

## Motivation

- In JET, ICRH can cause anisotropies of  $p_{\perp}/p_{\parallel} = 2.5$  [W Zwingmann *et al.* PPCF, 43(11):1441, 2001.]
- MAST has reached  $p_{\perp}/p_{\parallel} \approx 1.7$  [MJ Hole *et al.* PPCF, 53(7):074021, 2011.]

## 1. Equilibrium with flow, anisotropy

- Inclusion of anisotropy and flow in equilibrium MHD equations e.g. [R. Iacono, *et al.* Phys. Fluids B 2 (8). 1990]

$$\nabla \cdot (\rho \mathbf{v}) = 0, \quad \rho \mathbf{v} \cdot \nabla \mathbf{v} = \mathbf{J} \times \mathbf{B} - \nabla \cdot \bar{\mathbf{P}}, \quad \nabla \cdot \mathbf{B} = 0$$

$$\mu_0 \mathbf{J} = \nabla \times \mathbf{B}, \quad \nabla \times (\mathbf{v} \times \mathbf{B}) = 0,$$

$$\bar{\mathbf{P}} = p_{\perp} \bar{\mathbf{I}} + \Delta \mathbf{B} \mathbf{B} / \mu_0, \quad \Delta = \frac{\mu_0 (p_{\parallel} - p_{\perp})}{B^2}$$

- If **toroidal flow** only:  $\mathbf{v} = -R \phi'_E(\psi) \mathbf{e}_{\phi} = R \Omega(\psi) \mathbf{e}_{\phi}$   
 $\Rightarrow F(\psi) = RB_{\phi}(1-\Delta)$
- If two temperature Bi-Maxwellian model chosen

$$p_{\parallel}(\rho, B, \psi) = \frac{k_B}{m} \rho T_{\parallel}(\psi), \quad p_{\perp}(\rho, B, \psi) = \frac{k_B}{m} \rho T_{\perp}(B, \psi)$$

- Can write

$$W(\rho, B, \psi) = T_{\parallel} \ln \frac{T_{\parallel} \rho}{T_{\perp} \rho_0}, \quad H = W - \frac{1}{2} [R \phi'_E(\psi)]^2, \quad T_{\perp}(B, \psi) = \frac{T_{\parallel} B}{|B - T_{\parallel} \Theta(\psi)|}$$

$$\nabla \cdot \left[ (1-\Delta) \left( \frac{\nabla \psi}{R^2} \right) \right] = - \frac{\partial p_{\parallel}}{\partial \psi} - \rho H'(\psi) + \rho \frac{\partial W}{\partial \psi} - \frac{F'(\psi) F'(\psi)}{R^2 (1-\Delta)} + R^2 \rho \Omega(\psi) \Omega'(\psi)$$

$$\{H(\psi), T_{\parallel}(\psi), F(\psi), \Theta(\psi), \Omega(\psi)\}$$

- Implemented in EFIT TENSOR for equilibrium reconstruction [Fitzgerald, Appel, Hole, Nucl. Fusion 53 (2013) 113040]
- Implemented in HELENA-ATF for stability studies [Qu, Fitzgerald, Hole, PPCF 56 (2014) 075007]

## Constraining HELENA+ATF

Boundary  $(R, Z)$   
Profiles

$$VF2 = \frac{F(\psi)^2}{F(0)^2}, \quad VTE = \frac{T_{\parallel}(\psi)}{T_{\parallel}(0)}, \quad VOM2 = \frac{\Omega(\psi)^2}{\Omega(0)^2}, \quad VH = \frac{H(\psi)}{H(0)}, \quad VTH = \frac{\Theta(\psi)}{\Theta(0)}$$

$$H(\psi) = \frac{k_B}{m_i} T_{\parallel}(\psi) \ln \frac{T_{\parallel}(\psi) \rho}{T_{\perp} \rho_0} - \frac{1}{2} \Omega^2 R^2, \quad T_{\parallel}(\psi) = \frac{m p_{\parallel}}{k_B \rho}$$

$$F(\psi) = RB_{\phi}(1-\Delta), \quad \Theta(\psi) = \frac{m_i B}{k_B T_{\parallel}} \left( 1 - \frac{T_{\parallel}(\psi)}{T_{\perp}} \right), \quad \Omega(\psi) = v_{\phi} / R$$

Constants

$$B = \mu_0 \rho \frac{k_B a^2 T_{\parallel}(0)}{m_i \varepsilon^2 F(0)^2} = \frac{\beta(0)}{2}, \quad HOT = \frac{m_i H(0)}{k_B T_{\parallel}(0)}, \quad \varepsilon = \frac{a}{R_0} R_0 B_0$$

$$OMGOT = \frac{m_i R_0^2 \Omega(0)^2}{k_B T_{\parallel}(0)}, \quad THTOF = R_0 \Theta(\psi) \frac{T_{\parallel}(0)}{F(0)}$$

## Anisotropy Scans Methodology

Addition of physics produces free parameters that are either not fully constrained (or self-consistent), or it is the inferred quantities (e.g. q profile) that are effectively constrained.  
 $\Rightarrow$  choose appropriate constraints to resolve impact of pressure anisotropy of q profile.

- 1 Constrain to Grad-Shafranov solution
2. Select target  $\Theta(\psi)$
- 3 Constrain thermal energy  $W_{th}$  such that  $p = (2p_{\perp} + p_{\parallel})/3$ . Modify  $\beta$  to match  $W_{th}$  by iterating HELENA+ATF until  $\Delta W_{th} / W_{th} < \varepsilon_{th}$
4. Modify current profile to match q. Either

(a) Single - pass: Assume  $B_{\phi} \gg B_{\theta}$  and low  $\beta = 2\mu_0 p / B^2$

In the isotropic case:  $J_{\phi,i} = R \left( \frac{\partial p_i}{\partial \psi} \right)_B + \frac{1}{2\mu_0 R} \left( \frac{\partial (RB_{\phi,i})}{\partial \psi} \right)_B$

In the anisotropic case:  $J_{\phi,a} = R \left( \frac{\partial p_{\perp,a}}{\partial \psi} \right)_B + \frac{1}{2\mu_0 R} \left( \frac{\partial (RB_{\phi,a})}{\partial \psi} \right)_B$

Force  $J_{\phi,i} = J_{\phi,a}$ .  
 Integrate over  $\psi$ :  $F_a^2 = F_i^2 \frac{1-\Delta^2}{1-\frac{2}{3}\Delta} \approx F_i^2 \left( 1 - \frac{4}{3}\Delta \right)$

Rerun HELENA+ATF

(b) Iteration: In general the toroidal current can be written

$$J_{\phi} = - \frac{F(\psi) F'(\psi)}{(1-\Delta) R \mu_0} - \left[ T_{\parallel}'(\psi) + H'(\psi) - \left( \frac{\partial W}{\partial \psi} \right)_{\rho, B} \right]$$

Next, we assume:  $J_{\phi,a} + \frac{F_a(\psi) F_a'(\psi)}{(1-\Delta_a) R \mu_0} \approx J_{\phi,i} + \frac{F_i(\psi) F_i'(\psi)}{R \mu_0}$

Compute

$$\int_1^{\psi_n} (J_{\phi,a} (1-\Delta_a) R \mu_0 - J_{\phi,i} R \mu_0) (\psi_a - \psi_0) d\psi \approx [-F_a^2 + F_i^2] \quad (1)$$

$$= -\delta F^2(\psi_n) + \delta F^2$$

... and then update

$$F^2(\psi_n) \rightarrow F^2(\psi_n) - \delta F^2(\psi_n). \quad (2)$$

Rerun in HELENA+ATF and compute

$$\Delta q = \int_0^1 (q_{target} - q) d\psi_n \quad (3)$$

Iterate Eqs. (1) - (3), and the replacement for F until  $\Delta q < \varepsilon_q$

## 2. Ballooning modes in anisotropic plasmas

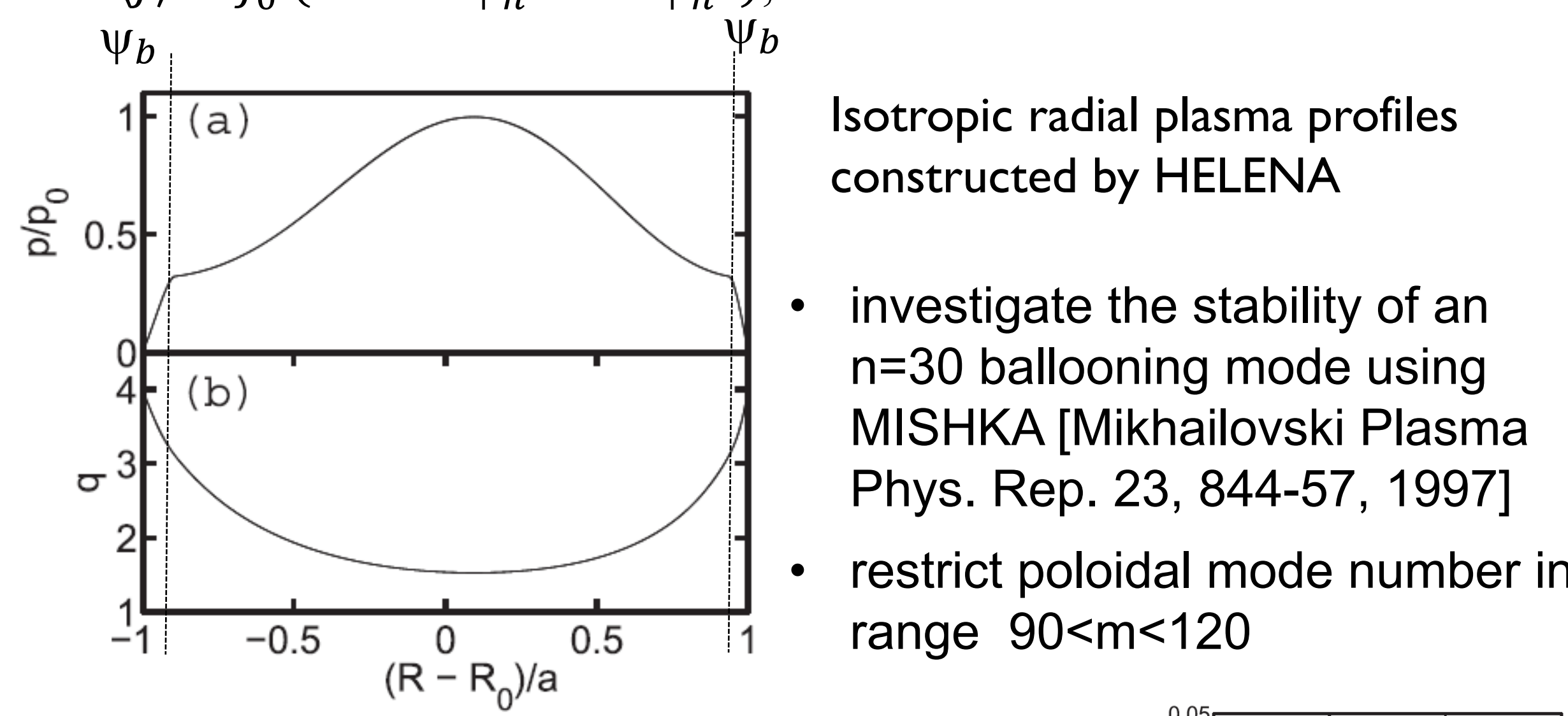
**Aim:** explore the impact of anisotropy on n=30 mode for the same q and  $J_{\phi}$  profile

- As in Huysmans *et al.* Phys. Plasmas 8 4292-305, 2001, assume a circular cross-section tokamak,  $R_0/a=4$ ,  $\beta_{pol}=1$ , and  $q_a=4$ .

$$p'(\psi) = p_0(1 - \psi_n), \quad \psi_n < \psi_b$$

$$p'(\psi) = p_0 \left( 1 - \psi_n + p_1 \left[ \frac{(\psi_n - \psi_b)^2 (3 - 2\psi_n - \psi_b)}{(1 - \psi_b)} \right]^{1/4} \right), \quad \psi_n > \psi_b$$

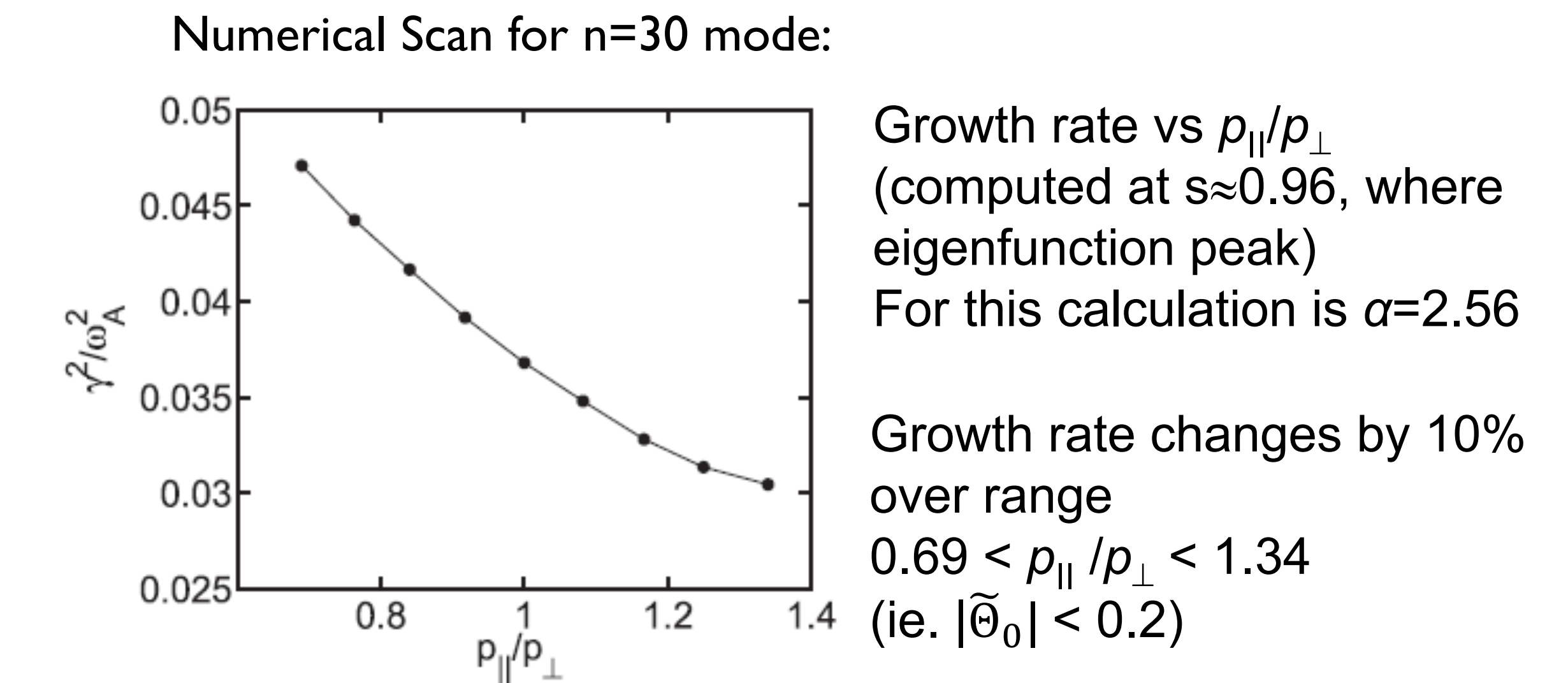
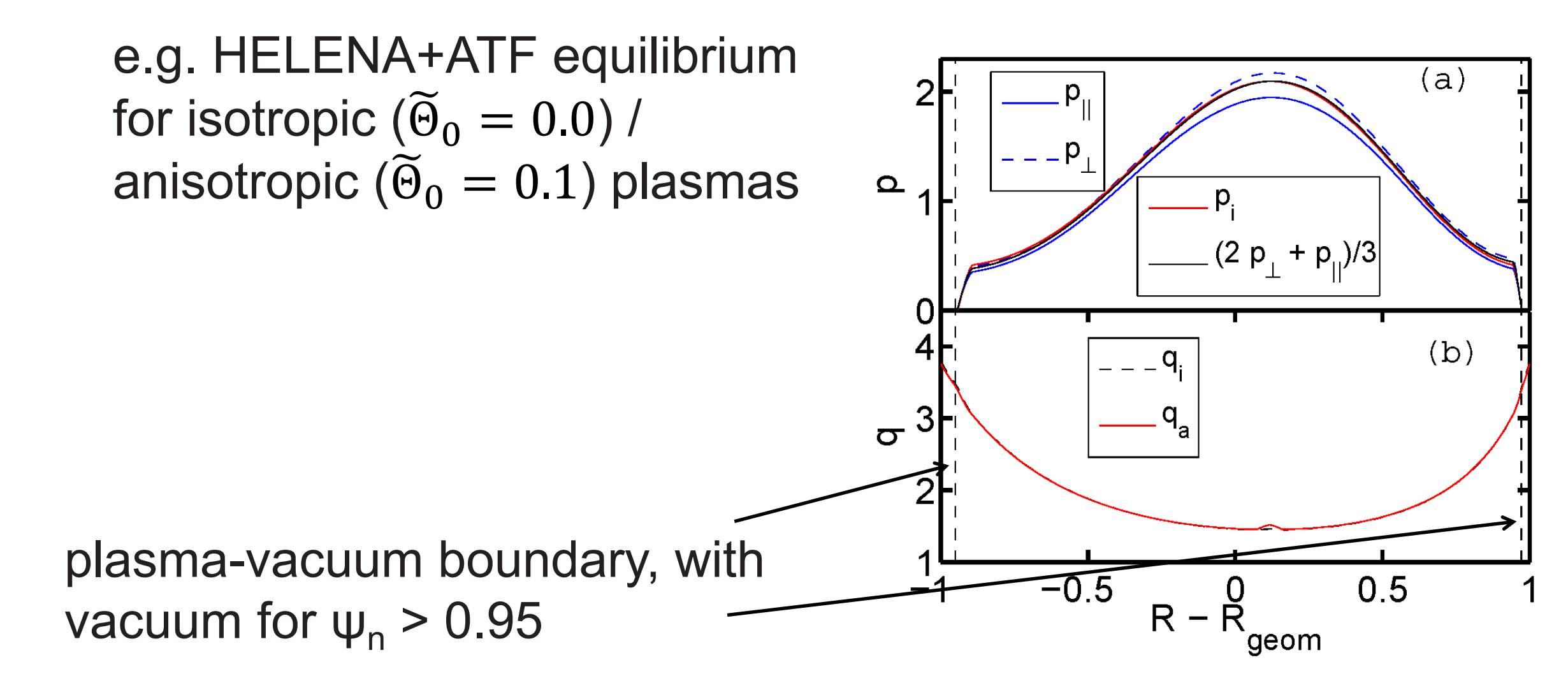
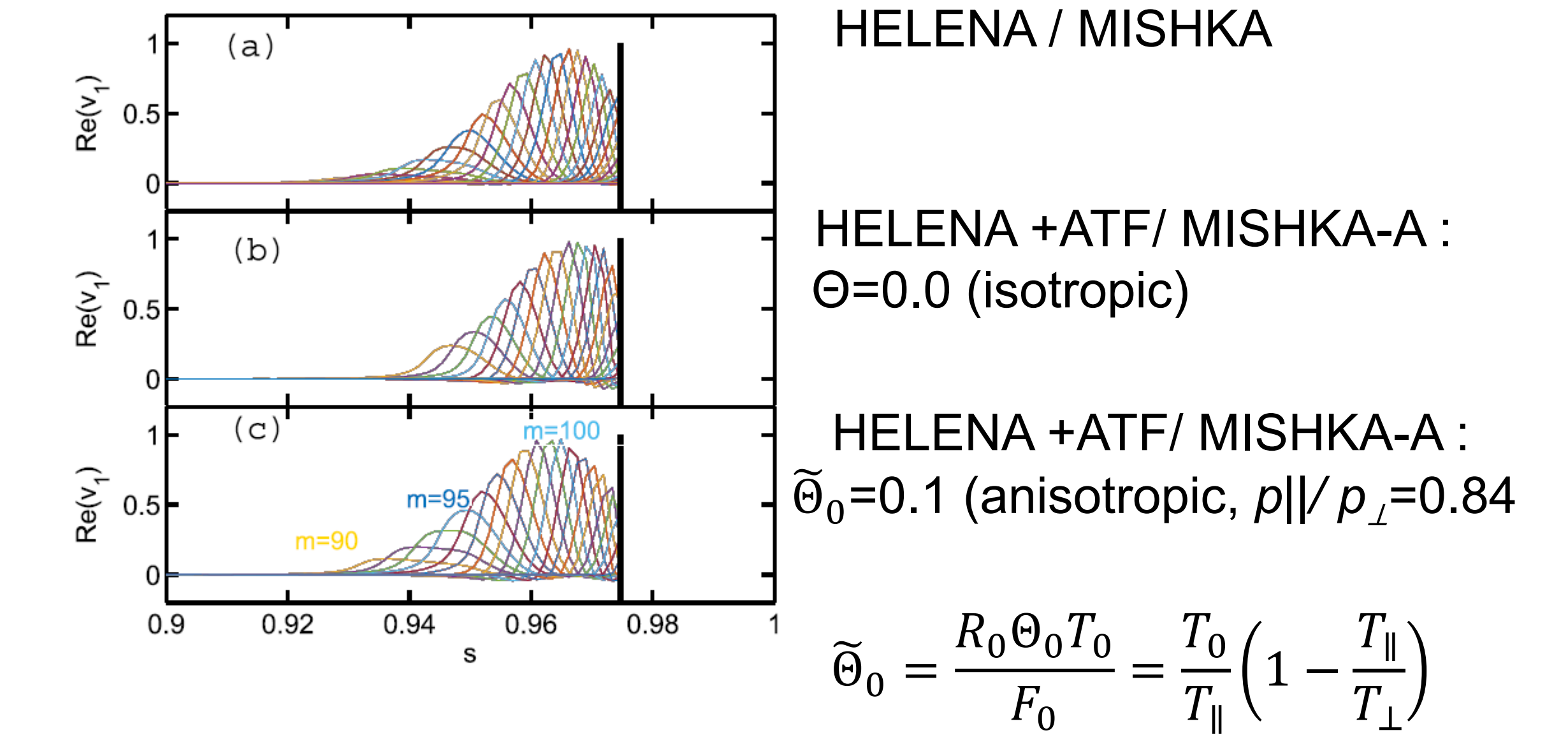
$$\langle J \rangle = J_0 (1 - 0.8\psi_n - 0.2\psi_n^2),$$



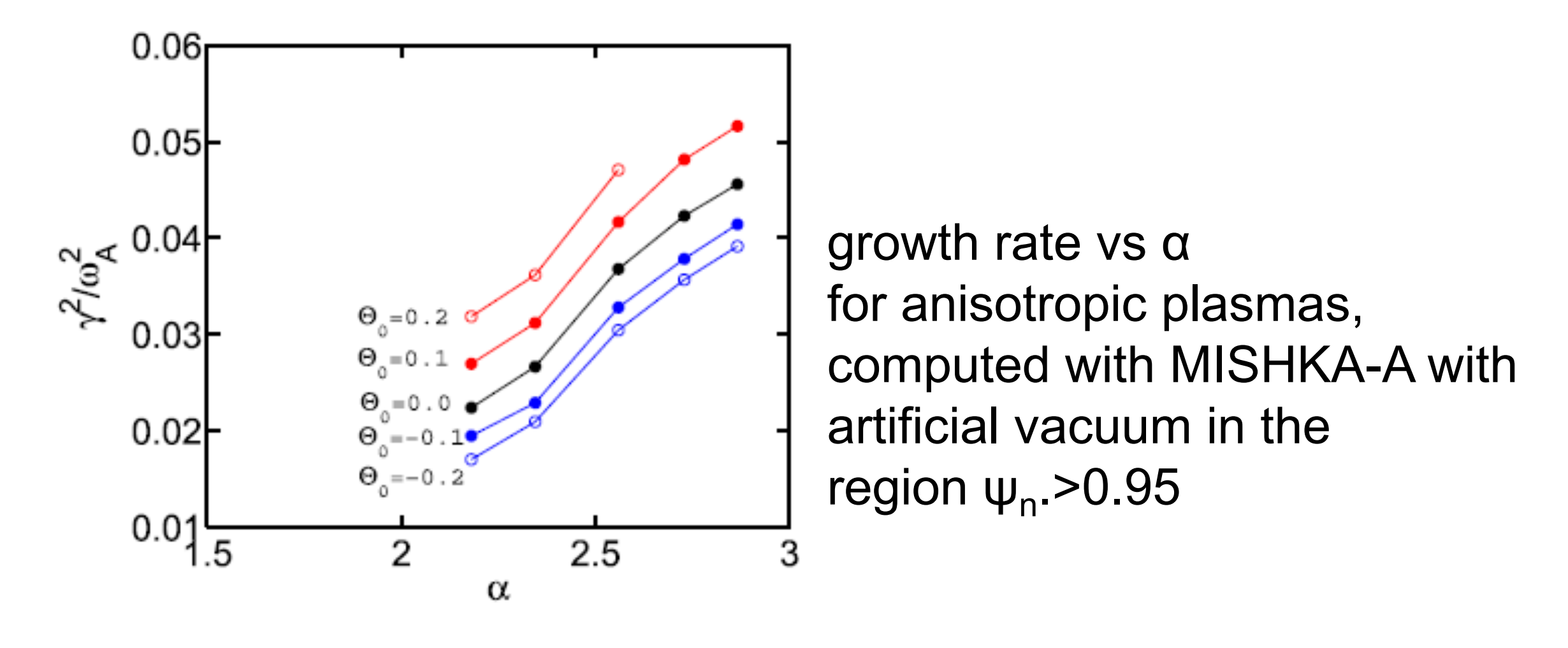
- investigate the stability of an n=30 ballooning mode using MISHKA [Mikhailovski Plasma Phys. Rep. 23, 844-57, 1997]
- restrict poloidal mode number in range  $90 < m < 120$
- Eigenmode benchmark
  - ✓ Mode structure matches
  - ✓  $\gamma^2/\omega^2$  versus  $\alpha$

**Including Anisotropy:** Employ Steps (1)-4(a)

- MISHKA-A [Qu *et al.* Plasma Phys. Control. Fusion 57 095005] extends MISHKA to plasmas with an anisotropy and flow
- MISHKA-A assumes a conformal wall at  $R_w = 1$ .
- Adapt HELENA equilibrium to include an "artificial vacuum": set the  $p'(\psi_n)=0$  for  $\psi_n > 0.95$



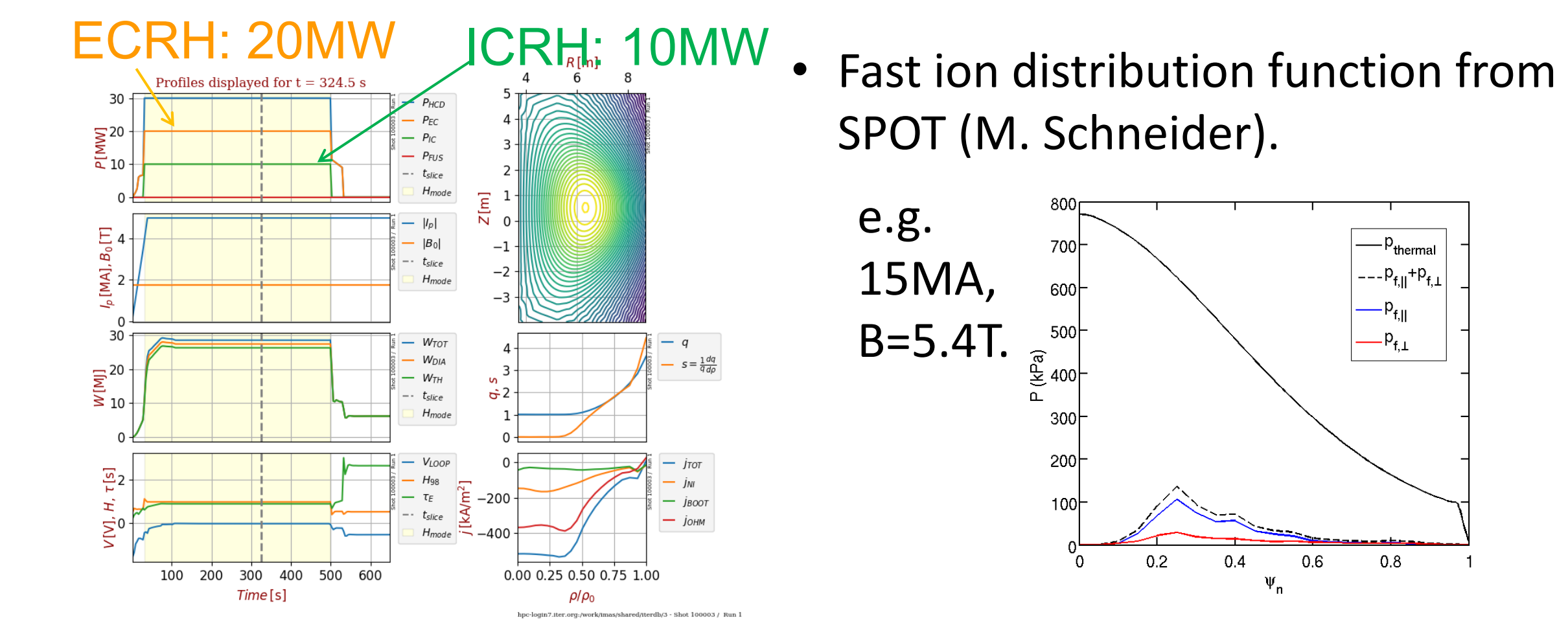
**Reason:** As  $T_{\perp}$  increases over  $T_{\parallel}$ ,  $p_{\perp}$  surfaces are displaced outboard to bad curvature region of an inward shift of surfaces stabilises the mode.



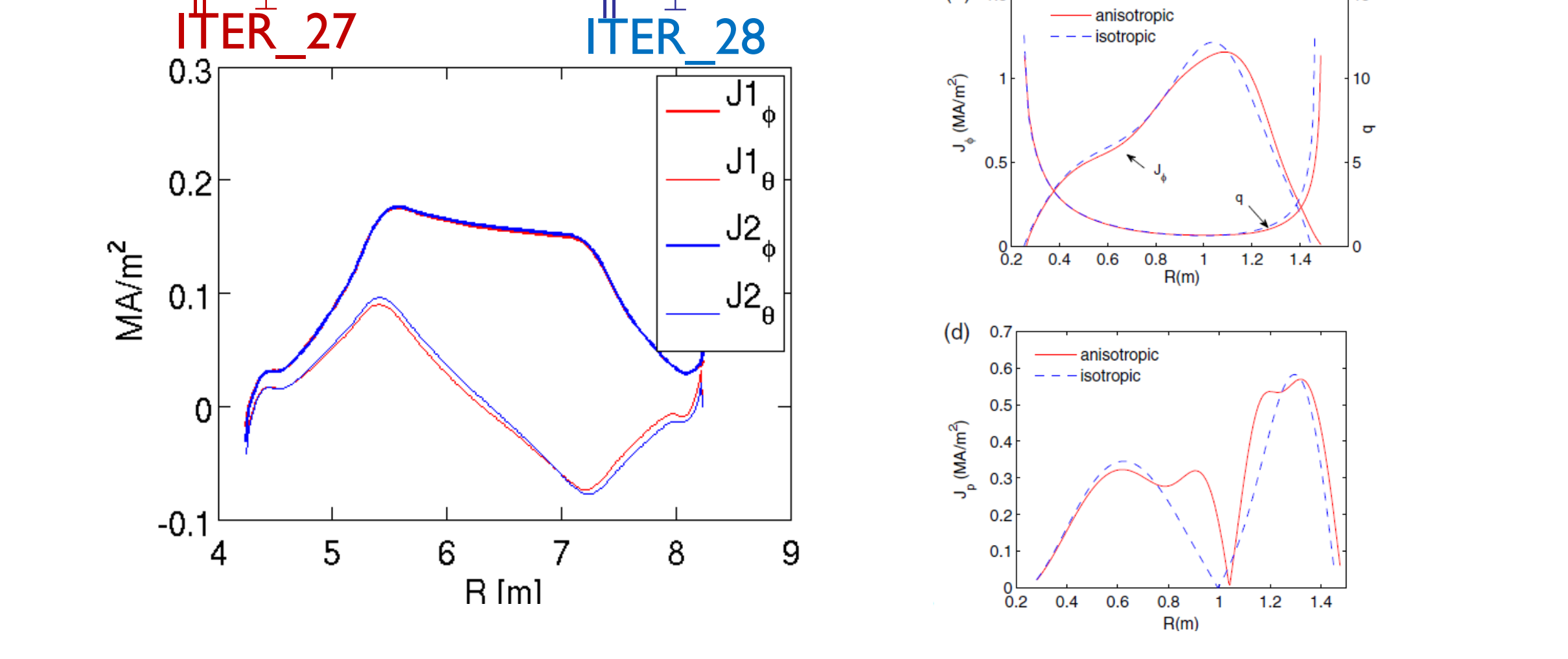
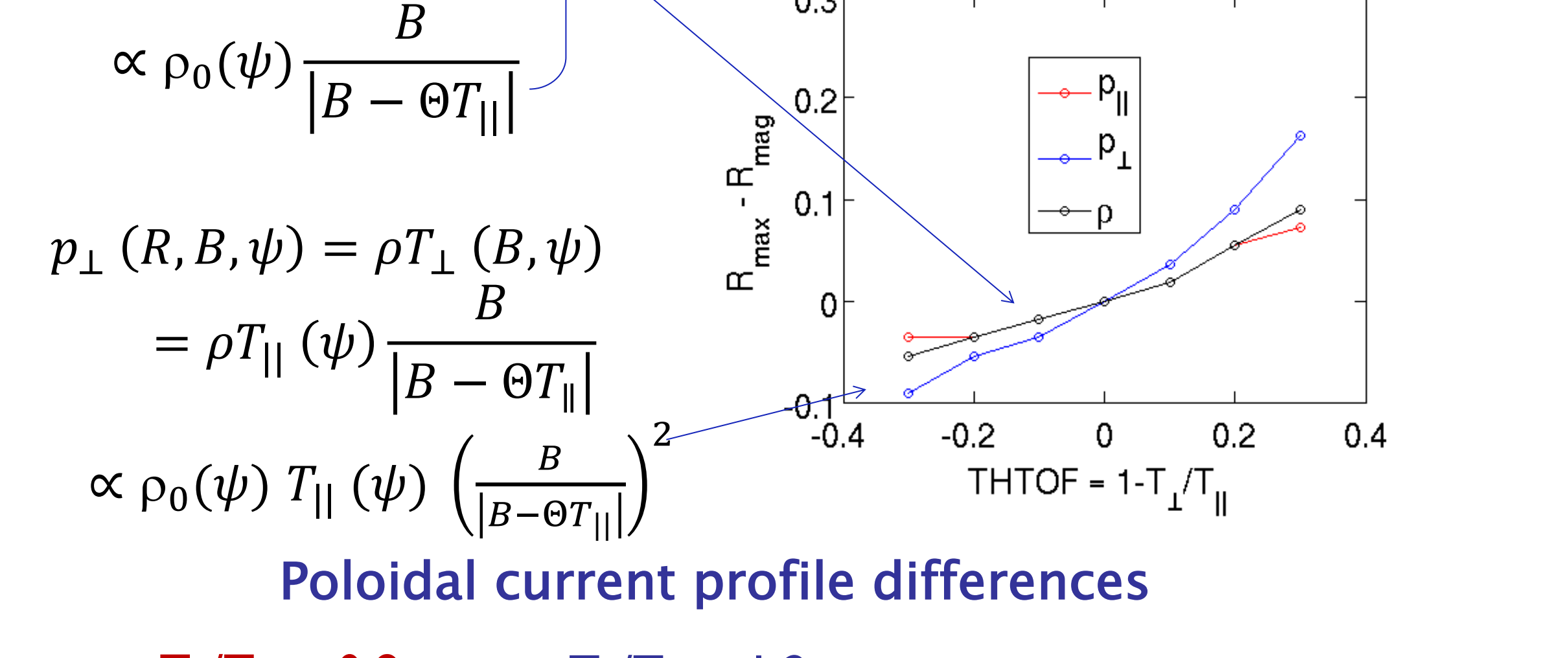
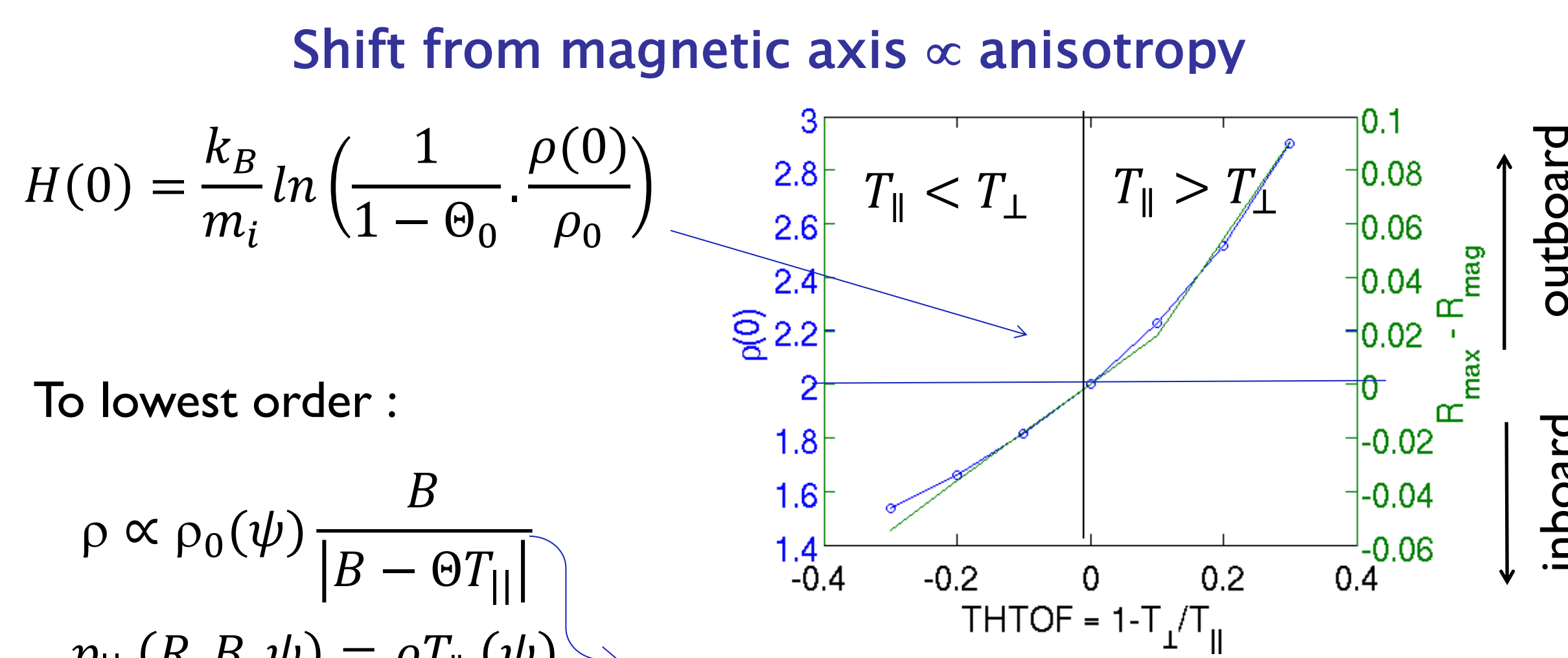
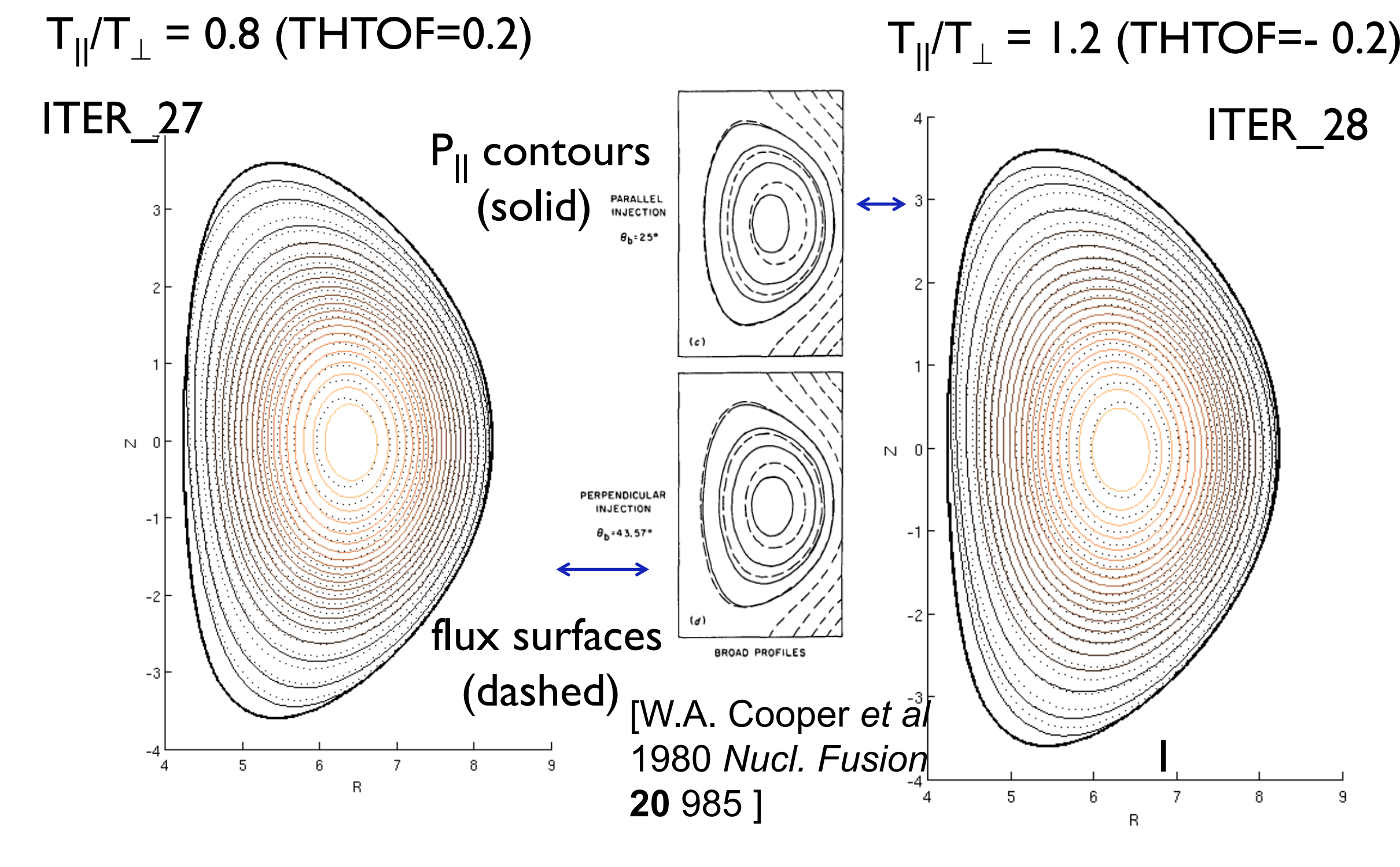
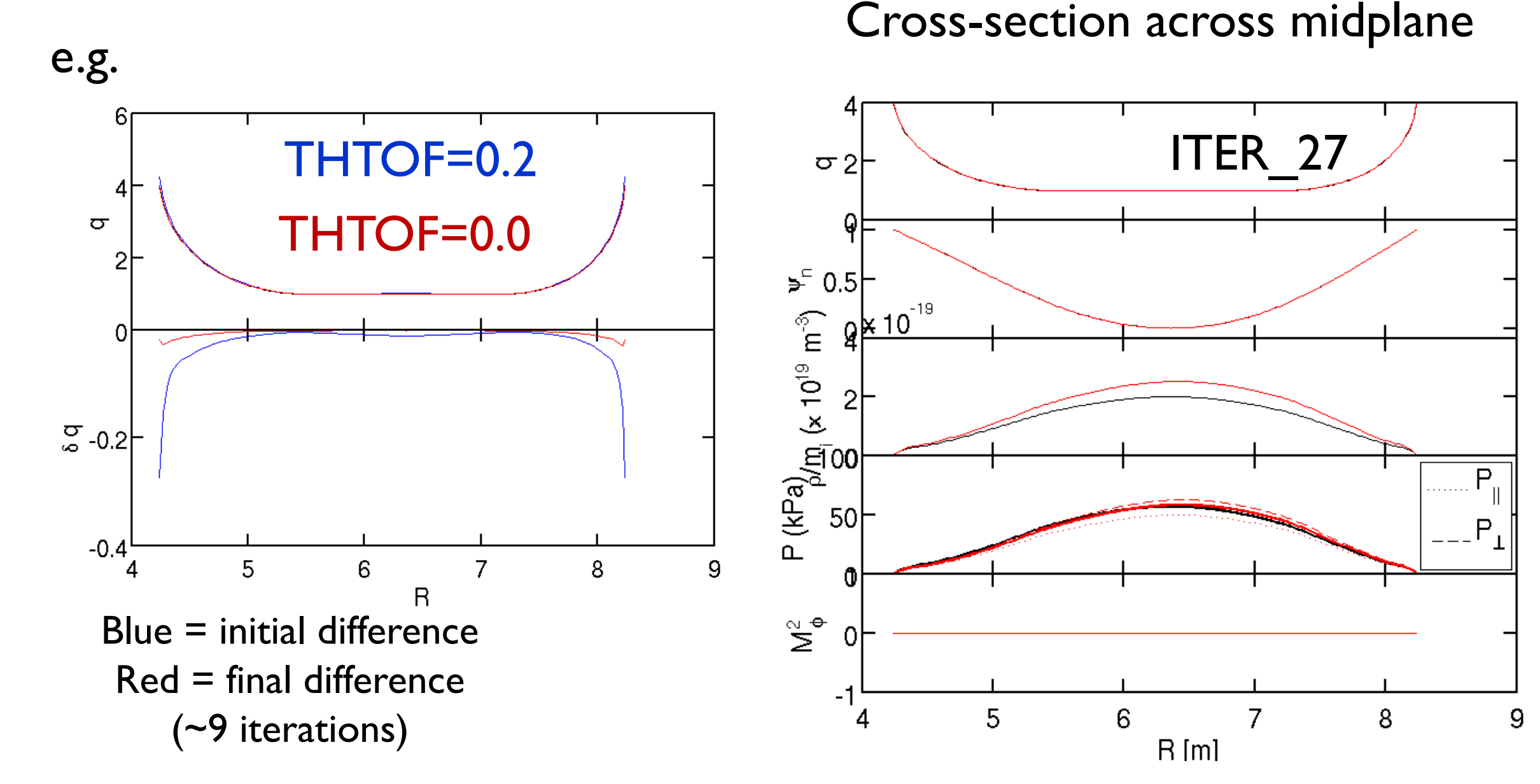
- Over parameter range explored growth rate has same dependence with  $\alpha$
- First step to (s,  $\alpha$ ) marginal stability boundaries with anisotropy.
- $\gamma^2/\omega^2$  increases with increasing  $p_{\perp}/p_{\parallel}$  (increasing  $\bar{\theta}_0 = T_0/T_{\parallel} (1 - T_{\perp}/T_{\parallel})$ ).
- Suggests increasing  $p_{\parallel}/p_{\perp}$  in the pedestal region might lead to higher ELM-free performance

## 3. ITER Pre-fusion power operation

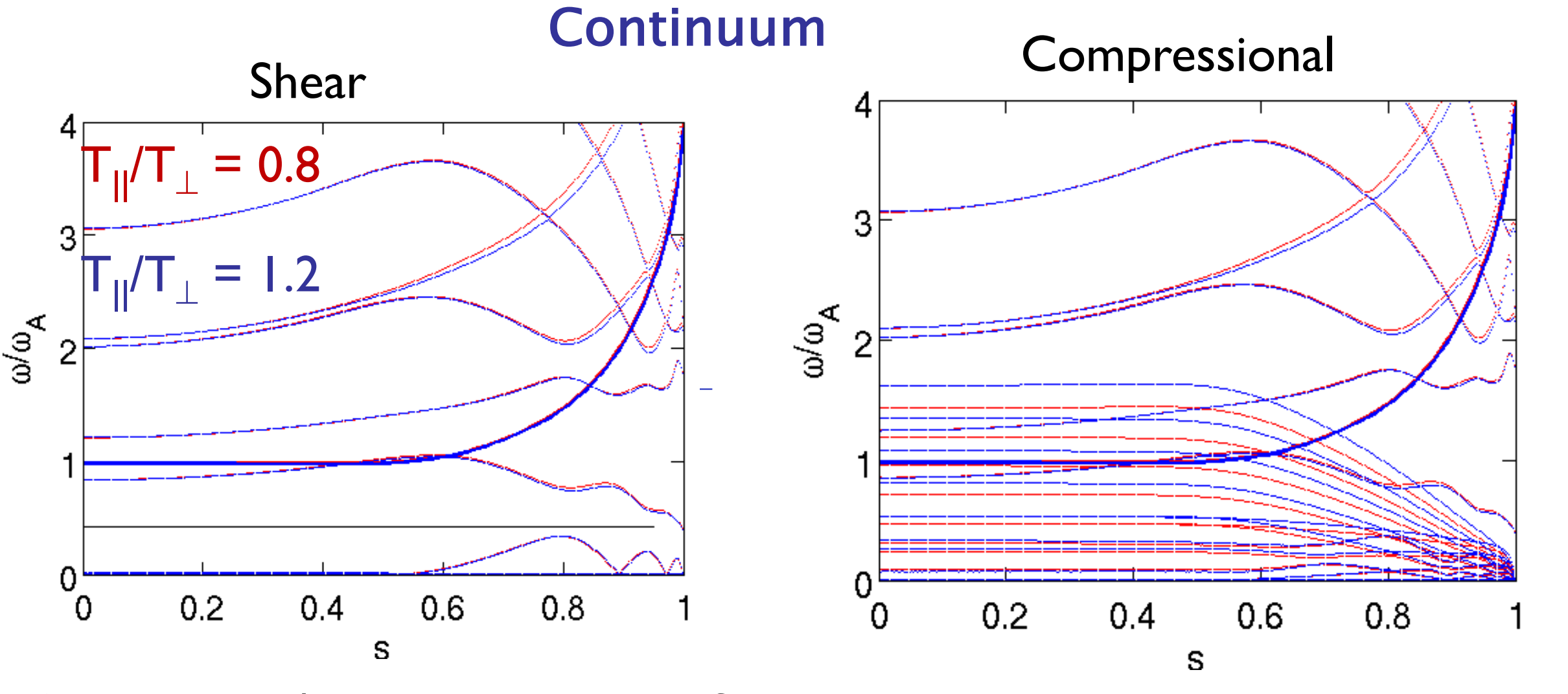
- Pre-fusion power operation H plasma, 5MA, B=1.8T (1/3 field) ICRH scenario: #100003, run 1 at 324s



- First Step: Vary central anisotropy  $THTOF = \frac{T_0}{T_{\parallel}} \left( 1 - \frac{T_{\perp}}{T_{\parallel}} \right)$
- Adjusted  $-0.2 < THTOF < 0.2$  to examine impact on equilibrium, continuum, and TAE mode structure.



Moderate for ITER, but larger for an ST (EFIT TENSOR with / without anisotropy) [Qu, Fitzgerald, Hole, PPCF 56 (2014) 075007]



- Density / pressure axis shift  $\propto$  anisotropy
- Little difference to shear continuum and gap eigenmodes
- Significant difference to compressional continuum,

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