Max-Planck-Institut für Plasmaphysik **Investigation of Ar and N seeded SOLPS 5.0** simulations for ASDEX Upgrade

F. Hitzler^{1,2*}, M. Wischmeier¹, F. Reimold³, D. P. Coster¹ and the ASDEX Upgrade Team⁴

¹Max-Planck-Institut für Plasmaphysik, 85748 Garching, Germany – ²Physik-Department E28, Technische Universität München, 85747 Garching, Germany – ³Max-Planck-Institut für Plasmaphysik, 17491 Greifswald, Germany – ⁴See the author list of "H. Meyer et al. 2019 Nucl. Fusion 59 112014"

1. Introduction: Power Exhaust

- Power exhaust is a key challenge in future tokamaks
- Divertor power loads need to be reduced significantly
- Controlled seeding of impurities like Ar or N
- → Radiative power dissipation: strongly reduced target temperatures and power fluxes
- Main task: maximizing the

4. SOLPS 5.0 Impurity Seeding Scans



As expected:

- Decreasing temperatures (due to radiative power dissipation)
- Stronger impact on the pedestal with Ar, despite rather low Ar core density $(\rightarrow \text{ see radiation efficiency})$
- Ar seeding scan: shifted density distribution at high seeding rates (LFS→HFS)
- No significant change in the N density distribution



radiative power dissipation and minimizing the impact on the confined plasma



 $q \approx 50 \, \mathrm{MW/m^2}$

(geom. effects /

∑a 525

⊢^e 500 [↓]

475

450

top

Ped.

2. Radiation Efficiency: Ar vs. N

- Ar rad. efficiency is higher in hot regions (\rightarrow SOL & core)
- N radiates more efficiently only below ~ 5eV (\rightarrow divertor)



3. SOLPS 5.0 Modeling

Simulations in this work:

- Computational grid based on AUG H-mode shot #29256 [5]
- Electron density at the midplane separatrix: $2.5 \cdot 10^{19} \text{ m}^{-3}$
- Input power (heat flux crossing the core boundary): 5MW
- No drifts terms are activated (challenging due to numerical instabilities)



Low seeding

66 % Ar

90

— 100 % Ar



Mixing Ar and N impurities

- Different mixing ratios: Γ_{Ar} : $\Gamma_{N} = 1:0, 2:1, 1:1, 1:2$ and 0:1
- Less pedestal top temperature drop with higher N fraction
- Lower pedestal top fuel dilution with higher Ar fraction
- Trade-off between pedestal top temperature drop and fuel dilution by mixing both impurities – further studies required to identify "optimum" ratio

5. Argon Impurity Transport & Divertor Retention

Forces acting on the impurities

66 % N

100 % N-

70

Low seeding

Far

SOL

Near

Friction force $F_{fr} \propto (u_{D^+} - u_{imp})$, equalizing impurity and main ion flows

80

Ped. top fuel dilution $\frac{n_{D^+}}{n_e}$ [%]

- Thermal force $F_{th} \propto \nabla T \rightarrow$ deviation of impurity flow from main ion flow
- ∇p and electrostatic forces negligible
- F_{th} induces an equivalent $F_{fr} \approx -F_{th}$ \rightarrow forces well balanced in steady state
- Impurity seeding modifies ∇T , and therefore, F_{th}
- Increasing F_{th} towards the inner divertor at increasing seeding levels





- Ar and N seeding, up to $1.8 \cdot 10^{21} \frac{e^{-1}}{s}$ (electron equivalent)
 - ► Seeding scans as pure "code experiments"

6. Expected Impact of Drifts

- Increased inner / outer target temperature asymmetry with hotter outer target & colder inner target (strongest impact at low densities and at the inner target) [6]
- Formation of the high-field side high density region [7]
- Ionization fronts shifted further away from the target [8]
- Poloidal particle flux in the SOL towards the outer divertor
- Impurity redistribution (as discussed in this contribution) possibly mitigated

7. Conclusions

Seeding scans:

- Lowest impact on the ped. top temperature with N
- Lower fuel dilution with Ar seeding
- Trade-off with mixed impurities further studies required to develop a rule of thumb for an "optimum" mixing ratio

SOL transport:

Inverted main ion flow patterns (due to modified ionization sources) and increasing thermal forces on the impurities





Main ion and impurity flows

- Inverted main ion plasma flow pattern at higher seeding levels
- Caused by modification of the deuterium ionization sources in the divertor regions
- Strongly modified impurity flow pattern (due to friction between main ions and impurities)
- Low seeding: no impurities can move from outer to inner divertor through the SOL
- High seeding: vice versa
- Strong & sudden* modification of the impurity density distribution at the transition *("sudden" in terms of the impurity seeding, i.e., as a function of the seeding level)





Reversed impurity flow & shifted density distribution

Divertor retention:

- Determined by the relative positions of the neutral impurity ionization front and the impurity stagnation point
- Both shifted away from target with increasing seeding
- Competition between both mechanisms
- Preliminary result: shift of ionization front dominates \rightarrow more leakage at higher impurity seeding levels

Impact of drifts:

Will be critical and might (quantitatively) alter the results

Impurity stagnation point

Ionization front position

- Lower $T \rightarrow$ shifted away from target
- More particles reach beyond impurity stagnation point & can escape
- Indicates increasing divertor leakage with higher seeding

 $u_{D^+} = -F_{th}/c_{fr}$ With $F_{th} \propto Z^2 \cdot \nabla T$ and $c_{fr} \propto \frac{nZ^2}{T^{3/2}}$ [9] $\Rightarrow F_{th}/c_{fr} \propto \frac{T^{3/2} \cdot \nabla T}{r}$

Stagnation point shifted away from the target with decreasing temperatures

 $= -c_{fr} (u_{D^+} - u_{imp})$

Indicates enhanced divertor retention with higher impurity seeding level

Competition between divertor retention and leakage due to shifted ionization front and stagnation point positions

References: [1] A. Loarte, et al., Nucl. Fusion 47, 2007, [2] M. Wischmeier, J. Nucl. Mater. 463, 2015, [3] A. Kallenbach, Plasma Phys. Control. Fusion 55, 2013, [4] H. P. Summers, The ADAS User Manual, 2004, [5] L. Xiang, Nucl. Mater. Energy 12, 2017, [6] L Aho-Mantila et al, Plasma Phys. Control. Fusion 59, 2017, [7] F. Reimold, Nucl. Mater. Energy 12, 2017, [8] I. Yu. Senichenkov, Plasma Phys. Control. Fusion, 2019, [9] P. C. Stangeby, The Plasma Boundary of Magnetic Fusion Devices, 2000

*Corresponding author: ferdinand.hitzler@ipp.mpg.de Third IAEA Technical Meeting on Divertor Concepts 2019 – Poster ID: 64



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