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Ion Heating by Nonlinear Landau Damping of High-n Toroidal Alfvén Eigenmodes in ITER

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Bulk ion heating from nonlinear Landau damping of high-n Toroidal Alfvén Eigenmodes (TAEs) excited by energetic ions is studied. Based on the nonlinear saturation level of high-n TAEs which is confirmed by numerical studies of a wave-kinetic equation incorporating the ion Compton scattering-induced spectral transfer and the particle trapping by the wave, we assess potentially beneficial effect of bulk ion heating, so-called alpha channeling for ITER DT plasmas using an integrated simulation. The result indicates that, in its optimistic limit, the anomalous ion heating due to high-n TAEs can compensate the heating reduction due to energetic particle loss.

Motivations

- For high-n TAEs ($n \ge 10$) in ITER, the nonlinear mode interaction such as ion Compton scattering (ICS, or nonlinear Landau damping) can be a dominant mechanism for the saturation of TAEs [1,2,3].
- Previous analytic studies derived the saturation level of the modes by ICS [1,3] and the consequent bulk ion heating rate [2].

3. Alpha-channeling via high-n TAEs in ITER

• EP transport and energy channeling by TAEs

$$\frac{\Delta E}{\Delta P_{\phi}} \simeq \frac{\omega}{n} = const.$$

$$P_{EP \to TAE} \equiv n_{EP} \left\langle -\frac{dE}{dt} \right\rangle_{EP} = n_{EP} \frac{\omega}{n} \left\langle -\frac{dP_{\phi}}{dt} \right\rangle_{EP} = \frac{\omega}{n} \frac{ZeBr}{q} \Gamma_{EP}$$
(5)
(6)

From the EP flux by [3], we can estimate the EP energy transfer to TAEs.

• We need to investigate how the multiple TAEs evolve dynamically and how significant the ion heating will be in ITER, which can be a candidate for the **alpha-channeling**.

2. Nonlinear evolution of high-n TAEs

• Wave-kinetic equation [1,2] $\frac{\partial}{\partial t} E_{\omega} = 2\gamma_{L}(\omega)E_{\omega} - \sum_{\omega'}M_{\omega,\omega'}E_{\omega'}E_{\omega} \qquad (1)$ $\frac{\partial}{\partial t}E_{ion} \simeq \sum_{\omega=\omega_{1}}^{\omega_{M}}\sum_{\omega'=\omega_{1}}^{\omega}\frac{\omega-\omega'}{\omega}M_{\omega,\omega'}E_{\omega'}E_{\omega} \qquad (2)$

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- γ_L > 0 modes in the gap grows linearly and transfer the energy to lower ω modes by ICS.
 γ_L < 0 modes near the lower continuum are linearly stable, but are driven nonlinearly by the energy transfer.
- This system exhibits a typical predator-prey (PP) behavior.



- PP oscillation of TAEs by ICS
- For a parabolic shape of the linear growth rate curve,
 Eq. (1) and (2) result in the PP oscillation of the
 multiple TAEs (Fig. 2 (b)).



Integrated simulation in ITER burning plasma



- The EPs are transported by TAEs, and the central plasma heating by EP drops about 1/3.

- The nonlinear transfer can sufficiently drive the linearly stable modes such as kinetic TAEs (Fig. 2 (c)).
- The transient peak level of the total mode energy

during the bursty behavior is much larger than the previously expected NL saturation level.



- Initially, the modes near the center of the gap are amplified.
- Then, the downward transfer by ICS excites the stable modes near the gap boundary.
- Finally, the low- ω modes are damped by the continuum.
- Since the gap structure varies radially, the mode amplitude evolves in real space as well (Fig. 3 (b)).
- It can lead to EP avalanches [4] which induce performance degradation and a wall damage.

- The ICS of high-n TAEs induces additional ion heating (alpha-channeling, Fig 6 (b)).
- The ion heating by alpha-channeling can compensate the collisional ion heating loss by the EP transport ($P_{TAE \rightarrow ion}/P_{EP,loss} > f_{i,coll} \sim 0.3$).

				Tab.1
[MW]	P _{abs}	P _{abs} ,elec	P _{abs} ,ion	Q
	$= P_{\alpha} + P_{NB} - P_{EP,loss}$			$= P_{fus} / P_{in}$
Case 1: Reference $(P_{EP,loss} = 0)$	150.1	109.3	40.8	10.1
Case 2: EP transport	121.7	91.0	30.7	8.79
Case 3: EP transport + channeling	128.5	87.4	41.2	9.64

- In Case 1, the EP energy is transferred to plasmas by coll. slowing down only ($Q \approx 10.1$).
- Considering the EP transport (Case 2), the reduced plasma heating by EPs decreases T_i and so does alpha heating further ($Q \approx 8.79$ with 25% reduction of ion heating).
- In Case 3, the ion heating fraction becomes higher by alpha-channeling, which compensates the ion heating loss by the EP transport ($Q \approx 9.64$).
- The alpha-channeling by TAEs is \sim 10MW, which is as influential as a single ICRH antenna.

Conclusion

- The NL evolution of high-n TAEs induced by ICS is investigated.
- The ICS shows that the modes evolve with PP bursts that can cause EP avalanches.
- If we consider a weak effect of particle trapping, the system becomes saturated.
- The ion heating during the spectral transfer, responsible for alpha-channeling, is modeled
- Mode saturation by considering the wave-trapping effect [5]



- When we take the wave-trapping effect into account, the system saturates after decaying oscillations (Fig. 4 (a)).
- The levels normalized to the analytic estimations converge into O(1).



(a)

 $\frac{\sum E_{\omega}}{E_{sat}} \quad 10$

 $\frac{P_{ion}}{P_{ion,sat}} \frac{30}{20}$

- The saturated energy spectrum (Fig. 4 (b)) shows different characteristics according to v_{d0} , but both are shifted downwards from the original linear growth curve.

for the predictive simulation.

- With the integrated modeling, it is found that the ion heating by the channeling can compensate the heating loss by the EP transport in the optimistic limit.
- It is still necessary to consider various factors such as GAM coupling [6], CAE-induced channeling [7] and cyclotron resonance of alpha particles [8] in the future.

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