

# model of interchange turbulent transport: on the correlation between scrape off layer width and core confinement in tokamaks

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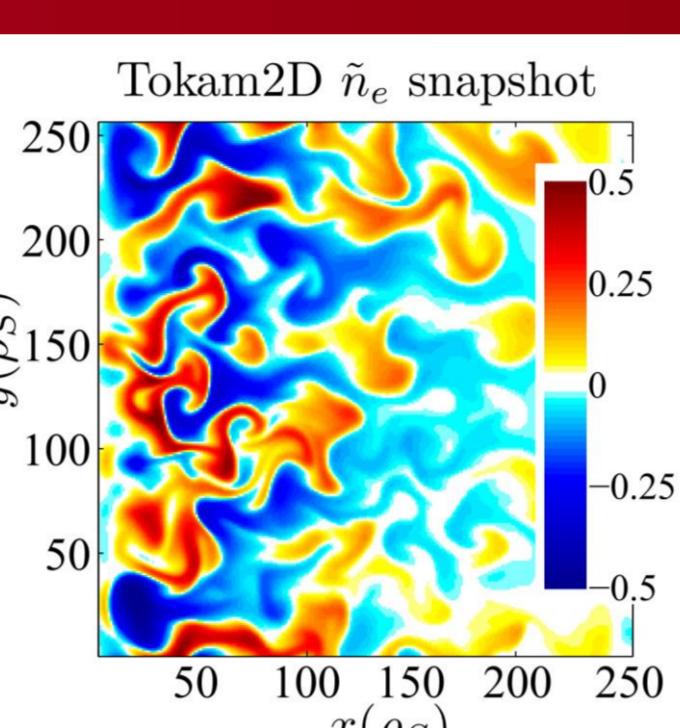
- Spectral filament paradigm: self consistent prediction of turbulent spectra ( $n_e, \Phi$ ) of flux driven 2D isothermal interchange turbulence
- Predicts macroscopic transport properties (flux, fluctuations) in scrape-off layer ( $\rightarrow$  SOL width  $\lambda_q$ ) but also in the core ( $\rightarrow$  confinement time  $\tau_E$ )
- SOL transport: quantitative agreement with Tore Supra data (circular plasmas), recovers impact of divertor leg length on SOL width ( $\lambda_q$ ) in TCV, recovers multi-machine parametric sensitivity of  $\lambda_q$
- Core confinement: recovers multi-machine parametric sensitivity of  $\tau_E$  and correlation between core confinement and heat flux width
- Offers a flexible paradigm for addressing the optimization of power exhaust versus core confinement (key: impact of geometry in model)

## 2D isothermal interchange model (TOKAM2D)

Y. Sarazin et al. Phys. Plasmas 5 (12) (1998) 4214–4228

$$\begin{aligned} \text{Flux driven (source } S\text{), control parameters: } g &\approx \frac{\rho_S}{R}, \sigma_{\parallel} \approx \frac{\rho_S}{q_{cyl} R} \\ d_t n + D\nabla^2 n = S - \sigma_{\parallel} n e^{\Lambda-\Phi} & \quad \text{curvature // loss rate} \\ d_t \omega + \nu \nabla^2 \omega = g \frac{\partial_y n}{n} - \sigma_{\parallel} (1 - e^{\Lambda-\Phi}) & \quad d_t = \partial_t + \{\Phi, \cdot\} \\ \omega = \nabla_{\perp}^2 \Phi & \quad \omega = \nabla_{\perp}^2 \Phi \end{aligned}$$

Intermittent filamentary transport, exponential SOL, etc



## Test isolated filament model

S. I. Krasheninnikov et al. J. Plasma Phys. 2008, 74, 679

Density filament: size  $\delta_f$  & normalized amplitude  $n_f$   
 $\rightarrow$  drift velocity ( $v_f$ ) predicted from vorticity balance

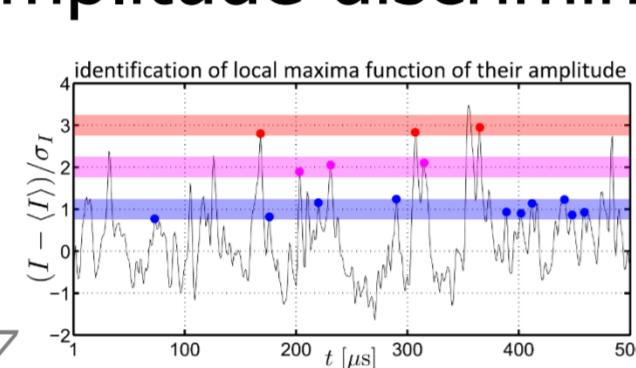
Hyp: well shaped filaments, not the case in simulations

Solution: construction of averaged filaments

$\rightarrow$  conditional average with event amplitude discrimination

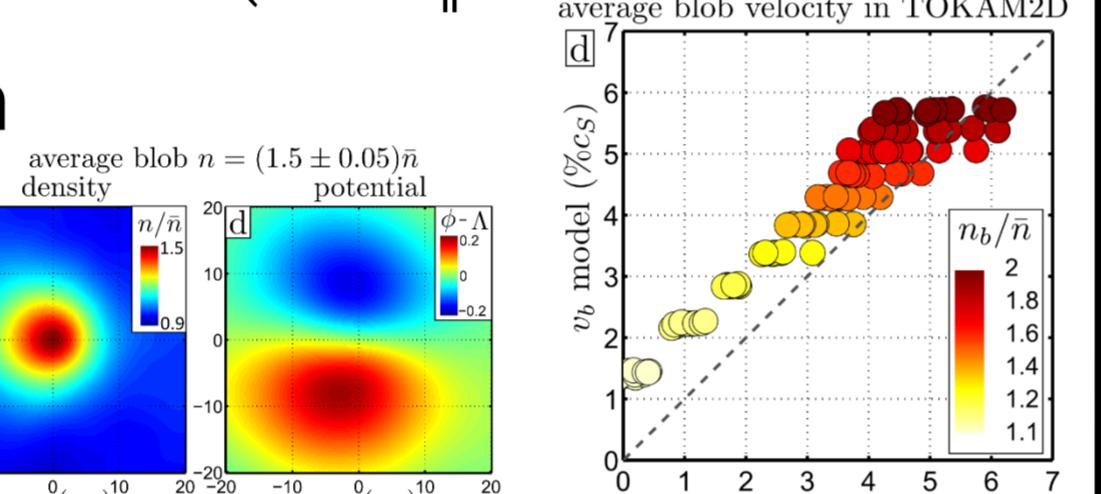
Isolated filament model works on  
TOKAM2D averaged filaments

N. Fedorczak et al. Contrib. Plasma Phys. 2018;58:471–477



$$v_f = \frac{1}{2} \sigma_{\parallel} \delta_b^3 \left[ \sqrt{1 + (\delta_0/\delta_b)^5} + 1 \right]$$

$$\delta_0 = \left( \frac{4g(n_b - 1)}{\sigma_{\parallel}^2} \right)^{\frac{1}{5}}$$



## Spectral filament (SF) model

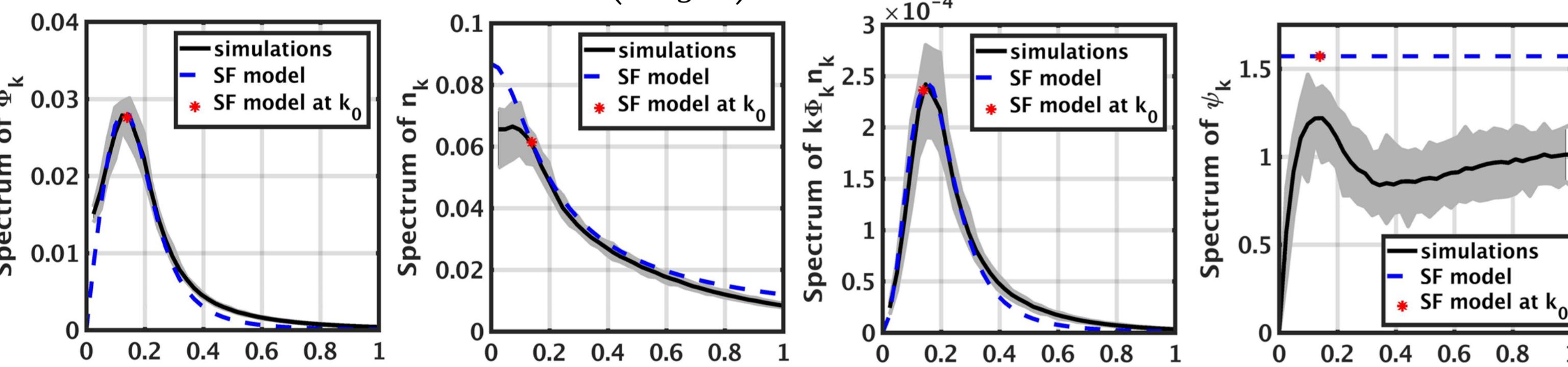
N. Fedorczak et al. Nuc. Mater. & Energy 19 (2019) 433–439

Principle: extend filament model to fluctuation spectra  $v_f(\delta_f) \rightarrow \Phi(k_{\theta})$  &  $n_b \rightarrow n(k_{\theta})$

Hyp: density & potential spectra in phase quadrature + time independent spectra (it is!)

- potential: include effective viscosity from dominant mode  $k^4 \Phi_{k_0} \Phi_k + \sigma_{\parallel} \Phi_k = g k n_k$
- density : mixing length + ad-hoc saturation at low  $k$   $\sqrt{k^2 + k_0^2} n_k = \lambda_n^{-1}$

$$\text{dominant wave number } k_0 = \left( \frac{1.45 \lambda \sigma_{\parallel}^2}{4g} \right)^{1/4}$$



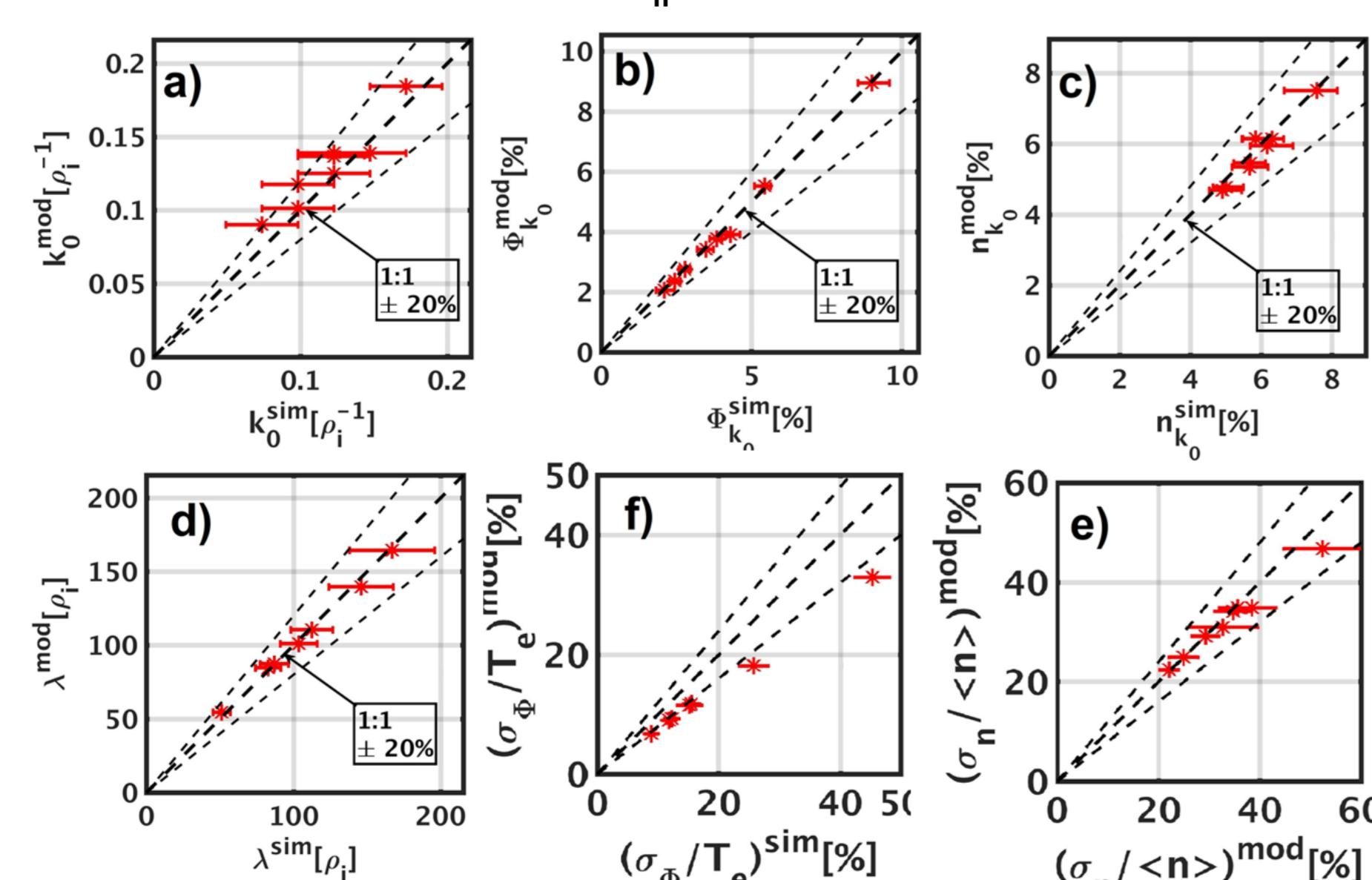
## Spectral integrals for macroscopic quantities:

- radial flux:  $\Gamma_r \equiv 8\pi^2 \int k \Phi_k n_k dk = 41 \lambda_n^{-7/4} g^{3/4} \sigma_{\parallel}^{-1/2} [\bar{n} c_s]$
- fluctuation levels  $\frac{\sigma_n}{\bar{n}} = 12.6 \lambda_n^{-9/8} g^{1/8} \sigma_{\parallel}^{-1/4}$  &  $\frac{\sigma_{\Phi}}{T_e} = 4.8 \lambda_n^{-7/8} g^{7/8} \sigma_{\parallel}^{-3/4}$
- scrape-off-layer width:  $\Gamma_r \equiv \lambda_n \sigma_{\parallel} \rightarrow \lambda_n = 3.9 g^{3/11} \sigma_{\parallel}^{-6/11} [\rho_L]$

$\rightarrow$  Quantitative agreement  
with TOKAM2D simulations  
(scan of  $g$  &  $\sigma_{\parallel}$ )

## Spectral filament model:

- stiff flux model  $\Gamma_r \propto \lambda^{-7/4}$
- not diffusive nor convective
- intermittency disappears in spectral paradigm
- key : model of  $g$  &  $\sigma_{\parallel}$



## SF model $\rightarrow$ prediction of heat flux width & confinement time

### Heat flux decay length

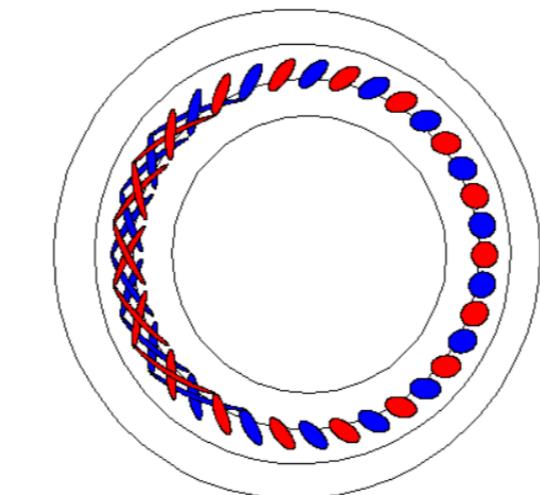
Hyp:  $\lambda_q = \alpha_q \lambda_n$ ,  $g = \rho_L/R$  &  $\sigma_{\parallel} = \frac{\rho_L}{L_{\parallel}} = \rho_L/\pi q_{cyl} R$ , 2-point model for  $T_{sep}$

$$\lambda_q = 50 \alpha_q q_{cyl}^{0.63} B^{[-0.65]} R^{0.25} \left( \frac{A}{Z} \right)^{0.33} P_{sep}[MW] [mm]$$

### Confinement time:

Hyp: density  $\rightarrow$  pressure & particle flux  $\rightarrow$  heat flux  
vorticity // dissipation still valid (HFS decorrelation of LFS mode)  
gradient scale length  $\lambda \approx a$

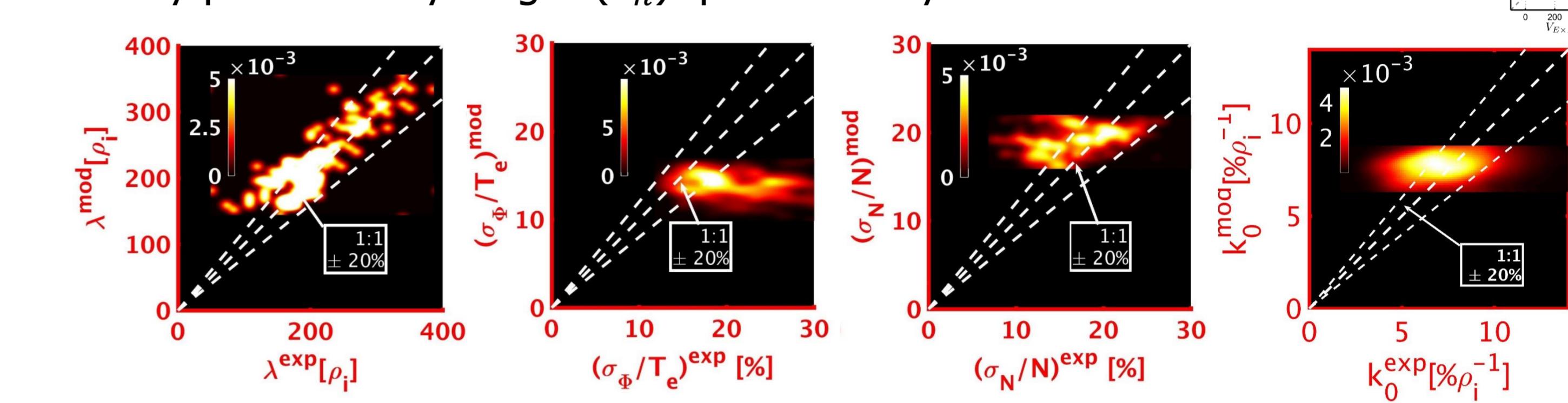
$$\tau_E = 0.007 \alpha_T I_P^{0.2} B^{0.6} P_{TOT}^{-0.6} n_e^{0.6} M^{0.2} R^{2.8} \left( \frac{a}{R} \right)^{1.9} [s]$$



## Model validated against broad SOL database from Tore Supra

>100 probe reciprocations with turbulence (1MHz  $\theta$ -rake) & profile (tunnel) collectors

- ExB velocities of averaged filaments match with isolated filament model
- fluctuation amplitudes ( $n_e$  &  $\Phi$ ) and main  $k_{\theta}$  well reproduced
- density profile decay length ( $\lambda_n$ ) quantitatively matched



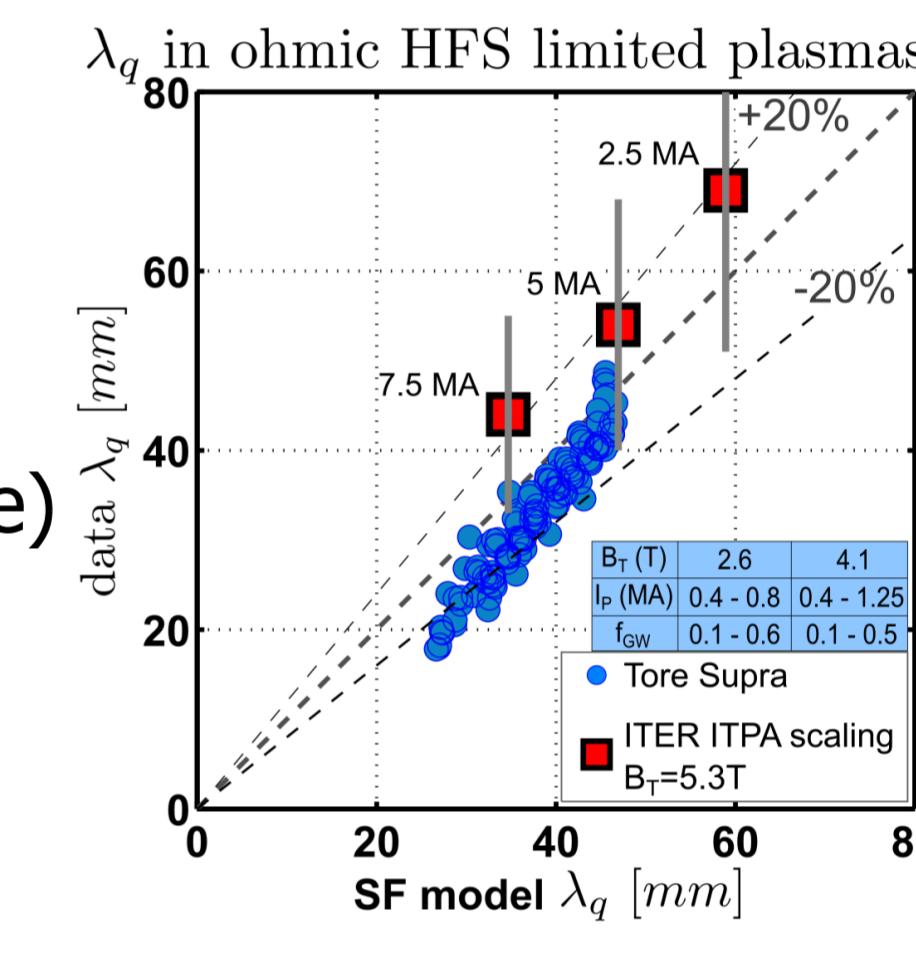
## Recovering multi-machine scaling of ITER startup heat flux width

N. Fedorczak et al. Nuc. Mater. & Energy 19 (2019) 433–439

Hyp:  $\lambda_q = 0.5 \lambda_n$  (correlation from Tore Supra database)

$$\lambda_q = 25 q_{cyl}^{0.63} B^{-0.65} R^{0.25} \left( \frac{A}{Z} \right)^{0.33} P_{sep}[MW] [mm]$$

- quantitative agreement with Tore Supra  $\lambda_q$  (ITER database)  
J. Gunn et al. J. Nucl. Mater. 438 (Suppl) (2013) S184–S188.
- quantitative agreement with multi-machine prediction of ITER startup SOL width  
J. Horacek et al. Plasma Phys. Controlled Fusion 58 (7) (2016) 074005.



## Recovering multi-machine scaling sensitivity of $\lambda_q$ in H-mode

Model:  $\lambda_q \propto R^{0.25} \dots$

- $\lambda_q \propto X^{3/4}$  for ITER & DEMO compared to  $B_{\theta}$  scaling

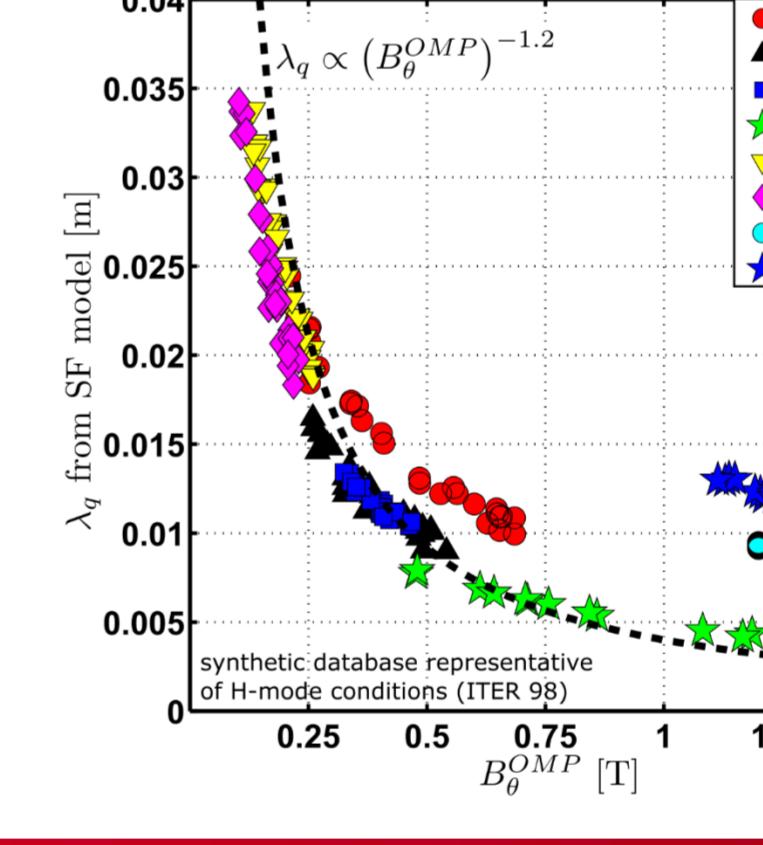
T. Eich et al. Nucl. Fusion 53 (9) (2013) 093013

- sounds with XGC1 results (different interpretation)

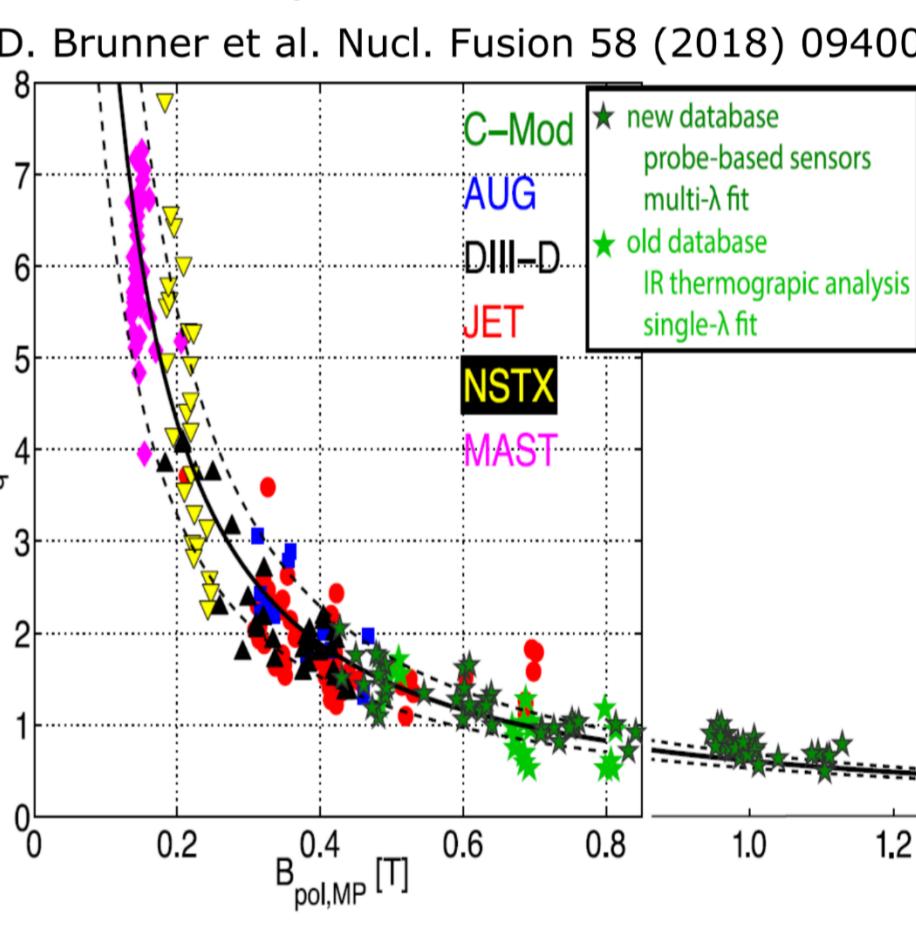
C.S. Chang et al. 2017 Nucl. Fusion 57 116023

### Heat flux decay length predicted by spectral filament model

versus poloidal magnetic field at outer midplane



### experimental data



## Recovering impact of divertor leg length on $\lambda_q$ in TCV L-mode

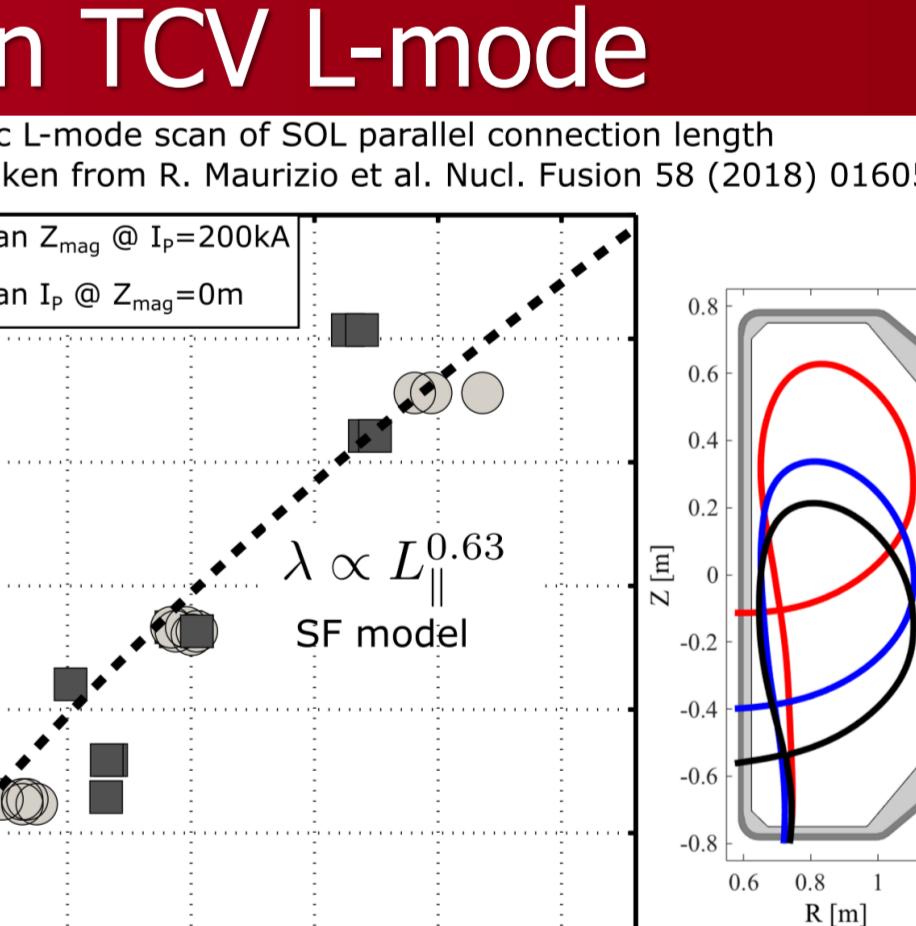
TCV ohmic pulses: scan of plasma current and divertor geometry gives same  $\lambda_q(L_{\parallel})$

R. Maurizio et al. Nucl. Fusion 58 (2018) 016052

Model:  $\lambda_q \propto L_{\parallel}^{0.63}$

- sensitivity recovered by the SF model
- 3D turbulent simulations show interchange along divertor leg, as the model considers

A. Gallo et al. Plasma Phys. Control. Fusion 60 (2018) 014007



## Recovering global multi-machine $\tau_E$ sensitivity & correlation with $\lambda_q$

Model :

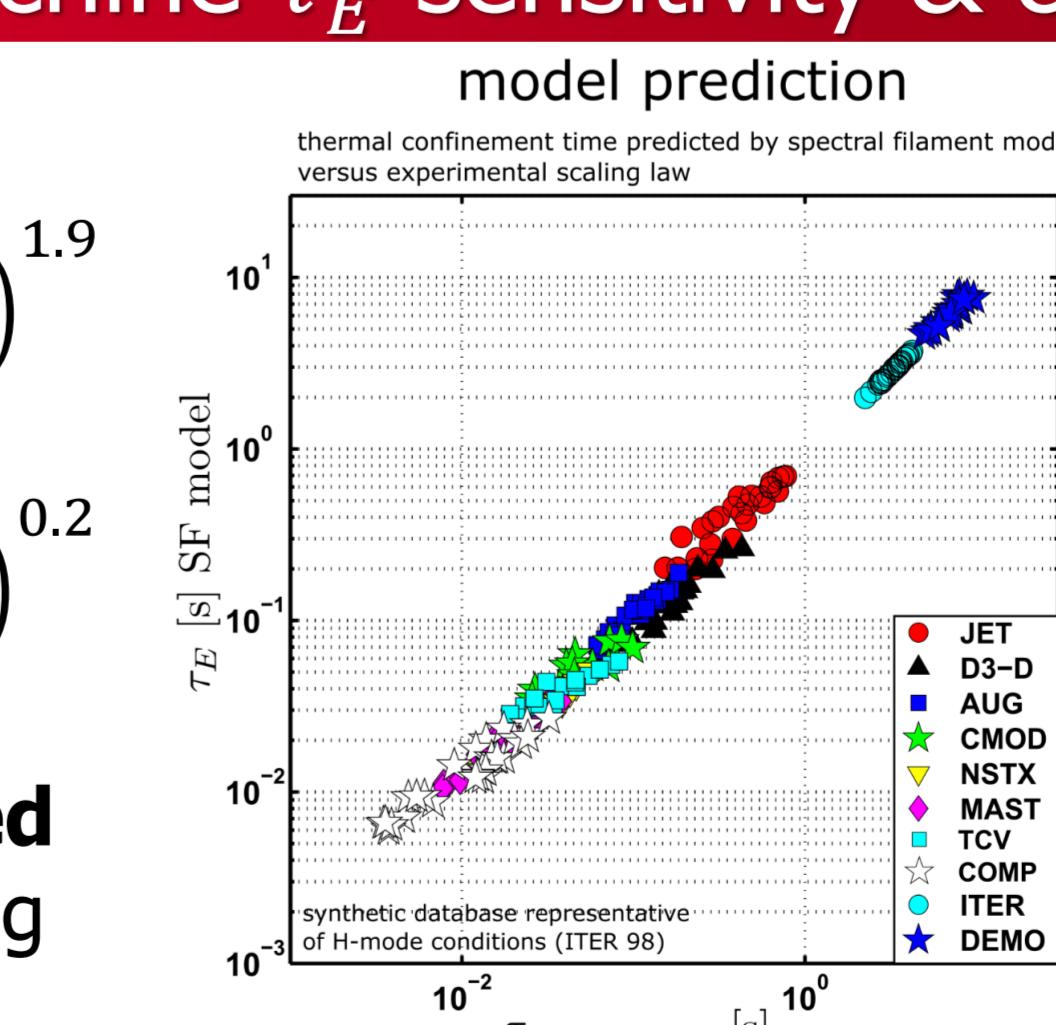
$$\tau_E \propto I_P^{0.2} B^{0.6} P_{TOT}^{-0.6} n_e^{0.6} M^{0.2} R^{2.8} \left( \frac{a}{R} \right)^{1.9}$$

Scaling law ITER 98:

$$\tau_E \propto I_P^{1.0} B^{0.1} P_{TOT}^{-0.6} n_e^{0.4} M^{0.2} R^{1.9} \left( \frac{a}{R} \right)^{0.2}$$

- significant difference in  $I_P$  &  $B$
- **global trend is well matched**

hints of correlations in exp. scaling



### experimental data

