

Modeling Snowflake Divertors in MAST-U Tokamak

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Magnetic equilibrium modeling using the FIESTA code shows that steady-state snowflake (SF) divertor (1) configurations can be created and maintained with the existing poloidal field coil set in the MAST-U tokamak. A full multi-fluid plasma transport model with a computational grid encompassing two poloidal magnetic field nulls, with charge-state-resolved carbon impurities sputtered at material boundaries, and fast convective SF plasma transport (mixing) has been applied to the standard and SF divertor configurations using the UEDGE code. Results show that the SF divertor 1) reduces peak heat loads to divertor plates by spreading particles and heat fluxes to additional strike points; 2) the additional strike points receive up to 10% of heat and particle fluxes (w.r.t. the outer strike point), with weak dependence on two-dimensional profiles of particle diffusivities and heat conductivities as well as SF geometry details; 3) SF divertors approach the outer strike point detachment conditions at lower upstream density w.r.t. the standard divertor. Experiments are planned to validate the model, as part of a broad advanced divertor research program in MAST-U (2). The SF divertor uses two nearby 1st order nulls separated by distance d_{xx} , for a larger region of low poloidal field B_p (PF) that modifies geometry and transport resulting in heat and particle flux sharing and reduction (1).

Snowflake magnetic equilibria created with a free-boundary Grad-Shafranov equilibrium code FIESTA using realistic currents I 4kA and existing MAST-U poloidal field coils show significantly increased near-separatrix connection length (cf. the standard divertor). A three-coil SF divertor algorithm previously developed in NSTX and DIII-D experiments (3) was used and improved with additional PF coils. Based on the theoretical SF criteria (1), families of up-down symmetric SF configurations were created (as, e.g., in Fig. 1) with $d_{xx} a(\lambda_q/a)^{1/3} = 0.15m$, where $a = 0.65m$ –plasma minor radius, $\lambda_q \sim 8mm$ –SOL heat flux width. The core kinetic profiles from the H-mode plasma model with modest shaping, $I_p = 1MA$ and $P_{NBI} = 2.5MW$, were used. Two main SF variants were investigated: the SF-plus configuration with a 2nd null in the private flux region (Fig. 1 (b)) and the SF-minus with a 2nd null in the low field side common flux region (Fig. 1 (c)). Flux tube connection length from outer midplane to outer strike points within 1-2 mm from the separatrix were up to 50% higher in both the SF-plus and SF-minus configurations (cf. standard divertor).

Two transport models in UEDGE were used over a broad range of core-boundary densities $2 \cdot 10^{19}m^{-3}$ to $2 \cdot 10^{20}m^{-3}$. In the standard transport model, SOL ion diffusivity and heat conductivity profiles were taken to match the MAST-U SOLPS/EIRENE model (4). A second model emulated the theoretically predicted fast plasma mixing (the “churning” mode) in the two-null SF region (5). The churning mode convective zone radius was estimated to be $r^* = 0.81a(B_p a/R)^{1/3} \sim 10cm$ for the SF divertor, and $r^* = 0.44B_p a^2/R < 1cm$ for the standard divertor. The enhanced SF transport was modeled as Gaussian profiles with $r^* = 10cm$ centered at the nulls (Fig. 1 (d), (e)), with peak values $140m^2/s$ (cf. $1m^2/s - 2m^2/s$ in the standard transport model). With the standard transport, peak heat fluxes at primary strike points SP1 and SP4 were 30-50% lower in the SF-plus and SF-minus configurations, cf. standard divertor. Heat flux profiles were considerably broader, due to diffusion into the secondary strike points and the increased connection length (Fig. 2 (a)). Primary strike point heat flux profiles in the SF-plus and SF-minus configurations with a small $d_{xx} = 2.5 - 3cm$ were similar. Heat fluxes showed weak dependence on wall albedo α_w (hydrogen atom wall reflection). However, heat and particle fluxes to secondary strike points SP2 and SP3 were found to be rather sensitive to the albedo. With $\alpha_w = 0.98$, additional strike points received heat and particle fluxes of about 5-10% of those in the outer strike point (Fig. 2 (c,d)), as compared to only 1-3% with $\alpha_w = 1$.

As the core-boundary (upstream) plasma density n_e was increased, primary strike points SP1 and SP4 in the SF divertors reached low temperatures (~ 1 eV) and low peak heat fluxes at lower n_e than the standard divertor (Fig. 2 (a)).

A complex interplay of SOL geometry effects, transport, and magnetic configurations was observed with the enhanced SF transport. The additional plasma mixing led to spreading of ions over a wider divertor region thereby reducing peak ion and neutral densities (Fig. 1 (d-e)). Heat flux profiles in all strike points were broadened (Fig. 2 (b-d)), peak heat fluxes were reduced. At higher core density ($10^{20}m^{-3}$), heat fluxes to additional strike points SP2 and SP3 were increased. With the enhanced SF transport (mixing), small differences between SF configurations with various inter-null distance $d_{xx} = 3 - 12cm$ were observed. Primary strike points SP1 and SP4 reached detachment conditions at lower core-boundary density n_e than with the standard transport assumptions.

In summary, significant progress in numerical modeling of SF divertor configurations was made in preparation for MAST-U experiments. The results show that heat and particle fluxes were spread among additional strike

points and reduced as compared to the standard divertor. The MAST-U modeling results are consistent with numerous findings in previous SF divertor experiments conducted in TCV, NSTX, and DIII-D tokamaks (1,3).

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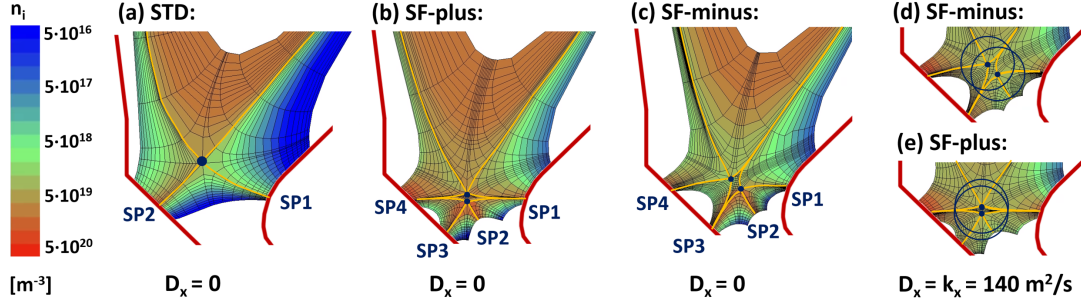


Figure 1. Divertor ion density distribution with standard transport profiles (a) the standard divertor; (b) SF-plus; (c) SF-minus configurations; and with fast plasma mixing in the (d) SF-plus; and (e) SF-minus configurations. The SF-plus configuration with $d_{xx}=2.5$ cm, and the SF-minus with $d_{xx}=5$ cm are shown. The core-boundary interface density was $n_e=6 \times 10^{19} m^{-3}$ (low), and albedo was 0.98 (high hydrogen absorption by the tokamak walls).

Figure 1: Image

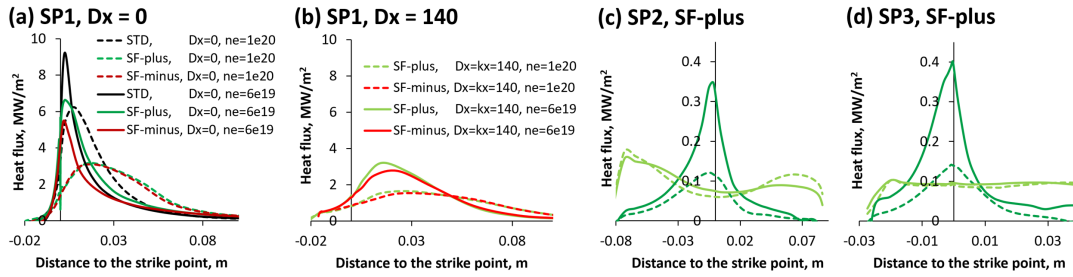


Figure 2. Modeled divertor heat flux profiles in the standard divertor, SF-plus and SF-minus configurations with standard and enhanced SF transport and at low and high densities. (a) In strike point SP1, standard transport, $n_e=6 \times 10^{19} m^{-3}$ (low) and $n_e=1 \times 10^{20} m^{-3}$ (high); (b) Same as in (a) but with enhanced SF transport; (c-d) SP2 and SP3 of the SF-plus, standard and enhanced transport, $n_e=6 \times 10^{19} m^{-3}$ and $n_e=1 \times 10^{20}$. Albedo value was 0.98 (high hydrogen absorption by the tokamak walls); carbon impurities were not modeled in the presented results.

Figure 2: Image

Country or International Organization

United States

Affiliation

Lawrence Livermore National Laboratory

Authors: Dr KHRABRYI, Aleksandr (LLNL); SOUKHANOVSKII, Vsevolod (Lawrence Livermore National Laboratory); UMANSKY, Maxim (Lawrence Livermore National Lab); HARRISON, James (CCFE); Dr MOULTON, David (CCFE); Mr ROGNLIEN, Thomas (Lawrence Livermore National Laboratory)

Presenter: Dr KHRABRYI, Aleksandr (LLNL)

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