

Isotope identity experiments in JET with ITER-like wall

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**See the author list of 'Overview of JET results for optimising ITER operation' by J. Mailloux et al. to be published in Nuclear Fusion Special issue: Overview and Summary Papers from the 28th Fusion Energy Conference (Nice, France, 10-15 May 2021)*

Isotope identity experiment



- Dimensionless expts test the invariance of plasma physics to changes in the dimensional parameters when the canonical dimensionless plasma physics parameters are conserved:

[Connor and Taylor, NF 1977]

$$\rho_i^* \sim \frac{\sqrt{m_i T_i}}{a B} \quad \beta \sim \frac{n T}{B^2} \quad v^* \sim \frac{n_e R Z_{eff}}{T_e^2} \quad q \sim \frac{B}{R j} \quad \text{Mach} \sim \frac{\sqrt{m_i} v_{tor}}{\sqrt{k T_e}}$$

- $\Omega_i \tau_{E,th} \sim \rho^{*\alpha_\rho} \beta^{-\alpha_\beta} v^{*\alpha_v} q^{-\alpha_q} M^{-\alpha_M} A^{-\alpha_A} Z^{-\alpha_Z} \dots$

[Luce, PPCF 2008]

$$A = m_i / m_p$$

$$\Omega_i = eB / A$$

- Plasma boundary effects (neutrals recycling, impurities...) not included → could potentially invalidate approach
- Isotope identity:** exploiting the change in isotope mass A , achieve discharges with matched dimensionless profiles in the same tokamak:

[Cordey et al., PPCF 2000]

$$\rightarrow B \ \& \ I_p \sim A^{3/4}; \quad n \sim A; \quad T \sim A^{1/2} \quad ; \quad \omega_{tor} \sim A^{-1/4}$$

Isotope identity not trivially expected a priori



- Although isotope mass appears explicitly only in ρ_i^* and Mach-number, changing A in experiment affects all plasma kinetic profiles:
 - Numerous plasma parameters and transport processes have isotope mass dependence [Weisen et al., J Plasma Phys 2020]
 - Operational effects are impacted by changes in isotope mass (NBI, RF heating)
 - → isotope identity not trivially expected a priori

- Isotope identity obtained in JET with C-wall in H and D with type I ELM H-mode
 - Profile similarity achieved over entire plasma radius [Cordey et al., PPCF 2000]



JET-ILW

- L-mode isotope identity in H and D
- Type I ELMy H-mode isotope identity in H and D
- Conclusions and outlook



*Isotope identity technique revisited in **JET-ILW**, with the addition of:*

- *Improved edge kinetic profiles → H-mode pedestal*
- *Sought similarity in Mach number*
- *Investigated role of v_{tor} and ExB shear on core transport*



Hydrogen and Deuterium L-mode

In JET-ILW:

- $\tau_{E,th} \sim A^{0.15 \pm 0.02} P_{abs}^{-0.63 \pm 0.02}$
- Edge particle and heat transport larger in H than D

[Maggi et al., PPCF 2018]

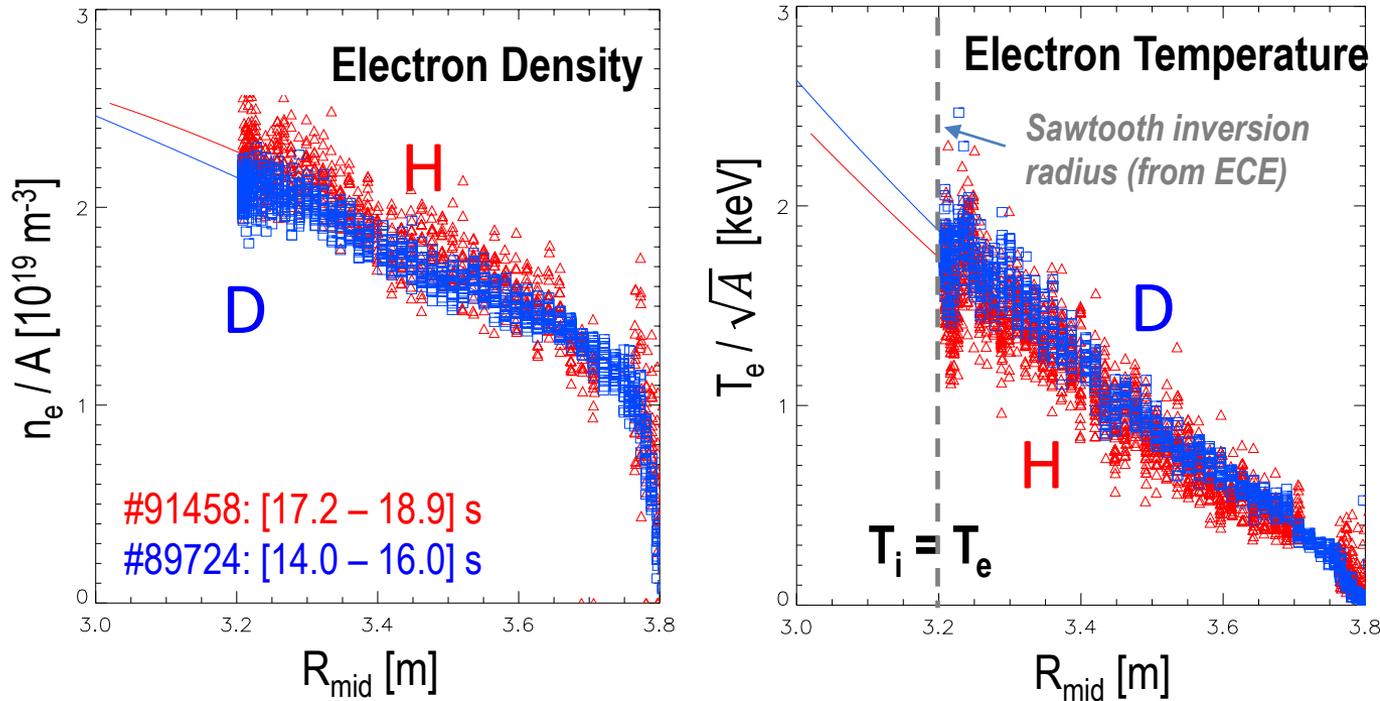
[Bonanomi et al., NF 2019]

H & D L-mode pair with good match of scaled profiles

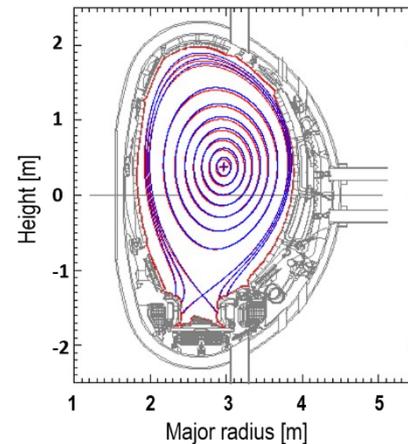


D: #89724 (3.0T/2.5MA) and **H: #91458 (1.74T/1.44MA)**
 Isotope purity $\geq 98\%$

Composite HRTS profiles over steady time window of discharge



Strike points on divertor vertical targets \rightarrow maximizes L-mode domain



- \rightarrow Similar density peaking in H and in D at same ρ^* , v^* , β , q
- Both pinch and NB particle source contribute to core n_e peaking (flux driven predictive modelling)

[Maggi et al., NF 2019]

Similar NB heating profiles in H and D

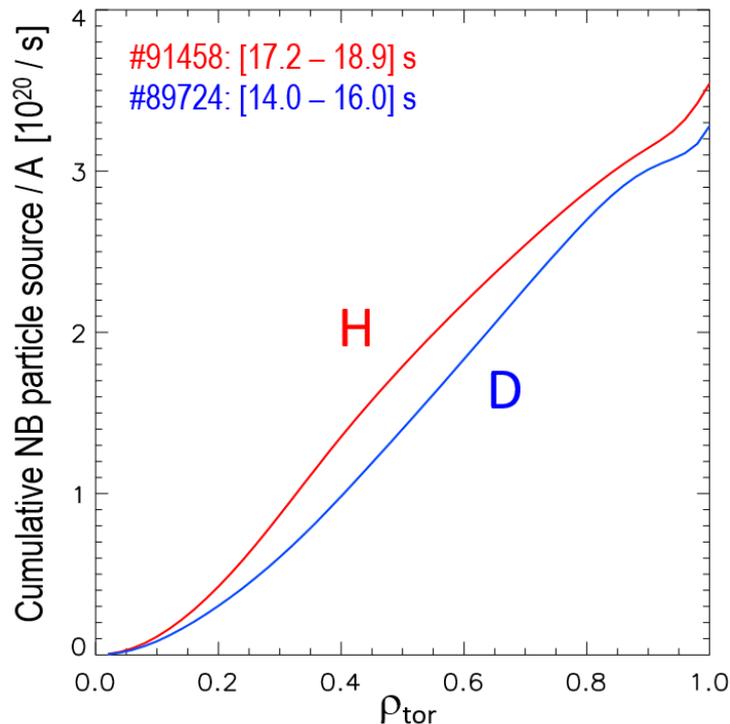
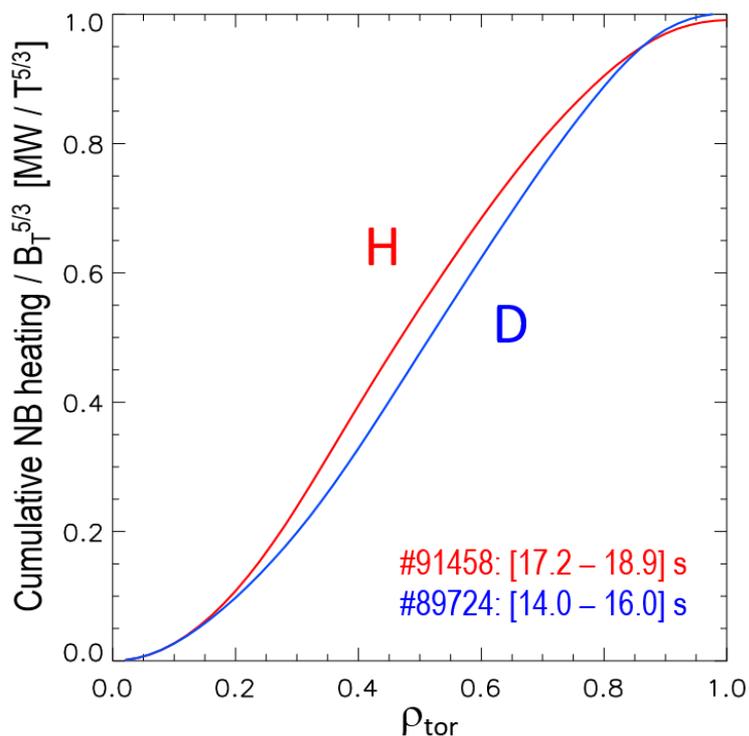


D: #89724 (3.0T/2.5MA) and

Beam energy: 82 – 91 keV; $P_{\text{NBI}} = 6.24$ MW

H: #91458 (1.74T/1.44MA)

Beam energy: 64 – 71 keV; $P_{\text{NBI}} = 2.56$ MW



TRANSP/NUBEAM, with $T_i = T_e$

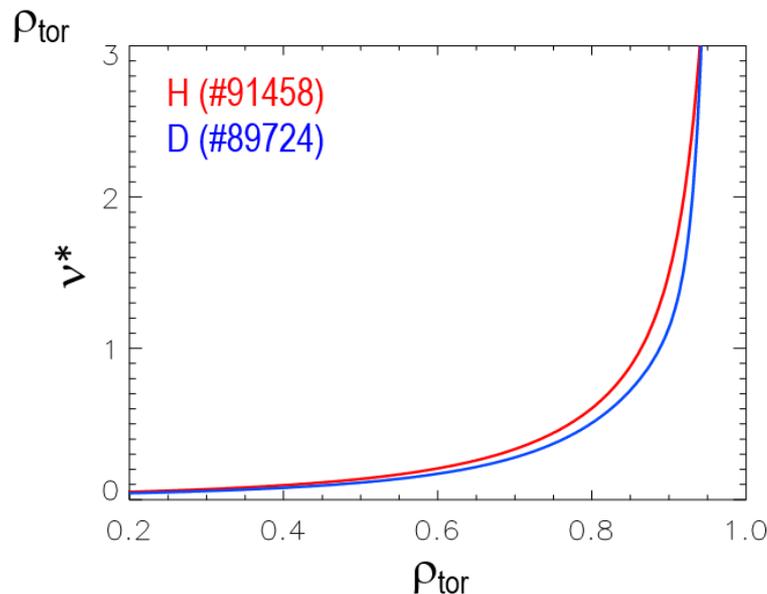
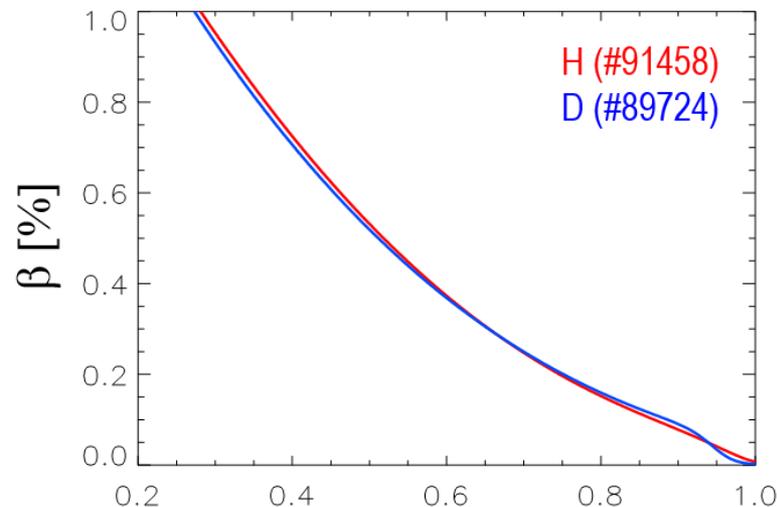
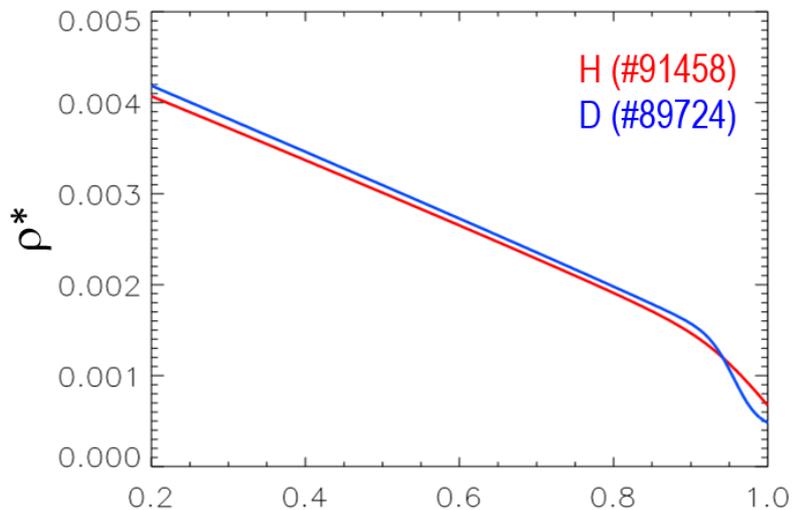
Similar NB heating profiles in H and D

Larger NB particle source in H plasma core

Dimensionless profiles matched in H and D pair



D: #89724 (3.0T/2.5MA) and H: #91458 (1.74T/1.44MA)



[Maggi et al., NF 2019]

H & D isotope identity achieved in confinement region



[Maggi et al., NF 2019]

Pulse # / Isotope	H: #91458	D: #89724
Time interval [s]	17.2 – 18.9	14.0 – 16.0
B_T [T] / I_p [MA] / q_{95}	1.74 / 1.44 / 3.4	2.95 / 2.46 / 3.4
P_{abs} [MW] ($\pm 10\%$)	2.56	6.24
$\tau_{E,th}$ [s] ($\pm 10\%$)	0.155	0.19
T_i / T_e	1.0	1.0
$P_{abs} / B_T^{5/3}$ [MW/T ^{5/3}]	1.02	1.03
Z_{eff} ($\pm 10\%$)	1.4	1.35
$\Omega_i \tau_{E,th}$ [T s]	0.27	0.28

Input power required for L-mode isotope identity: $P_{abs} \sim B_T^{5/3}$

$\Omega_i \tau_{E,th}$ identical in H and D

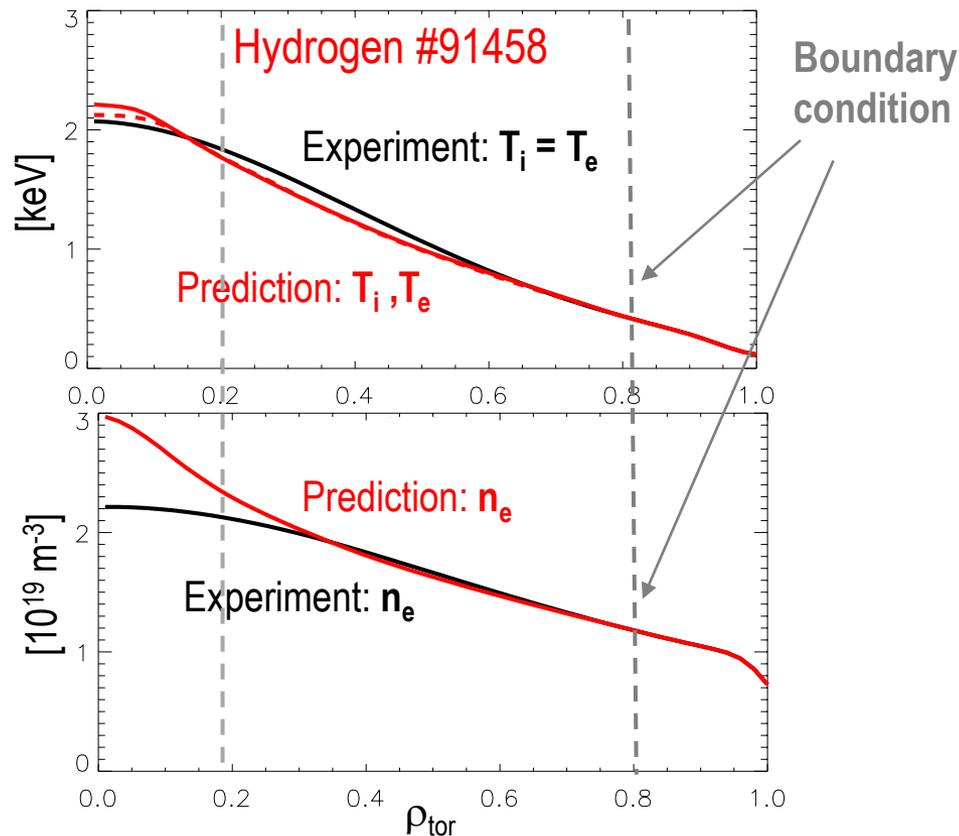
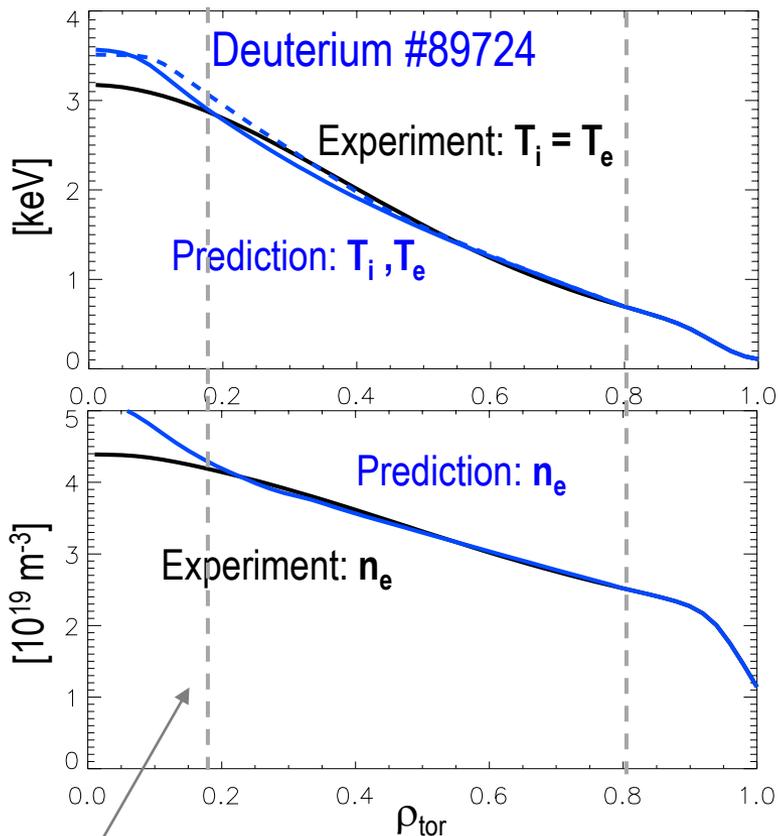
→ Confinement scale invariance principle satisfied in L-mode confinement region

- Scaled thermal energy confinement independent of isotope mass: $\Omega_i \tau_{E,th} \sim A^{0.05 \pm 0.1}$

Flux driven JETTO-TGLF (SAT1) predictive modelling



- Very good agreement with experiment for both isotopes:
- Stiff core heat transport offsets
 - Local gyro-Bohm scaling

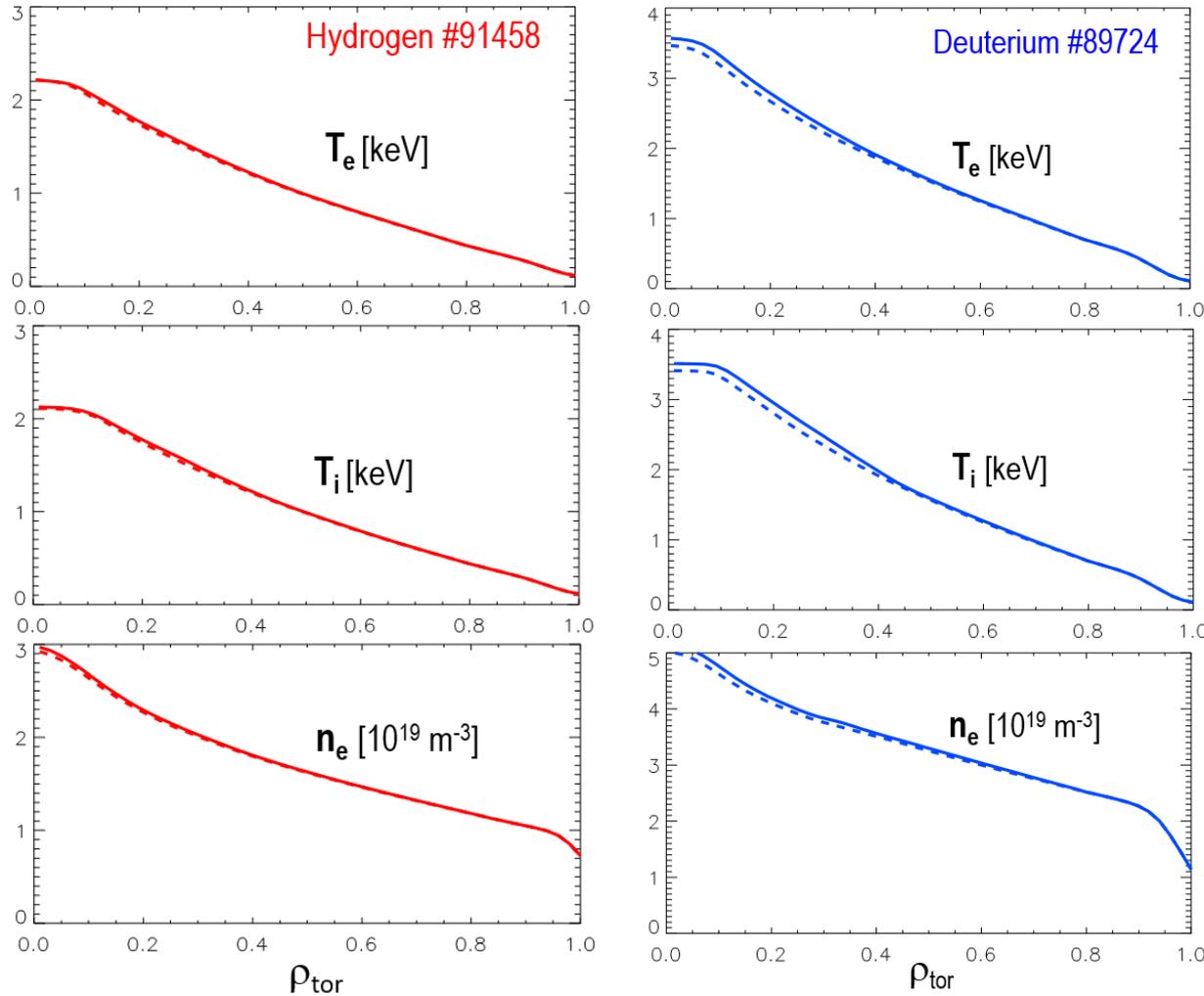


Sawtooth inversion radius

- ITG turbulence dominates for $\rho_{\text{tor}} < 0.8$ (TGLF)
- Effects of collisions, ExB shearing included
- No sawtooth model used

[Maggi et al., NF 2019]

Negligible effect of ExB shear on core transport



JETTO-TGLF:
 — with ExB shear
 - - - w/o ExB shear

v_{tor} and ExB shearing very low \rightarrow do not affect heat and particle transport of H and D L-mode



Hydrogen and Deuterium Type I ELMy H-mode at moderate beta

In JET-ILW:

- Energy, particle, momentum confinement $\sim A^{0.5}$
- Strong, favourable isotope mass dependence of energy confinement at pedestal
 - Primarily in particle channel (likely from inter-ELM particle transport)

[Maggi et al., PPCF 2018]

[Weisen et al., J Plasma Phys 2020]

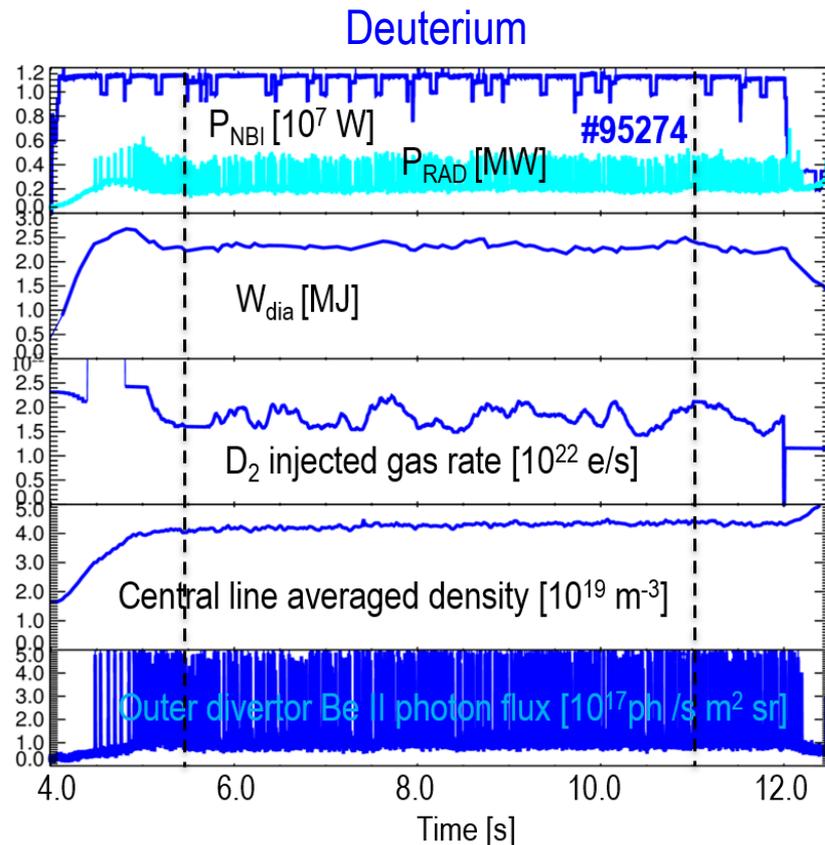
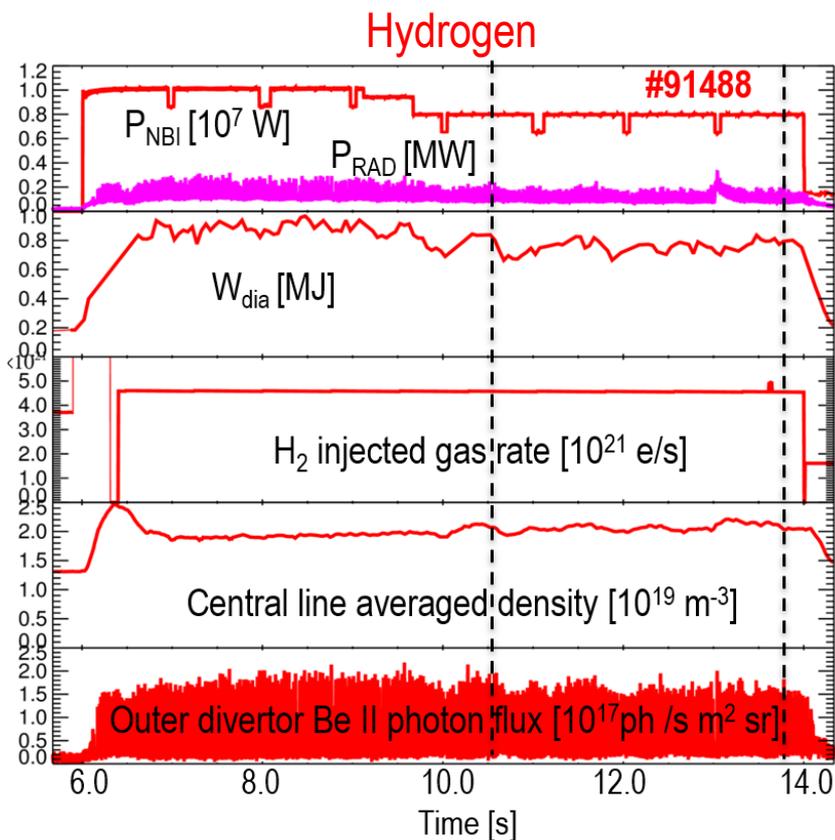
[Horvath et al., NF 2021]

[Schneider et al., this conference]

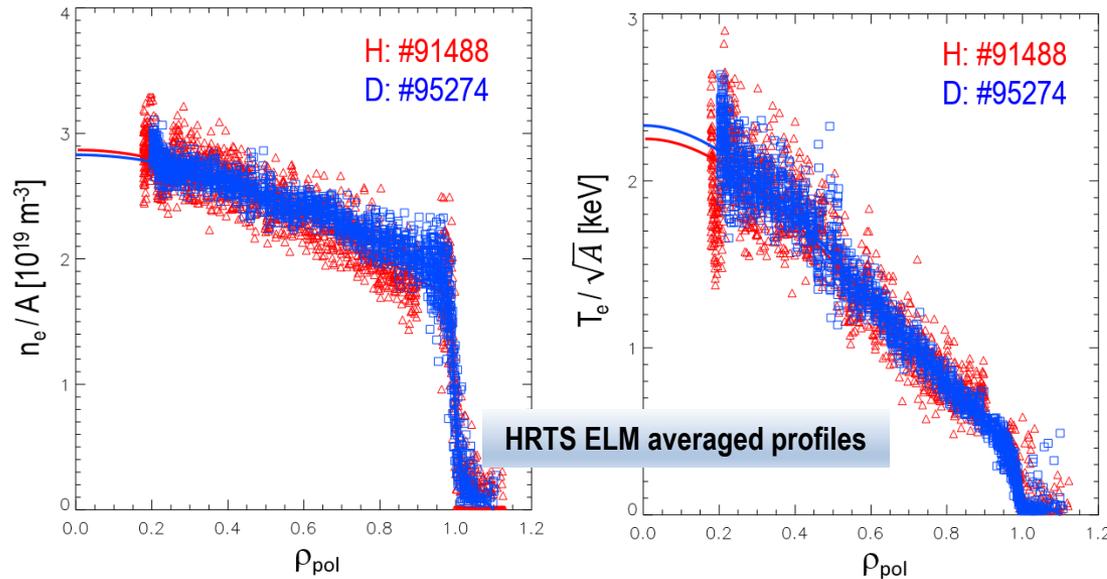
Type I ELMy H-mode pair in H and D



- Varied parameters:
 - Input NBI power and beam energies
 - Injected gas rate
 - f_{ELM} in D, via feedback control of D_2 gas rate \rightarrow same f_{ELM}/Ω_i in H and D \rightarrow match of $n_{e,PED}/A$

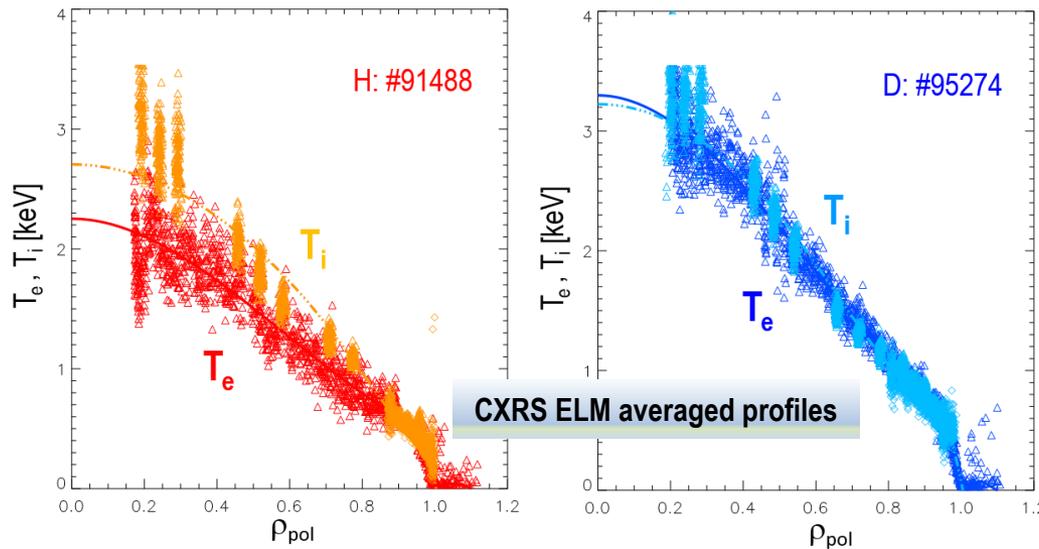
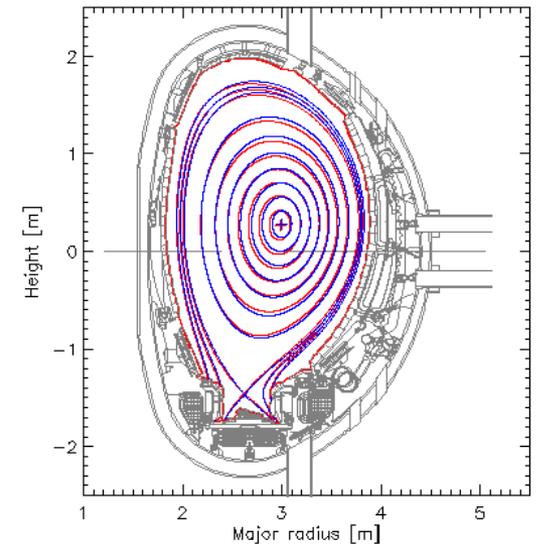


Similar ELM-averaged profiles in plasma core and pedestal



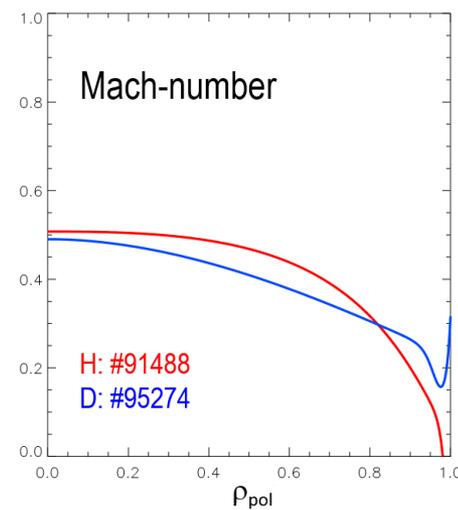
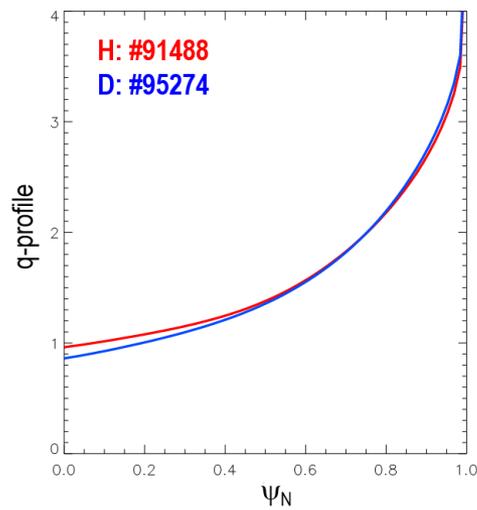
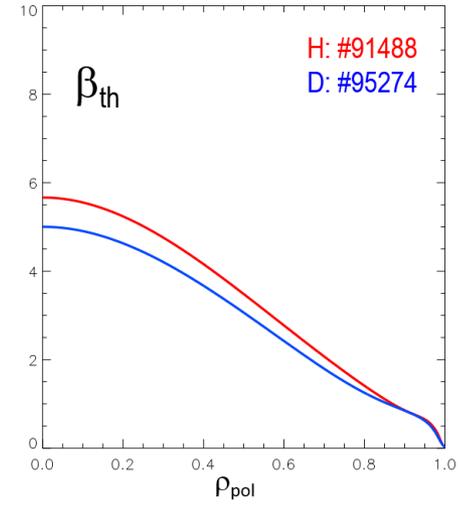
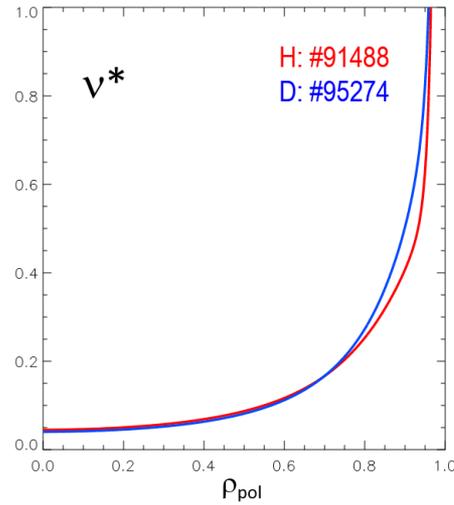
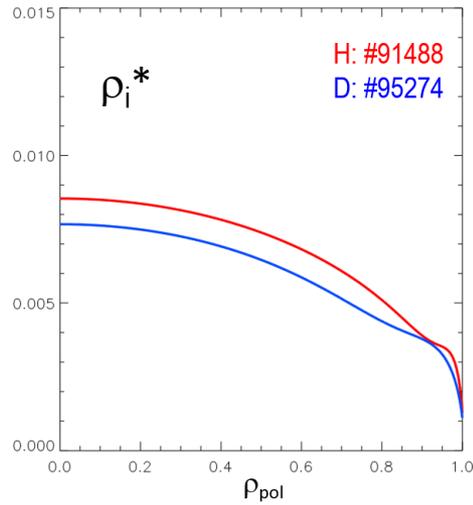
H: [10.5 – 13.8] s = $31 \times \tau_{E,\text{th}}$
D: [5.5 – 11.0] s = $31 \times \tau_{E,\text{th}}$

Same plasma shape - Strike points in divertor corner configuration for improved pumping and density control



- $T_i = T_e$ in D
- $T_i \sim 1.2 T_e$ in core plasma of H

Similar dimensionless profiles in H and D pair



$\Omega_i \tau_{E,th}$ not identical in H and D dimensionless pair



Pulse # / Isotope	H: #91458	D: #95274
Time interval [s]	10.5 – 13.8	5.5 – 11.0
B_T [T] / I_P [MA] / q_{95}	1.0 / 1.0 / 3.0	1.7 / 1.7 / 3.0
P_{abs} [MW] ($\pm 10\%$)	7.0	11.5
$\tau_{E,th}$ [s] ($\pm 10\%$)	0.105	0.175
T_i / T_e	1.0 – 1.2	1.0
P_{abs} / B_T [MW/T]	7.0	6.8
Z_{eff} ($\pm 10\%$)	1.4	1.4
$\Omega_i \tau_{E,th}$ [T s]	0.105	0.15
f_{ELM} / Ω_i [Hz / T]	54	54

Input power required for dimensionless pair: $P_{abs} \sim B_T$

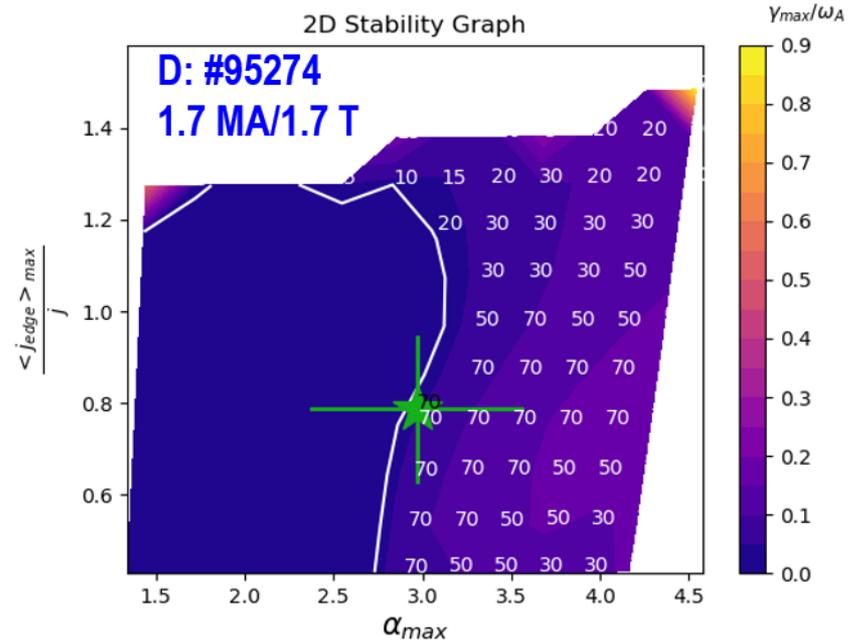
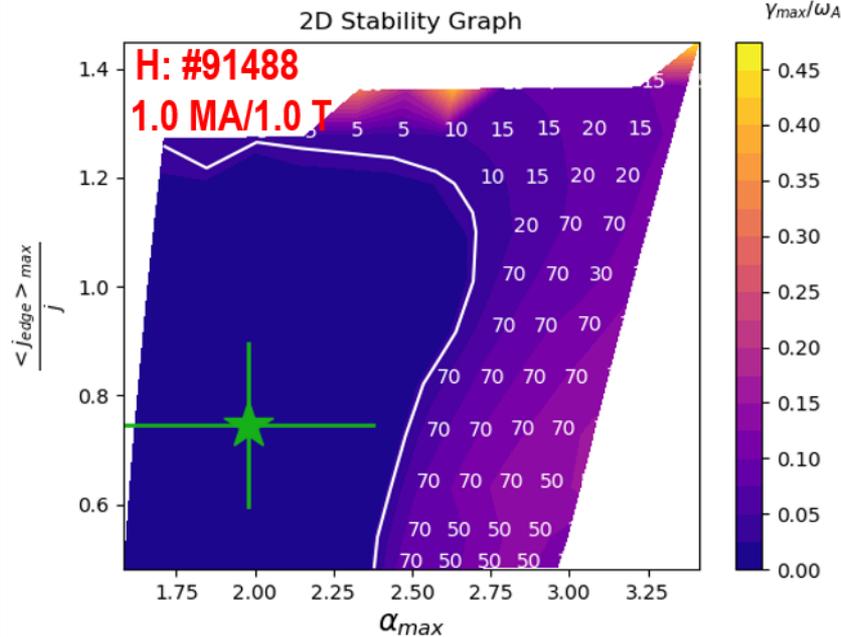
$\Omega_i \tau_{E,th}$ not identical in H and D

Confinement scale invariance principle not satisfied in type I ELMy H-mode at moderate beta

- Scaled thermal energy confinement time: $\Omega_i \tau_{E,th} \sim A^{0.51}$



Stability analysis with HELENA / ELITE



Although type I ELMy, H pedestal is found to be stable to P-B modes

See, e.g [Horvath et al., NF 2021]

D pedestal is at P-B boundary

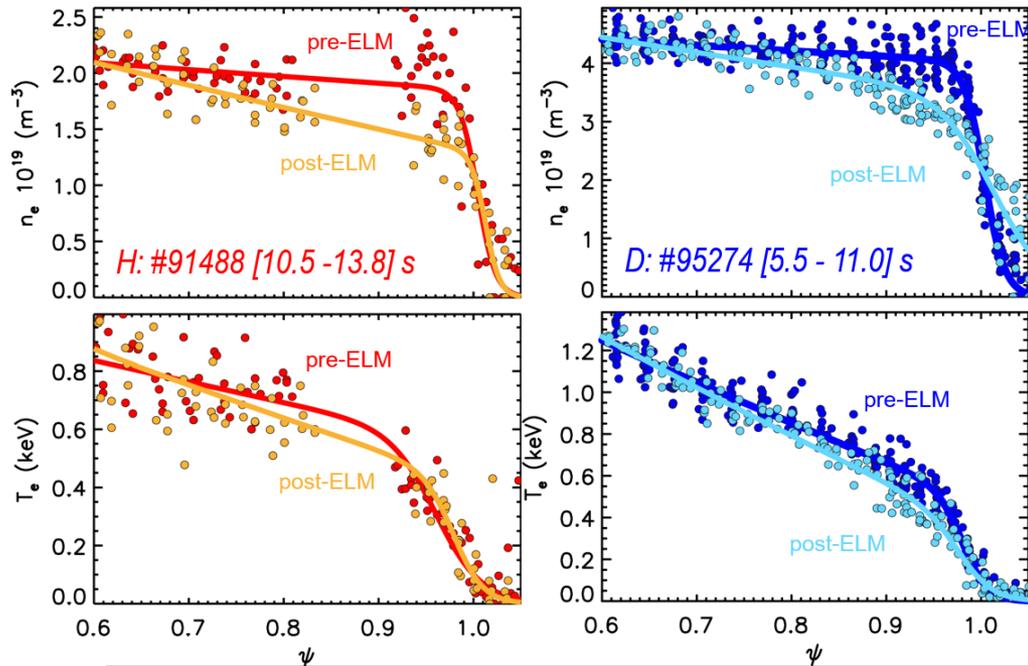
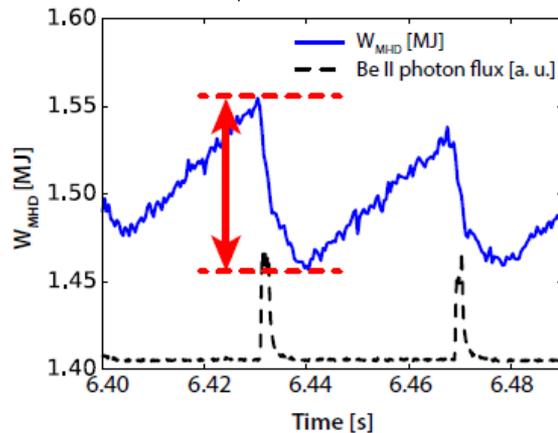
- Scaled pre-ELM pressure gradient not matched in H and D (larger in D) → **pedestal similarity not achieved for pre-ELM profiles**

Fractional pedestal energy and ELM losses similar in H and D

ELM energy losses ΔW_{ELM} derived by:

- Pre- and post-ELM HRTS profiles
- Drop in EFIT W_p at ELM crash

[Horvath et al., NF 2021]



ΔW_{ELM} primarily convective both in H and in D

Isotope / Pulse#	H / #91488	D / #95274
W_{PED} [kJ]	259.1	657
$W_{\text{PED}} / W_{\text{th}}$	0.35	0.33
$\Delta W_{\text{WP}}; \Delta W_{\text{HRTS}}$ [kJ]	36 ± 10 ; 50 ± 9	82 ± 15 ; 103 ± 12
$\Delta W_{\text{WP}} / W_{\text{PED}}; \Delta W_{\text{HRTS}} / W_{\text{PED}}$	0.14 ± 0.04 ; 0.19 ± 0.03	0.12 ± 0.02 ; 0.16 ± 0.02
$P_{\text{ELM (WP)}}$ [MW]	1.94 ± 0.54	3.77 ± 0.69
$P_{\text{ELM}} / P_{\text{abs}}$	0.28	0.33

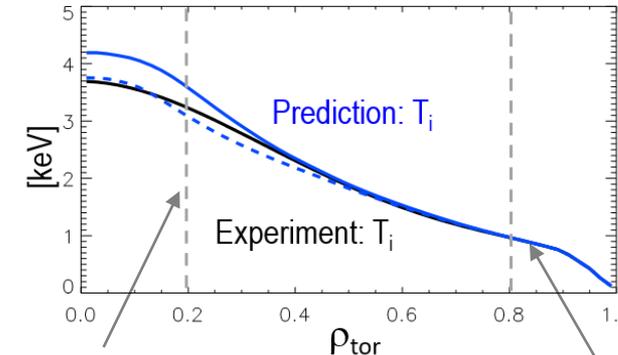
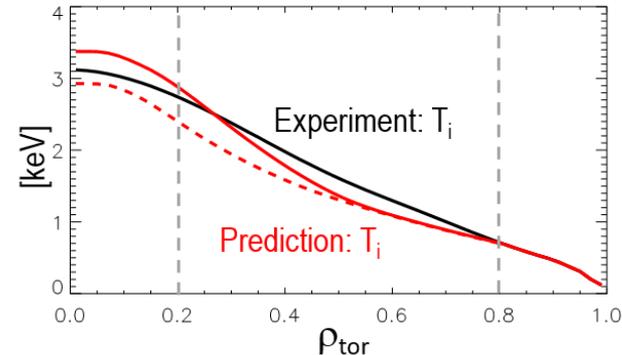
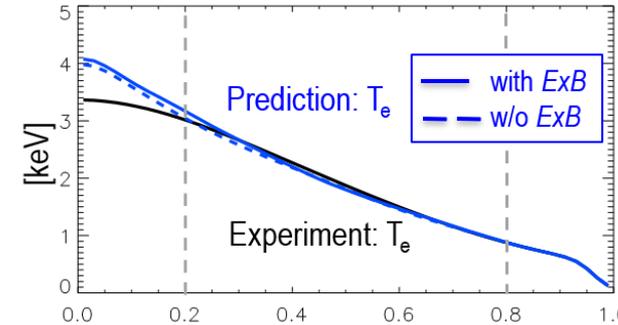
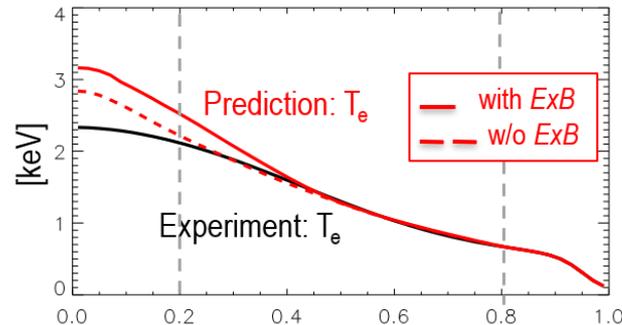
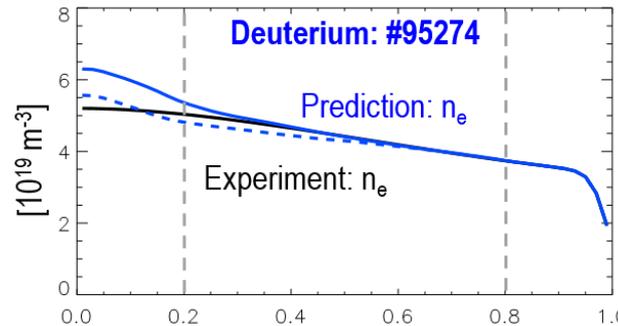
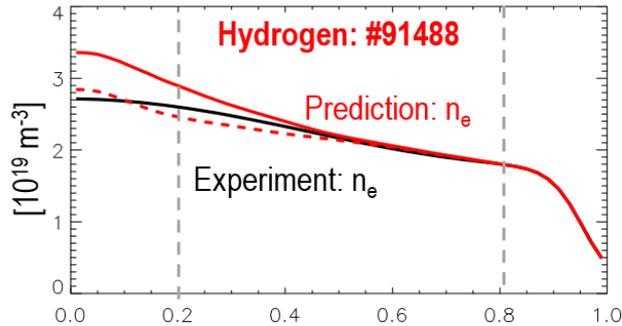
Similar fraction of pedestal to total stored energy in H & D

Similar fraction of P_{loss} carried away by ELMs in H & D

Flux driven TRANSP-TGLF (SAT1) predictive modelling



- Stiff core heat transport offsets local gyro-Bohm scaling of TGLF → consistent with increase of confinement with A, originating in pedestal region



- Weak effect of v_{tor} and ExB shearing on heat and particle transport
- Core ITG stabilization by fast ions pressure gradient not included
 - Expected to be weak, as f.i. content is low in these plasmas

Sawtooth inversion radius

Boundary condition



- Isotope identity achieved with H and D plasmas in JET-ILW in L-mode → invariance principle satisfied in core confinement region
- Similarity achieved for dimensionless ELM-averaged profiles with H and D in type I ELMy H-mode at moderate beta both in core and pedestal
 - but $\Omega_i \tau_{E,th}$ not identical → invariance principle not satisfied
 - scaled pre-ELM pressure gradient not similar in H and D → needs further investigation
- Sought similarity for ion ρ_i^* is (a posteriori) consistent with core transport dominated by ITG turbulence (ρ_e^* not matched) → approach likely not valid for plasmas dominated by electron and/or mixed ion-electron scales turbulence
- Predictive flux driven simulations of core transport in agreement with experiment: stiff core heat transport overcomes local gyro-Bohm scaling of TGLF, explaining
 - Lack of isotope mass dependence of core confinement in L-mode pair
 - Increase of confinement with A in H-mode pair, originating in pedestal region