

## First-time analysis of detached divertor conditions in RMP ELM suppressed H-mode plasma in ITER

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## Motivation: compatibility of detached divertor plasmas with RMPs for ELM control

- Control of divertor loads is required for next step magnetic fusion devices
  - $\rightarrow$  detached plasma operation
- Resonant magnetic perturbations (RMPs) will be used for ELM control in ITER
- A 3D plasma edge model (EMC3-EIRENE) is required to analyze the impact on divertor performance



- Staged Approach: Pre-Fusion Power Operation (PFPO): *P*<sub>SOL</sub> = 30 MW, *B*<sub>t</sub>/*I*<sub>p</sub> = 1.8 T/5 MA, *q*<sub>95</sub> = 3
- Focus on n = 3 RMP field with coil phasing optimized for ELM control



- 1 EMC3-EIRENE: 3D model for the (steady state) plasma boundary
- 2 Plasma response effects on the magnetic topology
- 3 Divertor performance with RMP application
- 4 Sensitivity on assumptions within the plasma response model

## EMC3-EIRENE extends the traditional (axisymmetric) framework for divertor perfomance analysis to 3 dimensions

Magnetic geometry is input for plasma boundary modeling



Boundary plasma is determined by particle, momentum, and energy balances

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 Magnetic geometry is input for plasma boundary modeling and can include plasma response effects (MARS-F, ...)



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## Resistive single fluid calculations (MARS-F) show strong screening of resonant field components

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## Resistive single fluid calculations (MARS-F) show strong screening of resonant field components

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- Ideal magneto-hydrodynamics (MHD) suggests screening of resonant fields
- Strong screening response is recoved in resistive single fluid (MARS-F) calculations
- But plasma response includes field amplification near separatrix
  - ightarrow important for divertor operation











Perturbed separatrix guides field lines from the bulk plasma to divertor targets

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Detached divertor conditions in RMP H-mode plasmas in ITER



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### EMC3-EIRENE simulations show heat flux peaking correlated W with radial connection of perturbed field lines

**•** Model parameters:  $\Gamma_{gas} = 3 \cdot 10^{22} \, \text{s}^{-1}$ ,  $P_{SOL} = 30 \, \text{MW}$ ,  $D_{\perp} = 0.3 \, \text{m}^2 \, \text{s}^{-1}$ ,  $\chi_{\perp} = 1 \, \text{m}^2 \, \text{s}^{-1}$ 



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RMPs: earlier onset of detachment in original strike zone (OSZ)

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## RMPs: significantly different exhaust characteristics at primary and secondary strike locations

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- Particle flux roll-over found in reference (unperturbed) configuration



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- Particle flux roll-over found in reference (unperturbed) configuration
- Early detachment of primary perturbed strike location
- Secondary perturbed strike location remains attached



### Leading role of $T_t$ facilitates parametrization of characteristic $\mathbb{W}$ curves for divertor operation



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Boundary condition sets link between particle and heat loads:

$$\Gamma_t = \frac{q_t}{\gamma T_t}, \qquad Q_t = q_t + \varepsilon \Gamma_t$$

Power losses in the divertor provide link to heat flux from bulk plasma:

$$q_t = (1 - f_{\text{cool}}) \cdot \underbrace{\mathbf{q}_{\parallel} \cdot B_t / B_u \cdot \sin \vartheta}_{= \widehat{\mathbf{q}}}$$

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Parametrization of f<sub>cool</sub>:

$$1 - f_{\text{cool}}(T_t) = \left[ A \left( 1 - e^{-T_t/T^*} \right)^{\alpha} \right]$$

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P.C. Stangeby PPCF 60 (2018) 044022

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• Parametrization of  $f_{cool}$ :



$$1 - f_{\text{cool}}(T_t) = A \left(1 - e^{-T_t/T^*}\right)^{\alpha}$$
 when onset happens  
how fast losses take over after that  
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 Toy model captures roll-over of unperturbed SP



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 Splitting of perturbed heat flux consistent with reduction of peak particle load by factor 2



Primary SP is on low T<sub>t</sub> branch ahead of seconday SP on same (type of) curve



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Primary SP is on low T<sub>t</sub> branch ahead of seconday SP on same (type of) curve
 Higher T<sub>t</sub> consistent with **deeper radial connection** from secondary SP to higher T<sub>u</sub>



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### Footprint size is sensitive to rotation profile in MARS-F



Amplification competes with screening, and depends on assumed rotation profile



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Heat loads may occur far outside of dedicated high heat flux region

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### Optimal ELM control phasings imply large footprint size



- ELM control optimized based on displacement near X-point
- Footprint size s: max. distance from original strike point with connection to  $\Psi_N < 1$



Reduced divertor closure at  $s \gtrsim 70 \text{ cm}$  (soft limit), and extension beyond dedicated high heat flux region (hard limit at full power?)  $\rightarrow$  reliable prediction of plasma response required!



- 1 An earlier transition to detachment is found at the OSZ with RMPs while the far SOL non-axisymmetric SP remains attached
  - $\rightarrow$  high  $T_t$  at secondary SP is problematic for extrinsic impurities required for dissipation at full power

- 2 Non-axisymmetric particle and heat loads during RMP application are sensitive to the **plasma response** (in particular the toroidal rotation used in MARS-F)
  - $\rightarrow$  Optimal coil phasing for ELM control not optimal for divertor operation? (rotation of RMPs possible but should be avoided)
  - ightarrow Reliable predictions for plasma response models required!

### **APPENDIX**



### Plasma response determines radial extent of perturbed SOL

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■ Screening response: smaller, non-overlapping island chains → reduced radial extent of perturbed SOL



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Perturbed SOL: field lines connect to divertor targets



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 $\blacksquare \ Screening \ response: \ smaller, \ non-overlapping \ island \ chains \\ \rightarrow \ reduced \ radial \ extent \ of \ perturbed \ SOL$ 

- Perturbed SOL: field lines connect to divertor targets
  - $\rightarrow$  higher edge temperature can be sustained with RMP screening



### Heat loads may occur far from original strike zone depending 🕅 on plasma response

**•** Model parameters:  $\Gamma_{gas} = 3 \cdot 10^{22} \, \text{s}^{-1}$ ,  $P_{edge} = 30 \, \text{MW}$ ,  $D_{\perp} = 0.3 \, \text{m}^2 \, \text{s}^{-1}$ ,  $\chi_{\perp} = 1 \, \text{m}^2 \, \text{s}^{-1}$ 



Optimal coil phasing for ELM control may not be optimal for divertor operation

### Screening competes with field amplification near separatrix





MARS-F plasma response includes both screening of resonances within the bulk plasma and amplification near the separatrix

### Absence of power and momentum losses confirm: secondary W SP (RMPs) remains attached

Evaluate power and momentum losses with respect to divertor entrance:

$$1 - f_{\rm cool} = q_t/q_u^{\rm (tot.)}$$

$$1 - f_{mom} = p_t^{(tot.)}/p_u^{(tot.)}$$



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### What are the implications of 2ndary SP for divertor operation?

- Initially even weaker  $T_t$  scaling than traditional attached conditions ( $T_t \sim \Gamma_t^{-1}$ ) because of increasing  $\hat{q}$ 
  - ature  $\rightarrow$  increasing  $q_t$  observed at EAST (J. Li, Nature Physics 9 817 (2013) arget temp
- $\blacksquare$  High  $T_t$  increases sputtering on tungsten plates
- Field lines connect deeper into the plasma
  - $\rightarrow$  even higher  $T_{\mu}$  at full power
- From divertor entrance (X-point) to targets: shorter connection length  $\rightarrow$  less dissipation possible (even if  $T_t$  can be brought down)

[V] [eV]



### Linearization of *S<sub>ee</sub>* provides stabilization at low divertor temperatures

First order Taylor expansion allows more accurate treatment of  $T_e$  dependence:

$$S_{ee}^{(j)} \approx S_{ee}(T_e^{(j-1)}) + \left(T_e^{(j)} - T_e^{(j-1)}\right) \left. \frac{dS_{ee}}{dT_e} \right|_{T_e^{(j-1)}}$$

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$$S_{ee}^{(j)} \approx S_{ee}(T_e^{(j-1)}) + (T_e^{(j)} - T_e^{(j-1)}) \frac{dS_{ee}}{dT_e} |_{T_e^{(j-1)}}$$

$$= \text{Source decomposition (linearization)} \\ \text{can be stabilizing if } S_{ee1} \leq 0: \\ S_{ee} \approx S_{ee0} + T_e \cdot S_{ee1} \\ \text{explicit method} \\ \text{implicit method}$$

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Particle balance (n: plasma density)

Y. Feng et al. PPCF 59 (2017) 034006

 $\nabla \cdot \left[ n u_{\parallel} \mathbf{e}_{\parallel} - D_{\perp} \mathbf{e}_{\perp} \mathbf{e}_{\perp} \cdot \nabla n \right] = S_{\rho}$ 

 $D_{\perp}$ : anomalous cross-field diffusion,

 $S_p$ : ionization of neutral particles



Particle balance (*n*: plasma density)

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 $\nabla \cdot \left[ n u_{\parallel} \mathbf{e}_{\parallel} - D_{\perp} \mathbf{e}_{\perp} \mathbf{e}_{\perp} \cdot \nabla n \right] = S_{\rho}$ 

Momentum balance ( $u_{\parallel}$ : fluid velocity parallel to magnetic field lines)

$$\nabla \cdot \left[ m_i n u_{\parallel}^2 \mathbf{e}_{\parallel} - \eta_{\parallel} \mathbf{e}_{\parallel} \mathbf{e}_{\parallel} \cdot \nabla u_{\parallel} - D_{\perp} \mathbf{e}_{\perp} \mathbf{e}_{\perp} \cdot \nabla \left( m_i n u_{\parallel} \right) \right] = -\mathbf{e}_{\parallel} \cdot \nabla n \left( T_e + T_i \right) + S_m$$

 $\eta_{\perp} = m_i n D_{\perp}$ : anomalous cross-field viscosity,  $S_m$ : interaction (CX) with neutral particles



Particle balance (*n*: plasma density)

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Energy balance  $(T_e, T_i:$  electron and ion temperature)

$$\nabla \cdot \left[ \frac{5}{2} T_{e} \left( n u_{\parallel} \mathbf{e}_{\parallel} - D_{\perp} \mathbf{e}_{\perp} \mathbf{e}_{\perp} \cdot \nabla n \right) - \left( \kappa_{e} \mathbf{e}_{\parallel} \mathbf{e}_{\parallel} + \chi_{e} n \mathbf{e}_{\perp} \mathbf{e}_{\perp} \right) \cdot \nabla T_{e} \right] = + K \left( T_{i} - T_{e} \right) + S_{ee} + S_{e,imp}$$

$$\nabla \cdot \left[ \frac{5}{2} T_{i} \left( n u_{\parallel} \mathbf{e}_{\parallel} - D_{\perp} \mathbf{e}_{\perp} \mathbf{e}_{\perp} \cdot \nabla n \right) - \left( \kappa_{i} \mathbf{e}_{\parallel} \mathbf{e}_{\parallel} + \chi_{i} n \mathbf{e}_{\perp} \mathbf{e}_{\perp} \right) \cdot \nabla T_{i} \right] = - K \left( T_{i} - T_{e} \right) + S_{ei}$$

 $\chi_e, \chi_i$ : anomalous cross-field transport  $S_{e...}$ : interaction with neutral particles and impurities



Particle balance (*n*: plasma density)

Y. Feng et al. PPCF 59 (2017) 034006

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Energy balance  $(T_e, T_i:$  electron and ion temperature)

$$\nabla \cdot \left[\frac{5}{2}T_{e}\left(nu_{\parallel}\mathbf{e}_{\parallel} - D_{\perp}\mathbf{e}_{\perp}\mathbf{e}_{\perp}\cdot\nabla n\right) - \left(\kappa_{e}\mathbf{e}_{\parallel}\mathbf{e}_{\parallel} + \chi_{e}n\mathbf{e}_{\perp}\mathbf{e}_{\perp}\right)\cdot\nabla T_{e}\right] = +K\left(T_{i} - T_{e}\right) + S_{e,imp}$$
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 $\chi_e, \chi_i$ : anomalous cross-field transport  $S_{e...}$ : interaction with neutral particles and impurities

### Energy losses are very sensitive at low temperature

EIRENE: kinetic transport of neutral particles

 $\rightarrow n_H,\ n_{H_2},\ n_{H_2^+}$ 

convolution of atomic and molecular processes



### Energy losses are very sensitive at low temperature





### Energy losses are very sensitive at low temperature







$$S_{ee}^{(j)} pprox S_{ee}ig(T_e^{(j-1)}ig)$$

Temperature dependence of S<sub>ee</sub> is not treated accurately enough at low T<sub>e</sub> relevant for detachment



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$$S_{ee}^{(j)} pprox S_{ee}ig(T_e^{(j-1)}ig)$$

Temperature dependence of S<sub>ee</sub> is not treated accurately enough at low T<sub>e</sub> relevant for detachment



More accurate treatment of T<sub>e</sub> dependence based on first order Taylor expansion:

$$S_{ee}^{(j)} \approx \left| S_{ee}^{(j-1)} - T_e^{(j-1)} \right| + \left( T_e^{(j)} - T_e^{(j-1)} \right) \left| \frac{dS_{ee}}{dT_e} \right|_{T_e^{(j-1)}}$$

H-plasma (PFPO, no seeded impurities), model parameters:

- $D_{\perp}~=~0.3\,{
  m m^2\,s^{-1}}$ ,  $\chi_{\perp}~=~1\,{
  m m^2\,s^{-1}}$
- Self-consistent particle balance  $\Gamma_{pump} = \Gamma_{gas} + \Gamma_{core}$  including recirculation below the dome
  - $\rightarrow$  Semi-transparent dome support (50 %) Pumping (0.72 %)
- Include N-N collisions (BGK), molecular assisted recombination (MAR)
- $\blacksquare$  Fueling scan ( $\Gamma_{gas}) \rightarrow$  evaluate neutral pressure & divertor loads



Core fuelling  $\Gamma_{core}$ 



Γαια

Gas puff

Continuous reduction of peak heat flux during gas puff scan



ОТ

Continuous reduction of peak heat flux during gas puff scan



Heat load profiles (Outer target)



Continuous reduction of peak heat flux during gas puff scan



Heat load profiles (Outer target)





- Continuous reduction of peak heat flux during gas puff scan
- A clear roll-over of peak particle flux is found by both codes in good agreement





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