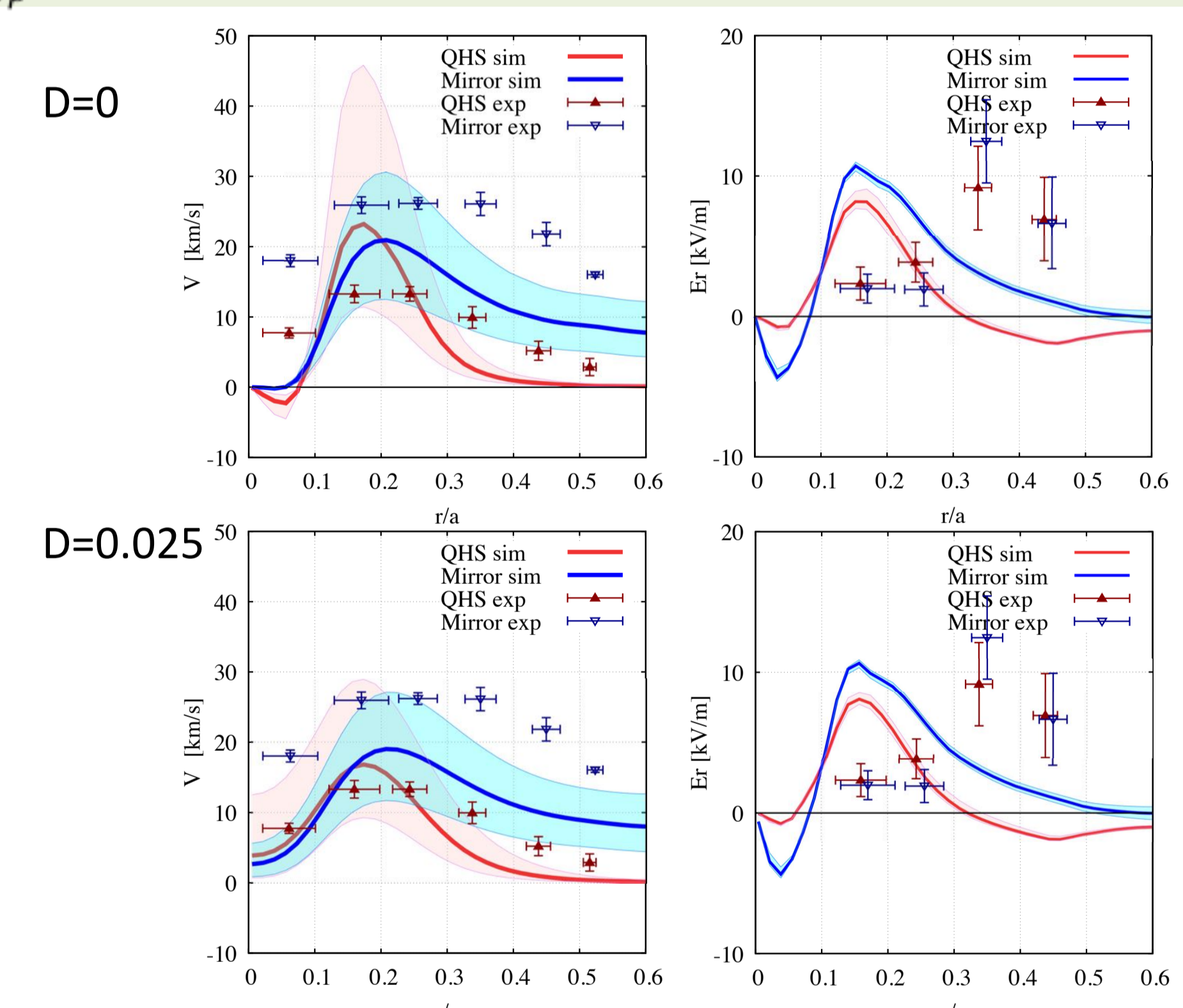
We solve the momentum balance equation and Ampere's law to evaluate the toroidal flow and the radial electric field.

$$m_i N_i \frac{\partial}{\partial t} (\mathbf{B} \cdot \mathbf{U}) = -\langle \mathbf{B} \cdot \nabla \cdot \Pi_i \rangle - m_i N_i v_{in} (\mathbf{B} \cdot \mathbf{U})$$

$$m_i N_i \frac{\partial}{\partial t} (\mathbf{B}_p \cdot \mathbf{U}) = -\frac{\sqrt{g} \mathbf{B}^c \mathbf{B}^\alpha}{c} (\mathbf{J}_{plasma} \cdot \nabla \rho) - \langle \mathbf{B}_p \cdot \nabla \cdot \Pi_i \rangle - m_i N_i v_{in} (\mathbf{B}_p \cdot \mathbf{U})$$

(+Diffusion term)

$$\frac{\partial}{\partial t} \frac{\partial \Phi}{\partial \rho} (\nabla \rho \cdot \nabla \rho) = 4\pi (\mathbf{J}_{plasma} \cdot \nabla \rho) + (\mathbf{J}_{ext} \cdot \nabla \rho)$$



- In the case of the completely helically symmetric configuration, the helical components cancel each other.
- There is a net force in the symmetry direction even in QHS config. because other small magnetic modes enhance the radial diffusion.
- The force in Mirror config. is about twice as large as that in QHS config. despite the less absorbed power.
- The obtained velocity has reasonable agreement with experiment, especially with the diffusion effect.

## Conclusion

- We have evaluated the  $J \times B$  and collisional forces and toroidal flows using GNET code in order to clarify the mechanism of the toroidal flow change in HSX and LHD.
- The obtained force by ECH is the same order as NBI force, and its direction agree with experiment observation.
- The evaluated toroidal flows have reasonable agreement, except for the co-NBI+ECH case of the LHD experiment.
- It indicates that ECH force would drive the toroidal flow.