

Local Current Injector Systems for Nonsolenoidal Startup in a Low Aspect Ratio Tokamak

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PEGASUS
Toroidal Experiment



Motivation: Intense Electron Current Sources Needed for Local Helicity Startup

- Significant progress with non-solenoidal startup of ST
 - Exploiting local helicity injection via current sources in plasma edge region
 - Technical attractiveness: can remove sources and anode after startup
 - Understanding of helicity balance and relaxation current limits guide hardware and operational changes
 - Helicity injection discharges couple to other current drive methods
- Tests and development on the Pegasus Toroidal Experiment
 - $A \sim 1$; $I_p = 0.1\text{-}0.3\text{ MA}$; $B_{\text{tf}} = 0.15\text{ T}$
 - $I_p \sim 0.17\text{ MA}$ using helicity injection and outer-PF rampup; $\sim 0.08\text{ MA}$ with HI only
 - Goal $\approx 0.3\text{-}0.4\text{ MA}$ non-solenoidal I_p to extrapolate to next level/NSTX
 - Issues in physics understanding: j_{edge} , Z_{inj} , confinement, etc.
- Exploitation of point-source helicity startup requires large-area sources of intense electron current
 - Developing understanding and designs of robust electron sources based on plasma arc sources
 - Exploring possibility of simpler large-area sources via gas-fed electrodes
 - Requires 2 kV, 15 kA programmable power systems



LOCAL HELICITY INJECTION OFFERS SCALABLE NONSOLENOIDAL STARTUP

- Inject Helicity for I_p startup using electron current source at the tokamak plasma edge

- I_p limited by available helicity drive, including PF induction. Helicity balance gives:

$$I_p \leq \frac{A_p}{2\pi R_0 \langle \eta \rangle} (V_{ind} + V_{eff}) \quad V_{eff} \approx \frac{A_{inj} B_{\phi, inj}}{\Psi_T} V_{bias}$$
$$I_p \leq \left[\frac{C_p}{2\pi R_{inj} \mu_0} \frac{\Psi_T I_{inj}}{w} \right]^{1/2}$$

- Max I_p set by relaxation to Taylor (constant λ) state:
- Helicity dissipation thru resistive losses in plasma

- Maximizing I_p requires

- Large helicity input rate: High A_{inj} , V_{inj}
- High relaxation limit: High I_{inj} , low w

A_p Plasma area

C_p Plasma circumference

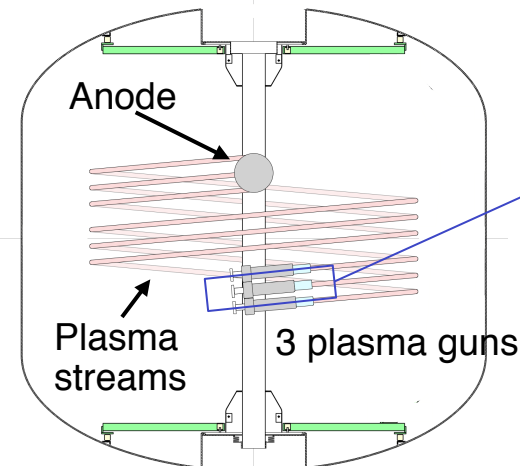
Ψ_T Plasma toroidal flux

w Edge current channel width

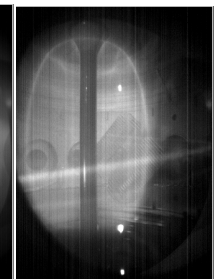
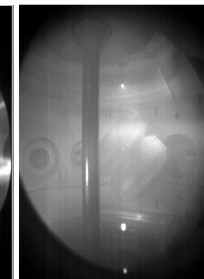
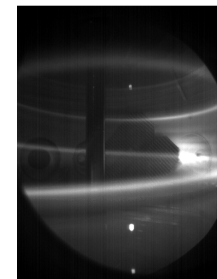


OUTER LFS INJECTION ADDS POLOIDAL INDUCTION TO HELICITY INJECTION

- Flexible geometry for injector locations
 - Outer midplane allows “port-plug” installation
- PF null via injection into helical (TF + PF) field; followed by relaxation to tokamak-like state
 - Rapid inward expansion and growth in I_p at low A
- Poloidal field induction adds to current growth



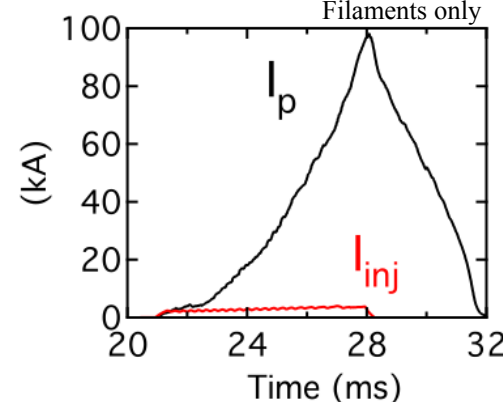
PEGASUS shot #40458: two midplane guns, outer-PF ram



$I_p = 2-3$ kA
Filaments only

$I_p = 42$ kA
Driven plasma

$I_p = 37$ kA
Guns off
Decaying

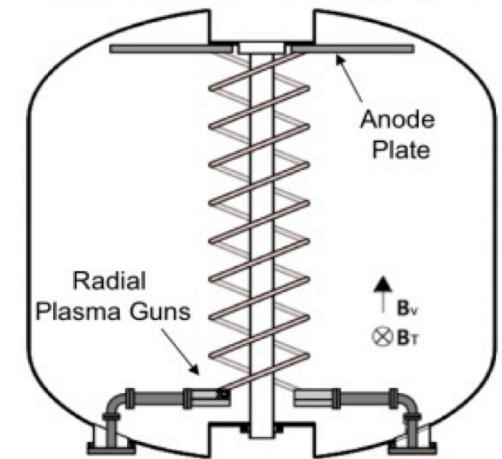




Inboard HFS Injection in Divertor Region Maximizes Helicity Input Rate

- HFS injection near centerstack maximizes helicity input rate
- Reduced plasma position control requirements
 - Static fields support easy control of position

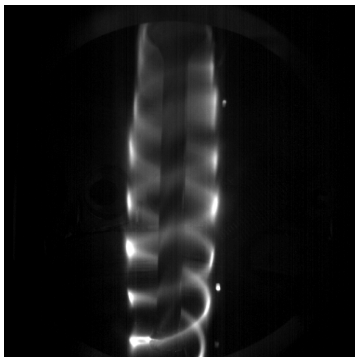
Inboard Divertor Gun Injection



$$R_{\text{gun}} = 16 \text{ cm}, Z_{\text{gun}} = -75 \text{ cm}$$



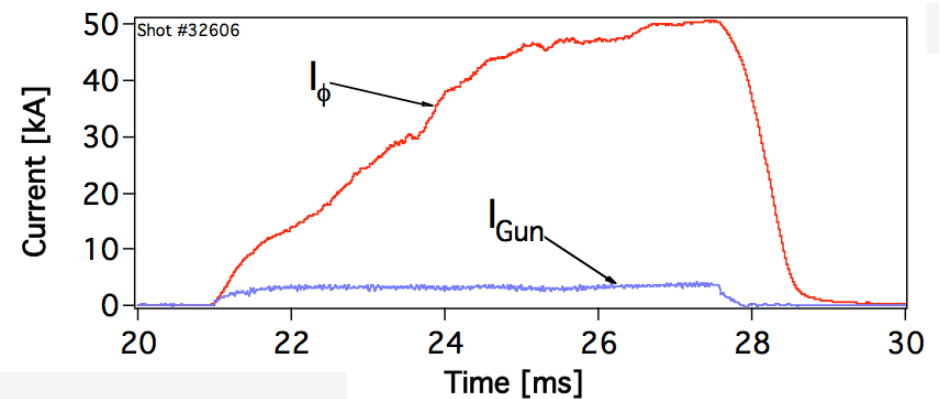
Increased I_{inj}
Reduced B_z



Current filaments



Relaxed tokamak





Plasma Arc Sources



Compact Plasma Arc Sources Provide Dense Plasma for Electron Current Extraction

- Plasma arc(s) biased relative to anode:

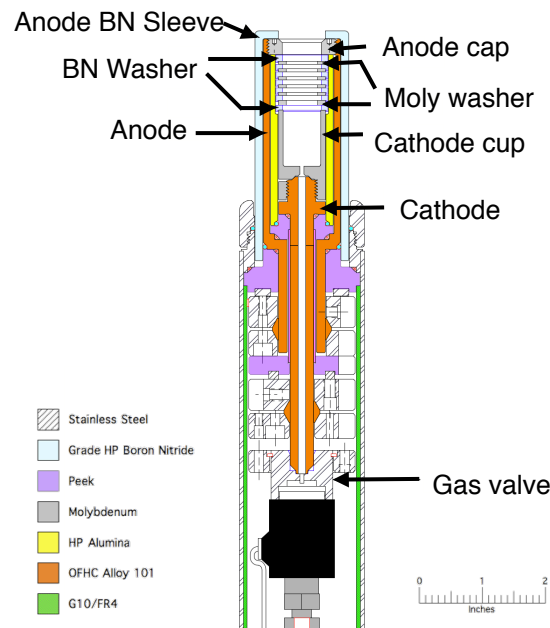
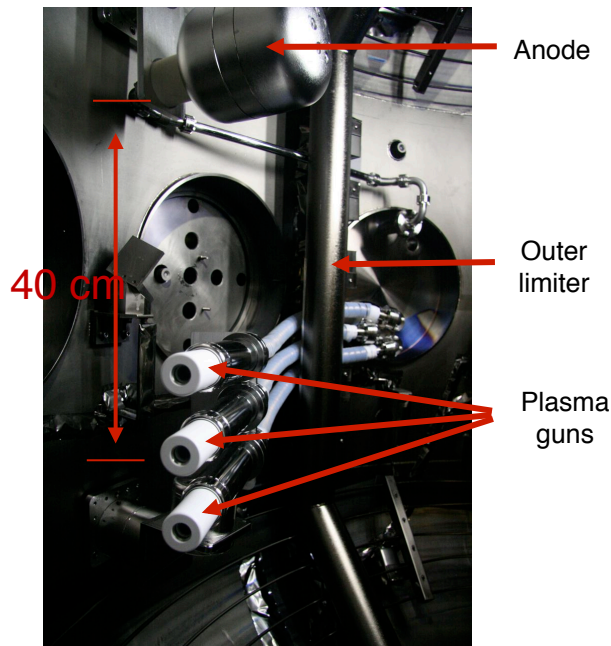
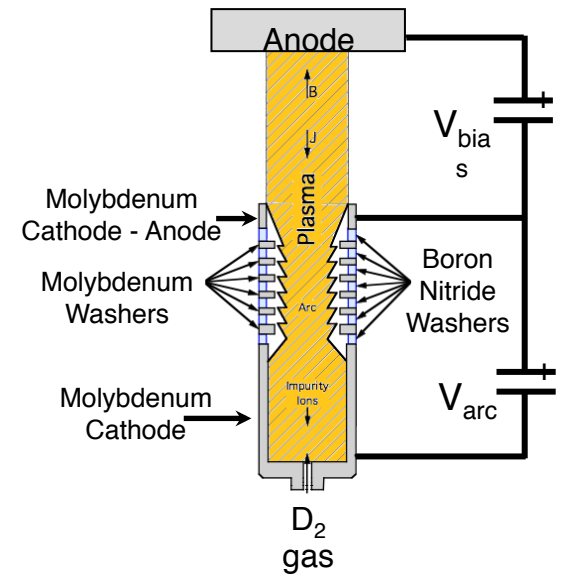
- Helicity injection rate:

$$\dot{K}_{inj} = 2V_{inj}B_NA_{inj}$$

V_{inj} - injector voltage

B_N - normal B field at gun aperture

A_{inj} - injector area



- Arc plasma fully ionized

- $N_e \sim 10^{20} \text{ m}^{-3}$

- $T_e \sim 10 \text{ eV}$

- $\text{Dia} = 1.6 \text{ cm}$

- $I_{arc} \sim 2 \text{ kA}$

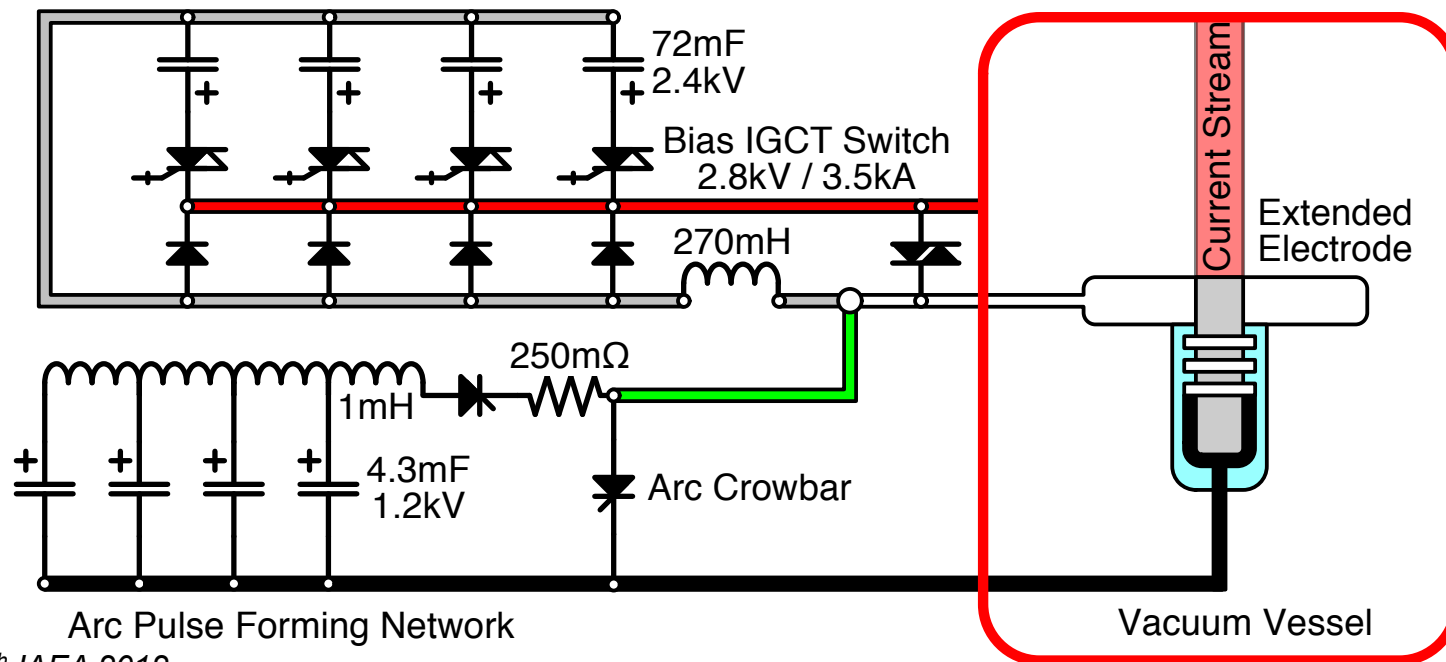
RJF 24th IAEA 2012

1 Fiksel, G, et. al., Plasma Sources Sci. & Tech. **5** (1996) 78.



Robust Switching Power Supplies Deployed for Arc & Injection

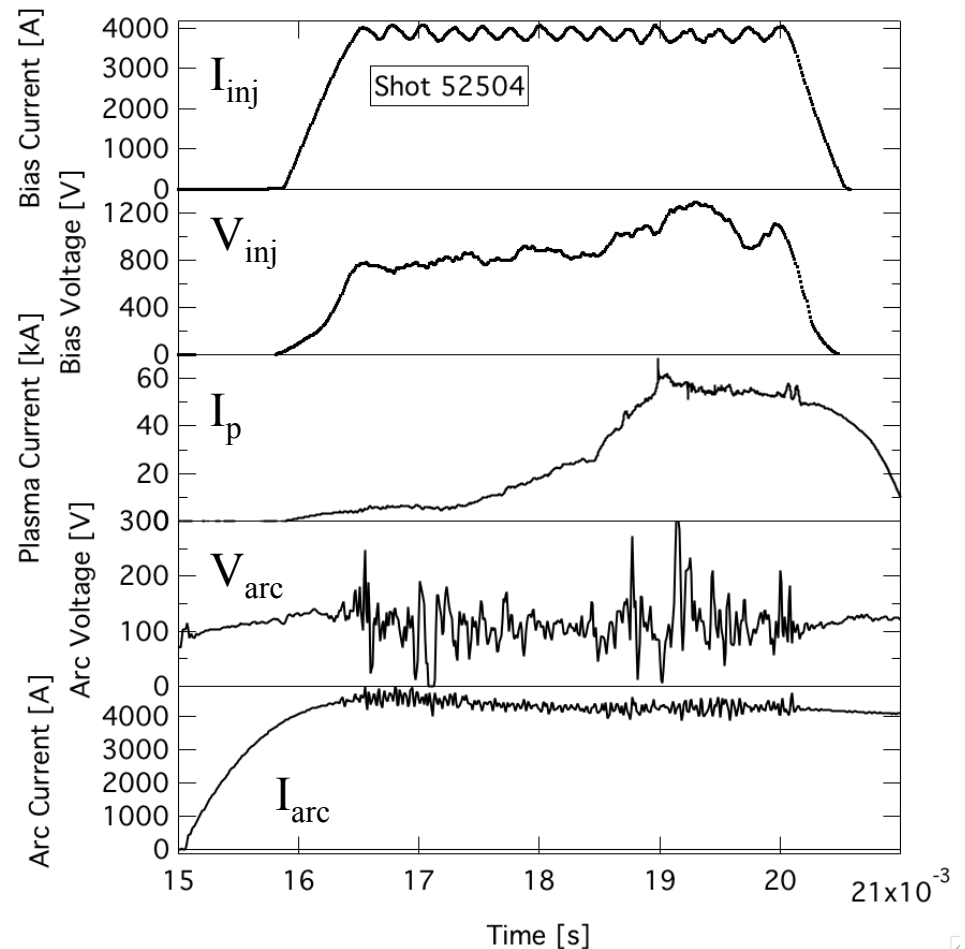
- Plasma Arc uses simple Pulse Forming Network
 - Once arc is established: $I_{\text{arc}} = 1\text{-}2 \text{ kA}$ @ $V_{\text{arc}} = 100\text{-}200 \text{ V}$
 - SCR terminates arc on demand
- Injection (Bias) circuit uses 4 IGCT switches in parallel
 - Total: $I_{\text{inj}} \leq 14 \text{ kA}$ @ $V_{\text{inj}} \leq 2.2 \text{ kV}$
 - Preprogrammed current control via stabilized PWM feedback controller
 - Series inductance stabilized, sometimes with parallel stabilizing capacitor and ballast resistor





Power Systems Provides Routine Programmable Injected Current and Helicity

- Injection circuit provides current feedback control
 - Impedance varies with resulting tokamak plasma so that V_{inj} varies through shot
 - Future upgrade: go to voltage feedback control
 - Active control of helicity injection rate
- Arc circuit fully ionizes injected gas
 - $I_{arc} \sim 2\text{-}4 \text{ kA}$ @ $V_{arc} \sim 150 \text{ V}$
 - With 1.6 cm diameter arc chamber, routine operation at 2 kA, with reduced lifetime at 4 kA
- Shot sequence
 - Inject gas flow into arc chamber
 - Strike Arc current; allow $\sim 1\text{ms}$ to establish arc
 - Extract I_{inj} ; usually with $I_{inj} < I_{arc}$





Arc Source Impedance



Predictive Impedance Models Required to Project to Future Startup Systems

- Current injector impedance is a critical parameter in local helicity injection startup
 - I_{inj} sets Taylor relaxation maximum I_p
 - V_{inj} sets effective V_{loop} for current drive
 - Impedance couples the two to define power requirements
- Double-sheath space-charge limits I_{inj} at low I_{inj} and V_{inj}

$$J_e = \frac{4}{9} \epsilon_o \sqrt{\frac{2e}{m_e}} \frac{V^{3/2}}{(\chi \lambda_{De})^2}$$

- At high I_{inj} ($> I_A$) and $V_{inj} > 10 \text{ kT}_e/e$, the Alfven-Lawson magnetic current limit dominates

$$I_{AL}^e = 1.65 \frac{4\pi m_e v_e}{e \mu_o} \equiv 1.65 I_A = 56 \sqrt{V_{inj}}$$

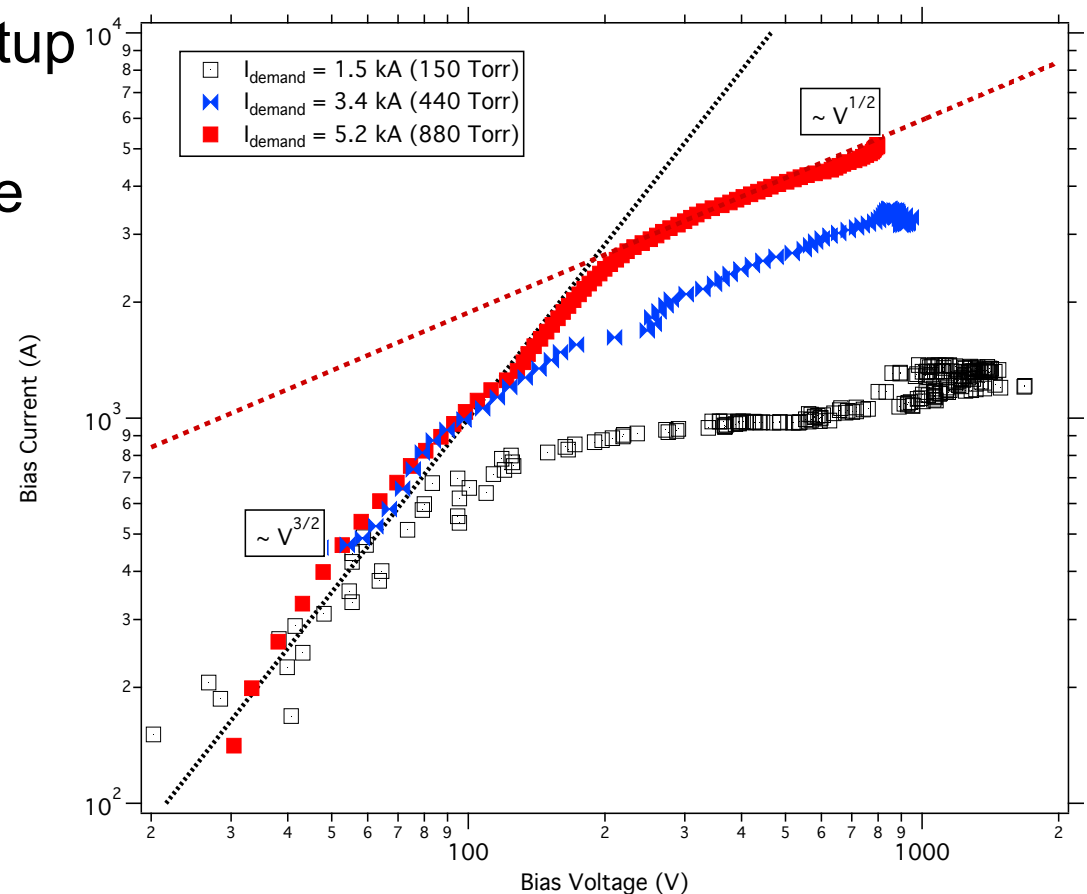
- For a uniform current density
 - Possible that sheath expansion also contributes in this region
- So far, these models and supporting evidence imply impedance determined by processes local to the injector and not the background plasma





Helicity Injection Process Governed by Space Charge and Magnetic Current Limits

- Arc source I-V characteristics obtained during plasma startup
- Double-sheath space-charge limits I_{inj} at low I_{inj} and V_{inj} : Initiation phase
 - $I_{inj} \sim n_e V^{3/2}$
- At high $I_{inj} > I_A$ and $V_{inj} > 10$ kT_e/e, the Alfven-Lawson magnetic current limit dominates
 - $I_{inj} \sim V^{1/2}$
 - Possible that sheath expansion also contributes here



I-V characteristics of arc plasma current injector for varied fueling rates.



Density Scaling in Injector Impedance May Reflect e⁻ Beam Profiles?

- I-V characteristics at varied fueling rates suggests a scaling with arc density

- Density variation may reflect changes in beam current density profile

- Alfven: uniform j with backward particle flow

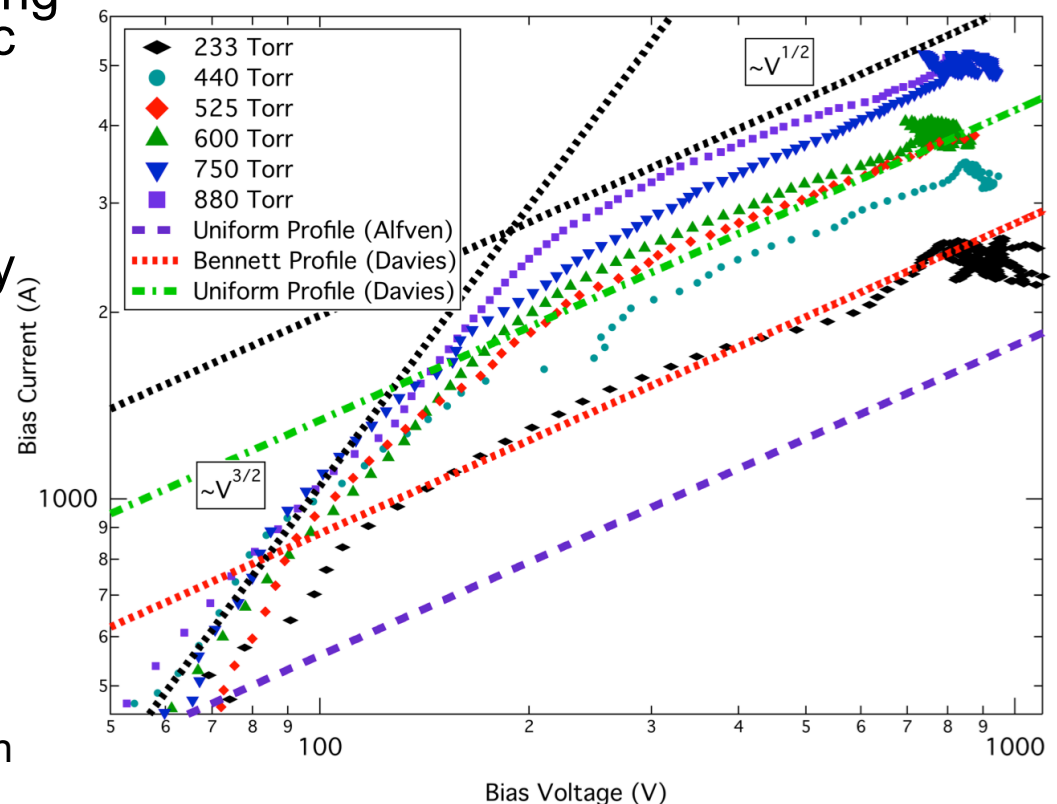
$$I_{AL}^e = 1.65 \frac{4\pi m_e v_e}{e\mu_o} \equiv 1.65 I_A = 56 \sqrt{V_{inj}}$$

- Davies: Uniform profile and Bennett profile for $j(r)$
 - Derived from energy conservation

$$I_{uniform}^e = 4.0 \frac{4\pi m_e v_e}{e\mu_o} = 134 \sqrt{V_{inj}}$$

$$I_{Bennett}^e = 2.59 \frac{4\pi m_e v_e}{e\mu_o} = 88 \sqrt{V_{inj}}$$

- Data shows inferred trends but detailed measurements needed



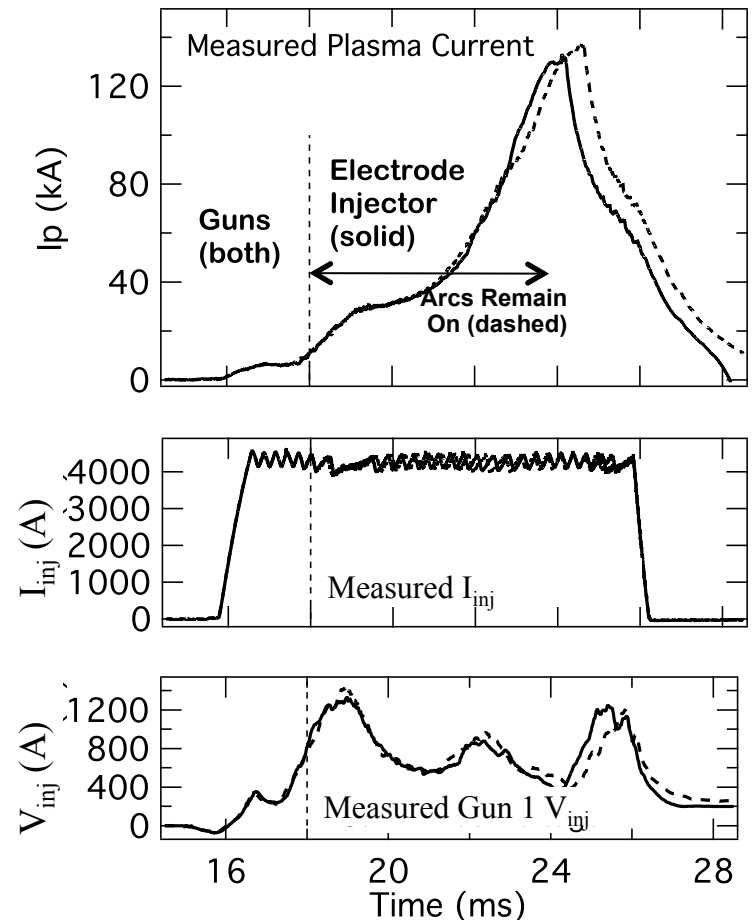


Current Injection via Gas-Fed Electrodes



Exploring Passive Injectors to Increase Helicity Injection Rates

- Maximizing Helicity (i.e., current drive) requires large area electron emitters
- Two possible paths
 - Large area active high-density plasma sources
 - Passive electron emission through driven electrodes
- To mitigate the effort in producing electron current, it is worthwhile to explore simple passive (i.e., no plasma arc) current sources
 - Form initial tokamak-like state with minimal active arc gun
 - Increase I_p with passive electrodes.
 - Critical feature is how to diffuse the current extracted from metallic electrode
- First tests were promising
 - Arc current cut off after relaxation and formation of tokamak-like state
 - Gas fueling through chamber continued
 - I_p rise is virtually the *same*, whether arc discharge or passive electrode provide the charge carriers
 - Suggests continuing development of electrode emitters

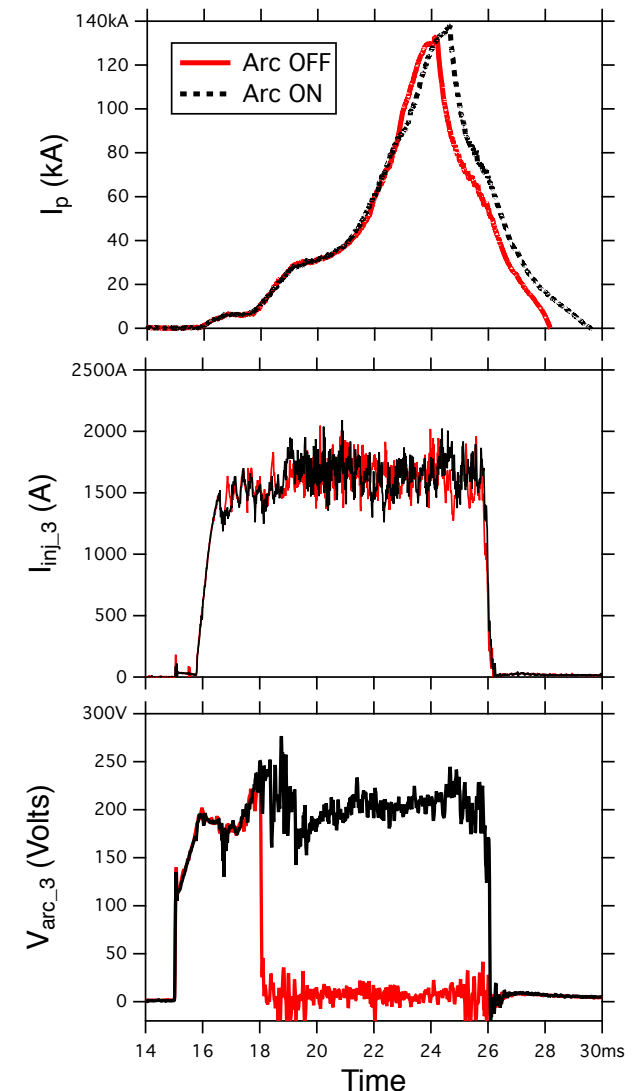
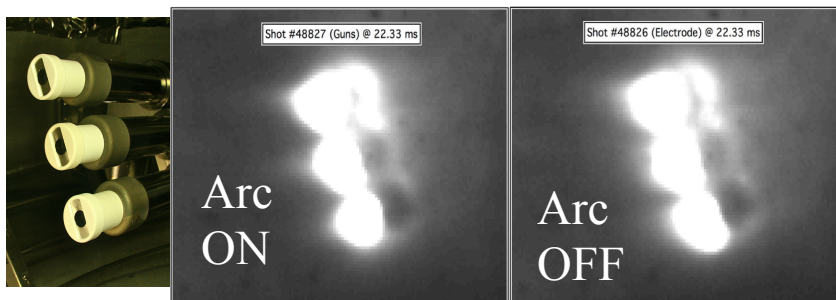




Identical Discharge Evolution Seen with Plasma Arc Turned Off

- Arc crowbarred out after tokamak discharge established to transition from arc plasma source to driven electrode system
 - Keep electrode widths narrow to maintain Taylor limit
 - Some limitations from PMI interactions at Mo/BN interface
- Demonstrated transition from active gun drive to passive electrode drive
 - Same extracted current whether arc is on or off, with same gas flow
 - Driven I_p virtually identical
 - Camera (low-res) images suggest similar current source regions

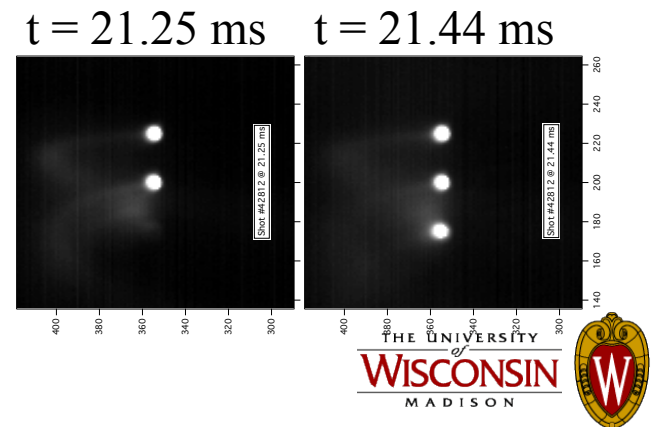
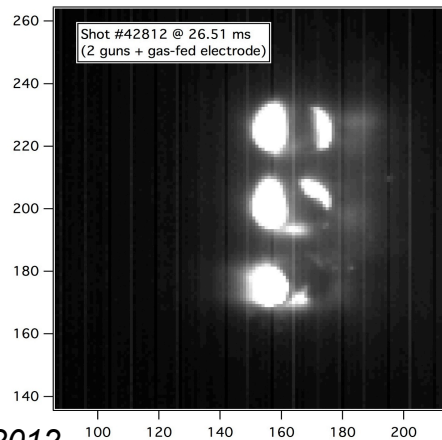
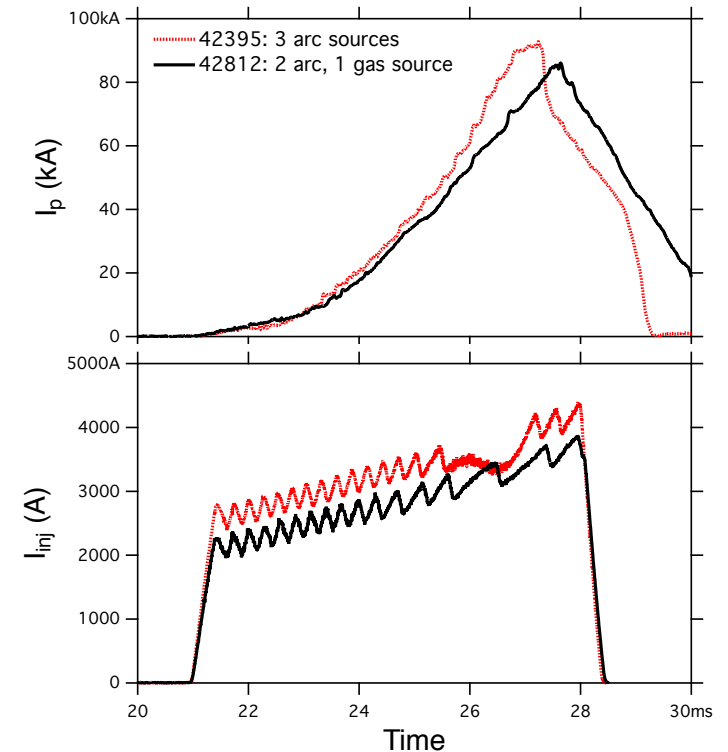
“Slot” Mo faces
with BN caps





Electrode with Integrated Gas Feed Behaves Similarly to Arc Source

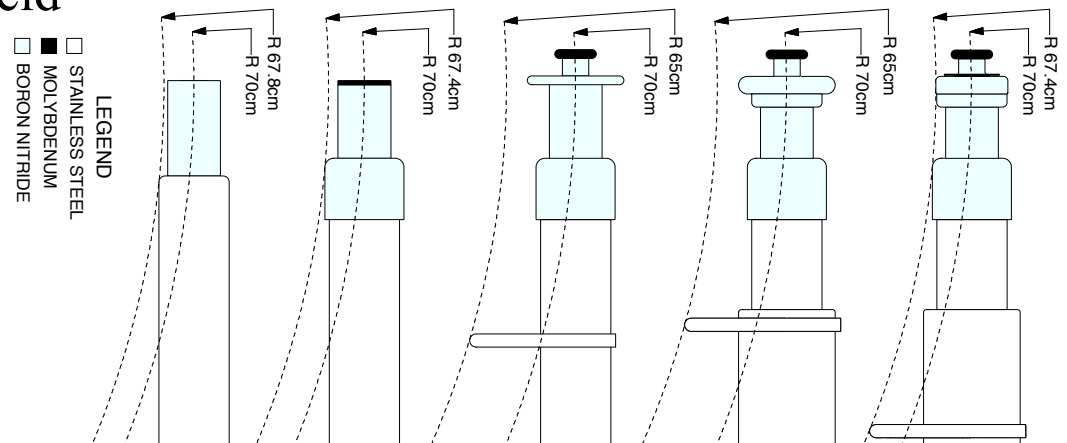
- Simple gas-fed electrode replaced a single arc source to test electrode concept
 - Passive electrode turns on spontaneously after 2 arc sources establish discharge
 - Discharge evolution to similar 3-arc source plasma
 - Suggests effective area of \sim size of gas source region
- Current \sim equally shared amongst 3 injectors





Large Area Electrode Development/Tests

- The integrated arc plasma and electrode system has evolved to provide mitigation against plasma interactions
 - BN “Bell” to avoid arc back to ground
 - Electrode raised above BN shield
 - PV-10 gas valve installed
 - Floating Mo shield plate behind electrode





Electrode Systems Evolved to Mitigate Deleterious Plasma-Material Interactions

- N dominant impurity with unprotected gun assembly

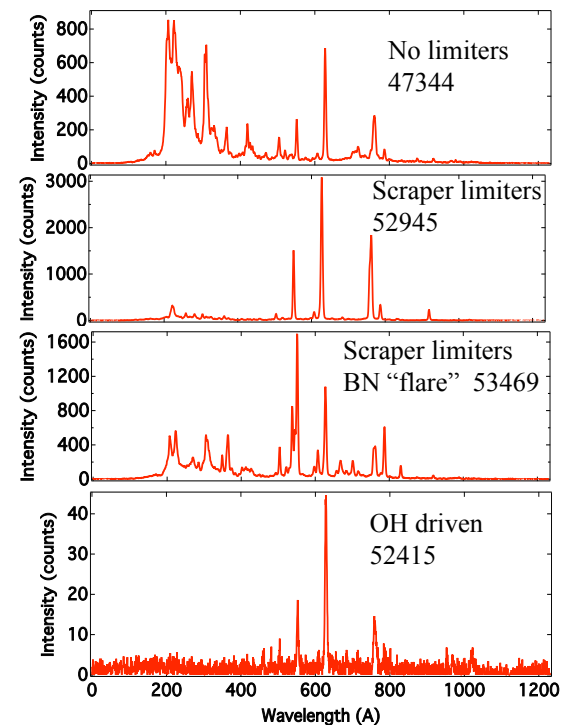
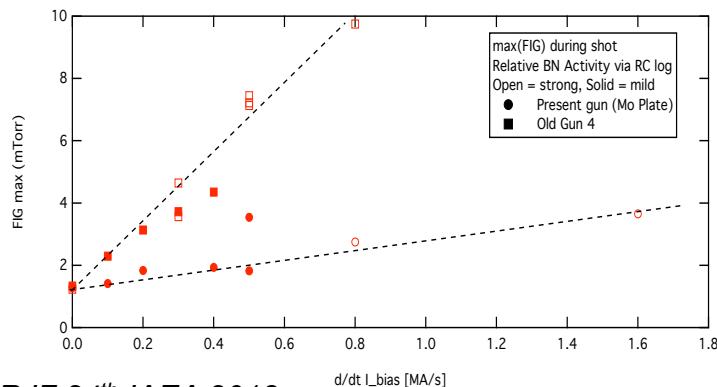
- $Z_{\text{eff}} \sim 2.2. \pm 0.8$ during; ≤ 1.4 after injection

- Local scraper limiters reduce N from unprotected gun case

- Also controls local edge N_e and injector impedance
 - O dominant impurity in OH and “well-behaved” helicity-driven plasmas

- Mo backing plate reduces BN interactions and undesired gas emission

- Arc-backs to limiter still occur at times





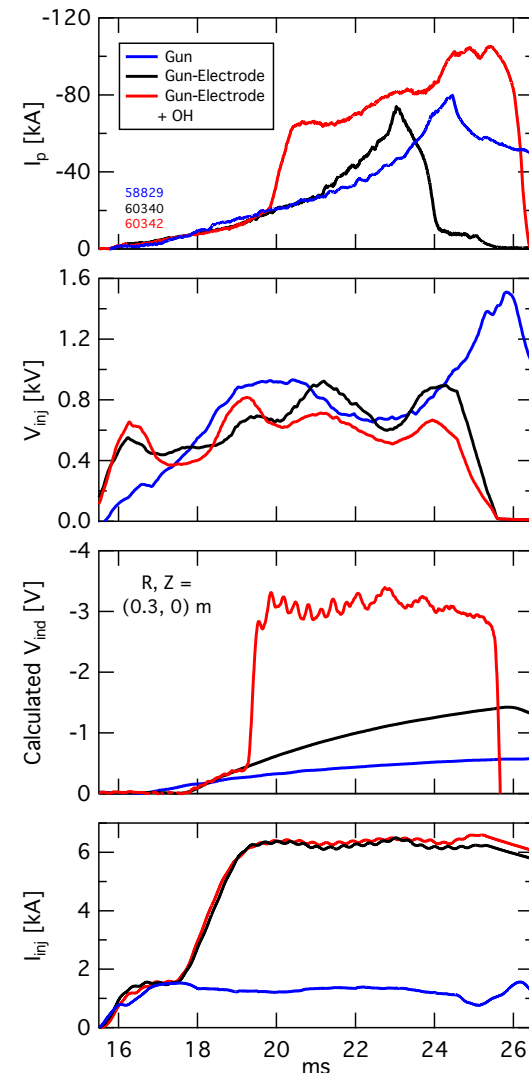
Extended Passive Electrode Tested as 1st Possibility for Large-Area Source

- Tests of current distribution on metallic electron emitter
 - No gas feed through electrode
 - Current in presence of plasma emitted from localized cathode spots
 - Similar to emission in vacuum
 - ~200-400 A per spot
 - For $I_{inj} \sim 6\text{kA}$, max effective area $< 0.3\text{ cm}^2 \sim A_{arc}/4$
 - Both single arc source and large passive electrode give similar I_p , well below the relaxation limit
 - Limit demonstrated with additional OH V_{loop}
- Need integrated gas fueling to spread I_{arc} across large area
 - Tests with single arc source cap underway to confirm and optimize

Single arc source with integrated large-area passive electrode



Small cathode spots emit current from simple metallic electrode





Summary: Significant Progress in Developing Edge Current Sources for Local Helicity Injection

- Miniature plasma arc sources provide local current sources for non-solenoidal startup of ST and other confinement devices
 - Very flexible geometry options
 - Can be combined with poloidal field induction when located in HFS region
 - Technical attractiveness: can remove sources and anode after startup
 - 1-2 kA/cm² available; low impurity content
- Arc source impedance, and helicity injection rate, appears to be governed by sheath effects and magnetic current limits
 - Further tests needed to understand apparent density scaling
- High helicity input requires large area current source and narrow current channel in edge region
 - Preliminary tests suggest gas-fed electrodes may be combined with arc sources to drive high I_p
 - Electrode design requires PMI mitigation techniques