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Max-Planck-Institut
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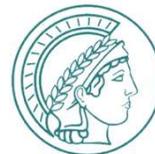
Electron heated plasmas in Wendelstein 7-X and ASDEX Upgrade; necessity to control turbulent transport

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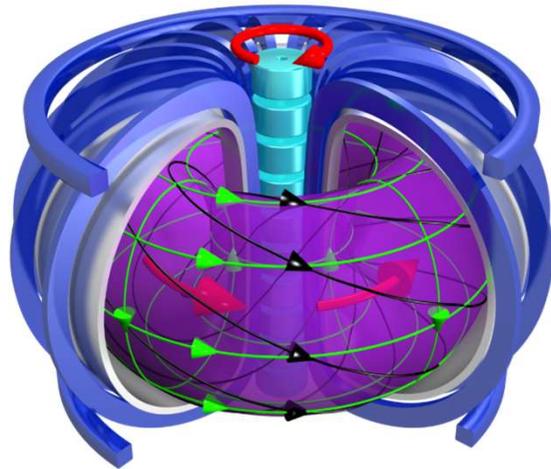
HELMHOLTZ
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 **EUROfusion**

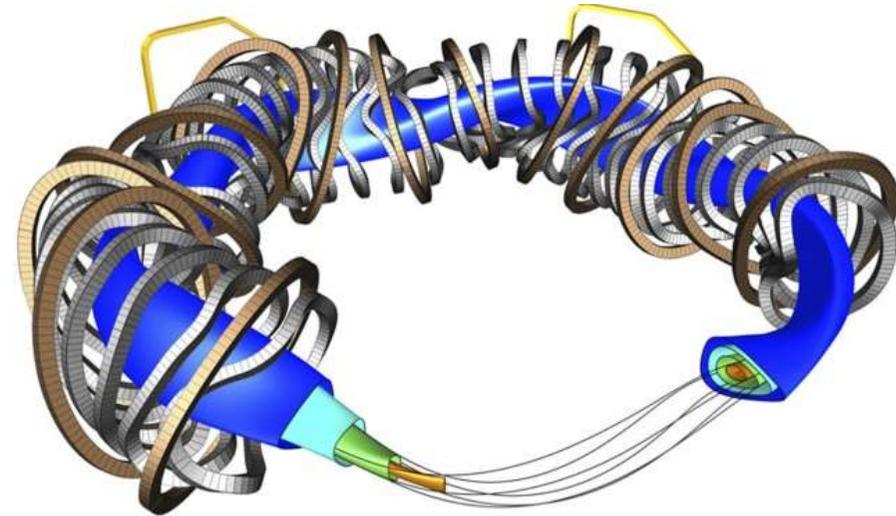
Magnetic confinement concepts

ASDEX Upgrade : Tokamak (2D)



- current $I_p \sim 1$ MA
- $a = 0.5$ m, $R = 1.6$ m, $B = 2.5$ T
- $P_{\text{ECRH}} = 5$ MW
- low neoclassical transport

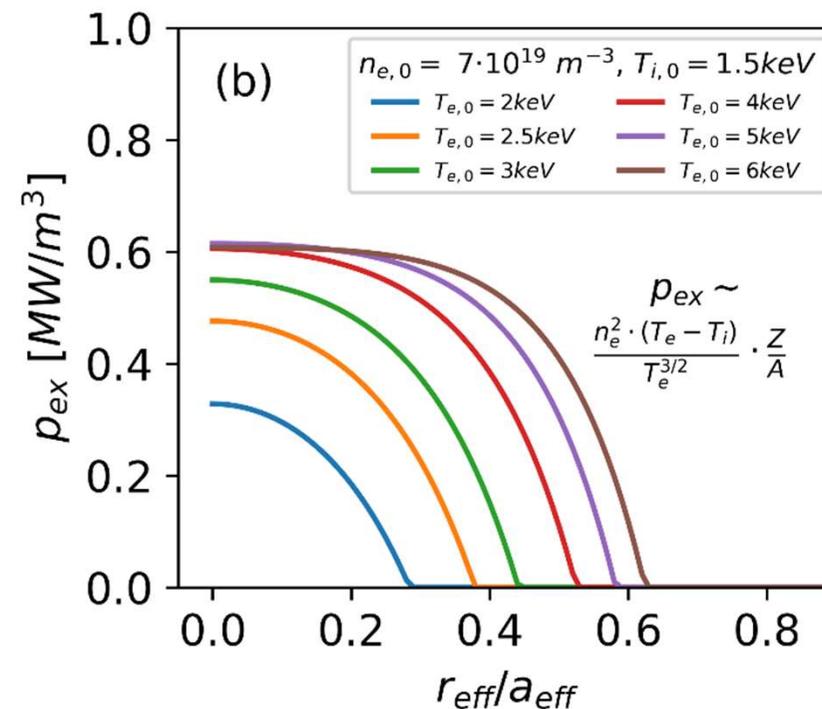
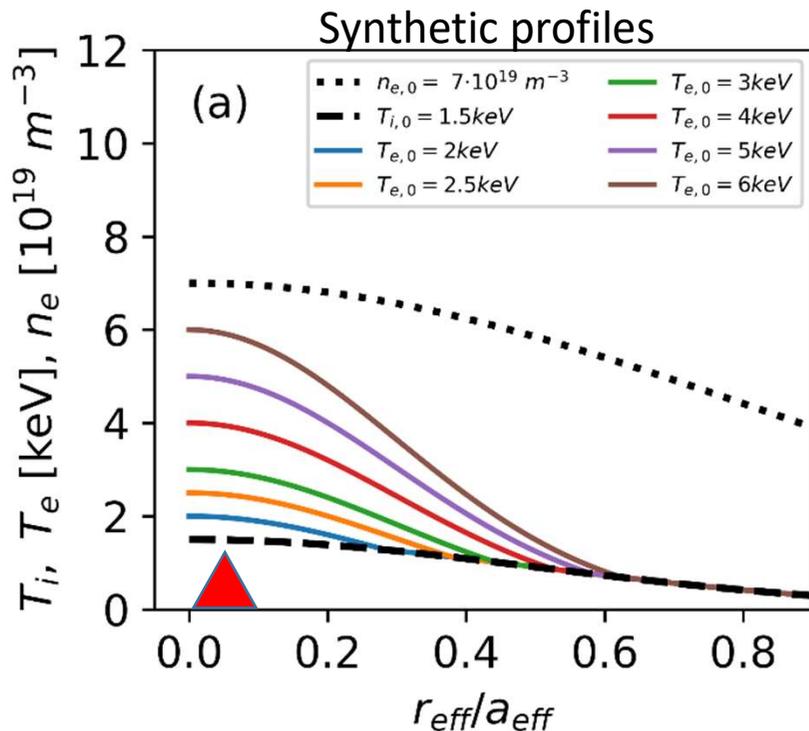
Wendelstein 7-X: Stellarator (3D)



- No induced plasmas current
- $a = 0.5$ m, $R = 5$ m, $B = 2.5$ T
- $P_{\text{ECRH}} = 7$ MW (15 MW envisaged)
- Optimised neocl. transport $\epsilon_{\text{eff}} < 1\%$

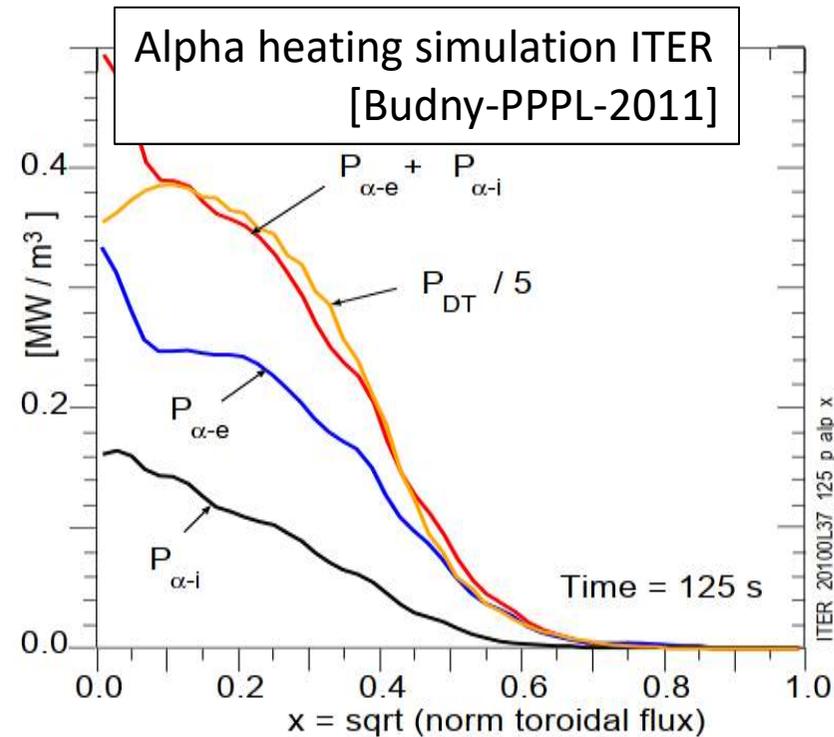
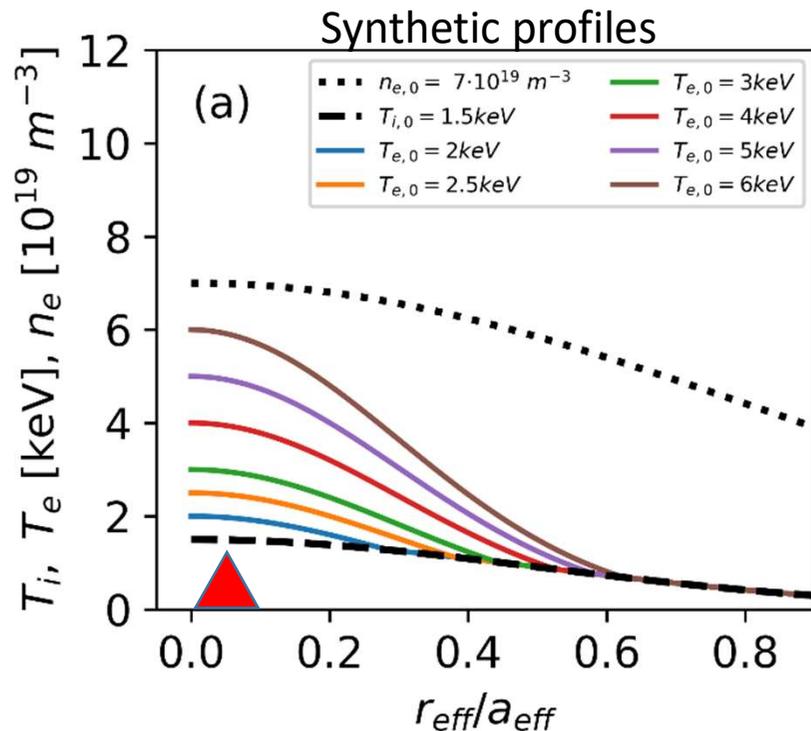
Electron heated plasma

- ECRH electron heating is very localised
- Ion heating (exchange power) is broad, $P_{ei} \sim \frac{n_e^2 (T_e - T_i)}{T_e^{3/2}} \cdot \frac{Z}{A}$



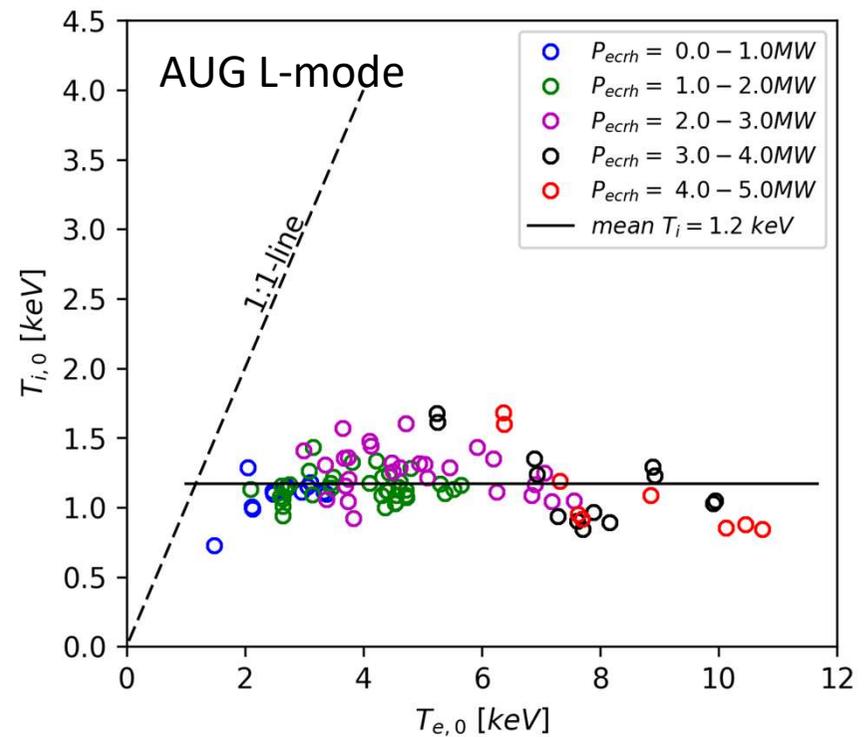
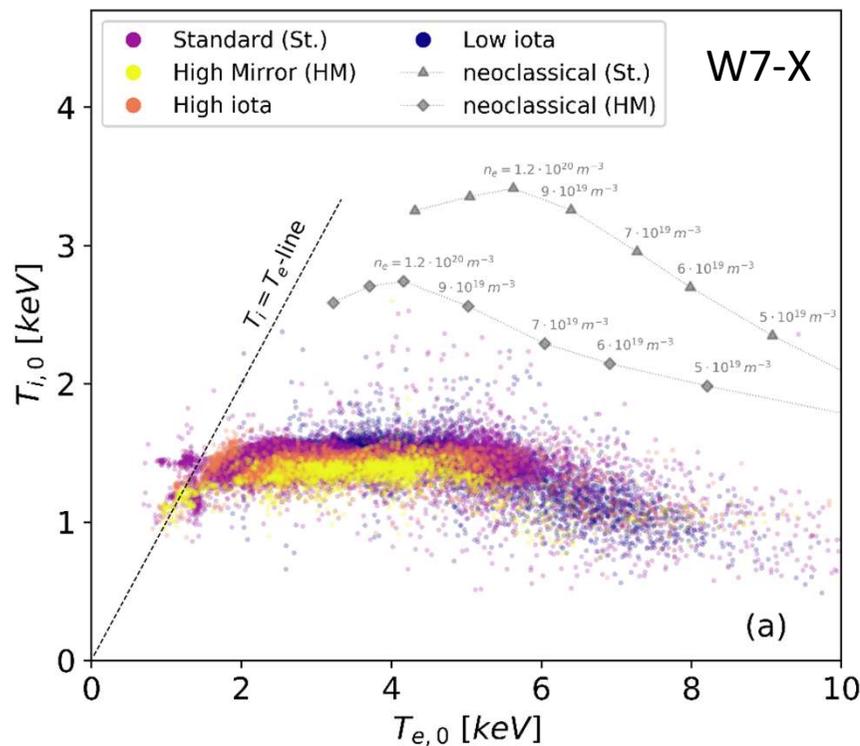
Electron heated plasma

- ECRH electron heating is very localised
 - Ion heating (exchange power) is broad
 - In a reactor alpha heating will predominantly heat electrons...
- e-heated plasma performance important to study!



Electron heated plasmas using ECRH

- **Ion temperature T_i -clamping in both devices: $T_i \leq 1.5$ keV**
 - Low exchange power with broad ion heating profile?
- **Under these conditions, turbulent transport is dominant**
 - Which type of turbulence is at play? Profile stiffness?



In electron heated L-mode plasmas of ASDEX Upgrade and Wendelstein7-X, the ion temperature clamps at $T_i \sim 1.5$ keV, independent of magnetic configuration:

- The ions are heated by energy exchange $p_{ei} \sim n_e^2 (T_e - T_i)/T_e^{3/2}$, which offers a broad ion heating profile only where $T_e/T_i > 1$.
- Heat transport due to ITG/TEM turbulence is strongly exacerbated as $T_e/T_i > 1$ even at small positive ε with $T_e/T_i = 1 + \varepsilon$
- The resulting “stiff” core transport leads to T_i -clamping.

To remedy the impact of clamping, tokamaks and stellarators either feature:

1. an **H-mode** with a strong **edge pedestal** to lift the “stiff” core T_i profiles
2. **suppressed or reduced core turbulence** by means of:
 - i. Fast ions stabilisation (tokamaks and stellarators)
 - ii. Enhanced density gradients (stellarators and negative-magn.-shear-tokamak)
 - iii. Designing turbulence-resilient-configurations with e.g. negative triangularity tokamaks or e.g. low elongation in stellarators.

Key findings of this work



In electron heated plasmas :

- T_i -clamping at ~ 1.5 keV occurs independent of magnetic configuration.
- Ion are heated through exchange power as $T_e/T_i > 1$.
- At the same time clamping puts a stronger constraint on T_i than profile stiffness as $T_e/T_i > 1$ exacerbates ITG/TEM turbulence.
- Therefore, in practical terms T_i -clamping leads to an apparent strongly “stiff” core transport in both AUG and W7-X

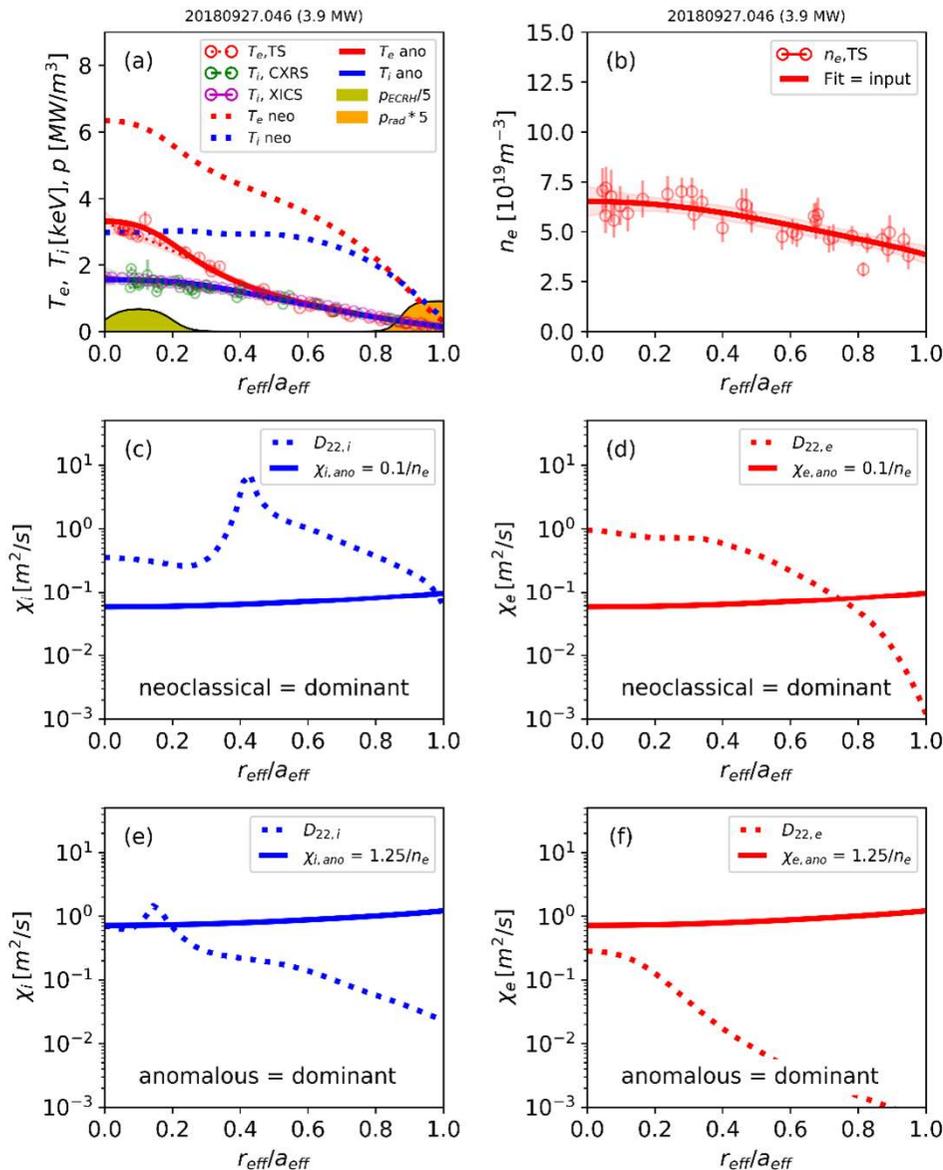
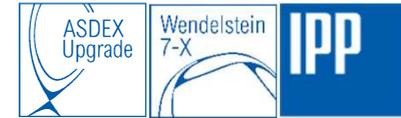
$T_e/T_i \leq 1$ reduces transport, but is not possible.

→ therefore T_i clamping is unavoidable in electron heated plasmas

To remedy the impact of clamping, tokamaks and stellarators either feature:

1. an H-mode with strong **edge pedestal** to lift the “stiff” core T_i profiles
2. **suppressed or reduced** core turbulent transport by means of:
 - i. Fast ions stabilisation (tokamaks and stellarators)
 - ii. Enhanced density gradients (stellarators and negative-magnetic-shear-tokamak)
 - iii. Negative triangularity (tokamaks only – exploratory)

W7-X: neoclassical transport overestimates performance



Wendelstein 7-X:

Comparing NC to experiment:

- $P_{\text{ECRH}} = 4\text{MW}$, $n_{e,0} = 7e19 \text{ m}^{-3}$,
- Central ECRH deposition
- Edge radiation assumed

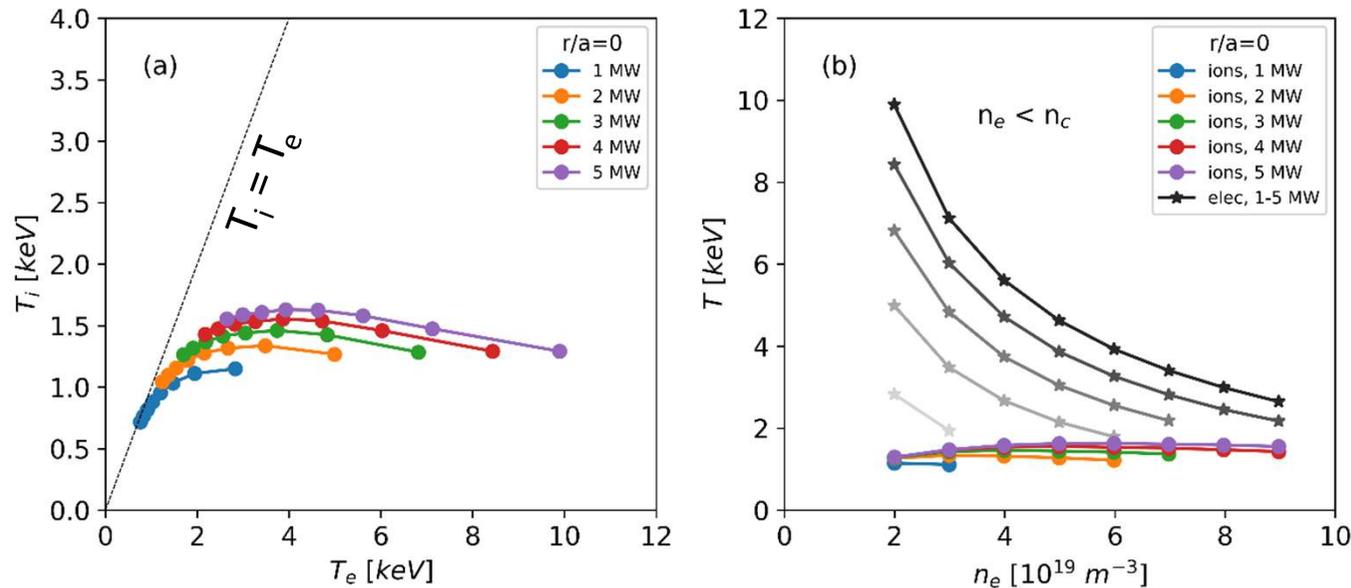
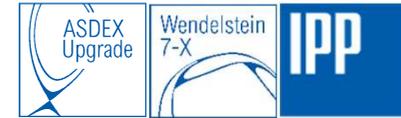
Experiments result:

$$T_{i,0} \sim 1.6 \text{ keV}$$

Neoclassical transport simulation:

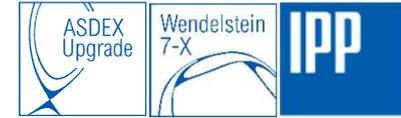
$$T_{i,0} \sim 3.0 \text{ keV}$$

configuration independent gyroBohm transport

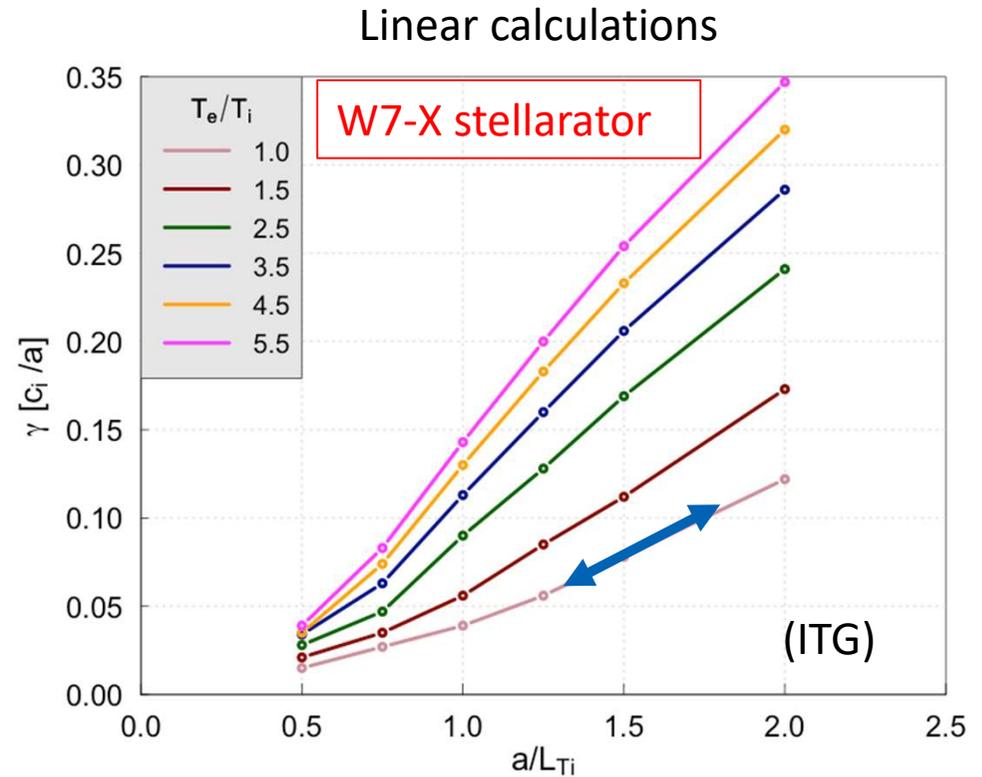
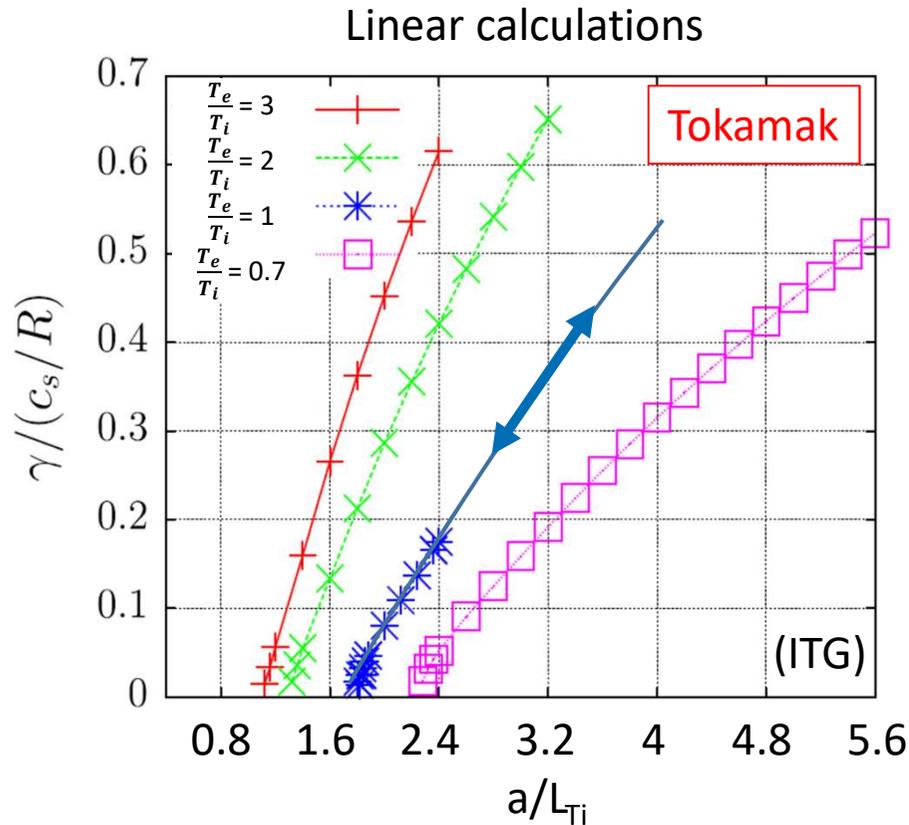


- NTSS simulations: $\chi_{turb} \gg \chi_{neo}$, Vary $P_{ECRH} = 1-5 \text{ MW}$ and $n_e = 10^{19} - 10^{20} \text{ m}^{-3}$
- GyroBohm ion diffusivity: $\chi_{i,gB} = \sqrt{\frac{m_i}{e}} \frac{T_i^{3/2}}{aB^2}$, (electrons: $\chi_{e,turb} = \text{constant}$)
- Tried in AUG and W7-X configurations \rightarrow clamping is configuration independent:
 - Due to broad ion heating profile (exchange power)
 - Large anomalous transport

Expectations for ion turbulent heat transport



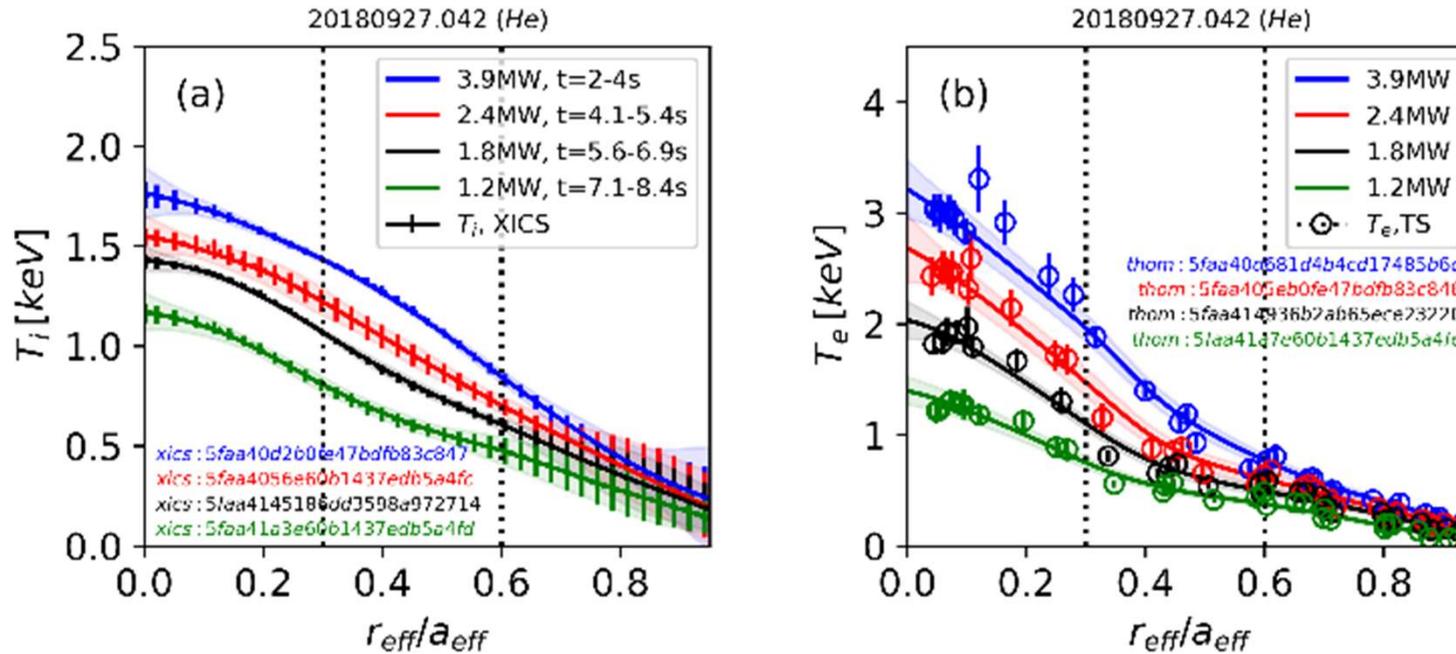
- Ion heat transport can be driven by Ion Temperature Gradient (ITG) or trapped electron modes (TEM)
- Both devices may also feature a strong T_e/T_i drive for ion transport



γ is growth rate of most unstable mode

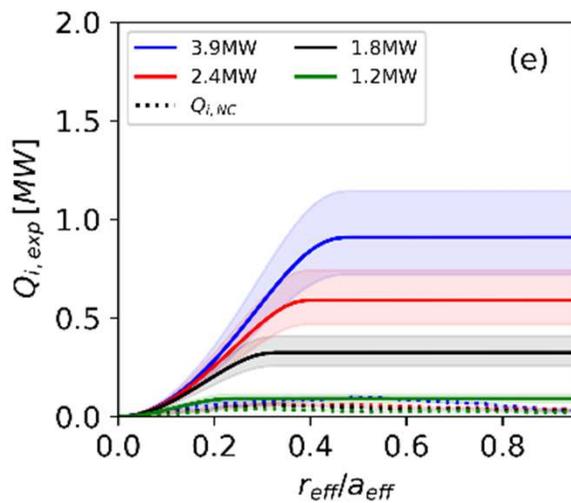
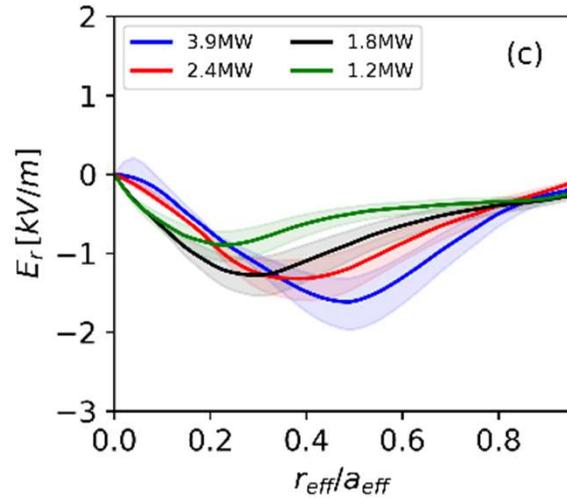
[Zocco, Xanthopoulos et al. JPP 84 (1), 2018]

Wendelstein 7-X power scan experiment



- Steps down $P_{ECRH} = 3.9 - 1.2$ MW
- T_e varies $3 \rightarrow 1.3$ keV
- T_i varies $1.7 \rightarrow 1.2$ keV
- n_e is virtually constant at $7 \cdot 10^{19} \text{ m}^{-3}$

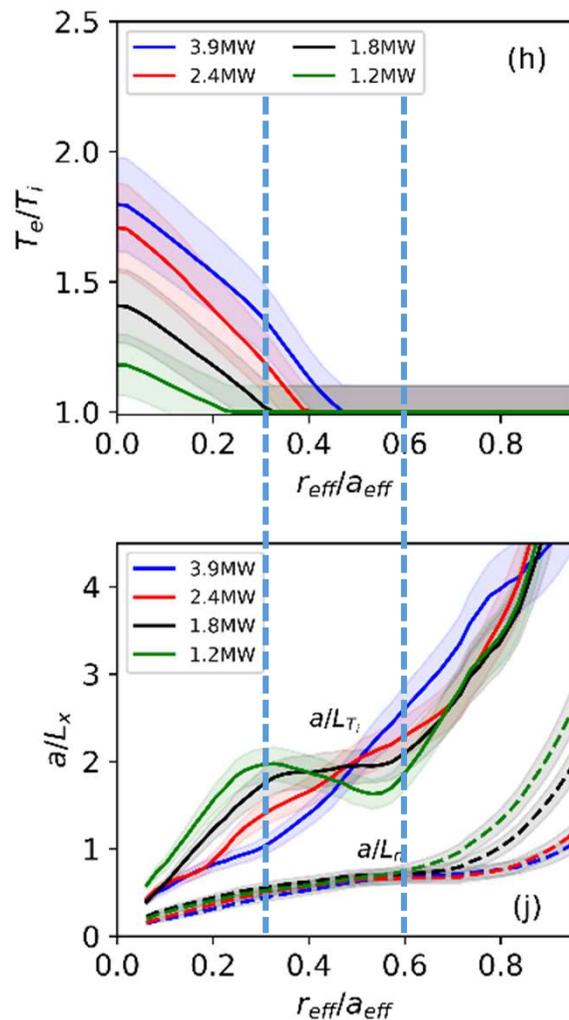
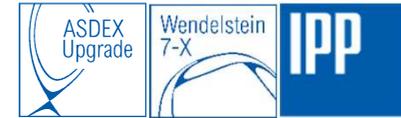
Wendelstein 7-X power scan experiment



With increasing P_{ECRH}	$r/a = 0.3$	$r/a = 0.6$
E_r	ion root	ion root
$Q_{i,\text{turb}} = Q_{i,\text{exp}} - Q_{i,\text{NC}}$	0.1 - 0.5 MW	0.1 - 1 MW
T_e/T_i		
a/L_{Ti}		
a/L_n		

- Avoid root transitions for accuracy of $Q_{i,\text{turb}}$
- Turbulent transport dominates: $Q_{i,\text{turb}} \gg Q_{i,\text{NC}}$

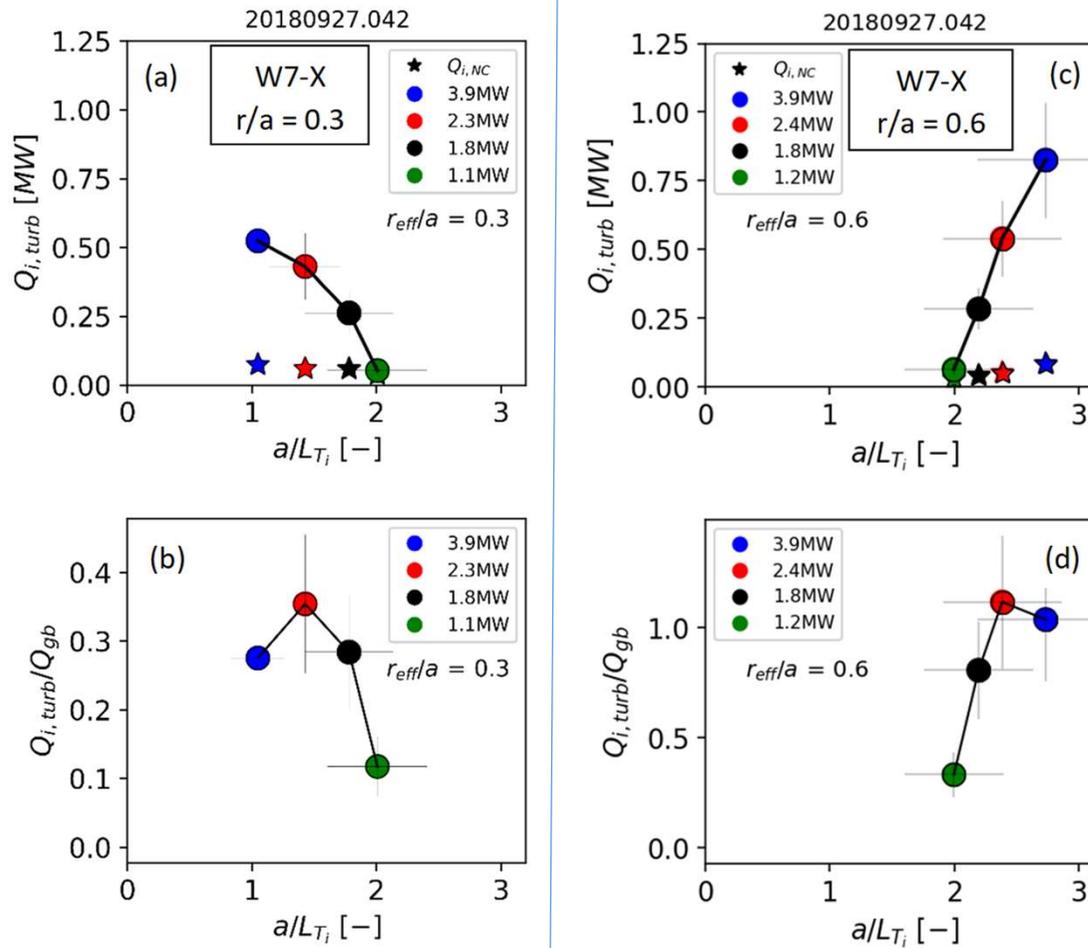
Wendelstein 7-X power scan experiment



With increasing P_{ECRH}	$r/a = 0.3$	$r/a = 0.6$
E_r	ion root	ion root
$Q_{i,turb} = Q_{i,exp} - Q_{i,NC}$	0.1 - 0.5 MW	0.1 - 1 MW
T_e/T_i	increases 1 - 1.3	constant 1
a/L_{Ti}	decreases 2 - 1	increases 1.5 \rightarrow 2.5
a/L_n	$a/L_n \ll a/L_{Ti}$	$a/L_n \ll a/L_{Ti}$

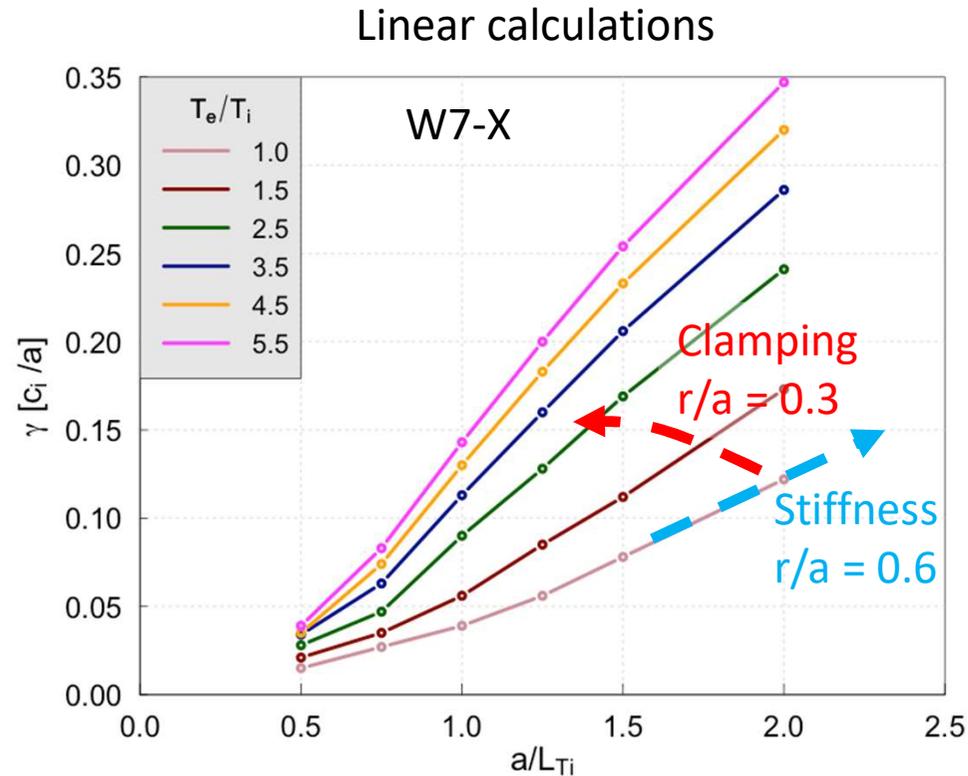
- Profiles of a/L_{Ti} cross over at $r/a \sim 0.5$
- This intersects with radius from where $T_e/T_i = 1$

Wendelstein 7-X power scan experiment



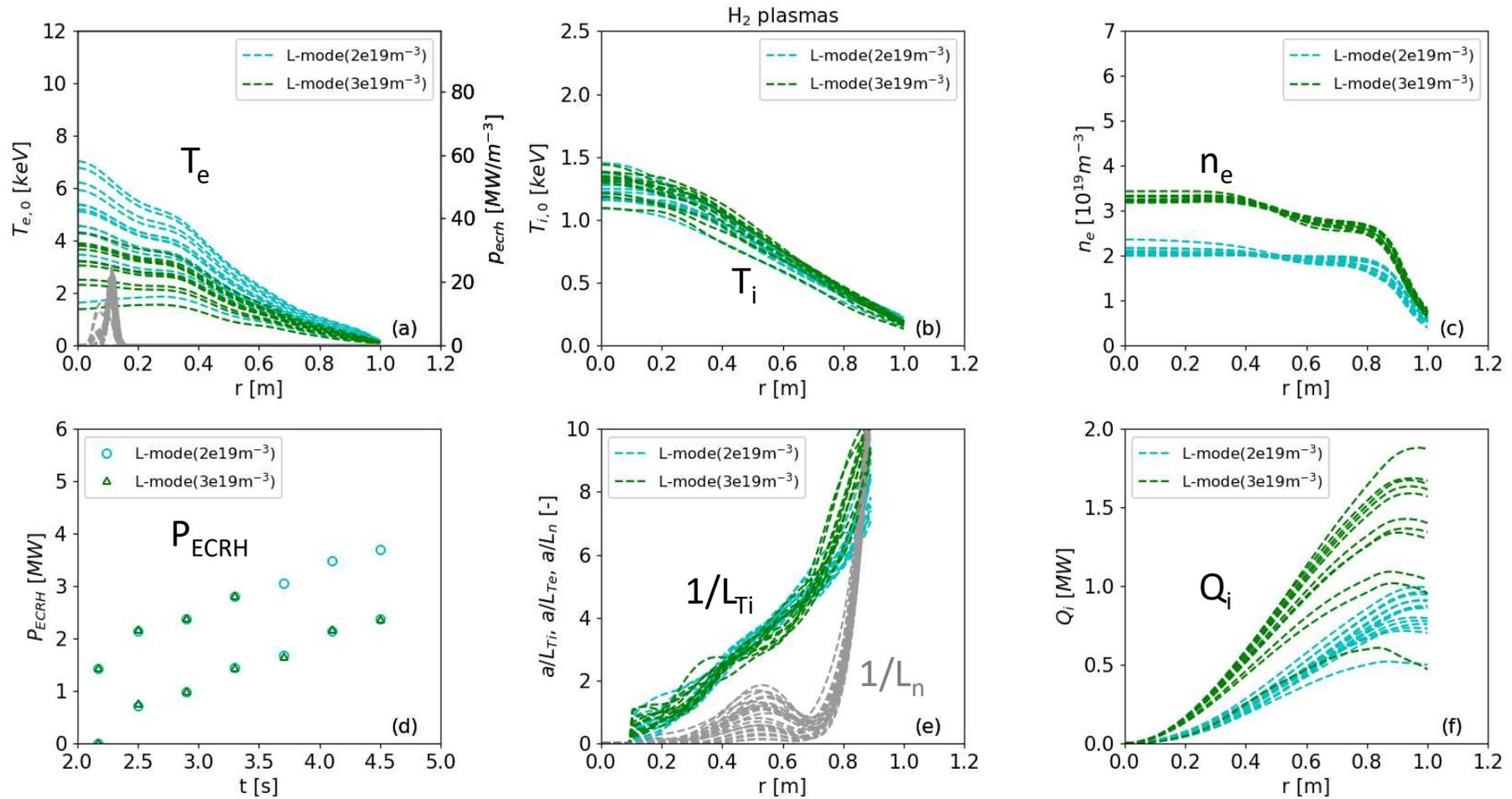
- At $r/a = 0.3$ we find a reduction of a/L_{T_i} with increasing $Q_{i,turb}$
- At $r/a = 0.6$, a/L_{T_i} increases with increasing $Q_{i,turb}$

Wendelstein 7-X power scan experiment



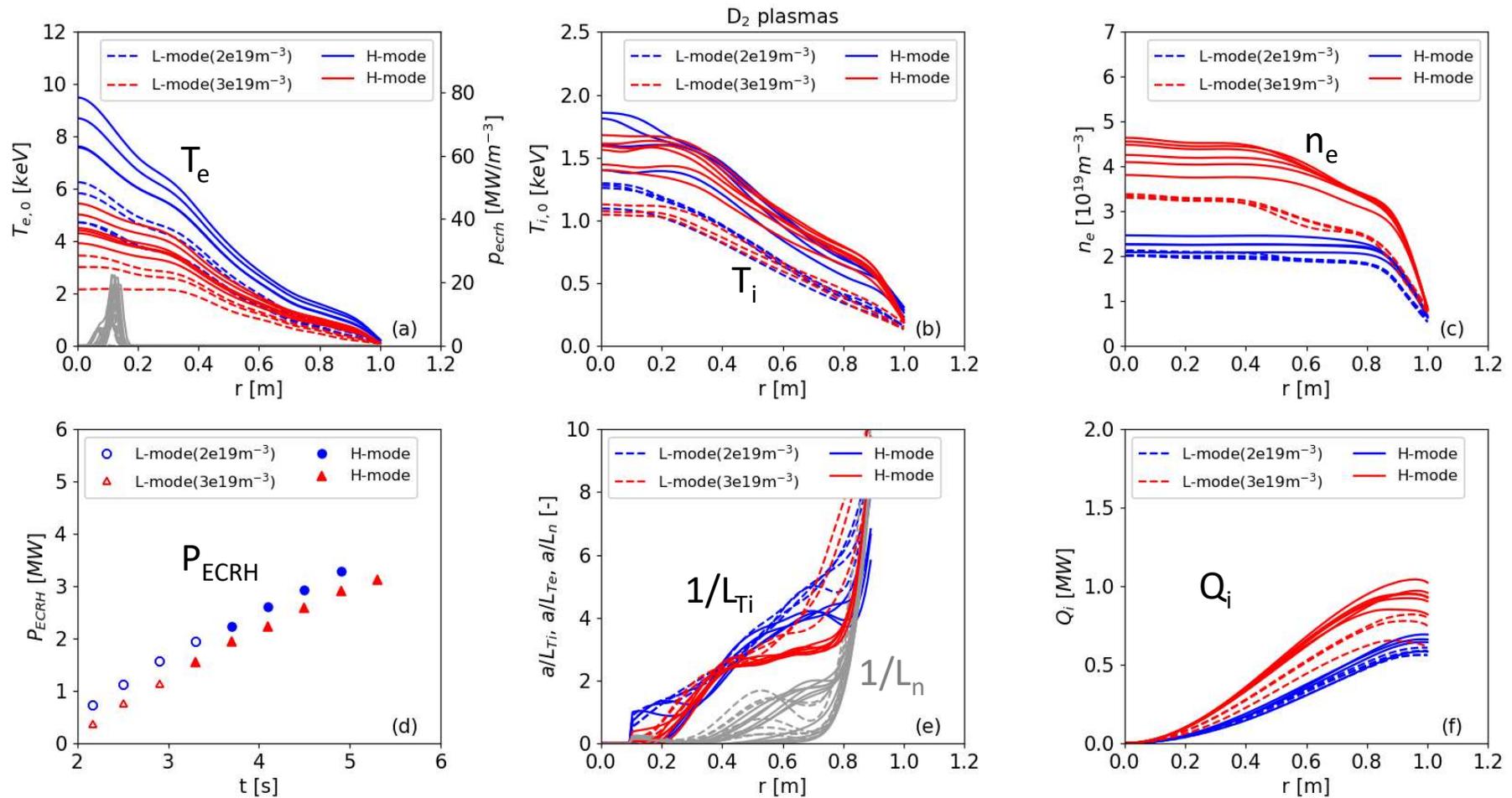
- **At $r/a = 0.3$** , both T_e/T_i and clamping increases (jump curves)
- **At $r/a = 0.6$** , as $T_e/T_i = 1$, the stiffness = constant (stay on curve)

AUG power scan experiment in H₂ (L-mode)



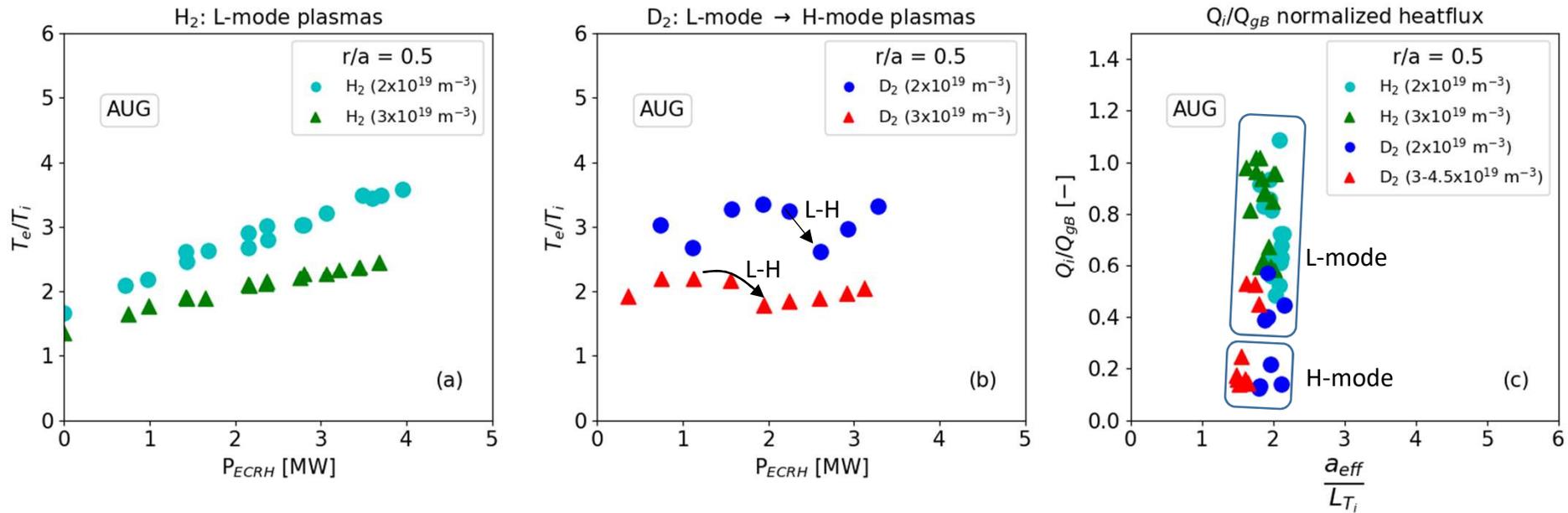
- Large variation of T_e , clamping of $T_i < 1.5$ keV in L-mode H₂ plasmas
- Q_i varied by factor 3

AUG power scan experiment in D₂ (H-mode)



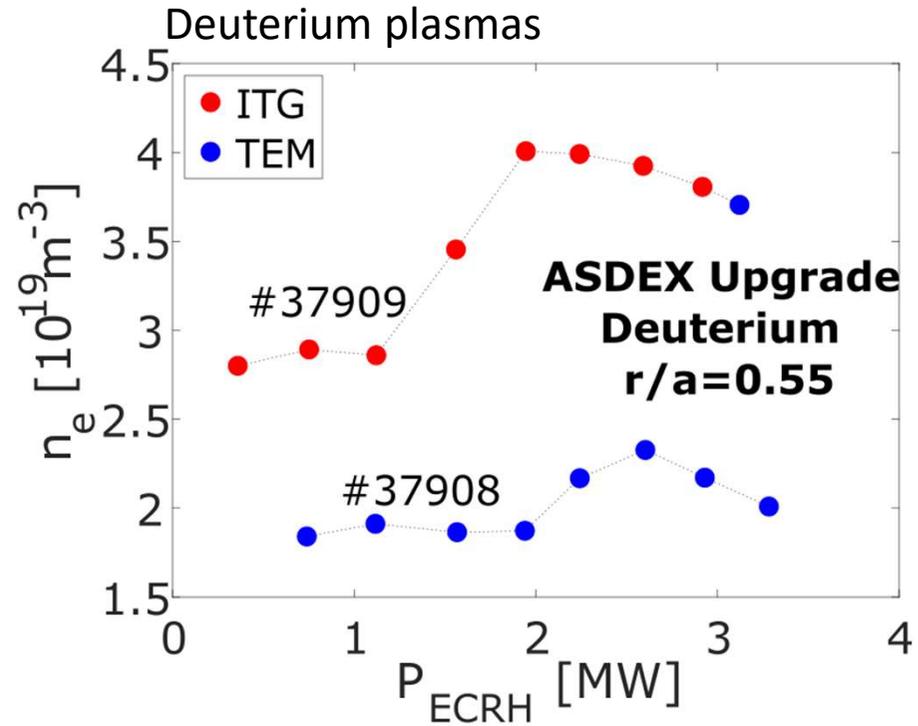
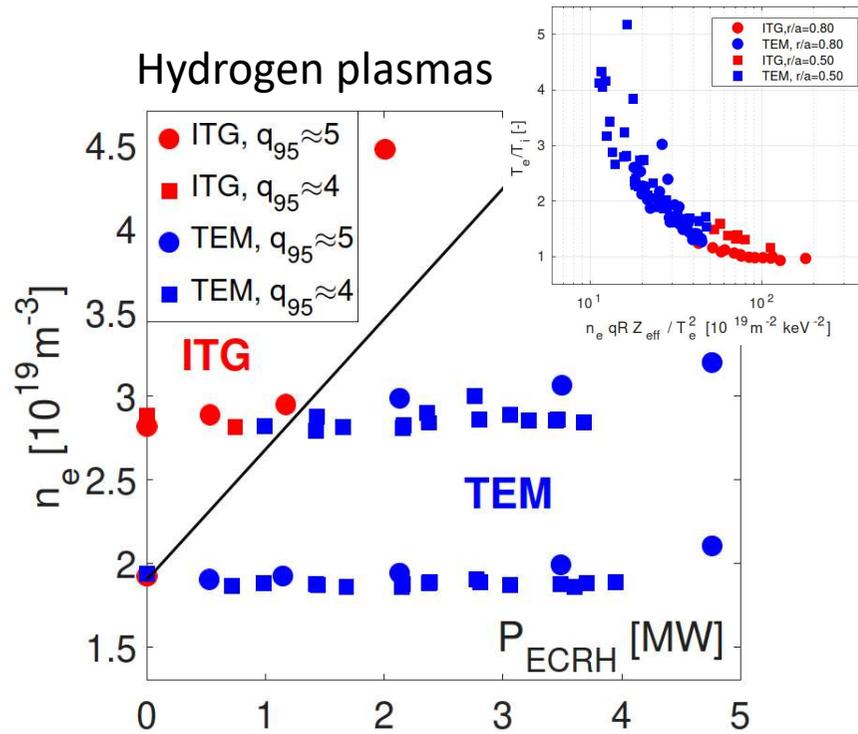
- Tokamak L-H transition can help enhance $T_{i,0}$ as pedestal lifts the core
- Although local gradients remain the same (or reduce somewhat)

Clamping due to turbulent transport



- Ion temperature gradient clamped at $a/L_{Ti} \sim 2$
- Variation of T_e/T_i from 1.5-3.5 enhances clamping (this is not infinite stiffness!)
- Issue for L-mode tokamak fusion reactor → impact on size requirement
- H-mode performance almost fully depends on the achievable pedestal

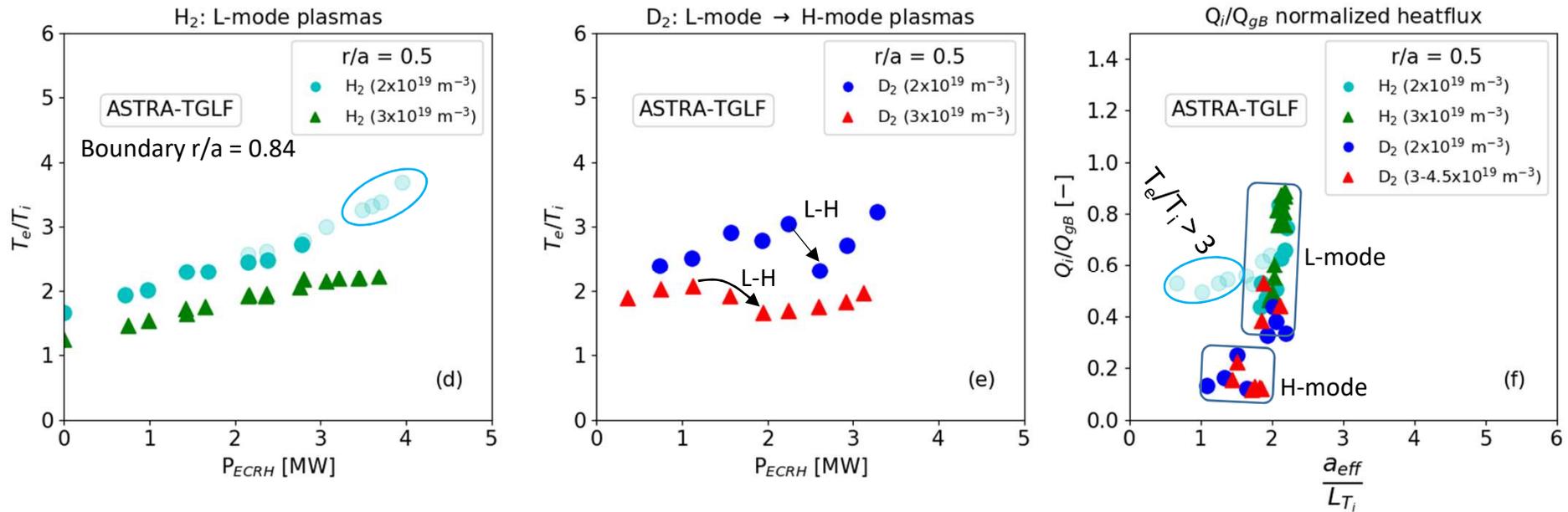
ASTRA-TGLF simulations: TEM or ITG dominant



Using ASTRA-TGLF simulations with quasi-linear gyrokinetic calculations

- Depending on collisionality either TEM or ITG is the dominant micro-instability
- Both instabilities feature enhanced stiffness with increasing T_e/T_i

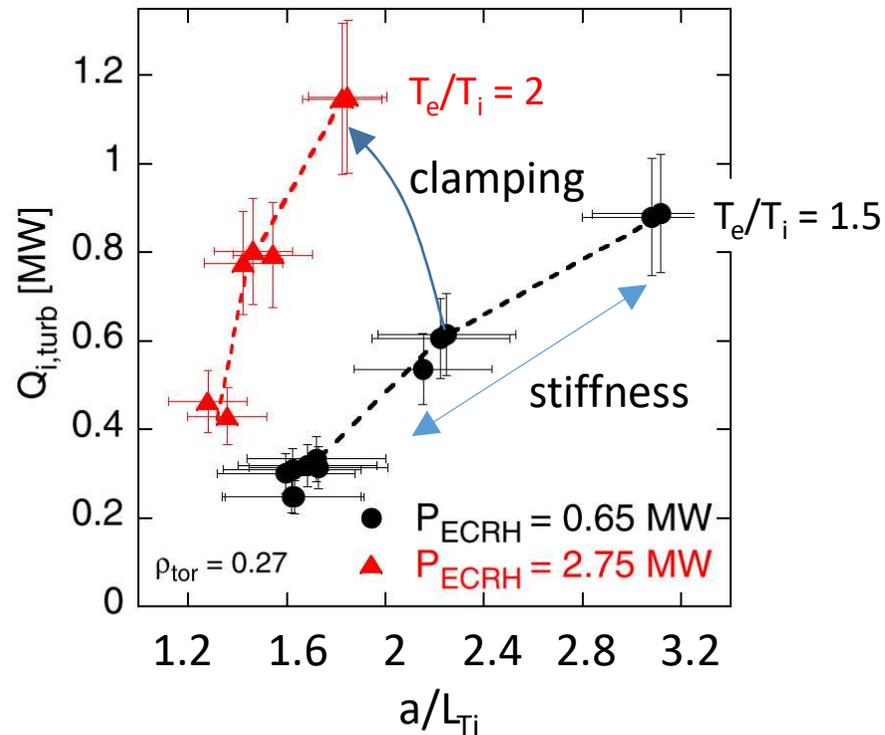
ASTRA-TGLF simulations reproduce clamping



- ASTRA-TGLF simulations show good quantitative agreement of core transport (pedestal from experiment):
 - With some systematic deviations
- L-mode plasmas: clamping effect is reproduced for $T_e/T_i < 3$
- H-mode plasmas: TGLF penalises T_e/T_i stronger than experiment

Clamping is not Stiffness

[F. Ryter et al 2019 Nucl. Fusion 59 096052]



Clamping >> Stiffness

In controlled profile experiment (NBI and ECRH deposition scan)

- with $T_e/T_i = \text{constant}$: much wider variation of $a_{eff}/L_{Ti} = 1.5 - 3$
- Here also, jumping curves when T_e/T_i changes enhances clamping.

T_i clamping is independent of configuration!



Strong T_i -clamping is equally found in AUG and W7-X electron heated plasmas.

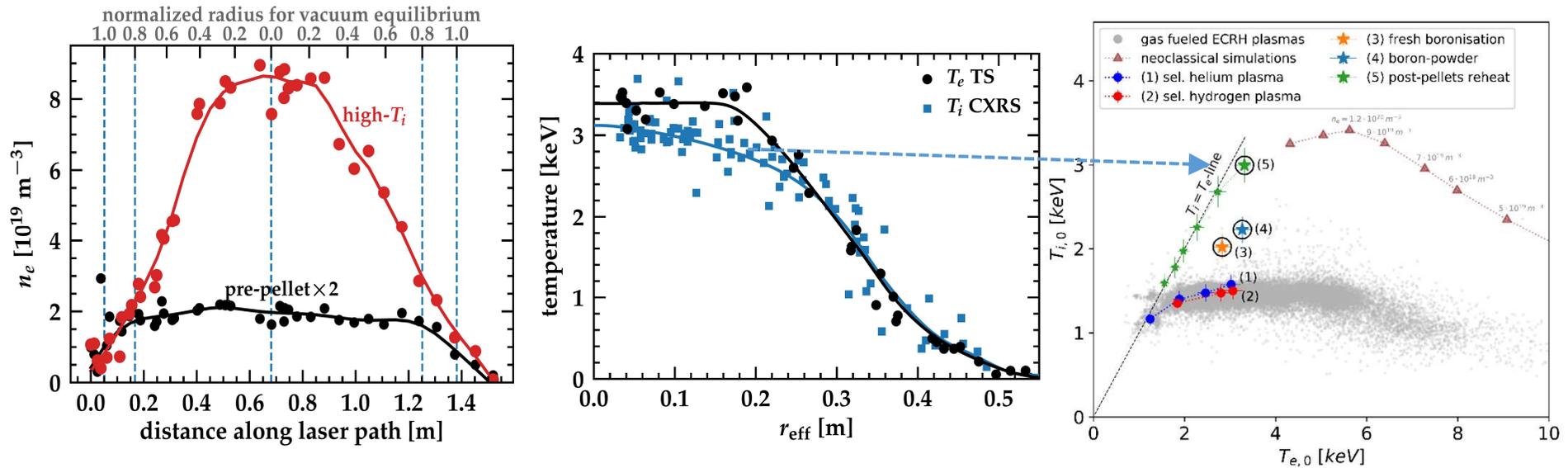
For tokamaks:

- Strong limitation on L-mode reactor (constraint on size!)
- H-mode operation may be a necessity (pedestal lifts core)
- Also beneficial effect of fast particle stabilisation and/or neg. triangularity concepts to be explored

For Stellarators (HELIAS):

- H-modes with good confinement still to be found. (J. Geiger this conference)
- Core turbulence suppression is possible through density gradient: “Stability valey” thanks to maximum-J configuration (T.S. Pedersen this conference)

Example of turbulence suppression scenario

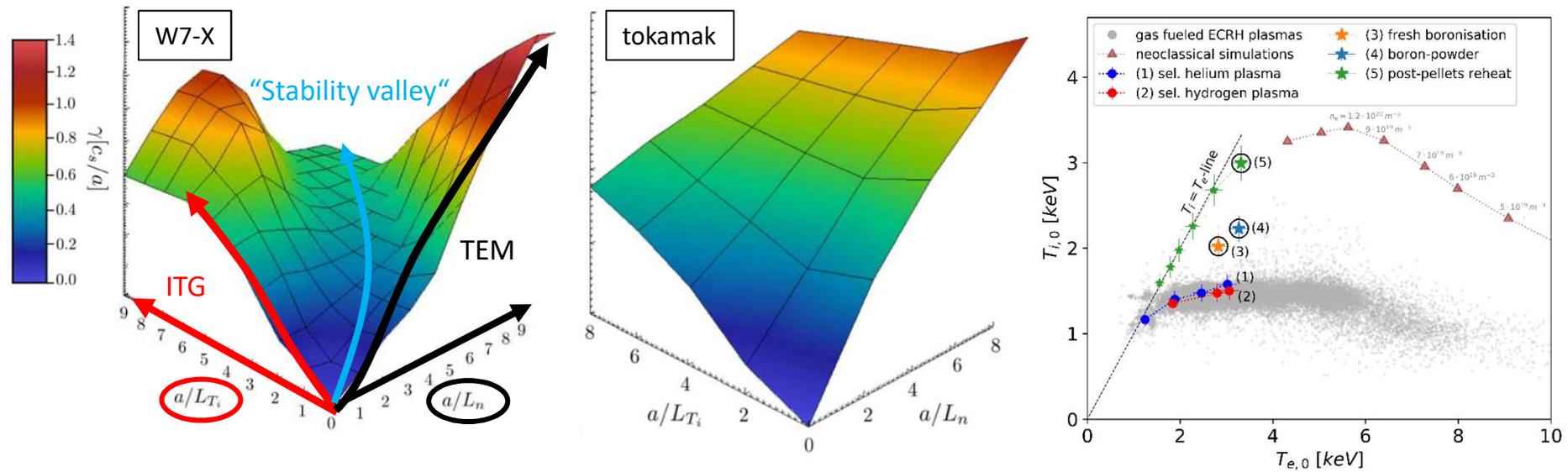


Post pellet experiments confirm reduced turbulence scenario:

- Enhanced density gradients $a/L_n \sim a/L_{T_i}$ help suppress ITG turbulence
- T_i clamping is (transiently) broken to get $T_i = 3$ keV

In W7-X we will attempt to make this a steady scenario and explore enhanced performance with new cryo pumps and a new continuous pellet injector

Example of turbulence suppressed scenario

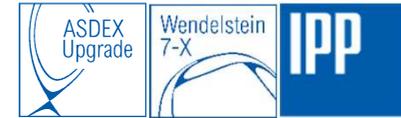


Wendelstein 7-X vs tokamak stability map

- Linear gyrokinetic simulations show **“Stability valley”**

Post pellet experiments confirms reduced turbulence scenario: $T_i = 3\text{keV}$

Key findings of this work



In electron heated L-mode plasmas of ASDEX Upgrade and Wendelstein7-X, the ion temperature clamps at $T_i \sim 1.5$ keV, independent of magnetic configuration:

- The ions are heated by energy exchange $p_{ei} \sim n_e^2 (T_e - T_i)/T_e^{3/2}$, which offers a broad ion heating profile only where $T_e/T_i > 1$.
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