Edge plasma turbulence modelling in realistic WEST geometry with the SOLEDGE3X code





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Introduction

A word about the numerics

Two types of fluid models are widely used to model edge plasma:

implicit/explicit time stepping. One uses structured grid aligned on magnetic flux surfaces + 2D transport codes such as SOLPS, UEDGE or SOLEDGE2D assuming plasma axi-symmetry and where cross-field turbulent transport is emulated by empirical diffusion. These codes account domain decomposition to treat X-points. The following discretization is used: Donat-Marquina fluxes + WENO interpolation for advection [JCP 125, 42-58 (1996)] for neutrals (fluid or kinetic) and impurities and are key tools to interpret experiments. They Günter scheme for parallel diffusion in the flux surface [JCP 209, 354-370 (2005)]

lack predictability due to their poor description of turbulence.

A mask function χ is used to define the wall location. Boundary conditions are imposed at the 3D turbulent codes such as BOUT++, GBS or TOKAM3X which are able to self-consistently model turbulent transport but often rely on idealistic plasmas. The description of realistic interface $\chi = 1 | \chi = 0$. 0.5

tokamak geometry, neutrals or impurities is not always addressed.

We propose here to capitalize on the experience in developing SOLEDGE2D and TOKAM3X at CEA to create a new code SOLEDGE3X integrating the advantages of the two codes and thus being able to simulate complex multi-species plasma in realistic geometry, either in transport mode (2D) or turbulent mode (3D).

Description of the model

SOLEDGE3X solves:

Mass balance [for each ion species (α)]:

$$\partial_t n_{\alpha} + \vec{\nabla} \cdot (n_{\alpha} \vec{v}_{\alpha}) = S_{\alpha}^n \quad \text{where} \quad \vec{v}_{\alpha} = v_{\parallel,\alpha} \vec{b} + \vec{v}_{E \times B} + \vec{v}_{\alpha}^{\star} - \frac{D_{\alpha} \nabla_{\perp} n_{\alpha}}{n_{\alpha}} + \vec{v}_{\alpha}^p$$

Parallel momentum balance [for each ion species (α)]:

 $\partial_t (m_\alpha n_\alpha v_{\parallel,\alpha}) + m_\alpha \vec{\nabla} \cdot (n_\alpha v_{\parallel,\alpha} \vec{v}_\alpha - v_\alpha n_\alpha \vec{\nabla}_\perp v_{\parallel,\alpha}) = -\nabla_\parallel p_\alpha - \vec{b} \cdot \vec{\nabla} \cdot \overline{\tau_\alpha}$ $+Z_{\alpha}en_{\alpha}E_{\parallel}+R_{\parallel,\alpha}+S_{\alpha}^{\Gamma}$

Energy balance [for each ion species (α) and electrons]:

 $\partial_t E_{\alpha} + \vec{\nabla} \cdot \left(E_{\alpha} \vec{v}_{\alpha} + p_{\alpha} v_{\parallel} \vec{b} + q_{\parallel,\alpha} \vec{b} - \frac{m_{\alpha}}{2} v_{\alpha} n_{\alpha} \vec{\nabla}_{\perp} v_{\parallel,\alpha}^2 - \chi_{\alpha} n_{\alpha} \nabla_{\perp} T_{\alpha} \right) =$ $Z_{\alpha}en_{\alpha}v_{\parallel,\alpha}E_{\parallel}+v_{\parallel,\alpha}R_{\parallel,\alpha}+Q_{\alpha}+S_{\alpha}^{E}$

Charge balance:

where
$$j_{\parallel} = -\sigma_{\parallel} \left(\nabla_{\parallel} \phi - \frac{\nabla_{\parallel} p_e}{n_e} - \frac{R_{\parallel,e}^T}{n_e} \right)$$
 $\vec{j}^* = \sum_{i,e} Z_{\alpha} n_{\alpha} \vec{v}_{\alpha}^*$ $\vec{j}^p = \sum_i Z_{\alpha} n_{\alpha} \vec{v}_{\alpha}^p$
with $n_{\alpha} \vec{v}_{\alpha}^p = -\partial_t \vec{\omega}_{\alpha} - \vec{\nabla} \cdot (\tilde{v}_{\alpha} \otimes \vec{\omega}_{\alpha})$ where $\vec{\omega}_{\alpha} = \frac{m_{\alpha}}{Z_{\alpha} B^2} \left(n_{\alpha} \vec{\nabla}_{\perp} \phi + \frac{1}{Z_{\alpha}} \vec{\nabla}_{\perp} p_{\alpha} + \frac{1}{Z_{\alpha}} \vec{\nabla}_{\perp} p_{\alpha}$

The following linear solver are used: LAPACK for small linear problems (1D)

- PASTIX (direct solver) for medium size problems (2D diffusion or 3D diffusion on small grids)
- AGMG (algebraic multigrid iterative solver) for 2D and 3D diffusion [agmg.eu]
- PETSC for 2D and 3D diffusion



An hybrid MPI/OpenMP parallelization is used to solve in parallel different sub-domains. An OpenMP layer is used to solve species in parallel.

Application to 2D transport modelling

SOLEDGE3X is used here as a transport code to help interpret experiments, namely WEST Ohmic shot #54057 at $B_T = 3,83T$, $I_p = 500kA$, $P_{Ohm} \approx 500kW$, $n_e(lid) = 2 \cdot 10^{19} m^{-2}$.

SOLEDGE3X numerical scheme is based on a 2nd order finite volume scheme with a mixed



The shot is simulated with the following parameters:

D = v	$\chi_e = \chi_i$	composition	R _n	Neutrals
$1m^2s^{-1}$	$2m^2s^{-1}$	D + O (2%)	98 %	EIRENE

0.3



- Electron density (from quasi-neutrality): $n_e = \sum Z_{\alpha} n_{\alpha}$
- **Parallel electron velocity (from parallel current):**

$$n_e v_{\parallel,e} = \sum_i Z_\alpha n_\alpha v_{\parallel,\alpha} - j_{\parallel}$$

Neutrals: EIRENE or fluid neutrals (diffusive)

In green are shown diffusive cross-field transport which can be used to emulate turbulence when the code is used as a transport code. In red are shown collisional closure terms (see below). In blue are shown terms not yet implemented.

Boundary conditions on the wall (Bohm-Chodura):

- Velocity: $|\vec{v}_{\alpha} \cdot \vec{n}_{wall}| \ge c_{s,\alpha} |\vec{b} \cdot \vec{n}_{wall}|$
- **Energy flux:** $\Phi_{E,\alpha} = \left(\gamma_{\alpha}T_{\alpha} + \frac{1}{2}m_{\alpha}v_{\parallel,\alpha}^{2}\right)\Phi_{n,\alpha}$ where γ_{α} is sheath transmission factor
- **Current:** $j = \left[\sum_{ions} Z_{\alpha} \Phi_{n,\alpha}\right] \times \left[1 \exp\left(\Lambda \frac{\phi}{T_{\alpha}}\right)\right]$

Multi-species Zhdanov closure



$$q_{\parallel,e} = -\kappa_e^0 T_e^{\frac{1}{2}} \nabla_{\parallel} T_e + 0.71 n_e T_e (v_{\parallel,e} - v_{\parallel,i}) \text{ or } R_{\parallel,e} = -0.71 n_e \nabla_{\parallel} T_e - 0.51 \frac{m_e}{\tau_e} n_e (v_{\parallel,e} - v_{\parallel,i})$$



• Code ability to simulate multi-component plasma with EIRENE coupling well tested □ Next step: validation with experimental data from WEST campaigns

Application to 3D turbulence modelling

First attempt to perform turbulence modelling in realistic WEST wall geometry. Machine size is reduced by a factor 2. Parallel resistivity is multiplied by 10. Grid resolution: $[N_{\varphi} \times N_{\theta} \times N_{\psi}] \sim [32 \times 560 \times 130] + \text{core cell} (1/4 \text{ torus})$



Transient showing interchange instability – Drive: energy source in the core cell

• Explicit expressions exist for friction forces for trace impurities [see Stangeby, Plasma boundary] of magnetic fusion devices (2000)]

□ Zhdanov closure [Transport processes in multicomponent plasmas (2002)] provides linear relation between temperature gradients, velocities and heat fluxes and friction forces. **Pros** &: No trace impurity assumption, no light species assumption

Cons (*): No explicit expression for forces and heat fluxes. Requires linear system inversion

 $\begin{cases} \frac{5}{2}n_{\alpha}k\nabla T_{\alpha} = \sum_{\beta} \left[\frac{5}{2}\frac{\mu_{\alpha\beta}}{m_{\alpha}}\bar{G}_{\alpha\beta}^{(2)}(\bar{w}_{\alpha} - \bar{w}_{\beta}) + \bar{G}_{\alpha\beta}^{(5)}\frac{\bar{h}_{\alpha}}{p_{\alpha}} + \bar{G}_{\alpha\beta}^{(6)}\frac{\bar{h}_{\beta}}{p_{\beta}} + \frac{\mu_{\alpha\beta}}{kT} \left(\bar{G}_{\alpha\beta}^{(9)}\frac{\bar{r}_{\alpha}}{p_{\alpha}} + \bar{G}_{\alpha\beta}^{(10)}\frac{\bar{r}_{\beta}}{p_{\beta}} \right) \right] \\ 0 = \sum_{\beta} \left[\frac{35}{2} \left(\frac{\mu_{\alpha\beta}}{m_{\alpha}} \right)^{2} \bar{G}_{\alpha\beta}^{(8)}(\bar{w}_{\alpha} - \bar{w}_{\beta}) + 7\frac{\mu_{\alpha\beta}}{m_{\alpha}} \left(\bar{G}_{\alpha\beta}^{(9)}\frac{\bar{h}_{\alpha}}{p_{\alpha}} + \bar{G}_{\alpha\beta}^{(10)}\frac{\bar{h}_{\beta}}{p_{\beta}} \right) + \frac{m_{\alpha}}{kT} \bar{G}_{\alpha\beta}^{(11)}\frac{\bar{r}_{\alpha}}{p_{\alpha}} + \frac{m_{\beta}}{kT} \bar{G}_{\alpha\beta}^{(12)}\frac{\bar{r}_{\beta}}{p_{\beta}} \right] \end{cases}$

Where $\bar{G}_{\alpha\beta}^{(n)}$ can be computed explicitly from local plasma parameters



Turbulent results remain preliminary Code scaling to higher resolution / real size WEST simulation ongoing Code validation with experiments on WEST to be carried out: special focus on core rotation measurements to validate the physical model on current balance.

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